

Definition 1.5.3

Relative Entropy

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Let $\rho_1, \rho_2 \in \mathcal{D}(S)$ such that $\text{ker}(\rho_2) \subseteq \text{ker}(\rho_1)$.

Then, the relative entropy of ρ_1 with respect to ρ_2 is

$$S(\rho_1 | \rho_2) := \text{Tr} \rho_1 (\log \rho_1 - \log \rho_2)$$

Lemma 1.5.4

$$S(\rho_1 | \rho_2) \geq 0$$

Proof:

1. if $\left[\sum_{\mathbb{R}} c_{\mathbb{R}} f_{\mathbb{R}}(x) g_{\mathbb{R}}(y) \geq 0 \right]$ for $c_{\mathbb{R}} \in \mathbb{R}$, $f_{\mathbb{R}}, g_{\mathbb{R}}: [a, b] \rightarrow \mathbb{R}$

then $\left[\sum_{\mathbb{R}} c_{\mathbb{R}} \text{Tr} (f_{\mathbb{R}}(A) g_{\mathbb{R}}(B)) \geq 0 \right]$ for A, B with spectrum in $[a, b]$

$$\sum_{\mathbb{R}} c_{\mathbb{R}} \text{Tr} (f_{\mathbb{R}}(A) g_{\mathbb{R}}(B)) = \sum_{\mathbb{R}} c_{\mathbb{R}} f_{\mathbb{R}}(a_{\mathbb{R}}) g_{\mathbb{R}}(b_{\mathbb{S}}) \text{Tr} (P_{\mathbb{R}} Q_{\mathbb{S}}) \geq 0$$

for $A = \sum_{\mathbb{R}} a_{\mathbb{R}} P_{\mathbb{R}}$ and $B = \sum_{\mathbb{S}} b_{\mathbb{S}} Q_{\mathbb{S}}$ (spectral representations)

- for $\eta(x) = -x \log x$, $\eta(x) - \eta(y) \leq \eta'(y)(x-y)$ (concavity)
 $x, y \in [0, 1]$

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$$\eta(y) - \eta(x) \geq -(1 + \log y)(y-x) \quad \text{implies}$$

$$\begin{aligned} \text{Tr}(\eta(\rho_2)) - \text{Tr}(\eta(\rho_1)) &\geq -\text{Tr}((1 + \log \rho_2)(\rho_2 - \rho_1)) \\ &= \text{Tr}(\eta(\rho_2)) + \text{Tr} \rho_1 \log \rho_2 \end{aligned}$$

$$\boxed{\text{Tr} \rho_1 (\log \rho_1 - \log \rho_2) = S(\rho_1 | \rho_2) \geq 0}$$

Proposition 1.5.3

Subadditivity of $S(\rho)$

Let $\rho_{12} \in \mathcal{D}(S_1 + S_2)$ be a bipartite state with marginal states

$\rho_1 = \text{Tr}_2 \rho_{12}$ and $\rho_2 = \text{Tr}_1 \rho_{12}$; then,

$$\boxed{S(\rho_{12}) \leq S(\rho_1) + S(\rho_2)}$$

Proof : $0 \leq S(\rho_{12} | \rho_1 \otimes \rho_2) = \text{Tr}(\rho_{12} (\log \rho_{12} - \log(\rho_1 \otimes \rho_2)))$

set $\rho_1 = \sum_{i \in I_1} r_{1i} P_{1i}$, $\rho_2 = \sum_{i \in I_2} r_{2i} P_{2i}$ (spectral representations)

$$\begin{aligned} \log(\rho_1 \otimes \rho_2) &= \sum_{i \in I_1, j \in I_2} \log(r_{1i} r_{2j}) P_{1i} \otimes P_{2j} \\ &= \sum_{i \in I_1, j \in I_2} (\log r_{1i} + \log r_{2j}) P_{1i} \otimes P_{2j} = \log \rho_1 \otimes \mathbb{1}_2 + \mathbb{1}_1 \otimes \log \rho_2 \end{aligned}$$

Then, $0 \leq \text{Tr}(\rho_{12} \log \rho_{12}) - \text{Tr}(\rho_{12} (\log \rho_1 \otimes \mathbb{1}_2)) - \text{Tr}(\rho_{12} \mathbb{1}_1 \otimes (\log \rho_2))$

$$0 \leq -S(\rho_{12}) + S(\rho_1) + S(\rho_2)$$

Corollary 1.5.1

$$S(\rho_{12}) = S(\rho_1) + S(\rho_2) \iff \rho_{12} = \rho_1 \otimes \rho_2$$

Proof $S(\rho_1 | \rho_2) = 0 \iff \rho_1 = \rho_2$

1.6. von Neumann entropy and bipartite entanglement

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Observation: consider two qubits in the entangled state

$$|\Psi_{\text{unif}}^{(2)}\rangle = \frac{|00\rangle + |11\rangle}{\sqrt{2}}, \quad P_{\text{unif}}^{(2)} = |\Psi_{\text{unif}}^{(2)}\rangle\langle\Psi_{\text{unif}}^{(2)}|$$

This pure state has zero entropy:

$$\boxed{S(P_{\text{unif}}^{(2)}) = 0}, \quad \text{more the less its two marginals}$$

$$\rho_1 = \text{Tr}_2 P_{\text{unif}}^{(2)} = \frac{1}{2} = \rho_2 = \text{Tr}_1 P_{\text{unif}}^{(2)} \quad \text{have both}$$

maximal entropies

$$\boxed{S(\rho_1) = S(\rho_2) = \log 2}$$

The whole (compound) system $S_1 + S_2$ can be perfectly known,

$$\boxed{S(P_{\text{unif}}^{(2)}) = 0}, \quad \text{without knowing anything of its parties,}$$

$$\boxed{S(\rho_1) = S(\rho_2) = \log 2}$$

Remark : This phenomenon is in contrast to the classical relations, $H(X, Y) \geq \max(H(X), H(Y))$, that says that the knowledge of the whole is always greater than the knowledge of its parts.
 This quantum occurrence $S(\rho_{12}) - S(\rho_1) \neq 0$ is due to entanglement.

Lemma 1.6.1

Let $|\psi_{12}\rangle$ be a separable state vector of a bipartite system $S_1 + S_2$. Then,

$S(\rho_{12}) = S(\rho_1) = S(\rho_2) = 0$, where $\rho_{12} = |\psi_{12}\rangle\langle\psi_{12}|$ and $\rho_1 = \text{Tr}_2 \rho_{12}$, $\rho_2 = \text{Tr}_1 \rho_{12}$.

Proof. $|\psi_{12}\rangle$ is separable iff $|\psi_{12}\rangle = |\psi_1\rangle \otimes |\psi_2\rangle$, then ρ_{12} , $\rho_1 = |\psi_1\rangle\langle\psi_1|$ and $\rho_2 = |\psi_2\rangle\langle\psi_2|$ are pure states with zero von Neumann entropy.

What about density matrices?

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Lemma 1.6.2 $\rho_{12} \in \mathcal{D}(S_1 + S_2)$ separable $\Rightarrow S(\rho_{12}) \geq \max\{S(\rho_1), S(\rho_2)\}$

Proof: assume $\rho_{12} = \sum_{i,j} d_{ij} \rho_i^{(1)} \otimes \rho_j^{(2)}$ so that

$$\rho_1 = \text{Tr}_2 \rho_{12} = \sum_i \left(\sum_j d_{ij} \right) \rho_i^{(1)}, \quad \rho_2 = \text{Tr}_1 \rho_{12} = \sum_j \left(\sum_i d_{ij} \right) \rho_j^{(2)}.$$

Set $d_2 = \dim(\mathcal{H}_2)$, then $\left| S(\rho_{12}, \rho_1 \otimes \frac{1}{d_2}) \geq 0 \right|.$

$$\begin{aligned} 0 &\leq S(\rho_{12}, \rho_1 \otimes \frac{1}{d_2}) = S\left(\sum_{i,j} d_{ij} \rho_i^{(1)} \otimes \rho_j^{(2)} \mid \sum_{i,j} d_{ij} \rho_i^{(1)} \otimes \frac{1}{d_2}\right) \\ &\leq \sum_{i,j} d_{ij} S(\rho_i^{(1)} \otimes \rho_j^{(2)} \mid \rho_i^{(1)} \otimes \frac{1}{d_2}) \quad (\text{Joint convexity of } S(\rho_1, \rho_2)) \\ &= \sum_{i,j} d_{ij} \text{Tr}_{12}(\rho_i^{(1)} \otimes \rho_j^{(2)} (\log \rho_i^{(1)} + \log \rho_j^{(2)} - \log \rho_i^{(1)} + \log d_2)) \\ &= \sum_{i,j} d_{ij} \text{Tr}_2 \rho_j^{(2)} \log \rho_j^{(2)} + \log d_2 = \log d_2 - \sum_{i,j} d_{ij} S(\rho_j^{(2)}) \end{aligned}$$

$$S(\rho_{12} | \rho_1 \otimes \frac{1}{d_2}) = \text{Tr}_{12} \rho_{12} (\log \rho_{12} - \log \rho_1 + \log \frac{1}{d_2}) = -S(\rho_{12}) + S(\rho_1) + \log d_2$$

$$\text{(from page 81)} \leq \log d_2 - \sum_{i,j} d_{ij} S(\rho_j^{(2)})$$

$$\text{Then, } S(\rho_{12}) \geq S(\rho_1) + \sum_{i,j} d_{ij} S(\rho_j^{(2)}) \geq S(\rho_1)$$

Remark: the entropy of a bipartite separable states
(uncertainty about the whole compound system)
is always larger than or equal to the entropy
of the marginal density matrices
(uncertainty about the constituent particles).

Question: is the von Neumann entropy an exhaustive entanglement witness? Namely, if $S(\rho_{12}) \geq \max\{S(\rho_1), S(\rho_2)\}$ is then ρ_{12} separable?

NO

Example 1.6.1 Horodecky's states

Let us consider the family of states $\rho_F \in M_d(\mathbb{C}) \otimes M_d(\mathbb{C})$ of the form

$\rho_F = aI + bP_{\text{unif}}^{(d)}$ where $P_{\text{unif}}^{(d)} = \frac{1}{d} \sum_{i,j=1}^d |i\rangle\langle j| \otimes |i\rangle\langle j|$, $\{|i\rangle\}_{i=1}^d$ ONB in \mathbb{C}^d

$\begin{cases} \text{Tr } \rho_F = 1 = ad^2 + b \\ F := \text{Tr}(\rho_F P_{\text{unif}}^{(d)}) = \text{Tr}(aP_{\text{unif}}^{(d)} + bP_{\text{unif}}^{(d)}) = a + b \geq 0 \end{cases} \Rightarrow \begin{cases} 1 = ad^2 + F - a \\ b = F - a \end{cases} \Rightarrow \begin{cases} a = \frac{1-F}{d^2-1} \\ b = \frac{Fd-1}{d^2-1} \end{cases}$

Positivity of ρ_F : $\text{Spectrum}(\rho_F) = \begin{cases} a \geq 0 \\ a + b = F \geq 0 \end{cases} \Rightarrow \mathbf{0 \leq F \leq 1}$

Partial transposition : $\rho_F = \frac{1-F}{d^2-1} + \frac{Fd^2-1}{d^2-1} P_{\text{sing}}^{(d)}$ sent into

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$$\rho_F^1 = \frac{1-F}{d^2-1} + \frac{Fd^2-1}{d^2-1} \frac{1}{d} V, \quad \forall |\psi \otimes \phi\rangle = |\phi \otimes \psi\rangle$$

$$\begin{aligned} \text{Spectrum}(\rho_F^1) &= \left\{ \frac{1-F}{d^2-1} + \frac{Fd^2-1}{d^2-1} \frac{1}{d} \right\} = \left\{ \frac{d-dF \pm Fd^2-1}{d(d^2-1)} \right\} = \left\{ \frac{\pm Fd(|d-1|) + |d-1|}{d(d+1)(d-1)} \right\} \\ &= \left\{ \frac{Fd+1}{d(d+1)}, \frac{1-Fd}{d(d-1)} \right\} \end{aligned}$$

ρ_F is Positive under Partial Transposition (PPT: see Section 1.4.)

if $1-Fd \geq 0 \Leftrightarrow 0 \leq F \leq \frac{1}{d}$.

If $F > \frac{1}{d}$ we know ρ_F is entangled.

For such a class of states it is also known that PPT is equivalent to separability

What about the von Neumann entropy?

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$$S(\rho_F) = -(d^2-1) \frac{1-F}{d^2-1} \log \frac{1-F}{d^2-1} - F \log F$$

$$\rho_F^{(1)} := \text{Tr}_2 \rho_F = d \frac{1-F}{d^2-1} \mathbb{1}_1 + \frac{F(d^2-1)}{d^2-1} \frac{1}{d} \mathbb{1}_1 = \frac{d^2 - Fd^2 + Fd^2 - 1}{d(d^2-1)} = \frac{\mathbb{1}_1}{d}$$

$$\rho_F^{(2)} := \text{Tr}_1 \rho_F = \frac{d(1-F)}{d^2-1} \mathbb{1}_2 + \frac{F(d^2-1)}{d^2-1} \frac{1}{d} \mathbb{1}_2 = \frac{\mathbb{1}_2}{d}$$

Choose $d=2$: $S(\rho_F) = -(1-F) \log \frac{1-F}{3} - F \log F$

$$S(\rho_F^{(1)}) = S(\rho_F^{(2)}) = \log 2$$

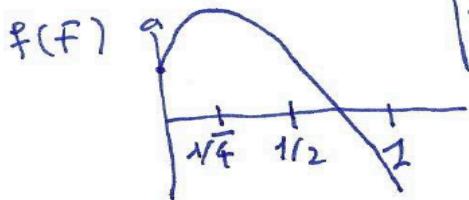
$$S(\rho_F) - S(\rho_F^{(1)}) = -(1-F) \log \frac{1-F}{3} - F \log F - \log 2 = f(F)$$

$$f(0) = \log 3 - \log 2 > 0$$

$$f(1) = -\log 2$$

$$f'(F) = \log \frac{1-F}{3F}$$

$$f''(F) = -\frac{1}{F(1-F)}$$



$$f\left(\frac{1}{2}\right) = \log \frac{3}{2} > 0$$

$\Rightarrow \exists F > \frac{1}{2}$ with $\begin{cases} \rho_F \text{ entangled} \\ \phi \\ S(\rho_F) > S(\rho_F^{(1)}) \end{cases}$