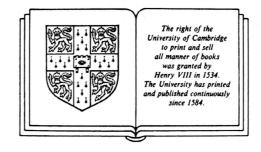
#### LEZIONE 5-6

# VORTEX ELEMENT METHODS FOR FLUID DYNAMIC ANALYSIS OF ENGINEERING SYSTEMS

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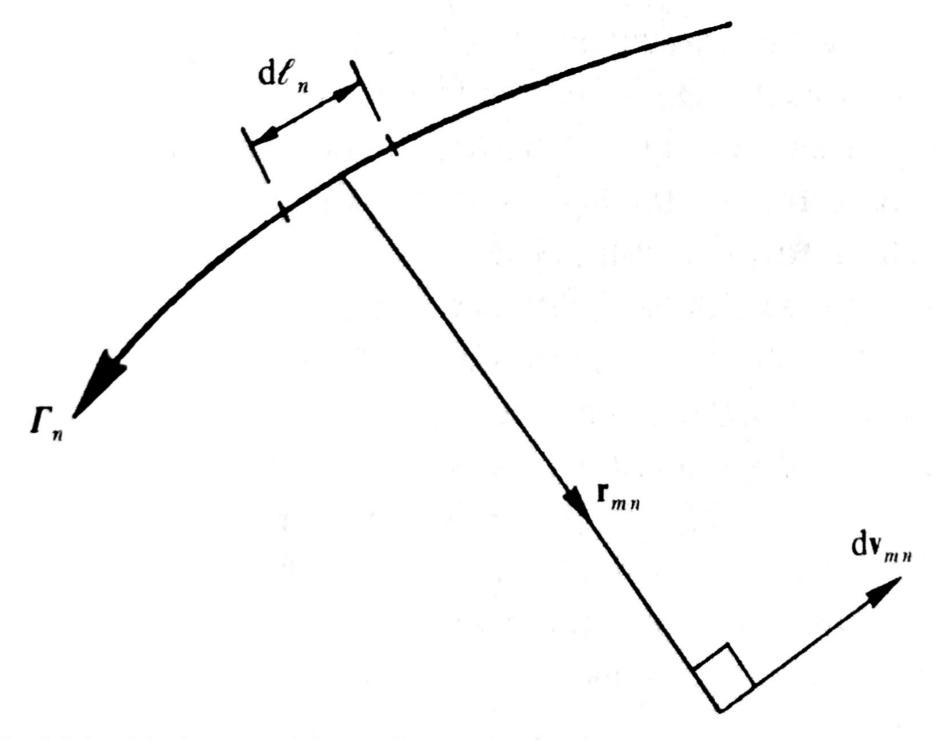


Fig. 1.2. Velocity induced by a line vortex element.

velocity induced at m by a small line vortex element at n of strength  $\Gamma_n$  per unit length\* and of length  $dl_n$  is given by the Biot-Savart law, namely, with reference to Fig. 1.2,

$$d\mathbf{v}_{mn} = \frac{\Gamma_n \, dl_n X \mathbf{r}_{mn}}{4\pi r_{mn}^3} \tag{1.8}$$

By taking the cross product of  $dv_{mn}$  with the unit vector  $i_m$  normal to the surface at m twice, we obtain the velocity parallel to the surface at m induced by the line vortex element. Thus

$$d\mathbf{v}_{smn} = \mathbf{i}_{m} X (d\mathbf{v}_{mn} X \mathbf{i}_{m})$$

$$= \frac{\mathbf{i}_{m} X ((\mathbf{\Gamma}_{n} X \mathbf{r}_{mn}) X \mathbf{i}_{m}) dl_{n}}{4\pi r_{mn}^{3}}$$
(1.9)

#### Physical significance of the surface vorticity model

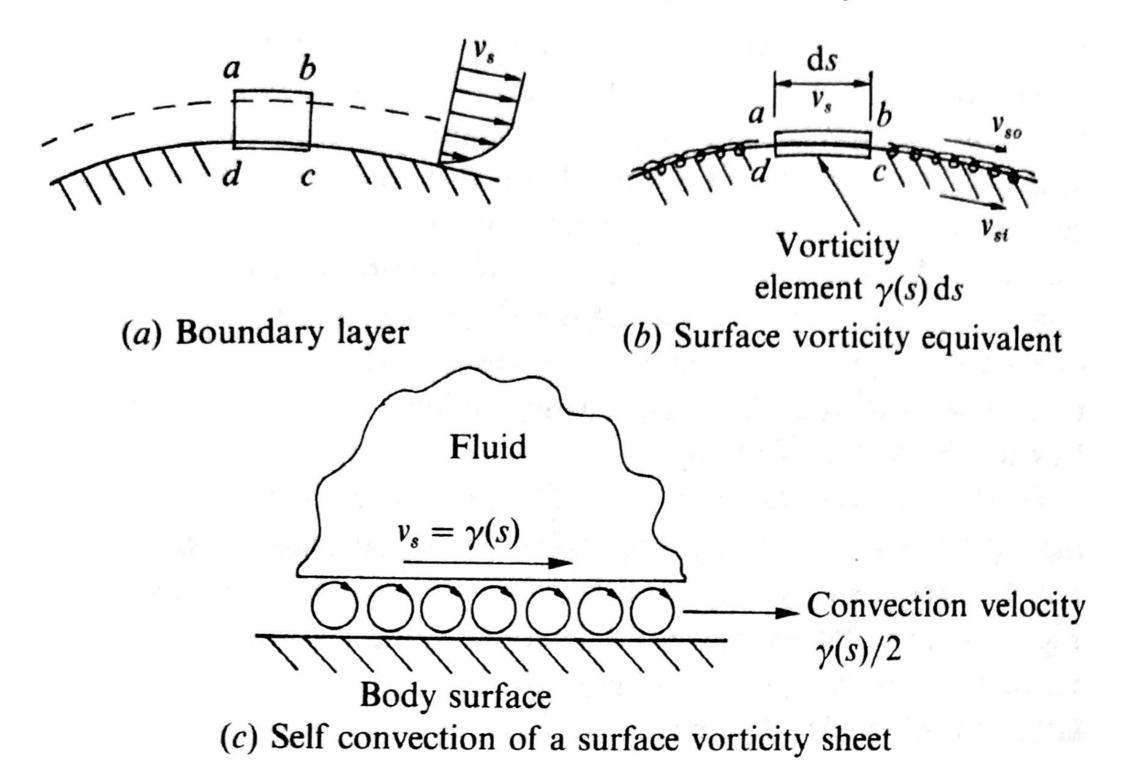


Fig. 1.3. Boundary layer and surface vorticity equivalent in potential flow.

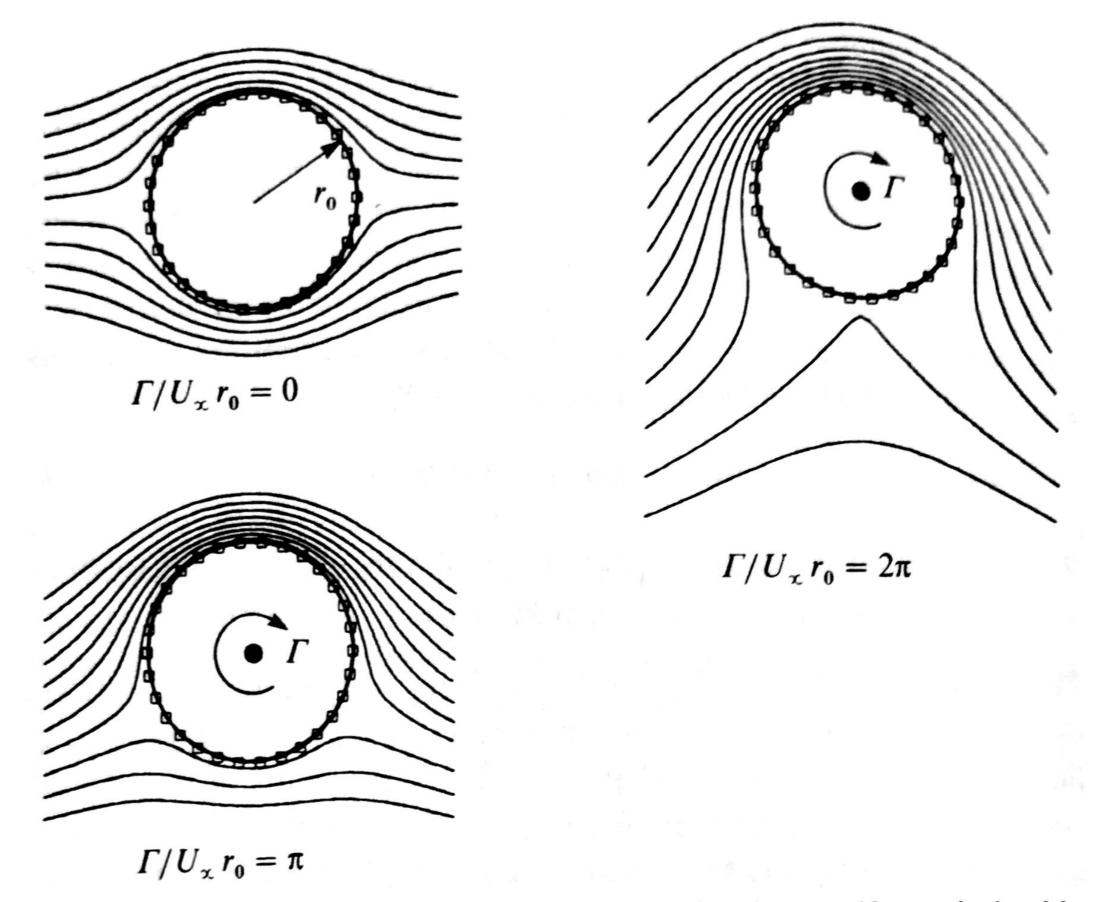


Fig. 2.1. Flow induced by cylinder with circulation in a uniform derived by the surface vorticity method.

through  $\alpha_{\infty}$  that this expression transforms to

$$v_s = 2W_\infty \sin(\theta - \alpha_\infty) + \frac{\Gamma}{2\pi r_0}$$
 (2.2a)

As shown by Glauert by integrating the surface pressure on the cylinder, a lift force L is generated in the direction normal to  $W_{\infty}$  given by the Magnus law.

$$L = \rho W_{\infty} \Gamma \tag{2.3}$$

Introducing the usual definition of lift coefficient

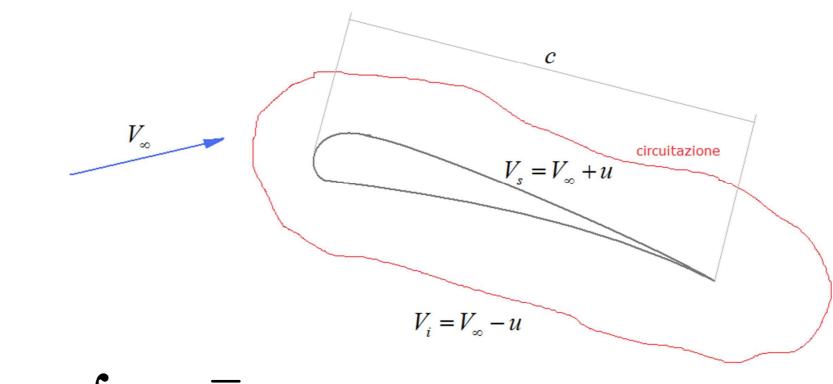
$$C_L = \frac{L}{\frac{1}{2}\rho W_{\infty}^2 \ell} \tag{2.4}$$

where  $\ell$  is a typical dimension of the body, in this case its diameter  $2r_0$ , we have

$$C_L = \Gamma/W_{\infty} r_0 \tag{2.5}$$

The dimensionless parameter  $\Gamma/U_{\infty}r_0$  previously referred to is in fact the lift coefficient with uniform stream  $U_{\infty}$ . To estimate the

# teorema di Kutta-Jukowsky

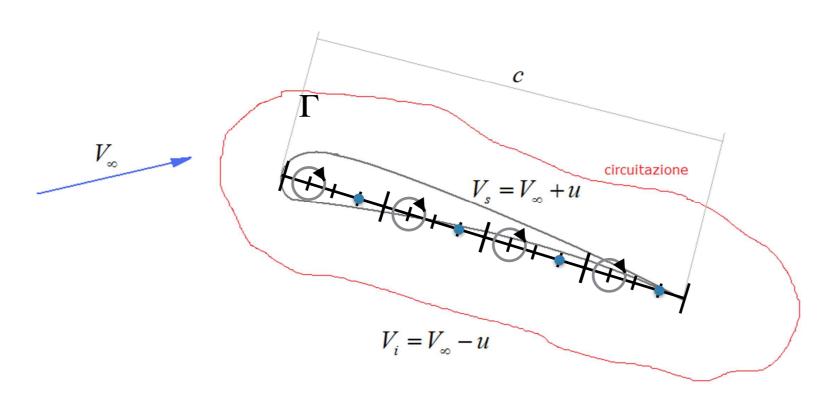


$$\Gamma = \oint_{\ell} \overline{v} \cdot d\overline{\ell} \qquad L = \rho \Gamma V_{\infty}$$

$$\Delta p = p_i - p_s = \frac{1}{2} \rho \left( V_s^2 - V_i^2 \right) = 2\rho u V_{\infty}$$

$$\Gamma = 2cu$$

# Metodo delle singolarità vorticose

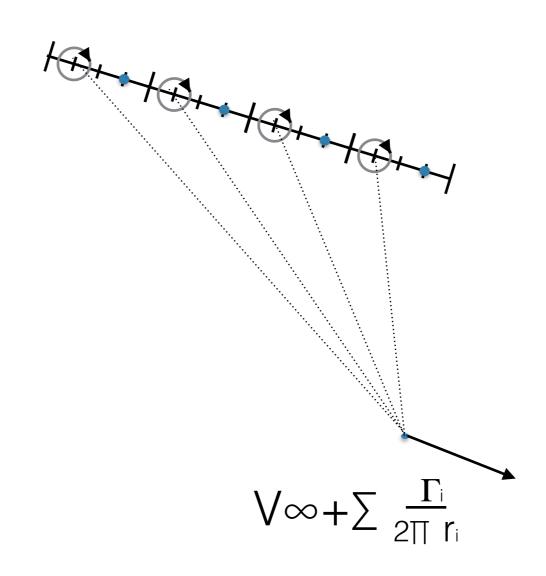


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# Metodo delle singolarità vorticose



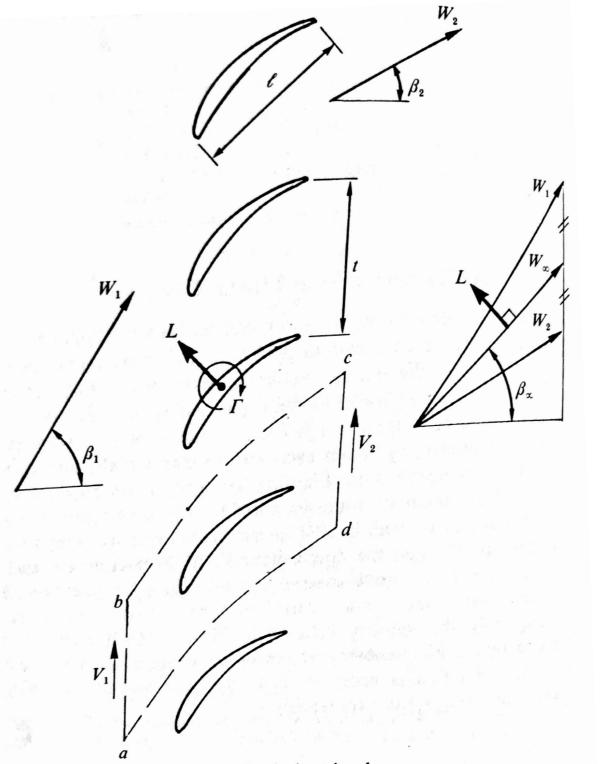


Fig. 2.12. Cascade geometry and velocity triangles.

#### Turbomachine linear cascades

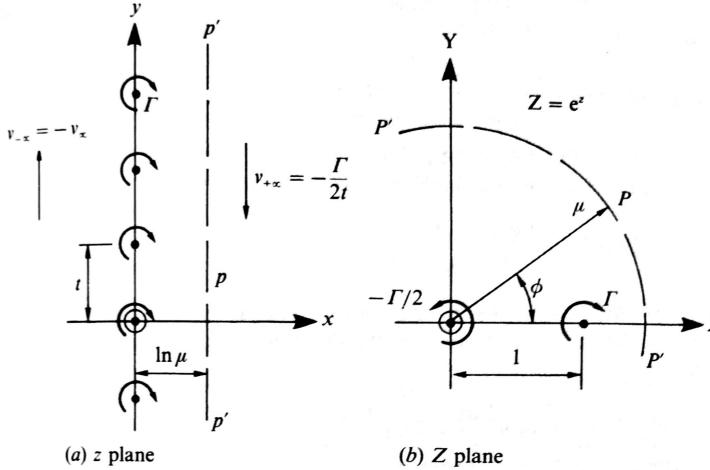


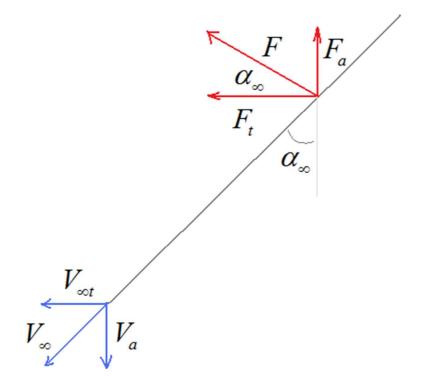
Fig. 2.13. Transformation of vortex array in z plane to vortex pair in Z plane.

# Schiere di pale

$$\Delta p_0 = 0$$
 ipotesi perdite nulle

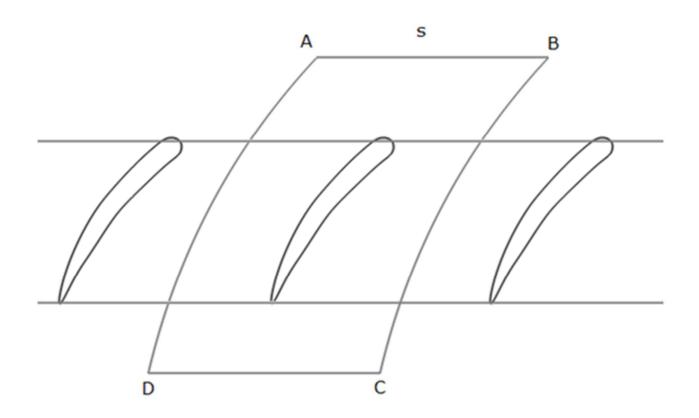
$$F_a = s\rho V_{\infty t} \left( V_{2t} - V_{1t} \right)$$
$$F_t = s\rho V_a \left( V_{1t} - V_{2t} \right)$$

$$\frac{V_{\infty t}}{V_a} = -\frac{F_a}{F_t} = \tan \alpha_{\infty}$$



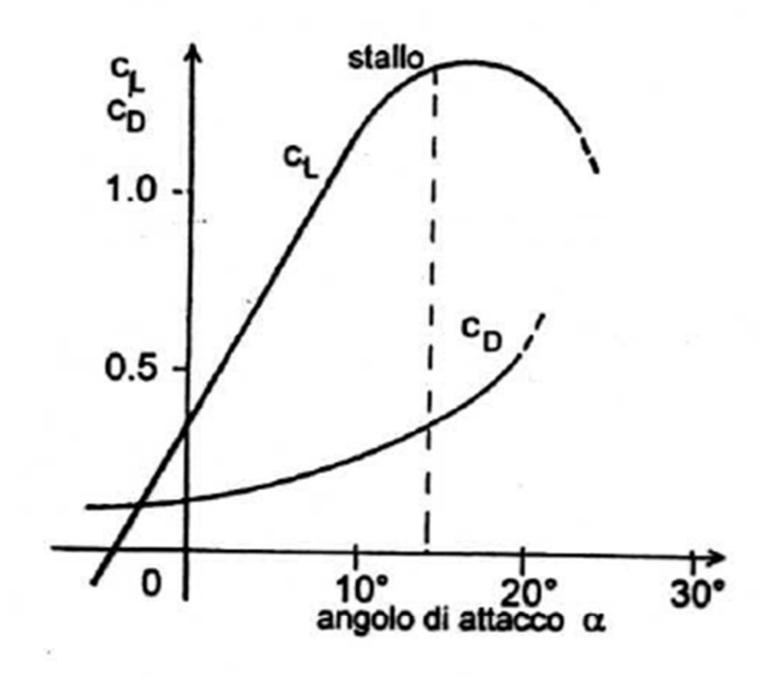
$$F = \frac{F_t}{\cos \alpha_{\infty}} = \rho \frac{V_a}{\cos \alpha_{\infty}} s(V_{1t} - V_{2t}) = \rho V_{\infty} s(V_{1t} - V_{2t}) = L$$

# Schiere di pale

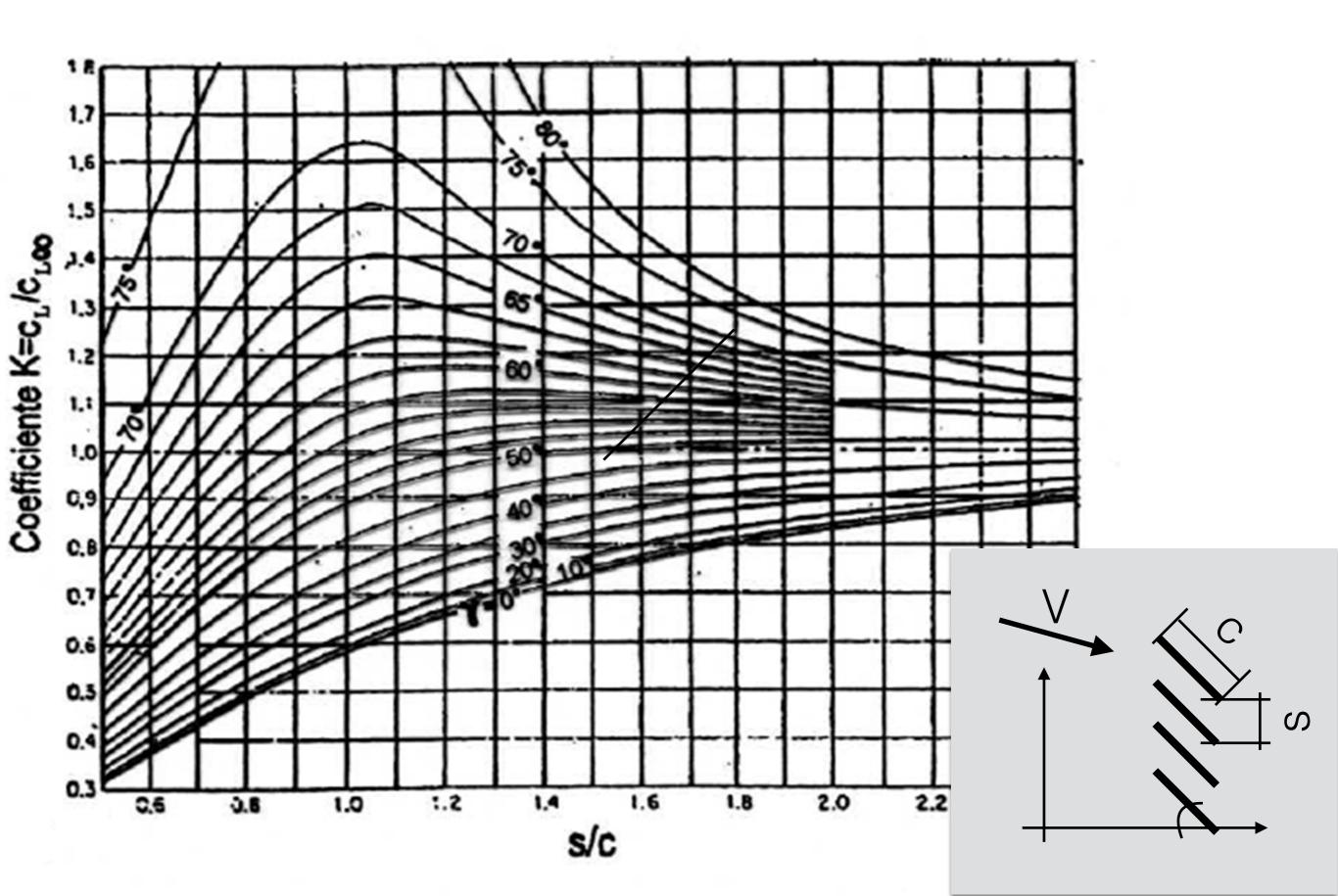


$$\Gamma = s(V_{1t} - V_{2t}) \longrightarrow L = \rho V_{\infty} \Gamma$$

#### Effetto schiera sulle prestazioni del profilo



# Effetto schiera sulle prestazioni del profilo



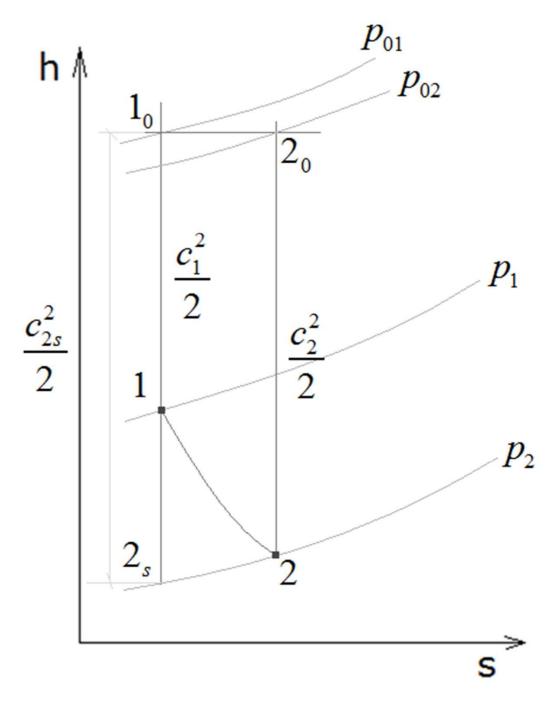
# ugelli e diffusori

#### Distinguiamo due casi:

- 1) Nell'elemento abbiamo un incremento di velocità a spese di una riduzione di pressione. Questi saranno gli *ugelli*.
- 2) Nell'elemento l'energia cinetica diminuisce ed aumenta la pressione. Questi saranno i *diffusori*.

# ugelli

$$\eta_{is} = \frac{h_1 - h_2}{h_1 - h_{2s}} = \frac{\frac{c_2^2}{2} - \frac{c_1^2}{2}}{\frac{c_{2s}^2}{2} - \frac{c_1^2}{2}} = \frac{c_2^2 - c_1^2}{c_{2s}^2 - c_1^2}$$



# ugelli

(Ma < 0,3)

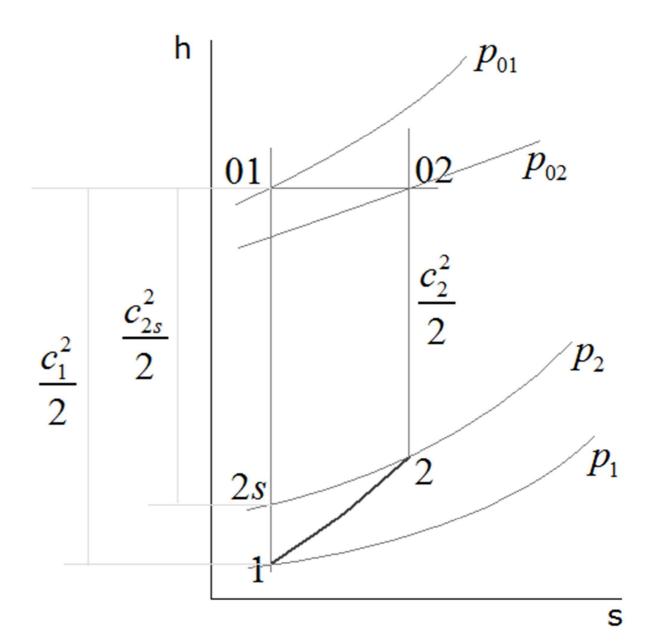
$$p_{01} = p_1 + \frac{1}{2}\rho c_1^2 \quad \Rightarrow \quad c_1^2 = \frac{2}{\rho}(p_{01} - p_1)$$

$$p_{01} = p_2 + \frac{1}{2}\rho c_{2s}^2 \quad \Rightarrow \quad c_{2s}^2 = \frac{2}{\rho}(p_{01} - p_2)$$

$$p_{02} = p_2 + \frac{1}{2}\rho c_2^2 \quad \Rightarrow \quad c_2^2 = \frac{2}{\rho}(p_{02} - p_2)$$

$$\eta_{is} = \frac{p_{02} - p_2 - (p_{01} - p_1)}{p_{01} - p_2 - (p_{01} - p_1)} = \frac{p_{02} - p_2 - (p_{01} - p_1)}{p_1 - p_2} = 1 - \frac{\Delta p_0}{p_1 - p_2}$$

$$\eta_{is} = \frac{h_{2s} - h_1}{h_2 - h_1} = \frac{c_1^2 - c_{2s}^2}{c_1^2 - c_2^2}$$



(Ma < 0,3)

$$\eta_{is} = \frac{p_{01} - p_1 - (p_{01} - p_2)}{p_{01} - p_1 - (p_{02} - p_2)} = \frac{p_2 - p_1}{p_2 - p_1 + (p_{01} - p_{02})} = \frac{1}{1 - \frac{\Delta p_0}{p_2 - p_1}}$$

coeff. recupero di pressione:  $c_p = \frac{p_2 - p_1}{p_{01} - p_1}$ 

(Ma < 0,3)

Legame tra  $c_p$  e  $\eta_{is}$ 

$$\eta_{is} = \frac{p_2 - p_1}{p_2 - p_1 + (p_{01} - p_{02})}$$

$$\frac{1}{\eta_{is}} = \frac{p_2 - p_1 + (p_{01} - p_{02})}{p_2 - p_1} = \frac{p_{01} - p_1 - (p_{02} - p_2)}{p_2 - p_1} = \frac{1}{c_p} - \frac{p_{02} - p_2}{p_2 - p_1}$$

$$c_{pi} = \frac{p_2 - p_1 + (p_{01} - p_{02})}{p_{01} - p_1}$$

(Ma < 0,3)

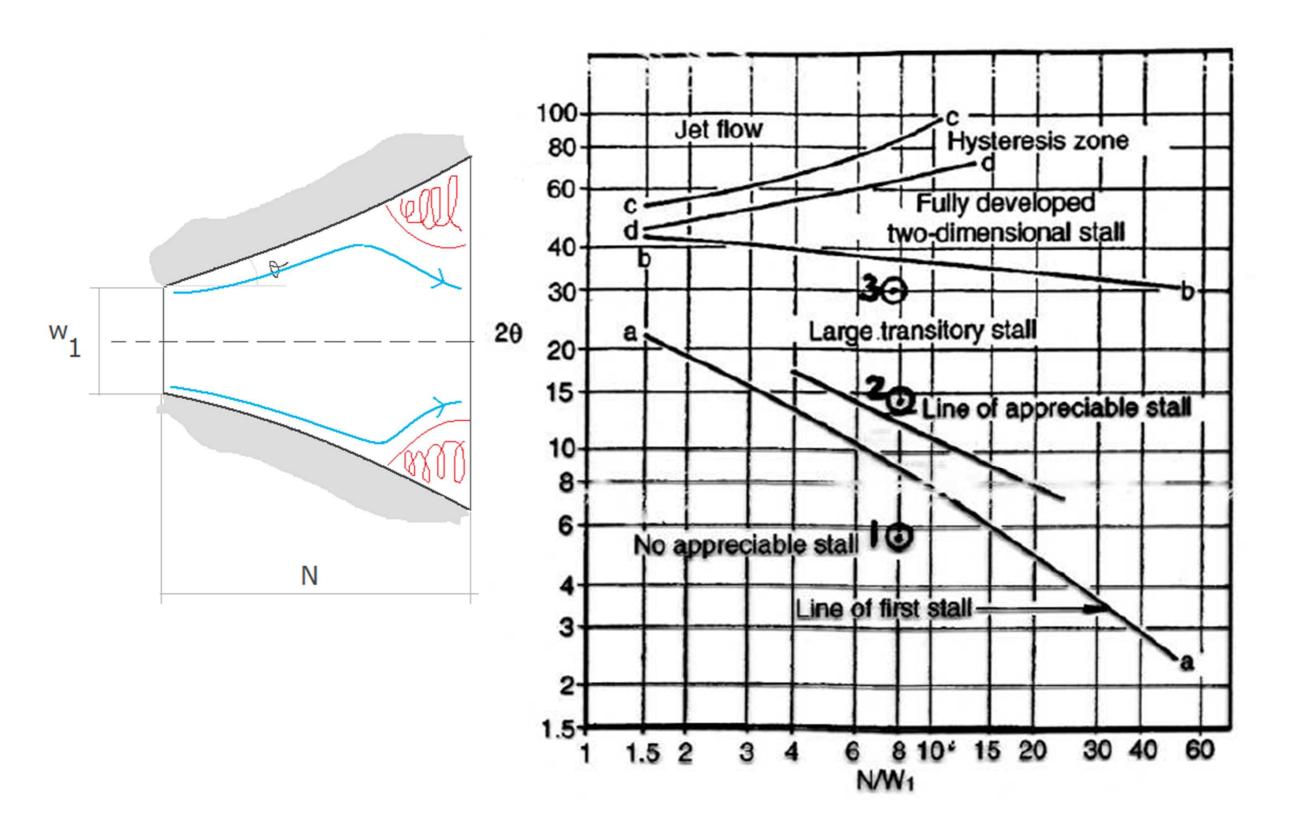
$$p_2 = p_{02} - \frac{1}{2}\rho c_2^2$$

$$p_1 = p_{01} - \frac{1}{2}\rho c_1^2$$

$$c_{pi} = \frac{c_1^2 - c_2^2}{c_1^2} = 1 - \left(\frac{c_2}{c_1}\right)^2 = 1 - \left(\frac{A_1}{A_2}\right)^2 = 1 - \frac{1}{A_R^2}$$

Legame tra  $c_p$ ,  $\eta_{is}$  e  $c_{pi}$ 

$$\frac{c_p}{c_{pi}} = \frac{p_2 - p_1}{p_{01} - p_1} \cdot \frac{p_{01} - p_1}{\left(p_2 - p_1\right) + \left(p_{01} - p_{02}\right)} = \eta_{is}$$



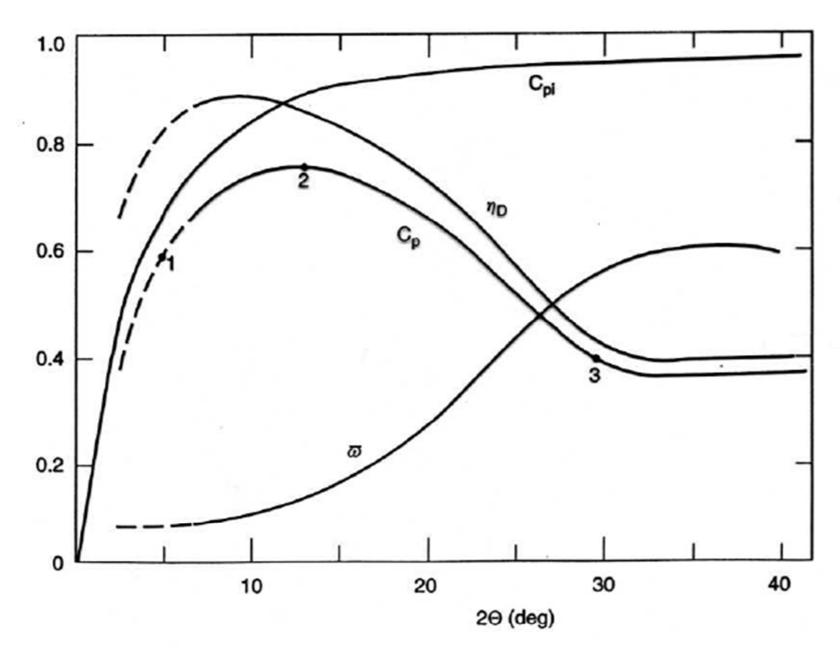


Fig. 2.16. Typical diffuser performance curves for a two-dimensional diffuser, with  $L/W_1 = 8.0$  (adapted from Kline et al. 1959).

https://www.youtube.com/watch?
v=JhlEkEk7igs&list=PL0EC6527BE871ABA3&index=8