

# Exercises QFT II — 2018/2019

## Problem Sheet 4

### Problem 7: Full scalar propagator in terms of 1PI 2pt function

Consider the full propagator, i.e. the two-point Green's function  $\tilde{G}^{(2)}(p)$  in momentum space, for a scalar self-interacting theory (keep the interaction Lagrangian generic, but proportional to a coupling constant  $\lambda \ll 1$ ). We have seen that

$$-i\tilde{G}^{(2)}(p) = \left(\tilde{\Gamma}^{(2)}(p)\right)^{-1}, \quad (1)$$

where  $\tilde{\Gamma}^{(2)}(p)$  is the Fourier transform of  $\Gamma^{(2)} \equiv \frac{\delta^2\Gamma}{\delta\psi^2}$ .

In perturbation theory one can express  $\tilde{\Gamma}^{(2)}(p)$  as the sum (we neglect the  $i\epsilon$  term)

$$\tilde{\Gamma}^{(2)}(p) = p^2 - m^2 - i\lambda\Pi. \quad (2)$$

- Justify the expression (2), starting from the fact that  $\tilde{G}^{(2)}(p) = \frac{i}{p^2 - m^2}$  at leading order in  $\lambda$ .
- By using (2) and (1), show that the full propagator (2pt connected Green's function) can be expressed as

$$\tilde{G}^{(2)}(p) = \frac{i}{p^2 - m^2} + \frac{i}{p^2 - m^2} \lambda\Pi \frac{i}{p^2 - m^2} + \frac{i}{p^2 - m^2} \lambda\Pi \frac{i}{p^2 - m^2} \lambda\Pi \frac{i}{p^2 - m^2} + \dots \quad (3)$$

[Use the fact that  $\frac{1}{1-X}$  can be expressed as the series  $\sum_{n=0}^{\infty} X^n$ . In this exercise,  $X$  will be proportional to  $\lambda$ .]

- Interpret the relation (3) diagrammatically (the full propagator should be now written in terms of the 1PI 2pt function and the free propagator). [The proper vertices are  $i\tilde{\Gamma}^{(n)}$ ].

## Problem 8: Physical mass and two point function

At the first order in perturbation theory, the two-point function of  $\lambda\phi^4$ -theory is

$$\tilde{\Gamma}^{(2)}(p) = p^2 - m^2 + \frac{\lambda m^2}{32\pi^2} \left( \frac{1}{2-\omega} + \psi(2) - \ln \left( \frac{m^2}{4\pi\mu^2} \right) + \mathcal{O}(2-\omega) \right) \quad (4)$$

In the *minimal subtraction scheme* ( $MS$ ), the renormalized mass  $m_r$  is defined through the relation

$$m_r^2 = m^2 - \frac{\lambda m^2}{32\pi^2} \frac{1}{2-\omega} \quad (5)$$

1. Write  $\tilde{\Gamma}^{(2)}(p)$  in terms of  $m_r$  instead of  $m$ . [Invert the expression (5) and then substitute the found result for  $m^2$  into (4).]
2. Find the expression of the physical mass  $m_{ph}$  in terms of  $m_r$ .  
[Solve the equation  $\tilde{\Gamma}^{(2)}(p^2 = m_{ph}^2) = 0$ .]
3. Write  $\tilde{\Gamma}^{(2)}(p)$  in terms of  $m_{ph}$  instead of  $m_r$  and show that the dependence on  $\mu$  disappears.
4. Repeat points 1,2,3 by using a different subtraction scheme, i.e. the *modified minimal subtraction scheme* ( $\overline{MS}$ ):

$$m_r^2 = m^2 - \frac{\lambda m^2}{32\pi^2} \left( \frac{1}{2-\omega} + \psi(2) \right) \quad (6)$$