Corso di Laurea in Fisica – UNITS ISTITUZIONI DI FISICA PER IL SISTEMA TERRA

Born of the Wave Equation

FABIO ROMANELLI

Department of Mathematics & Geosciences

University of Trieste

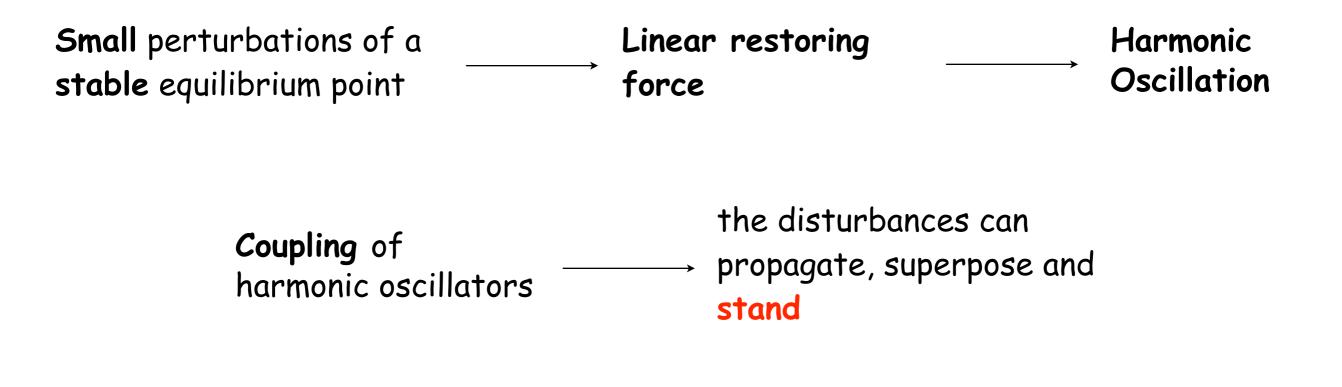
romanel@units.it

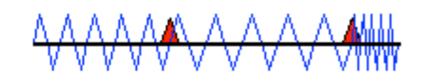
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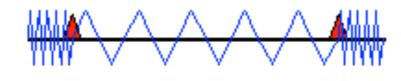










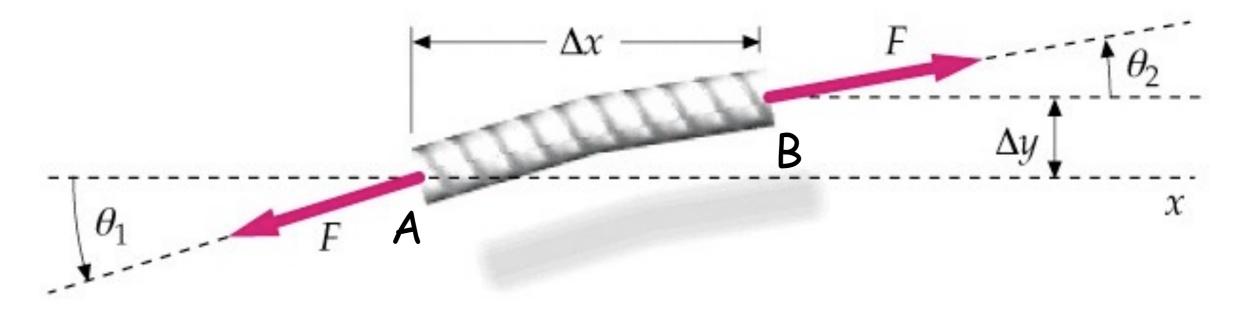




Normal modes of the system







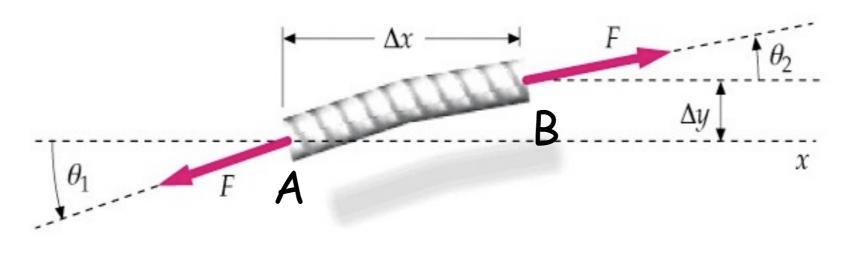
Consider a small segment of string of length Δx and tension F

The ends of the string make small angles θ_1 and θ_2 with the x-axis.

The vertical displacement Δy is very small compared to the length of the string







Resolving forces vertically

From small angle approximation $sin\theta \sim tan\theta$

$$\Sigma F_{y} = F \sin \theta_{2} - F \sin \theta_{1}$$
$$= F (\sin \theta_{2} - \sin \theta_{1})$$
$$\Sigma F_{y} \approx F(\tan \theta_{2} - \tan \theta_{1})$$

The tangent of angle A (B) = pendence of the curve in A (B) given by $\frac{\partial Y}{\partial x}$





$$\therefore \Sigma F_{y} \approx F\left(\left(\frac{\partial y}{\partial x}\right)_{B} - \left(\frac{\partial y}{\partial x}\right)_{A}\right)$$

We now apply N2 to segment

$$\Sigma \mathbf{F}_{\mathbf{y}} = \mathbf{ma} = \mu \Delta \mathbf{x} \left(\frac{\partial^2 \mathbf{y}}{\partial \mathbf{t}^2} \right)$$

$$\mu \Delta \mathbf{x} \left(\frac{\partial^2 \mathbf{y}}{\partial t^2} \right) = \mathbf{F} \left(\left(\frac{\partial \mathbf{y}}{\partial \mathbf{x}} \right)_{\mathbf{B}} - \left(\frac{\partial \mathbf{y}}{\partial \mathbf{x}} \right)_{\mathbf{A}} \right)$$

$$\frac{\mu}{F} \left(\frac{\partial^2 \gamma}{\partial t^2} \right) = \frac{\left[\left(\frac{\partial \gamma}{\partial x} \right)_B - \left(\frac{\partial \gamma}{\partial x} \right)_A \right]}{\Delta x}$$



as $\Delta x \rightarrow 0$



$$\frac{\mu}{F} \left(\frac{\partial^2 y}{\partial t^2} \right) = \frac{\left[\left(\frac{\partial y}{\partial x} \right)_B - \left(\frac{\partial y}{\partial x} \right)_A \right]}{\Delta x}$$

The derivative of a function is defined as

$$\begin{pmatrix} \frac{\partial f}{\partial x} \end{pmatrix} = \lim_{\Delta x \to 0} \frac{\left[f(x + \Delta x) - f(x) \right]}{\Delta x}$$

If we associate $f(x+\Delta x)$ with $(\partial y/\partial x)_B$ and f(x) with $(\partial y/\partial x)_A$

$$\frac{\mu}{F} \left(\frac{\partial^2 \gamma}{\partial t^2} \right) = \frac{\partial^2 \gamma}{\partial x^2}$$

This is the linear wave equation for waves on a string



Consider a solution of the form $y(x,t) = A \sin(kx-\omega t)$

$$\frac{\partial^2 \gamma}{\partial t^2} = -\omega^2 A \sin(kx - \omega t) \qquad \frac{\partial^2 \gamma}{\partial x^2} = -k^2 A \sin(kx - \omega t)$$

If we substitute these into the linear wave equation

$$\frac{\mu}{F}(-\omega^2 A \sin(kx - \omega t)) = -k^2 A \sin(kx - \omega t)$$
$$\frac{\mu}{F}\omega^2 = k^2$$

and, using $\omega^2/k^2 = F/\mu = v^2$, i.e. $v = \omega/k$

$$\frac{\partial^2 \gamma}{\partial x^2} = \frac{1}{v^2} \frac{\partial^2 \gamma}{\partial t^2}$$

 $v=\sqrt{F/\mu}$

General form of LWE





A general property of waves is that the speed of a wave depends on the properties of the medium, but is independent of the motion of the source of the waves.

Consider a wave moving along a rope experimentally we find

(i) the greater the tension in a rope the faster the waves propagate

(ii) waves propagate faster in a light rope than a heavy rope

ie v \propto tension (F) and v \propto 1/mass

known as Mersenne's law



Mersenne's law





L'Harmonie Universelle (1637)

This book contains (Marine) Mersenne's laws which describe the frequency of oscillation of a stretched string.

This frequency is:

a) Inverse proportional to the length of the string (this was actually known to the ancients, and is usually credited to Pythagoras himself).
b) Proportional to the square root of the stretching force, and

c) Inverse proportional to the square root of the mass per unit length.

HARMONIE VNIVERSELLE, CONTENANT LA THEORIE ET LA PRATIQUE DE LA MUSIQUE,

Oùil eft traité de la Nature des Sons, & des Mouuemens, des Confonances, des Diffonances, des Genres, des Modes, de la Composition, de la Voix, des Chants, & de toutes fortes d'Instrumens Harmoniques.

Par F. MARIN MERSENNE de l'Ordre des Minimes.



A PARIS, Chez SEBASTIEN CRAMOISY, Imprimeur ordinaire du Roy, ruë S. Iacques, aux Cicognes.

M. DC. XXXVI. Auce Prinilege du Roy, & Approbation des Docteurs.





Earlier we introduced the concept of a wavefunction to represent waves travelling on a string.

All wavefunctions y(x,t) represent solutions of the LINEAR WAVE EQUATION

The wave equation provides a complete description of the wave motion and from it we can derive the wave velocity

The most general solution is, for 1D homogeneous medium,

$$y(x,t)=g(x+vt)+f(x-vt)$$



D'Alembert's solution





D'Alembert (1747) "Recherches sur la courbe que forme une corde tendue mise en vibration" (Researches on the curve that a tense cord forms [when] set into vibration), Histoire de l'académie royale des sciences et belles lettres de Berlin, vol. 3, pages 214-219.

D'Alembert (1750) "Addition au mémoire sur la courbe que forme une corde tenduë mise en vibration," Histoire de l'académie royale des sciences et belles lettres de Berlin, vol. 6, pages 355-360.

y(x,t)
$$\rightarrow$$
 y(ξ,η) with ξ=x-vt, η=x+vt

$$y_{x} = \frac{\partial y}{\partial x} = y_{\xi}\xi_{x} + y_{\eta}\eta_{x} = y_{\xi} + y_{\eta}; \quad y_{xx} = \frac{\partial}{\partial x}(y_{x}) = y_{\xi\xi} + 2y_{\xi\eta} + y_{\eta\eta}, \quad y_{tt} = v^{2}(y_{\xi\xi} - 2y_{\xi\eta} + y_{\eta\eta})$$
$$\Rightarrow y_{\xi\eta} = \frac{\partial^{2}y}{\partial\xi\partial\eta} = \frac{\partial}{\partial\xi}\left(\frac{\partial y}{\partial\eta}\right) = 0$$
$$y = h(\xi) + g(\eta) \implies y(x, t) = h(x - vt) + g(x + vt)$$

and if the initial conditions are y(x,0)=f(x) and initial velocity=0

$$\mathbf{y}(\mathbf{x},\mathbf{t}) = \frac{1}{2} \Big[\mathbf{f}(\mathbf{x} - \mathbf{v}\mathbf{t}) + \mathbf{f}(\mathbf{x} + \mathbf{v}\mathbf{t}) \Big]$$

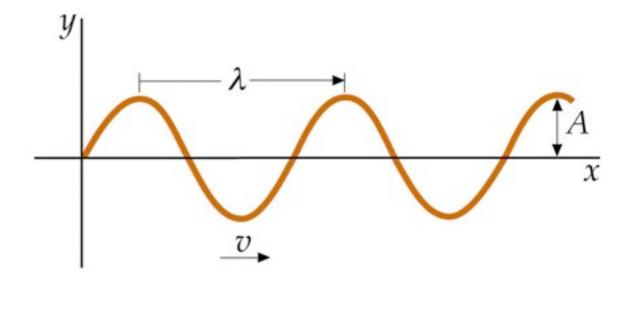
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Wave equation





A harmonic wave is sinusoidal in shape, and has a displacement y at time t=0 $y = A sin\left(\frac{2\pi}{\lambda}x\right)$



A is the **amplitude** of the wave and λ is the **wavelength** (the distance between two crests);

if the wave is moving to the right with speed v, the wavefunction at some t is given by:

$$y = Asin\left[\frac{2\pi}{\lambda}(x - vt)\right]$$





Time taken to travel one wavelength is the period T

Velocity, wavelength and period are related by

$$\mathbf{v} = \frac{\lambda}{T} \quad \text{or} \quad \lambda = \mathbf{v}T$$

$$\mathbf{v} = \mathbf{A}\sin\left[2\pi\left(\frac{\mathbf{x}}{\lambda} - \frac{\mathbf{t}}{T}\right)\right]$$

The wavefunction shows the periodic nature of y:

at any time t y has the same value at x, x+ λ , x+2 λ

and at any x y has the same value at times t, t+T, t+2T.....





It is convenient to express the harmonic wavefunction by defining the wavenumber \mathbf{k} , and the angular frequency $\boldsymbol{\omega}$

where
$$k = \frac{2\pi}{\lambda}$$
 and $\omega = \frac{2\pi}{T}$
 $\therefore y = A \sin(kx - \omega t)$

This assumes that the displacement is zero at x=0 and t=0. If this is not the case we can use a more general form

$$y = A\sin(kx - \omega t - \phi)$$

where $\boldsymbol{\varphi}$ is the phase constant and is determined from initial conditions

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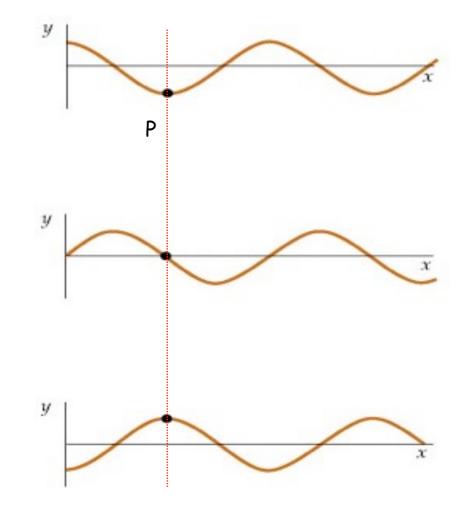


The wavefunction can be used to describe the motion of any point P.

If
$$y = A sin(kx - \omega t)$$

Transverse velocity
$$v_y$$

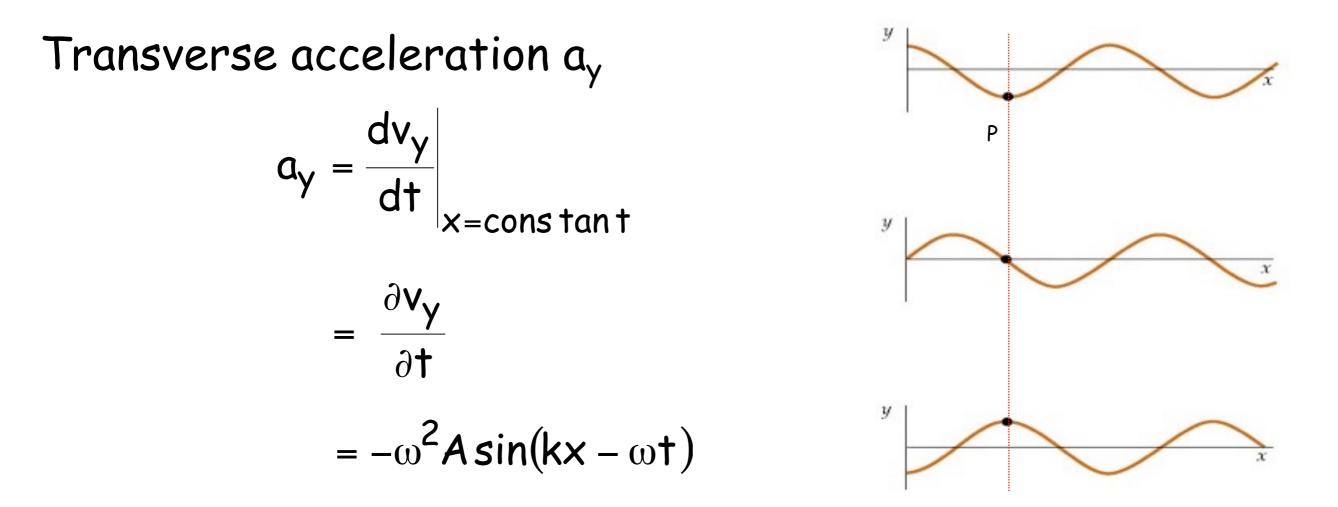
 $v_y = \frac{dy}{dt}\Big|_{x=cons tan t}$
 $= \frac{\partial y}{\partial t}$
 $= -\omega A cos(kx - \omega t)$



which has a maximum value, $(v_y)_{max} = \omega A$, when y = 0







which has a maximum absolute value, $(a_y)_{max} = \omega^2 A$, when t=0

NB: x-coordinates of P are constant

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A harmonic wave on a rope is given by the expression y(x,t) = 10 sin(2x - 5t)

where the amplitude is in mm, $\,k$ in rad m^{-1}, and ω in rad s^{-1}

(a) Determine the velocity and acceleration for each element of the rope.

(b) What are the maximum values of the acceleration and velocity ?

(c) Is the displacement +ve or -ve at x=1m and t=0.2s?





(a) Determine the velocity and acceleration for each element of the rope.

Generally $y(x,t) = A \sin(kx - \omega t)$ \therefore $v_y = -\omega A \cos(kx - \omega t)$ y(x,t) = 10 sin(2x - 5t) $\therefore v_v = -5 \times 10 \cos(2x - 5t)$ $v_v = -50\cos(2x-5t)$ Generally $y(x,t) = A \sin(kx - \omega t)$ \therefore $a_v = -\omega^2 A \sin(kx - \omega t)$:. $a_v = -5^2 \times 10 \sin(2x - 5t)$ $a_v = -250 \sin(2x - 5t)$





(b) What are the maximum values of the acceleration and velocity ?

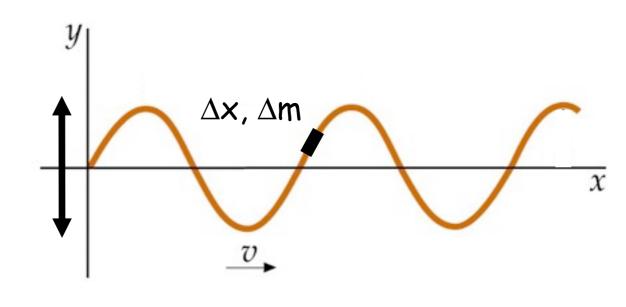
$$\begin{pmatrix} a_{y} \end{pmatrix}_{max} = \omega^{2}A \qquad \qquad \begin{pmatrix} v_{y} \end{pmatrix}_{max} = \omega A \\ \begin{pmatrix} a_{y} \end{pmatrix}_{max} = 5^{2} \times 10 \qquad \qquad \begin{pmatrix} v_{y} \end{pmatrix}_{max} = 5 \times 10 \\ \begin{pmatrix} a_{y} \end{pmatrix}_{max} = 250 \text{ mms}^{-2} \qquad \qquad \begin{pmatrix} v_{y} \end{pmatrix}_{max} = 50 \text{ mms}^{-1}$$





Consider a harmonic wave travelling on a string.

Source of energy is an external agent on the left of the wave which does work in producing oscillations.



Consider a small segment, length Δx and mass Δm .

The segment moves vertically with SHM, frequency $\boldsymbol{\omega}$ and amplitude A.

Generally
$$E = \frac{1}{2}m\omega^2 A^2$$





$$\mathsf{E} = \frac{1}{2}\mathsf{m}\omega^2 \mathsf{A}^2$$

If we apply this to our small segment, the total energy of the element is 1 < 1 < 2 < 2

$$\Delta \mathsf{E} = \frac{1}{2} (\Delta \mathsf{m}) \omega^2 \mathsf{A}^2$$

If μ is the mass per unit length, then the element Δx has mass $\Delta m = \mu \Delta x$ $\Delta E = \frac{1}{2}(\mu \Delta x)\omega^2 A^2$

If the wave is travelling from left to right, the energy ΔE arises from the work done on element Δm_i by the element Δm_{i-1} (to the left).





Similarly Δm_i does work on element Δm_{i+1} (to the right) ie. energy is transmitted to the right.

The rate at which energy is transmitted along the string is the power and is given by dE/dt.

If $\Delta x \rightarrow 0$ then Power = $\frac{dE}{dt} = \frac{1}{2}(\mu \frac{dx}{dt})\omega^2 A^2$ but dx/dt = speed \therefore Power = $\frac{1}{2}\mu \omega^2 A^2 v$





Power =
$$\frac{1}{2}\mu \omega^2 A^2 v$$

Power transmitted on a harmonic wave is proportional to

(a) the wave speed v
(b) the square of the angular frequency ω
(c) the square of the amplitude A

All harmonic waves have the following general properties:

The power transmitted by any harmonic wave is proportional to the square of the frequency and to the square of the amplitude.





 Small perturbations of a stable equilibrium point
 Linear restoring force
 Harmonic Oscillation

 Coupling of harmonic oscillators
 The disturbances can propagate, superpose and stand
 Harmonic Oscillators

WAVE: organized propagating imbalance, satisfying differential equations of motion

$$\frac{\partial^2 \gamma}{\partial x^2} = \frac{1}{v^2} \frac{\partial^2 \gamma}{\partial t^2}$$

General form of LWE