Exercises QFT II — 2018/2019

Problem Sheet 12

Problem 23 - Anomaly and Ward Identities I

Consider the free Lagrangian

$$
\mathcal{L} = \bar{\psi}(i \ \partial - m)\psi \,,\tag{1}
$$

and the following currents:

$$
j^{\mu} = \bar{\psi}\gamma^{\mu}\psi \qquad j_5^{\mu} = \bar{\psi}\gamma^{\mu}\gamma_5\psi \qquad (2)
$$

1. Using the equation of motion, compute the classical value for $\partial_\mu j^\mu$ and $\partial_\mu j^\mu_5$ $\frac{\mu}{5}$. What happens when $m = 0$?

Consider the following (Fourier transform of the) Green function:

$$
T_{\mu\nu\lambda} = i \int d^4x \, d^4y \, d^4z \, e^{ik_1x + ik_2y - iqz} \, \langle 0|Tj_\mu(x)j_\nu(y)j_\lambda^5(z)|0\rangle \tag{3}
$$

- 2. Draw the (triangle) Feynmann diagrams that describe this Green function.
- 3. Using Feynman rules, show that

$$
T_{\mu\nu\lambda}(m) = i \int \frac{d^4p}{(2\pi)^4} (-) \text{tr} \frac{i}{\not p - m} \gamma_\lambda \gamma_5 \frac{i}{\not p - \not q - m} \gamma_\nu \frac{i}{\not p - \not k_1 - m} \gamma_\mu + \begin{pmatrix} k_1 \leftrightarrow k_2 \\ \mu \leftrightarrow \nu \end{pmatrix}
$$
 (4)

Consider also

$$
T_{\mu\nu} = i \int d^4x \, d^4y \, d^4z \, e^{ik_1x + ik_2y - iqz} \langle 0|Tj_\mu(x)j_\nu(y)P(z)|0\rangle \qquad \text{with } P = \bar{\psi}\gamma_5\psi \tag{5}
$$

4. Show that plugging the classical results for $\partial_\mu j^\mu$ and $\partial_\mu j^\mu_5$ $\frac{\mu}{5}$ into (3), one obtains:

$$
\text{AWI}: \ q^{\lambda} T_{\mu\nu\lambda} = 2m T_{\mu\nu} \qquad \qquad \text{VWI}: \ k_1^{\mu} T_{\mu\nu\lambda} = k_2^{\nu} T_{\mu\nu\lambda} = 0 \tag{6}
$$

These are called the Axial Ward Identity and the Vector Ward Identity.

5. Draw the Feynman diagrams for $T_{\mu\nu}$ and show that:

$$
T_{\mu\nu}(m) = i \int \frac{d^4p}{(2\pi)^4} (-) \text{tr} \frac{i}{\not p - m} \gamma_5 \frac{i}{\not p - \not q - m} \gamma_\nu \frac{i}{\not p - \not k_1 - m} \gamma_\mu + \begin{pmatrix} k_1 \leftrightarrow k_2 \\ \mu \leftrightarrow \nu \end{pmatrix}
$$
(7)

6. Prove that this expression is convergent. [Write down explicitly the fermion propagator and use the properties of the traces of the gamma matrices with γ_5 .

Let us consider the Axial Ward Identity and see if the classical result is the same as the quantum one.

7. Using $\oint \gamma_5 = \gamma_5(\not p - \oint -m) + (\not p - m)\gamma_5 + 2m\gamma_5$, show that

$$
q^{\lambda}T_{\mu\nu\lambda} = 2mT_{\mu\nu} + R_{\mu\nu}^{1} + R_{\mu\nu}^{2}
$$
\n(8)

where

$$
R_{\mu\nu}^{1} = \int \frac{d^{4}p}{(2\pi)^{4}} [f(p - k_{2}) - f(p)] \quad \text{with} \quad f(p) = \text{tr} \frac{1}{p - m} \gamma_{5} \gamma_{\nu} \frac{1}{p - k_{1} - m} \gamma_{\mu} \quad (9)
$$

and $R_{\mu\nu}^2$ is equal to $R_{\mu\nu}^1$ with $k_1 \leftrightarrow k_2$ and $\mu \leftrightarrow \nu$.

8. Prove that $R^1_{\mu\nu}$ is linearly divergent. Then we need to regularize it.

Problem 24 - Anomaly and Ward Identities II

Pauli-Villars regularization: We introduce a regulator that preserves the Vector WI, while it violates the Axial WI (there exsists no regulator that preserves both). We define:

$$
T_{\mu\nu\lambda}^{\text{reg}} \equiv T_{\mu\nu\lambda}(m) - T_{\mu\nu\lambda}(m = M) \quad \text{and} \quad (10)
$$

we take the difference before doing the integral. Then the physical Green function is defined as: $T_{\mu\nu\lambda}^{\text{phys}}$ $\lim_{M\to\infty} T_{\mu\nu}^{\text{reg}}$ µνλ

- 1. Prove that $T^{\text{reg}}_{\mu\nu\lambda}$ is finite.
- 2. Prove that now $R^{\text{1reg}}_{\mu\nu}$ is finite and actually zero. [Hint: Show first that $\int_{-\infty}^{+\infty} dx$ [$f(x + a) - f(x)$] = 0 if $\int_{-\infty}^{+\infty} dx f(x)$ is convergent.]
- 3. Prove that the Vector WI is satisfied, i.e. $k_1^{\mu}T_{\mu\nu\lambda}^{\text{reg}} = 0 = k_2^{\nu}T_{\mu\nu\lambda}^{\text{reg}}$.
- 4. Show that the Axial WI is now:

$$
q^{\lambda}T_{\mu\nu\lambda}^{\text{phys}} = 2mT_{\mu\nu}(m) - \lim_{M \to \infty} 2MT_{\mu\nu}(M)
$$
\n(11)

5. Show that $T_{\mu\nu}(M)$ can be written as

$$
T_{\mu\nu}(M) = \int \frac{d^4p}{(2\pi)^4} 2 \int_0^1 dx \int_0^{1-x} dy \frac{4iM \epsilon_{\mu\nu\alpha\beta} k_1^{\alpha} k_2^{\beta}}{[p^2 - 2p \cdot \tilde{k}_1 - \tilde{M}^2]^3} + \left(\begin{array}{c} 1 \leftrightarrow 2\\ \mu \leftrightarrow \nu \end{array}\right) \tag{12}
$$

with appropriate \tilde{k}_1 and \tilde{M} . Determine \tilde{k}_1 and \tilde{M}_1 explicitly.

6. Use the relation

$$
\int \frac{d^n p}{(p^2 - 2p \cdot \tilde{k} - \tilde{M}^2)^A} = i^{1 - 2A} \pi^{n/2} \frac{\Gamma(A - \frac{n}{2})}{\Gamma(A)} \frac{1}{(\tilde{k}^2 + \tilde{M}^2)^{A - n/2}}
$$
(13)

to compute the integral (12).

7. Perform the limit $M \to \infty$ to show that

$$
q^{\lambda}T_{\mu\nu\lambda}^{\text{phys}} = 2mT_{\mu\nu}(m) - \frac{1}{2\pi^2} \epsilon_{\mu\nu\alpha\beta} k_1^{\alpha} k_2^{\beta} . \tag{14}
$$