

1) $\vec{v} = \frac{1}{\hbar} \vec{\nabla}_{\vec{k}} \mathcal{E}$; seguendo il suggerimento al punto seguente, calcoliamo $\vec{\nabla}_{\vec{k}} \left[\mathcal{E}(\vec{k}) (1 + 2\alpha \mathcal{E}(\vec{k})) \right] = \vec{\nabla}_{\vec{k}} \left(\frac{\hbar^2 \vec{k}^2}{2m^*} \right)$

che dà: $\vec{\nabla}_{\vec{k}} \mathcal{E} [1 + 2\alpha \mathcal{E}] = \frac{\hbar^2 \vec{k}}{m^*}$ (eq 2)

Però dunque esplicitare $\mathcal{E}(\vec{k})$ dall'eq. (1):

$$2\mathcal{E}^2 + \mathcal{E} - \frac{\hbar^2 \vec{k}^2}{2m^*} = 0 \Rightarrow \mathcal{E}_{1,2}(\vec{k}) = \frac{-1 \pm \sqrt{1 + 2\alpha \frac{\hbar^2 \vec{k}^2}{m^*}}}{2\alpha}$$

Per identificare la soluzione che si riduce alla banda parabolica quando $\alpha \rightarrow 0$, considero lo sviluppo di Taylor rispetto a α :

$$\mathcal{E}_{1,2}(\vec{k}) \approx \frac{-1 \pm \left(1 + \frac{1}{2} 2\alpha \frac{\hbar^2 \vec{k}^2}{m^*} \right)}{2\alpha} = \begin{cases} + \frac{\hbar^2 \vec{k}^2}{2m^*} & \leftarrow \text{parabola} \\ -\frac{1}{\alpha} - \frac{\hbar^2 \vec{k}^2}{2m^*} & \text{(divergente)} \end{cases}$$

Quindi va considerata la soluz con '+'.
Dall'eq. (2):

$$\boxed{\vec{v} = \frac{1}{\hbar} \vec{\nabla} \mathcal{E} = \frac{\hbar^2 \vec{k}}{m^*} \frac{1}{1 + 2\alpha \mathcal{E}(\vec{k})} = \frac{\hbar^2 \vec{k}}{m^*} \frac{1}{\sqrt{1 + 2\alpha \frac{\hbar^2 \vec{k}^2}{m^*}}}}$$

2) $\boxed{g(\mathcal{E})} = 2 \int \frac{1}{4\pi^2} \frac{d\vec{k}}{|\vec{\nabla} \mathcal{E}|}$ → elemento di linea su un cerchio (perché $\mathcal{E}(\vec{k}) = \mathcal{E}(|\vec{k}|)$, quindi le surf. a energie costanti sono cerchi)
curva a $\mathcal{E}(\vec{k}) = \text{cost.} = \mathcal{E}$

$$= 2 \int \frac{1}{4\pi^2} \frac{k d\theta}{\frac{\hbar^2 k}{m^*} [1 + 2\alpha \mathcal{E}]^{-1}} = 2 \cdot 2\pi \cdot \frac{1}{4\pi^2} \frac{m^* [1 + 2\alpha \mathcal{E}]}{\hbar^2 [1 + 2\alpha \mathcal{E}]}, \mathcal{E} \geq 0$$

dall'eq. (2) che si riduce a $g(\mathcal{E}) = \frac{m^*}{\hbar^2 \pi}$ per $\alpha \rightarrow 0$ ($\mathcal{E} \geq 0$)

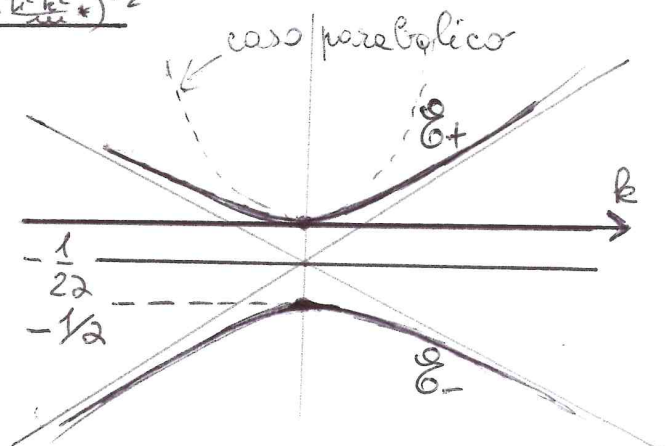
3) grafico di $\mathcal{E}(\vec{k})$: $\mathcal{E}(\vec{k}) = \frac{-1 \pm (1 + 2\alpha \frac{\hbar^2 \vec{k}^2}{m^*})^{1/2}}{2\alpha}$

Identifichiamo gli asintoti:

$$\mathcal{E}(\vec{k}) \underset{|\vec{k}| \rightarrow +\infty}{\sim} \frac{-1 \pm \sqrt{2\alpha \frac{\hbar^2}{m^*}} k}{2\alpha}$$

chiamo $\mathcal{E}_0(\vec{k})$

$$\text{Ho: } \mathcal{E}_0(\vec{k}=0) = -\frac{1}{2\alpha}$$



Per $\alpha \rightarrow 0$ la curva $\mathcal{E}(\vec{k})$

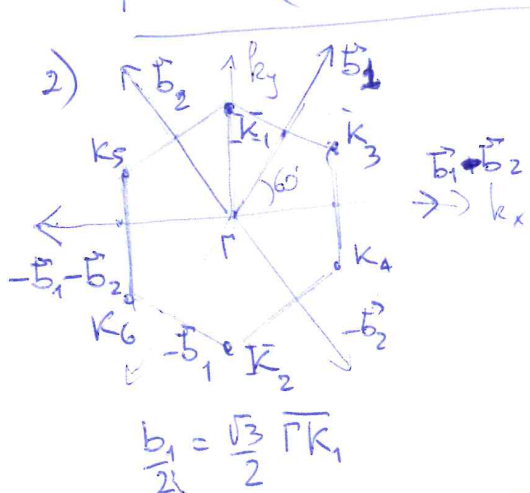
da iperbole diventa parabola $\Rightarrow g(\mathcal{E})$ è diversa

$$a = |\vec{a}_1| = |\vec{a}_2|$$

$$\vec{a}_{1,2} = \left(\pm a \frac{\sqrt{3}}{2}, \frac{a}{2} \right) \Rightarrow \begin{cases} \vec{b}_1 \perp \vec{a}_2 \Rightarrow \\ \vec{b}_1 \text{ inclinato di } 60^\circ: \vec{b}_1 = b \left(\frac{1}{2}, \frac{\sqrt{3}}{2} \right) \\ \vec{b}_2 \perp \vec{a}_1 \Rightarrow \\ \vec{b}_2 \text{ inclinato di } 120^\circ: \vec{b}_2 = b \left(-\frac{1}{2}, \frac{\sqrt{3}}{2} \right) \end{cases}$$

$$\vec{b}_1 \cdot \vec{a}_1 = 2\pi \Rightarrow a \frac{\sqrt{3}}{2} b \frac{1}{2} + a \frac{\sqrt{3}}{2} b = 2\pi \Rightarrow ab \frac{\sqrt{3}}{2} = 2\pi \Rightarrow b = \frac{4\pi}{\sqrt{3}} \frac{1}{a}$$

$$\Rightarrow \vec{b}_{1,2} = \left(\pm \frac{2\pi}{\sqrt{3}} \frac{1}{a}, \frac{2\pi}{a} \right) \rightarrow b = |\vec{b}_{1,2}| = \frac{2\pi}{a} \left(\frac{1}{3} + 1 \right)^{\frac{1}{2}} = \frac{4\pi}{\sqrt{3}} \frac{1}{a}$$



$$\vec{K}_{1,2} = \left(0, \pm \frac{4\pi}{3} \frac{1}{a} \right)$$

$$\vec{K}_{3,4} = \left(\frac{\vec{b}_1 - \vec{b}_2}{2} \right) = \left(\frac{b}{2}, \frac{b}{2} \right) = \left(\frac{2\pi}{\sqrt{3}a}, \frac{2\pi}{3a} \right)$$

$$= \left(\frac{2\pi}{\sqrt{3}a}, \frac{2\pi}{3a} \right) \quad \pm, \pm \text{ tutte le 4 combinazioni}$$

$$3) E(\vec{K}_{1,2}) = \pm \gamma \sqrt{1 + 4 \cos\left(\frac{\sqrt{3}a \cdot 0}{2}\right) \cos\left(\frac{4\pi a \cdot 1}{3a \cdot 2}\right) + 4 \cos^2\left(\frac{4\pi a \cdot 1}{3a \cdot 2}\right)}$$

$$= \pm \gamma \left(1 + 2 \cos\left(\frac{2\pi a}{3a}\right) \right) = \pm \gamma \left(1 + 2 \cdot \left(-\frac{1}{2}\right) \right) = 0$$

$$E(\vec{K}_{3,4}) = \pm \gamma \sqrt{1 + 4 \cos\left(\frac{\sqrt{3}a \cdot \frac{2\pi}{\sqrt{3}a}}{2}\right) \cos\left(\frac{2\pi a}{3a \cdot 2}\right) + 4 \cos^2\left(\frac{\pi}{3}\right)}$$

$$= \pm \gamma \sqrt{1 - 4 \cos\frac{\pi}{3} + 4 \cos^2\frac{\pi}{3}}$$

$$= \pm \gamma \left(1 - 2 \cos\frac{\pi}{3} \right) = \pm \gamma \left(1 - \frac{1 \cdot 2}{2} \right) = 0$$

$E(\vec{K}_i) = 0$ e l'energia è simmetriche rispetto a $E=0$
(le bande sono)

4) Per la simmetria \Rightarrow 1 e⁻ per atomo \Rightarrow $E_F = 0$

5) $k_x = 0 \Rightarrow$

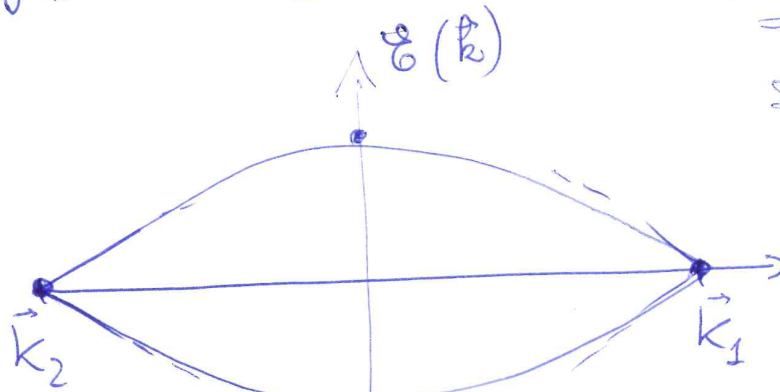
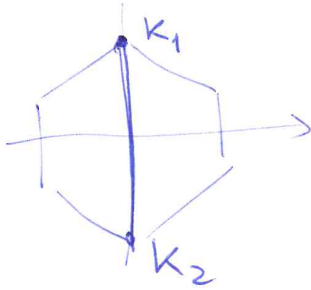
$$\mathcal{E}(0, k_y) = \pm f \sqrt{1 + 4 \cos\left(\frac{k_y a}{2}\right) + 4 \cos^2\left(\frac{k_y a}{2}\right)}$$

$$= \pm f \left| 1 + 2 \cos\left(\frac{k_y a}{2}\right) \right| \rightarrow \text{periodo } \frac{k_y a}{2} = 2\pi$$

$$\Rightarrow k_y = \frac{4\pi}{a}$$

Se vado da $-\frac{4\pi}{3a}$ a $+\frac{4\pi}{3a}$

NON FACCI
TUTTO IL
PERIODO



$$\mathcal{E}\left(0, \pm \frac{4\pi}{3} \frac{1}{a}\right) = \pm f \left| 1 + 2 \cos\left(\frac{2}{3} \frac{4\pi}{3} \frac{1}{a} \frac{1}{2}\right) \right| = \pm f \left(1 + 2 \cos\left(\frac{4\pi}{3}\right) \right) = 0$$

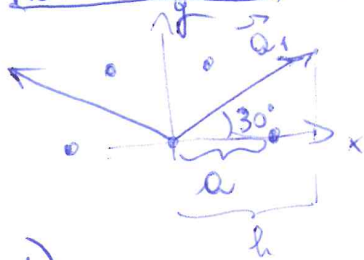
$$\mathcal{E}(0, 0) = \pm f |1 + 2| = \pm 3f \quad \begin{matrix} \text{max!} \\ \text{min} \end{matrix}$$

$$\frac{\partial}{\partial k_y} \mathcal{E}(0, k_y) = \pm f \frac{2(-1)a}{2} \sin \frac{k_y a}{2} = \mp a f \sin \frac{k_y a}{2}$$

$$\lim_{k_y \rightarrow \pm \frac{4\pi}{3} \frac{1}{a}} \frac{\partial}{\partial k_y} \mathcal{E}(0, k_y) = \mp a f \sin \frac{2}{3} \frac{4\pi}{3} \frac{1}{a} \frac{1}{2} = \mp a f \frac{\sqrt{3}}{2} \neq 0$$

\Rightarrow derivata $\neq 0 \Rightarrow$ comportamen. lineare

ALTERNATIVA: chiamiamo "a" la dist NN



$$\begin{cases} \frac{3}{2}a = h \\ h = a_1 \frac{\sqrt{3}}{2} \end{cases} \Rightarrow \boxed{a_1 = \frac{2}{\sqrt{3}} \frac{3}{2}a = \sqrt{3}a}$$

$$\boxed{\vec{a}_{1,2} = \left(\pm \frac{3}{2}a, \frac{\sqrt{3}}{2}a \right)}$$

1)

$$\vec{b}_1 = b \left(\frac{1}{2}, \frac{\sqrt{3}}{2} \right); \quad \vec{b}_1 \cdot \vec{a}_1 = 2\pi \Rightarrow \frac{3}{4}ab + \frac{3}{4}ab = 2\pi \Rightarrow ab = 2\pi \frac{2}{3} = \frac{4\pi}{3}$$

$$\Rightarrow b = \frac{4\pi}{3a} \Rightarrow \boxed{\vec{b}_{1,2} = \frac{2\pi}{3a} \left(\pm \frac{1}{2}, \frac{\sqrt{3}}{2} \right)}$$

$$\vec{b}_1 \cdot \vec{a}_1 = \frac{3}{4}a \cdot \frac{2\pi}{3a} + \frac{\sqrt{3}}{4}a \cdot \frac{2\pi}{3} \frac{\sqrt{3}}{2} = 2\pi = \vec{b}_2 \cdot \vec{a}_2$$

$$\vec{b}_1 \cdot \vec{a}_2 = \vec{b}_2 \cdot \vec{a}_1 = 0$$

2) "vecchio a" \rightarrow ^{poi} $\sqrt{3}a$

quindi $\vec{k}_{1,2} = \left(0, \pm \frac{4\pi}{3\sqrt{3}a} \right)$

$$\vec{k}_{3,4} = \left(\pm \frac{2\pi}{3a}, \pm \frac{2\pi}{3\sqrt{3}a} \right)$$

Altri punti di conseguenza ...