

# Image Processing for Physicists

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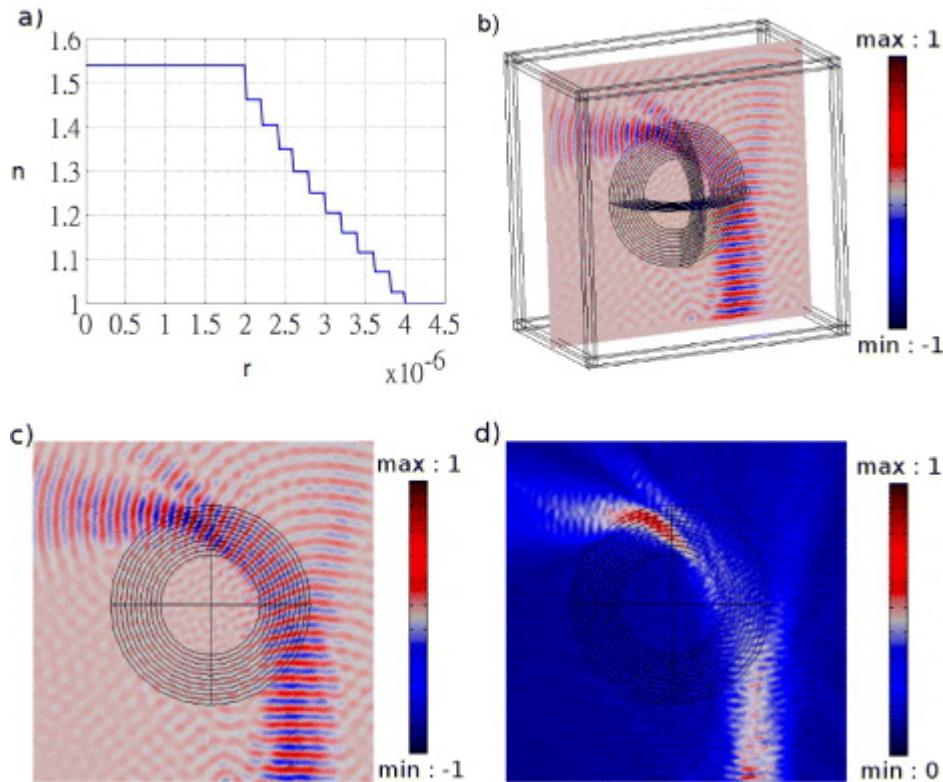
# Overview

- Propagation modelization
- Wave propagation:
  - Near-field regime
  - Far-field regime

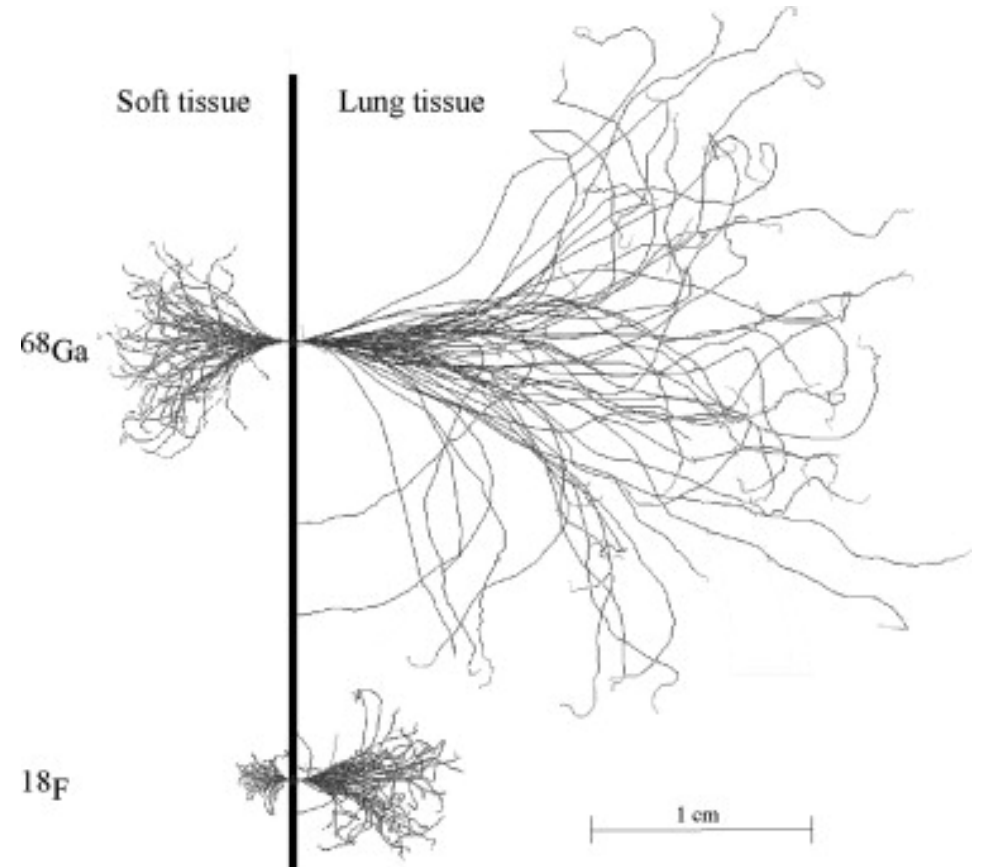
# Propagation modeling

- Motivations:

## 1. Validation



Finite element simulation of an electromagnetic field in a dielectric



Monte Carlo simulation of positrons trajectories resulting from  $^{68}\text{Ga}$  and  $^{18}\text{F}$  decay.

sources: T.M. Chang *et al.* New J. Phys. (2012)  
A. Sanchez-Crespo, Appl. Rad. Isotopes (2012)

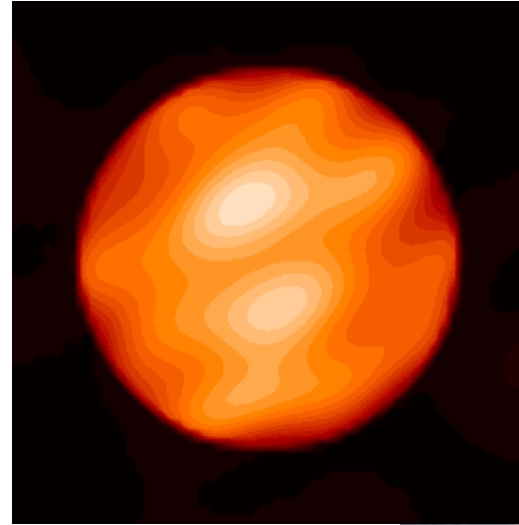
# Propagation modeling

- Motivations:

## 2. Inversion



Image reconstruction from sound wave propagation (ultrasonography)



The surface of Betelgeuse reconstructed from interferometric data (IOTA)



sources: wikipedia

Haubois *et al.* *Astronom. & Astrophys.* (2009)



# Propagation modeling

- Particles
  - Model particle tracks (rays) through different media
  - Model may include: refraction, force fields, particle decay and interactions
  - Not included: diffraction
- Wave
  - Model the interaction of a field with a medium
  - Can be very complicated → approximations are needed

# Propagation modeling

Starting point: Helmholtz equation

- for EM field: neglect polarization (scalar wave approximation)
- for electron wave, assume high energy electrons

Maxwell eq.  $\rightarrow \nabla^2 \psi + \frac{n^2}{c^2} \frac{\partial^2}{\partial t^2} \psi = 0$

index of refraction (spatially dependent)

$\psi$ : complex-valued field (scalar: no polarization)

$n$ : index of refraction

$c$ : speed of light

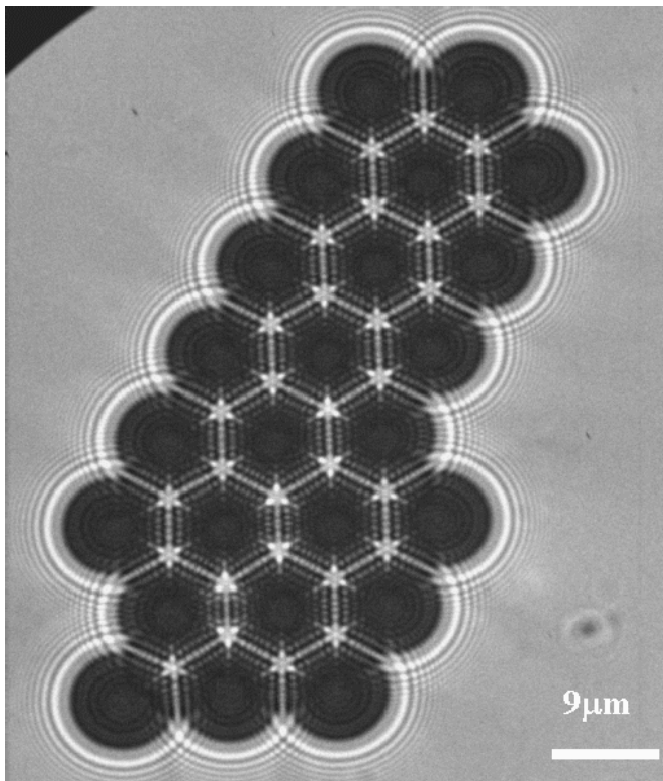
constant  $n \rightarrow$  plane wave solution  $\psi = \psi_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$

$k^2 = n^2 \frac{\omega^2}{c^2}$  "dispersion relation"

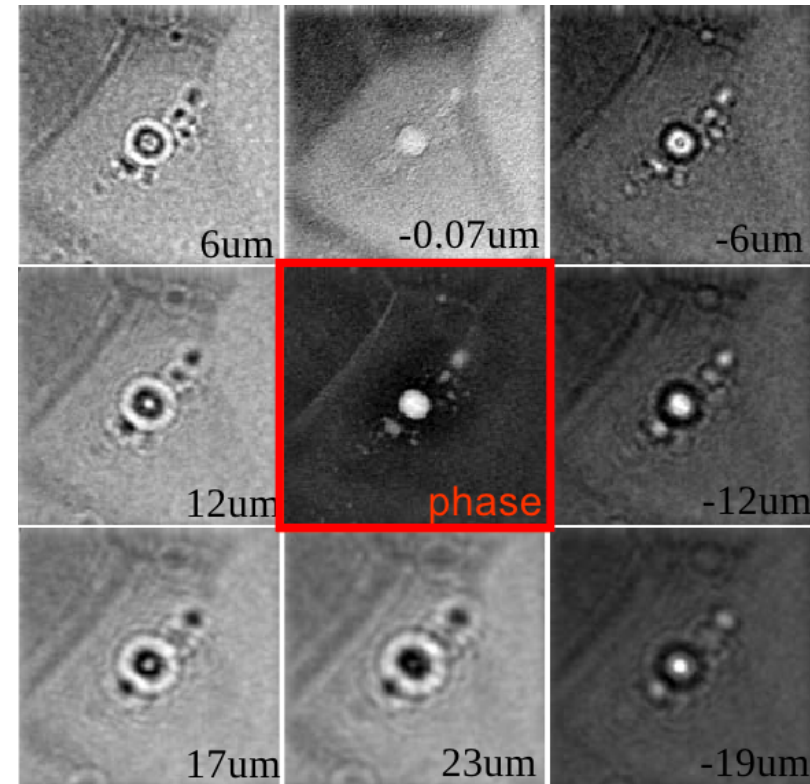
# Propagation modeling

- Useful to:
  - better understand optical systems
  - understand diffraction, holography, phase contrast, interferometry, ...

X-ray hologram



TEM through-focus series



sources: Mayo *et al.* Opt. Express (2003)  
<http://www.christophtkoch.com/Vorlesung/>

# The physics of propagation

In free space ( $n=1$ ) General solution

$$\psi(\vec{r}, t) = \sum_{\omega} \sum_{\vec{k}} A_{\omega \vec{k}} e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

$\uparrow$   $|\vec{k}| = \frac{\omega}{c}$

$$|\vec{k}| = \frac{2\pi}{\lambda}$$
$$"u = \frac{1}{P}"$$

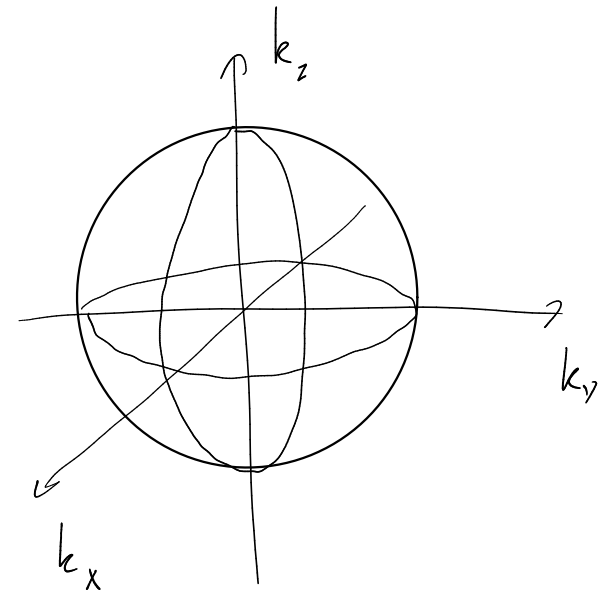
Commonly: fix  $\omega$  and solve monochromatic case

$$\rightarrow \psi(\vec{r}) = \sum_{\vec{k}} A_{\vec{k}} e^{i\vec{k} \cdot \vec{r}} \quad \leftarrow \text{is this a Fourier transform?}$$

$\uparrow$   $|\vec{k}| = \frac{\omega}{c}$

$$k_x^2 + k_y^2 + k_z^2 = \frac{\omega^2}{c^2}$$

"Ewald" sphere



# The physics of propagation

## "Angular spectrum" representation

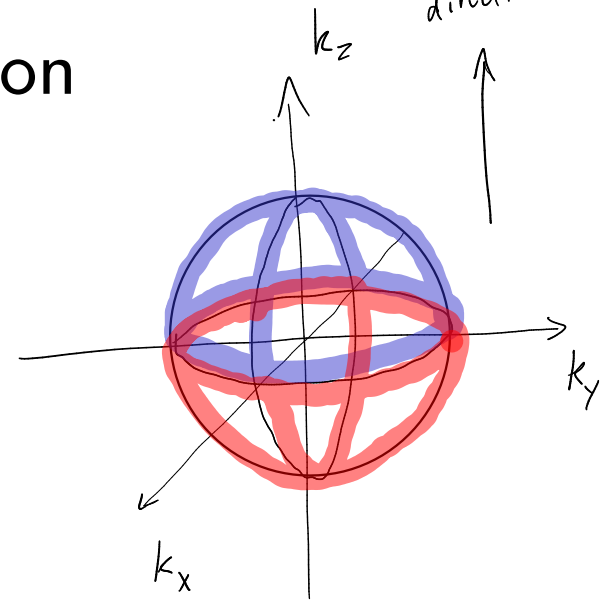
$$\frac{\omega^2}{c^2} = k^2 = k_z^2 + k_x^2 + k_y^2$$

$$k_z = \begin{matrix} + \\ - \end{matrix} \sqrt{k^2 - k_x^2 - k_y^2}$$

$$\psi(\vec{r}) = \sum_{k_x k_y} A_{k_x k_y}^+ e^{i(k_x x + k_y y + \sqrt{k^2 - k_x^2 - k_y^2} z)}$$

~~$$+ A_{k_x k_y}^- e^{i(k_x x + k_y y - \sqrt{k^2 - k_x^2 - k_y^2} z)}$$~~

← forward propagation



$$\psi(\vec{r}_\perp, z) = \sum_{\vec{k}_\perp} A_{\vec{k}_\perp} \underbrace{\exp(i\vec{k}_\perp \cdot \vec{r}_\perp)}_{\text{2D Fourier transform}} \exp(i\sqrt{k^2 - k_\perp^2} z)$$



# Forward propagation

Angular spectrum  
 $k_x = "k \sin \theta"$

$$z=0: \psi(\vec{r}_\perp; z=0) = \sum_{\vec{k}_\perp} A_{\vec{k}_\perp} \exp(i\vec{k}_\perp \cdot \vec{r}_\perp)$$

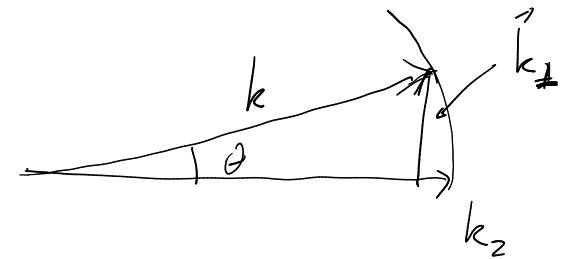
$$\Rightarrow A_{\vec{k}_\perp} = \mathcal{F} \{ \psi(\vec{r}_\perp; z=0) \}$$

Often:  $|\vec{k}_\perp| \ll k$  (small angle diffraction)

$$\sqrt{k^2 - k_\perp^2} \approx k \sqrt{1 - \frac{k_\perp^2}{k^2}} \approx k \left( 1 - \frac{k_\perp^2}{2k^2} \right)$$

"paraxial approximation"

$$= k - \frac{k_\perp^2}{2k}$$



parabolic approximation  
of a sphere

$$\exp(i\sqrt{k^2 - k_\perp^2} z) \approx \exp(ikz) \exp\left(\frac{-iz k_\perp^2}{2k}\right)$$

"Fresnel propagator"

# Forward propagation

$$\psi(r_{\perp}; z) = \mathcal{F}^{-1} \left\{ \mathcal{F} \left\{ \psi(r_{\perp}; z=0) \right\} \exp(-2\pi i \lambda z u^2) \right\}$$

The trick for numerical implementations:

F.T.  $e^{i2\pi x u}$

D.F.T.  $e^{i2\pi \frac{n \cdot k}{N}}$  ↑ integers

$$x = n \cdot \Delta x$$

$$u = k \cdot \Delta u$$

$$\Delta x \Delta u = \frac{1}{N}$$

e.g.  $\Delta x, N$  are known  
 $\hookrightarrow \Delta u = \frac{1}{\Delta x N}$

sampling in real and Fourier space.

$$\exp(-\pi i \lambda z u^2) \rightarrow \exp(-\pi i \lambda z \cdot k^2 \frac{1}{\Delta x^2 N^2})$$

$$\exp\left(-\pi i \frac{\lambda}{\Delta x} \frac{z}{\Delta x} \left(\frac{k}{N}\right)^2\right)$$

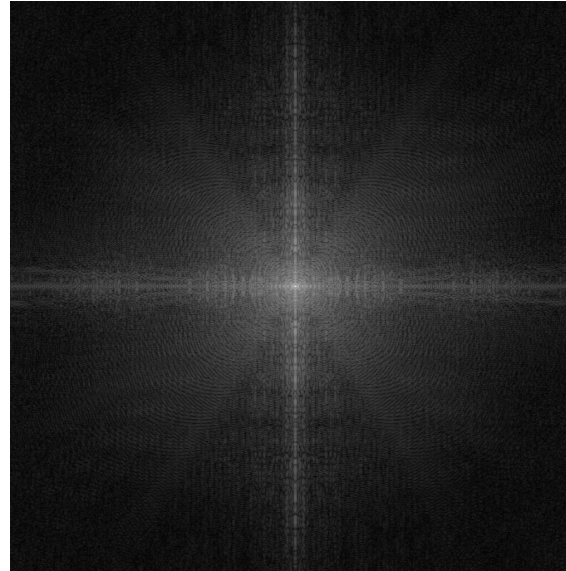
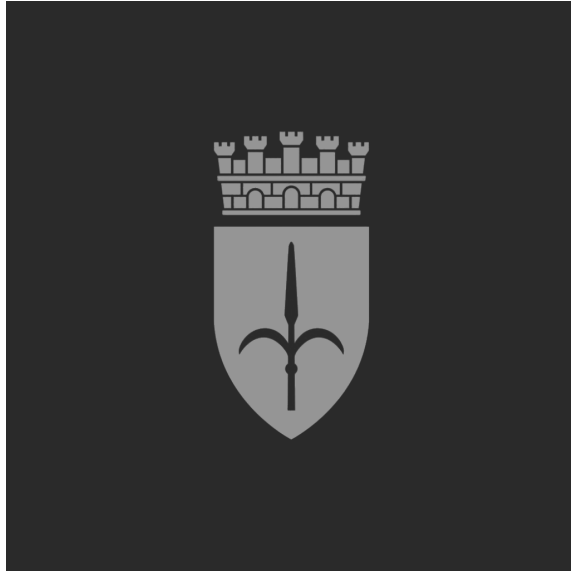
$\frac{k}{N}$  output of

not  $\frac{2\pi}{\lambda}$  here!  
 numpy. fft, fftfreq

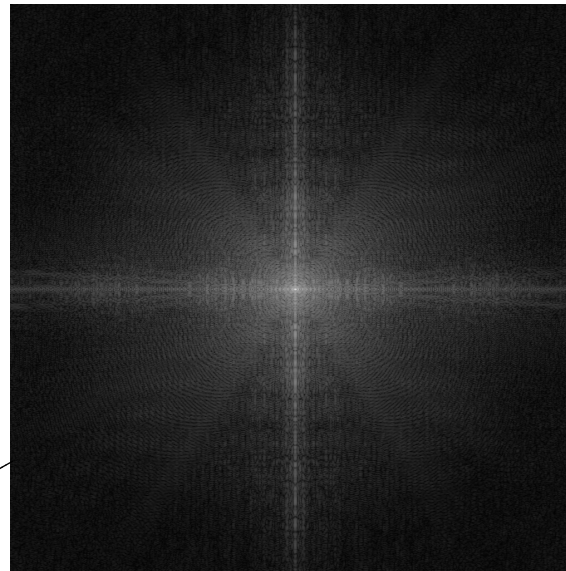
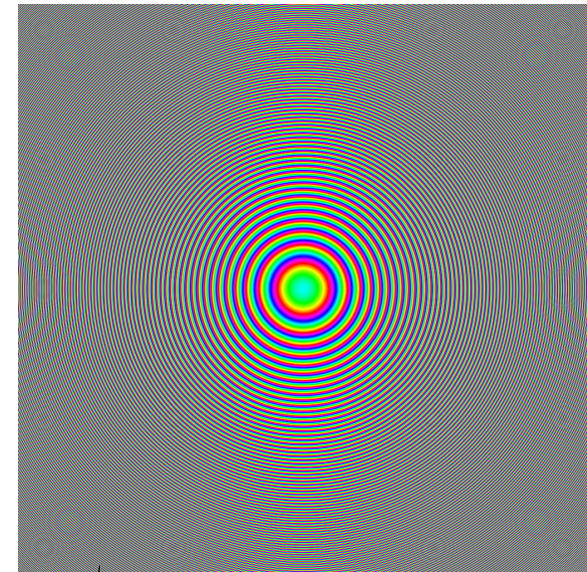
# Forward propagation

A numerical recipe

$$\psi(\vec{r}_\perp; z=0)$$



$$\times \exp\left(\frac{-iz|k_\perp|^2}{2k}\right)$$



# Near field, far field

$$\psi(r_{\perp}; z) = \mathcal{F}^{-1} \left\{ \mathcal{F} \left\{ \psi(r_{\perp}; z=0) \right\} \exp \left( - \frac{iz k_{\perp}^2}{2k} \right) \right\} \frac{2\pi}{\lambda}$$

$$\psi(r_{\perp}; z) = \mathcal{F}^{-1} \left\{ \mathcal{F} \left\{ \psi(r_{\perp}; z=0) \right\} \exp(-\pi i \lambda z \vec{u}^2) \right\}$$

"F<sup>-1</sup>{F.G}" = f \* g

$$\vec{k}_{\perp} = \vec{u} \cdot 2\pi$$

$e^{i k x}$

notation  
 $(\vec{r}_{\perp} \rightarrow \vec{r})$

$$= \mathcal{P}_z(\vec{r}_{\perp}) * \psi(\vec{r}_{\perp}; z=0)$$

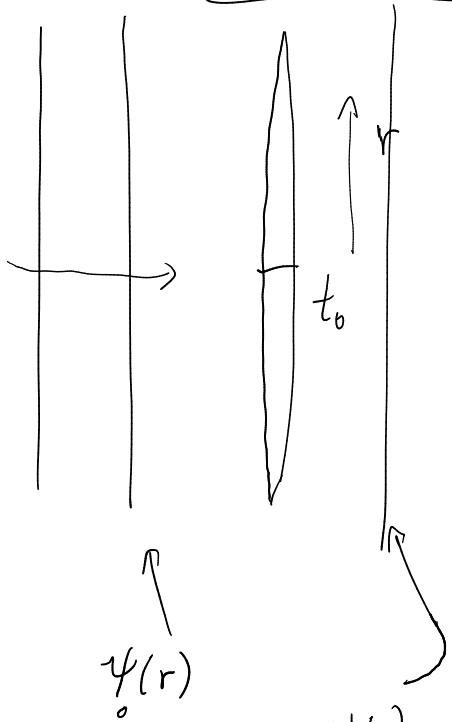
$$\mathcal{F}^{-1} \left\{ \exp(-\pi i \lambda z \vec{u}^2) \right\}$$

$$\psi(\vec{r}, z) = \frac{-i 2\pi}{\lambda z} \int d^2 r' \psi(r', z=0) \exp \left( \frac{ik(r-r')^2}{2z} \right)$$



"Fresnel - Huygens integral"

# Back focal plane of a lens



$$t(r) = t_0 - \alpha r^2 \quad (\text{model for thin lens thickness profile})$$

Passed the lens:

$$\phi(r) = \frac{2\pi}{\lambda} (n-1) t(r)$$

$$= k(n-1)t_0 - k(n-1)\alpha r^2$$

$$\frac{1}{2f} = (n-1)\alpha$$

$$\Psi(r; z) = \frac{-ik}{z} \int d^2 r' \underbrace{\psi_0 e^{-ik(n-1)\alpha r'^2}}_{\text{exit wave}} e^{\frac{ik(\vec{r}-\vec{r}')^2}{2z}}$$

(A lens acts as a Fourier transform operator)

$$= \frac{-ik}{z} \int d^2 r' \psi_0 \exp\left(ik\left[-\frac{1}{2f} r'^2 + \frac{r^2}{2z} + \frac{r'^2}{2z} - \frac{\vec{r}\cdot\vec{r}'}{z}\right]\right)$$

$$= \frac{-ik}{z} e^{\frac{ikr^2}{2z}} \int d^2 r' \psi_0 \exp\left(\frac{ik}{z} \vec{r}\cdot\vec{r}'\right) \exp\left(\frac{ikr'^2}{2}\left(\frac{1}{z} - \frac{1}{f}\right)\right)$$

$\underbrace{\frac{ik}{z} \vec{r}\cdot\vec{r}'}_{\vec{u}}$

$$\Psi(r; z=f) = \frac{ik}{z} e^{\frac{ikr^2}{2z}} \int \{\psi_0(r)\} \left(\vec{u} = \frac{k\vec{r}}{z}\right)$$

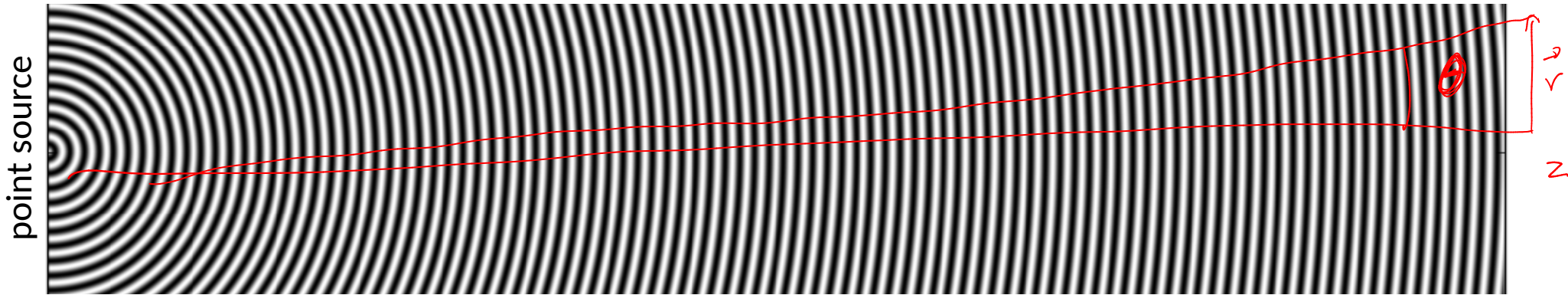


# Plane waves, point sources

Far-field propagation (Fraunhofer regime):

$$\Psi(r; z \rightarrow \infty) = \frac{ik}{z} e^{i\frac{kr^2}{2z}} \mathcal{F}\{\psi_0(r)\} \left( \vec{u} = \frac{kr^2}{z} \right)$$

" $z \sin \theta$ "

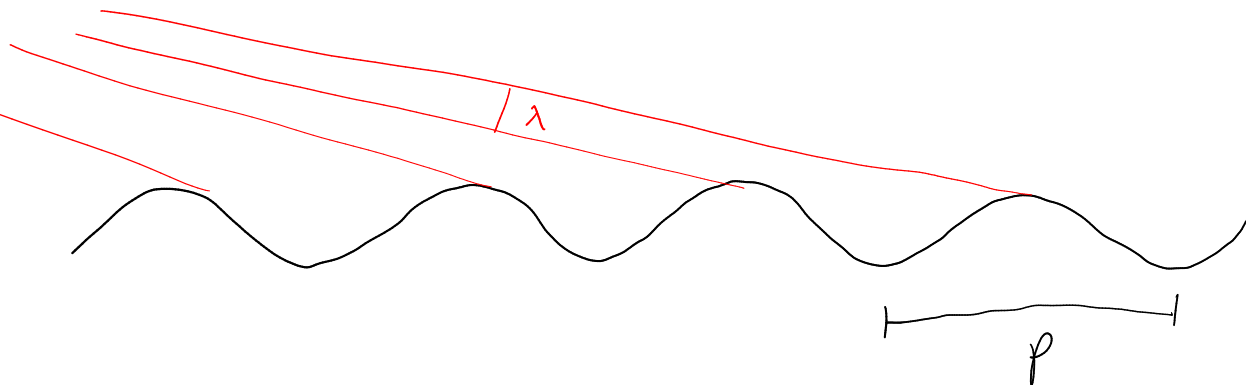
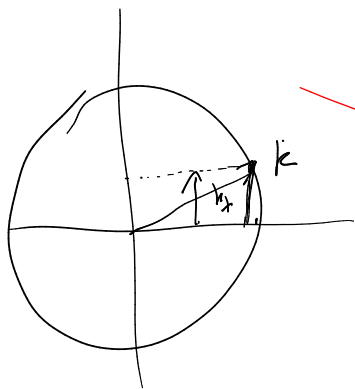


circular waves  
evanescent waves  
contact region

parabolic waves  
near field  
Fresnel region

carrier wave  $\lambda$   $H$   
MWW

plane waves  
far field  
Fraunhofer region

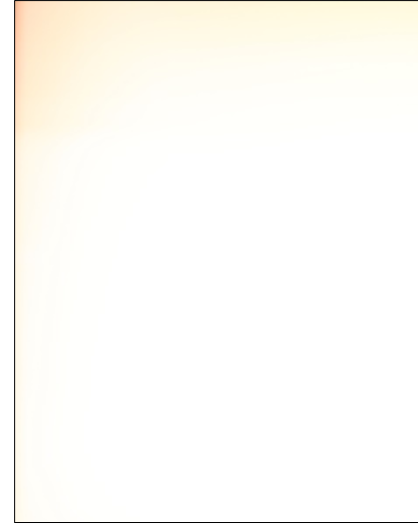


oscillation of given  
spatial  
frequency  
in initial  
image

# Why optical elements?



with objective lens

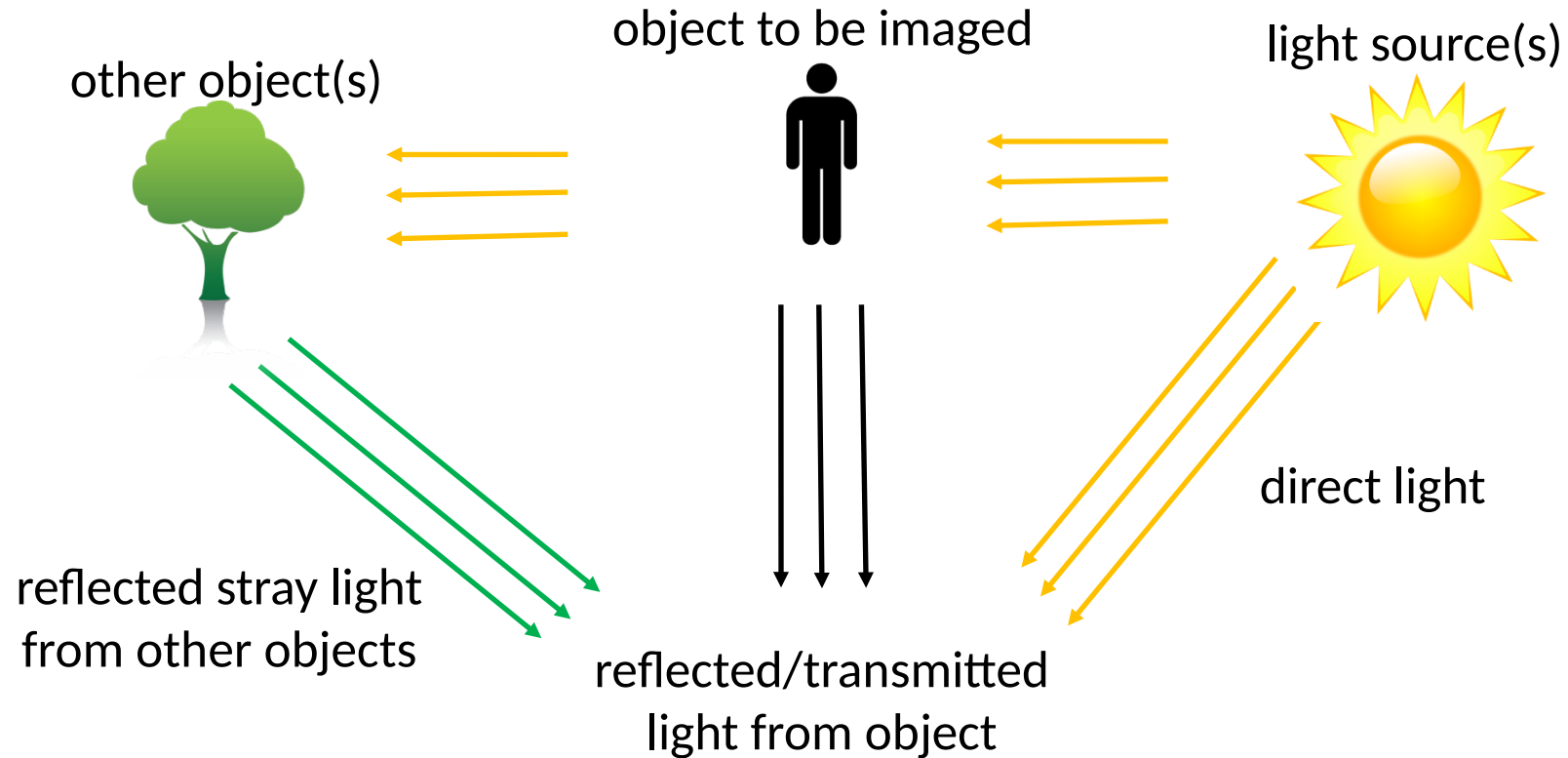


without objective lens



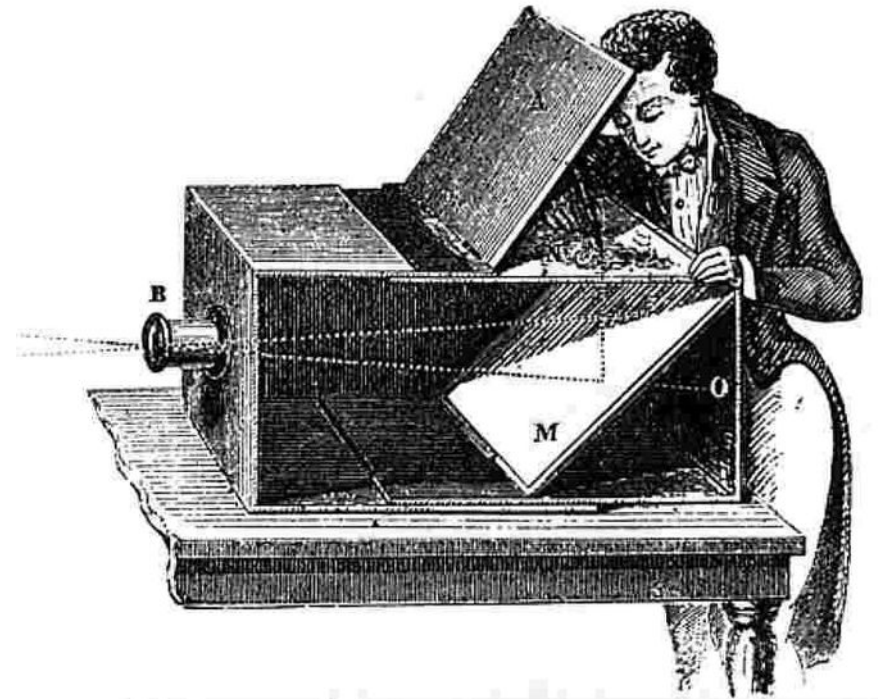
# Why optical elements?

- Information from many sources overlaps in detector plane
- Need models to understand image forming systems



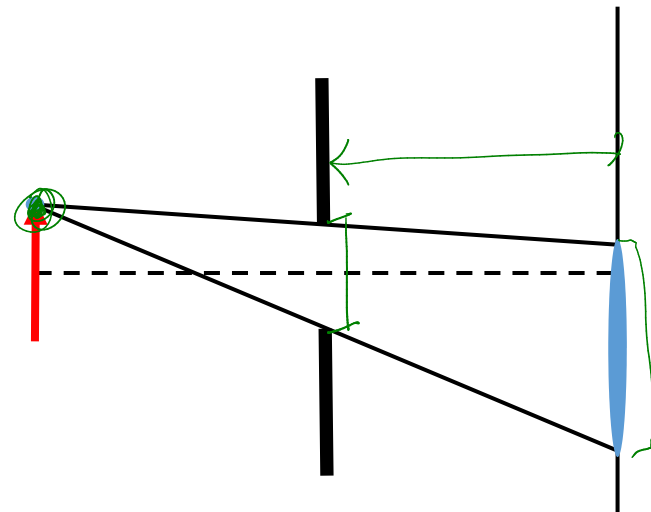
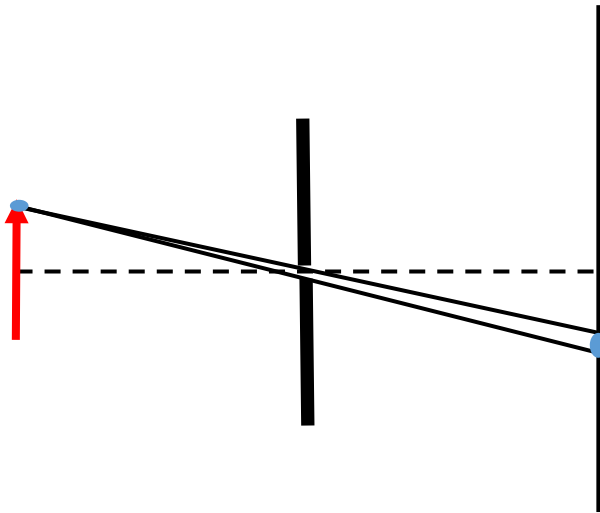
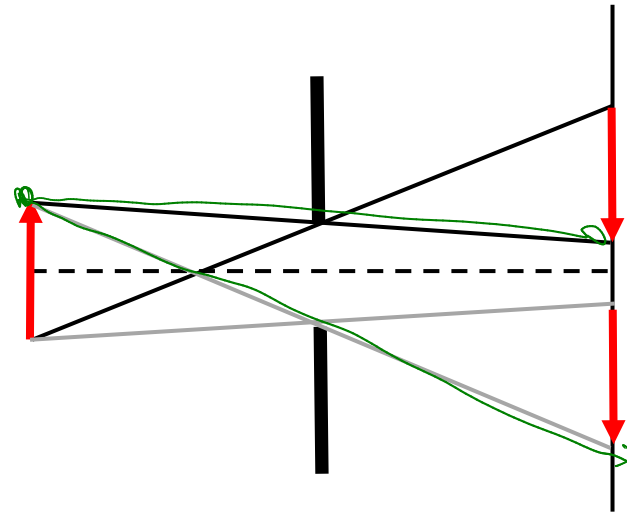
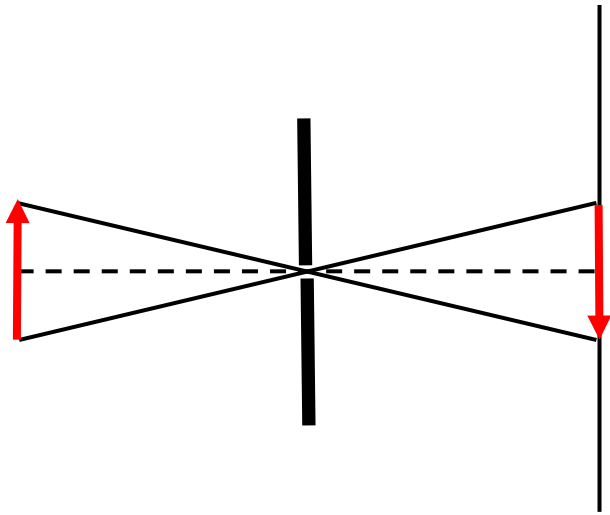
# Pinhole camera model

camera obscura



# Pinhole camera model

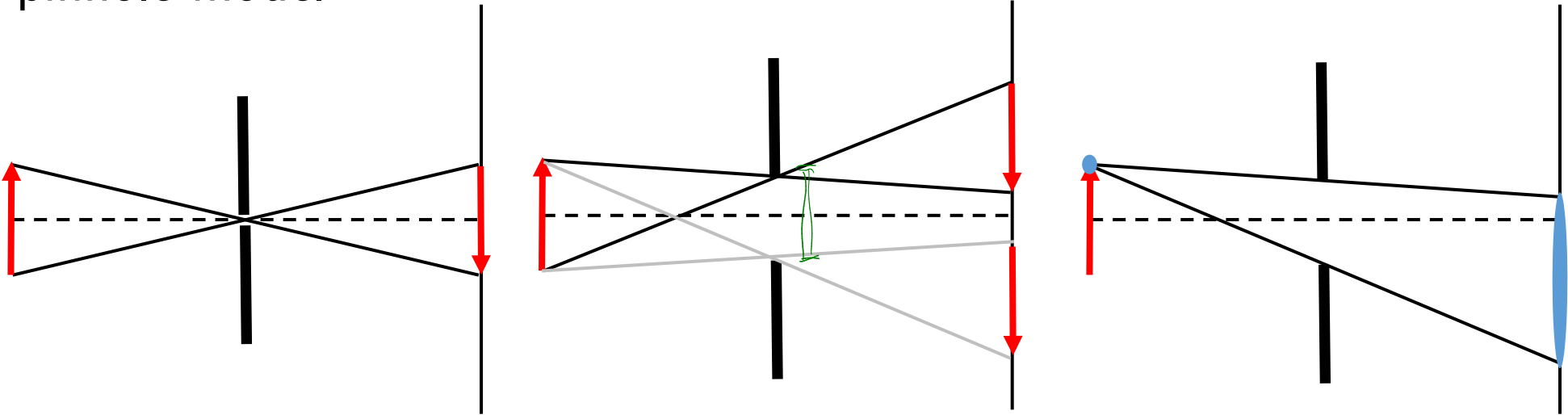
PSF determined by aperture width





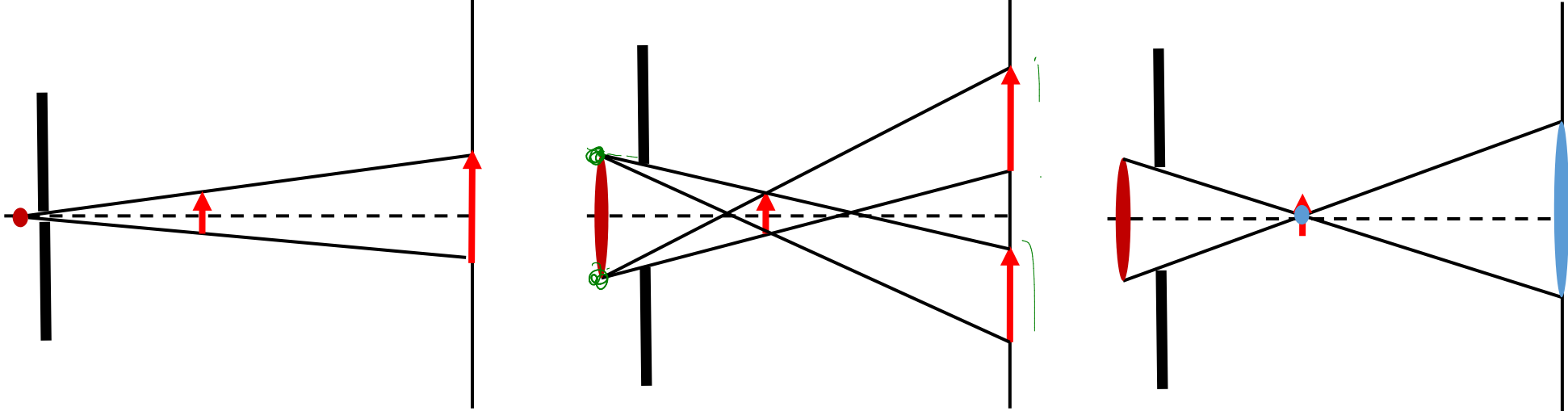
# Projection model

pinhole model

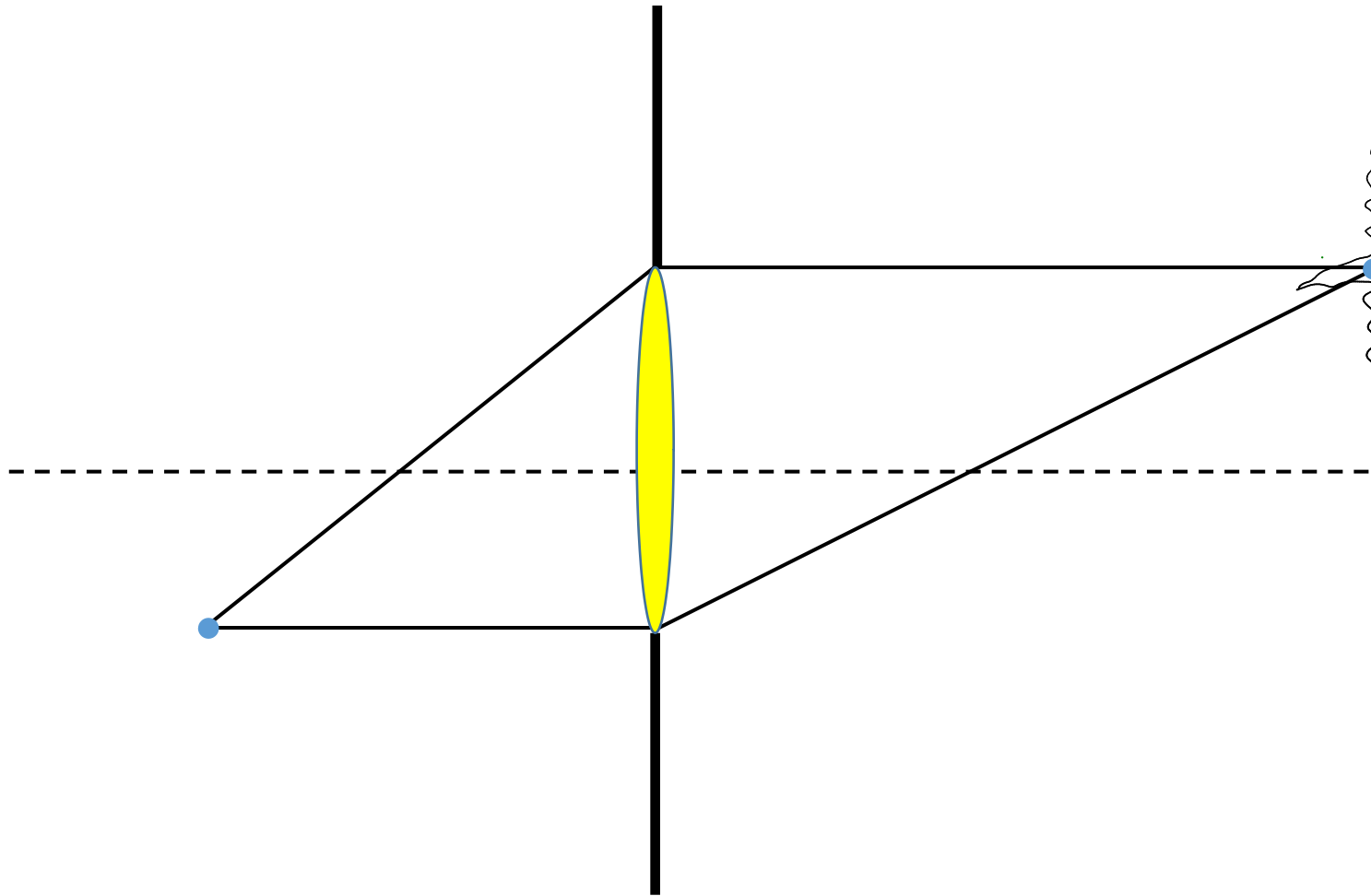


projection model

$PSF \leftrightarrow$  source size



# Lens camera model

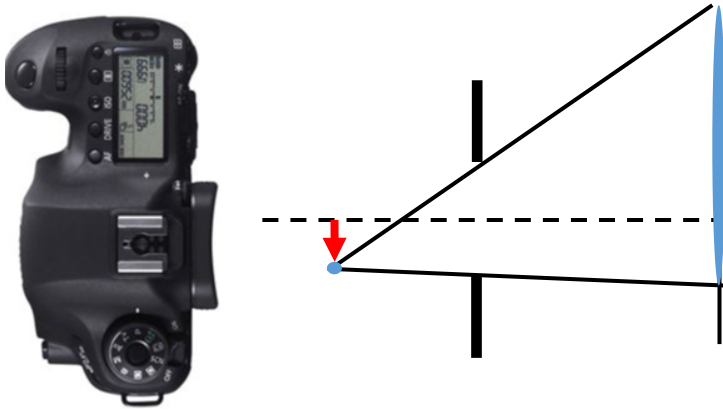


PSF:  
diffraction

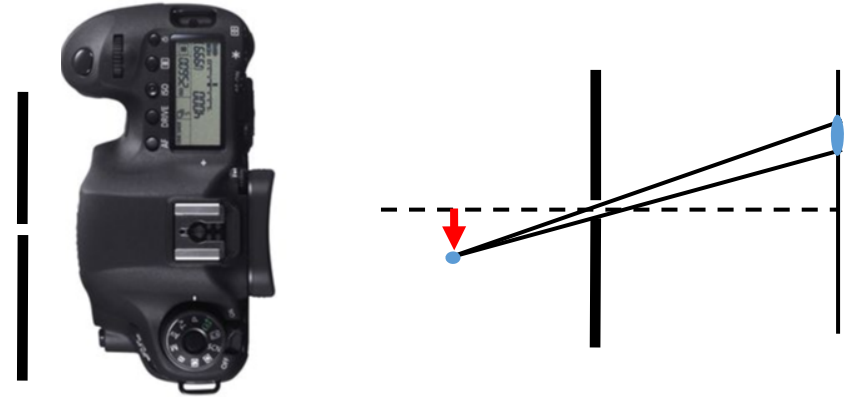
↖ PSF  
related to  
F.T. of  
aperture

# Lens camera model

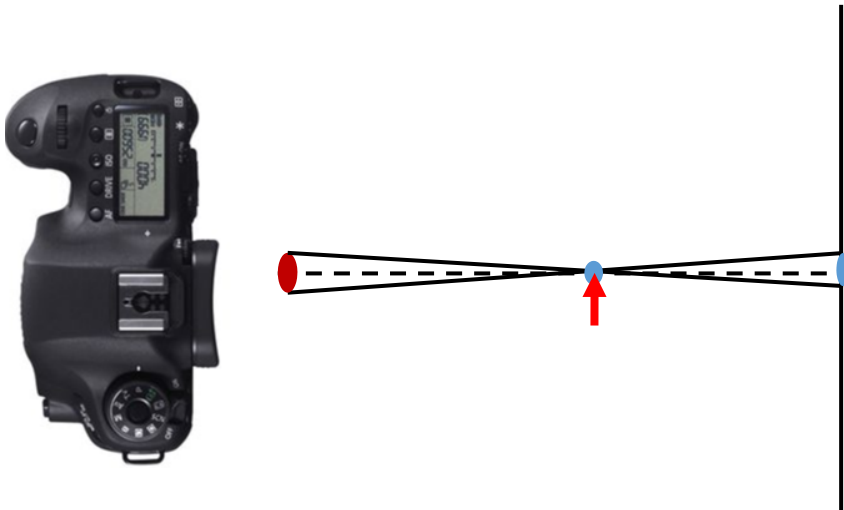
lensless model



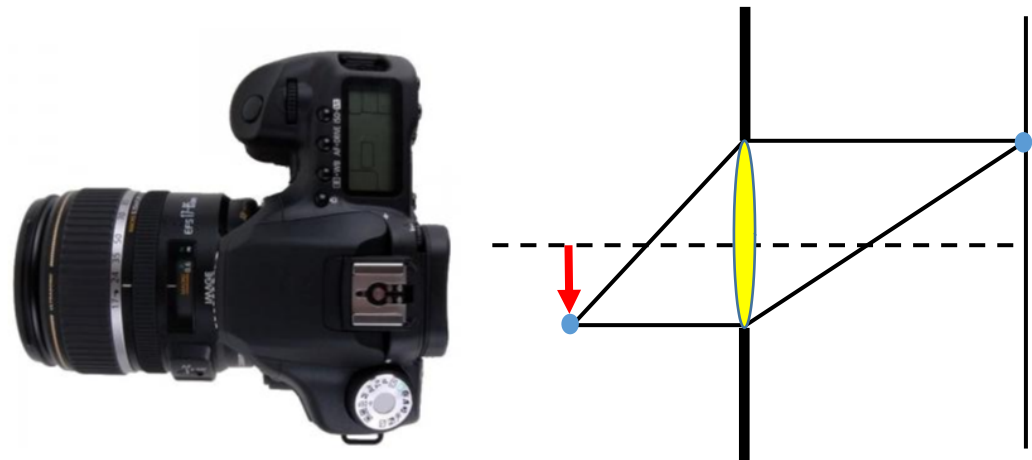
pinhole camera model



projection model



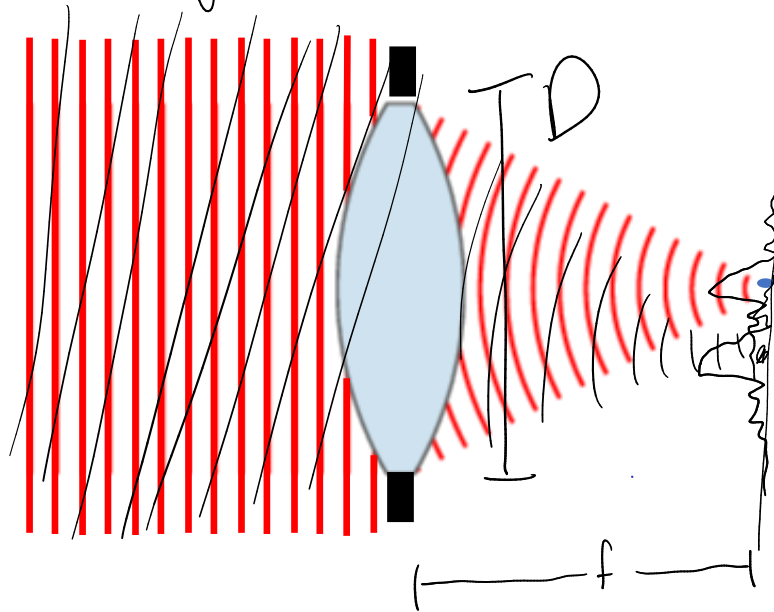
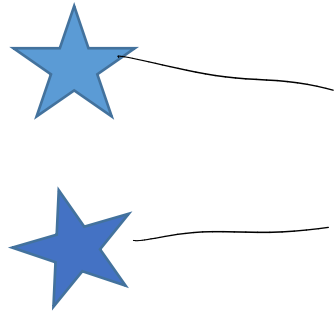
lens camera model



# Diffraction-limited imaging systems

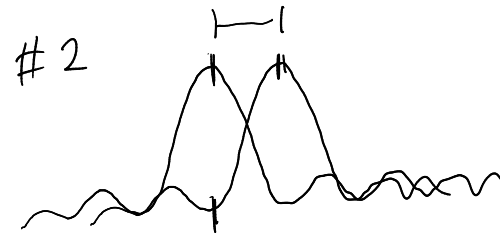
- Rayleigh criterion

*e.g. astronomy*

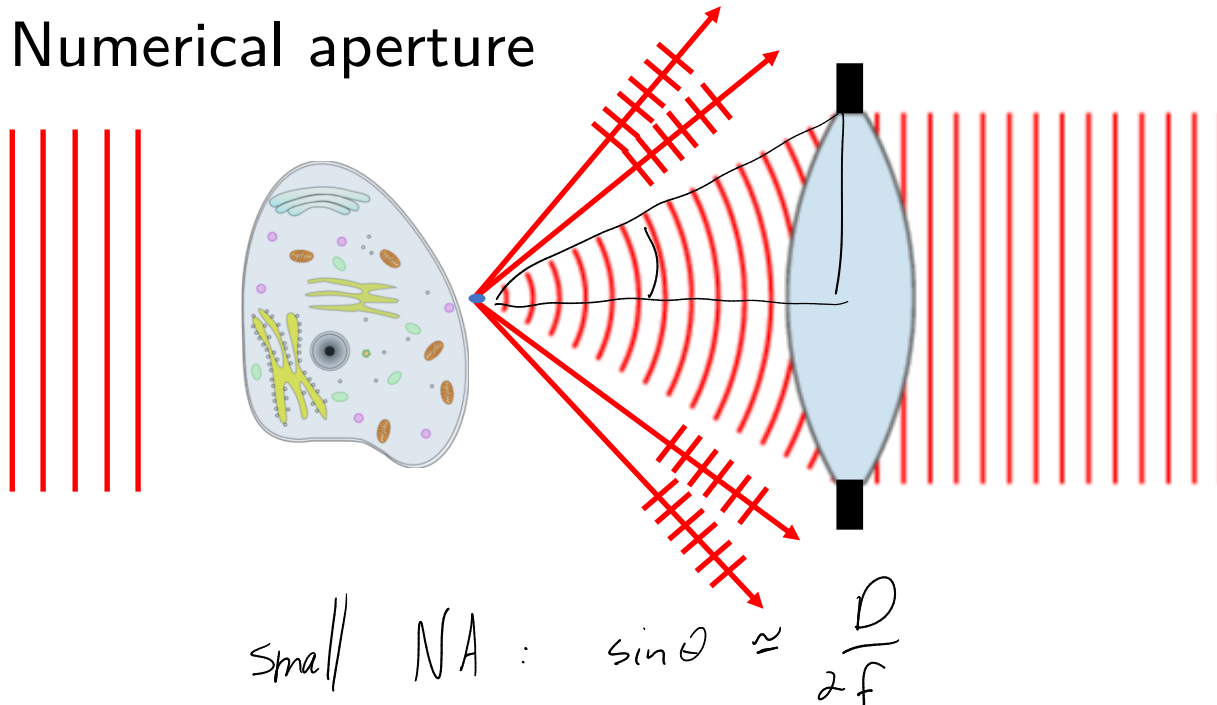


PSF: Airy "disc"

Rayleigh criterion:  
two points are resolvable  
if maximum of #1  
aligned with 1<sup>st</sup> minimum  
of #2



- Numerical aperture



Small NA:  $\sin \theta \approx \frac{D}{2f}$

$$d_{\min} = \frac{f \cdot \lambda \cdot 1.22}{D}$$

NA: numerical aperture

$$NA = n \cdot \sin \theta$$

$$d_{\min} = 1.22 \frac{\lambda}{2NA}$$

# Scanning systems

## Transmission

- **Scanning Transmission Electron Microscopy**
- **Scanning Transmission X-ray Microscopy**
- ...

## Indirect (reflection, scattering, fluorescence, ...)

- **Laser Scanning Confocal Microscopy**
- **Scanning Electron Microscopy**
- **X-ray Fluorescence Microscopy**
- **PhotoEmission Electron Microscopy**
- ...

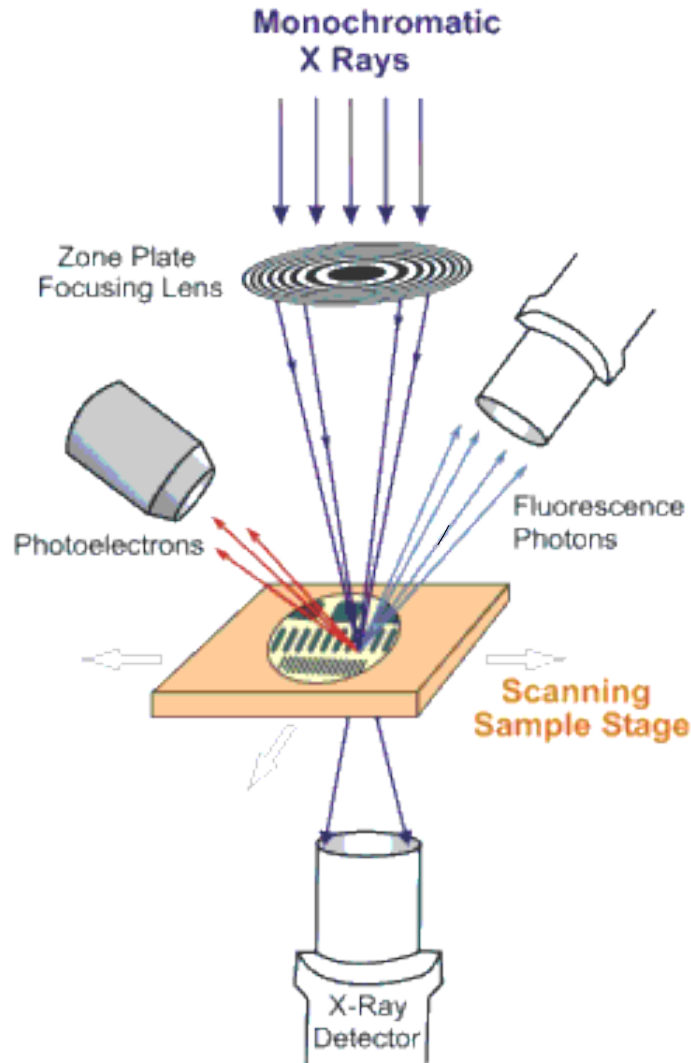
## Physical probe

- **Atomic Force Microscopy**
- **Scanning Tunneling Microscopy**
- ...

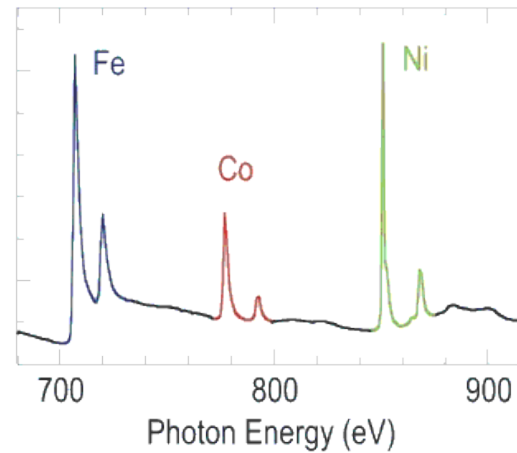


# Scanning transmission X-ray microscopy

Scanning Transmission X-ray Microscopy  
STXM

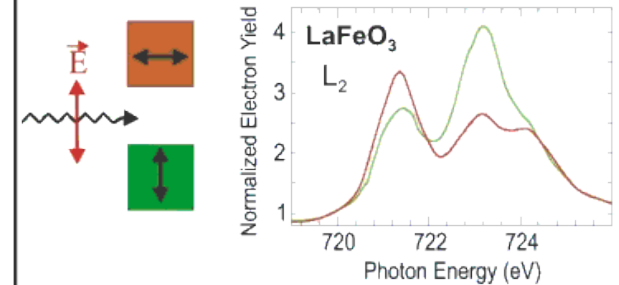


Tune x-ray **energy**  
for elemental specificity

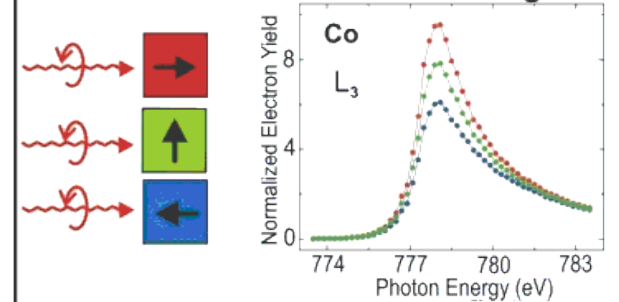


Tune x-ray **polarization**  
for magnetic specificity

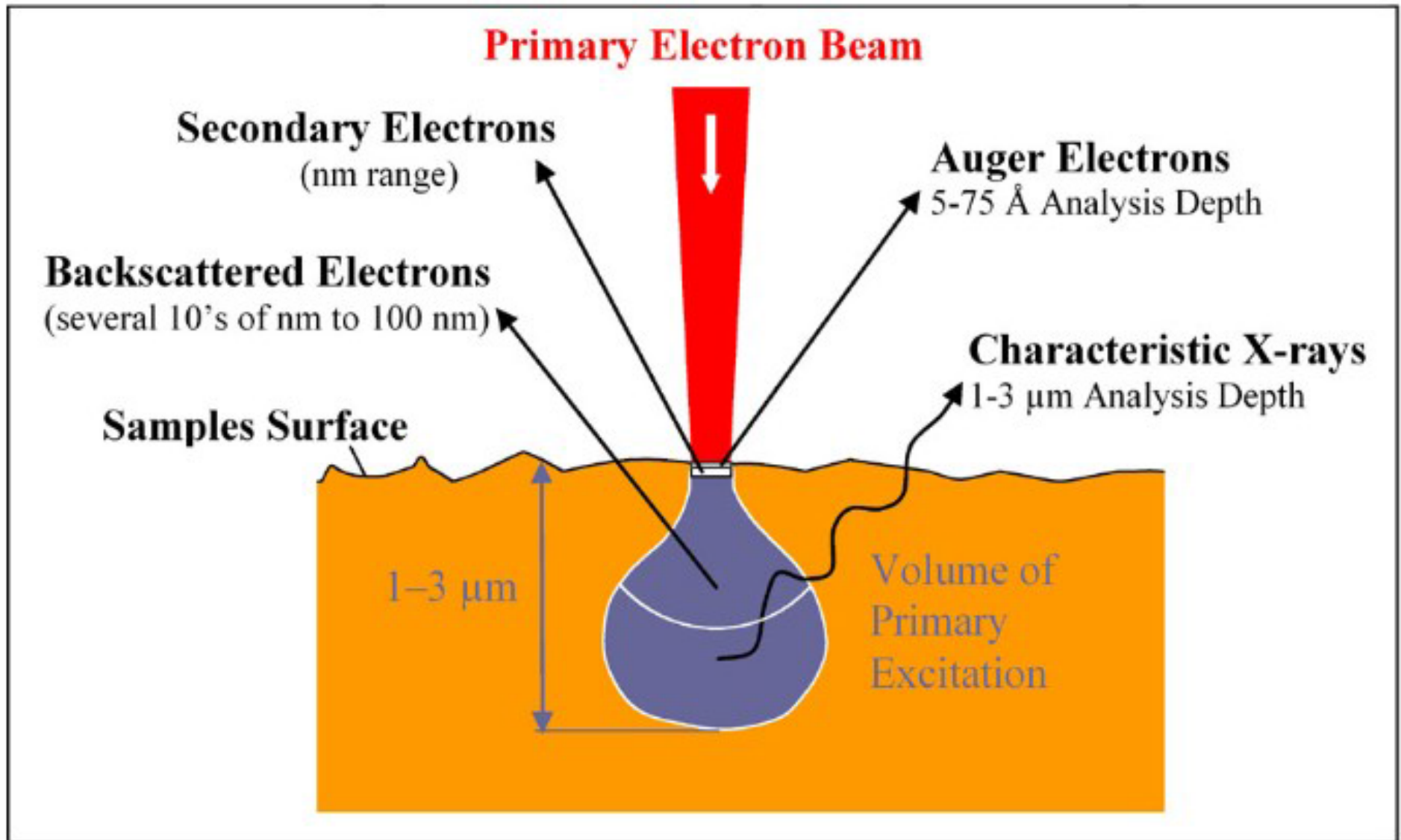
Linear Dichroism - Antiferromagnets



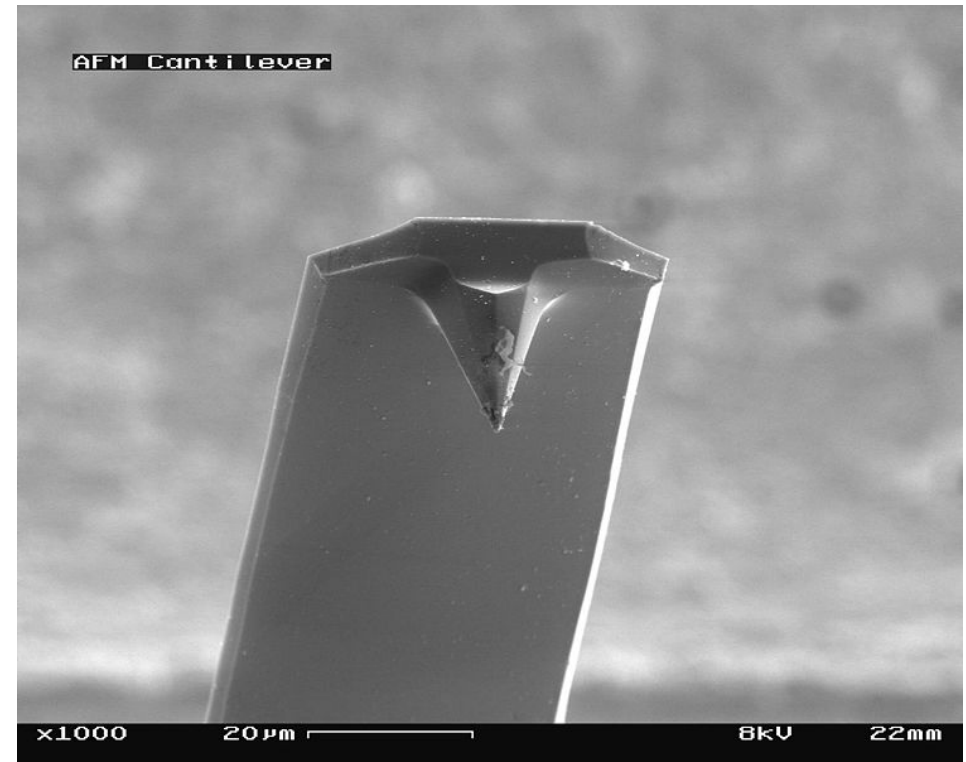
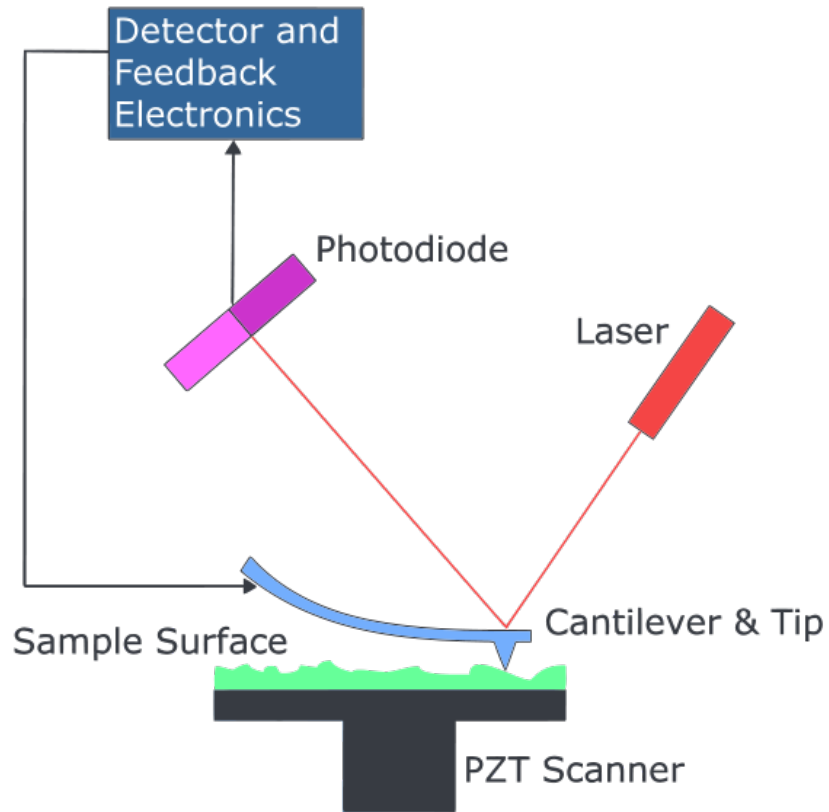
Circular Dichroism - Ferromagnets



# Scanning electron microscopy



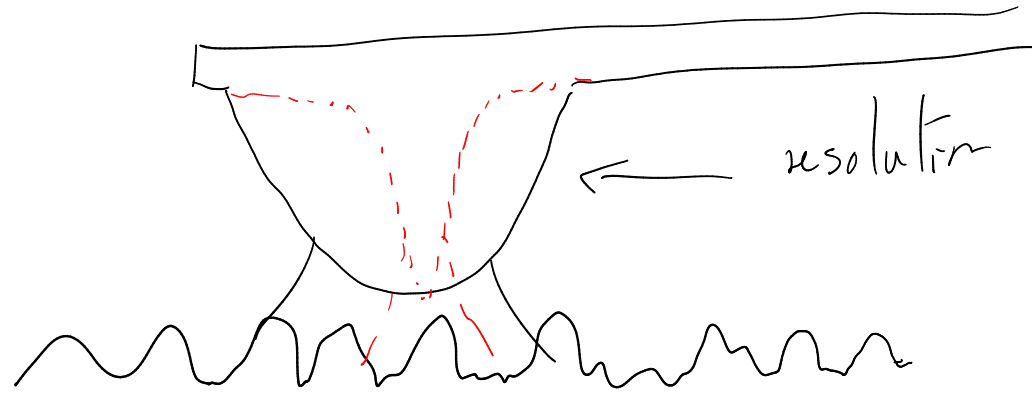
# Atomic force microscopy



# Resolution in scanning systems

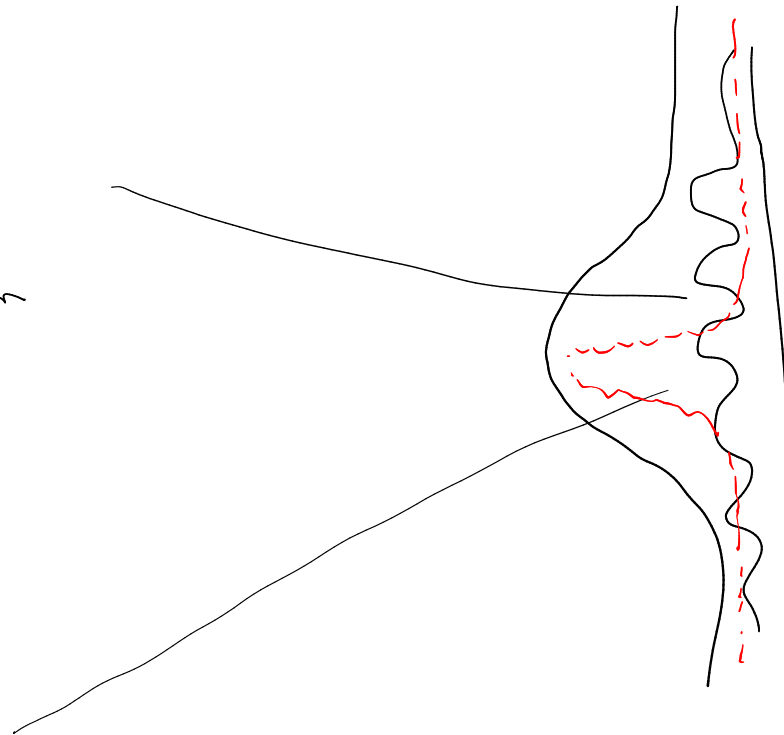
Resolution mainly limited by probe size

e.g. AFM



resolution limited  
by probe size

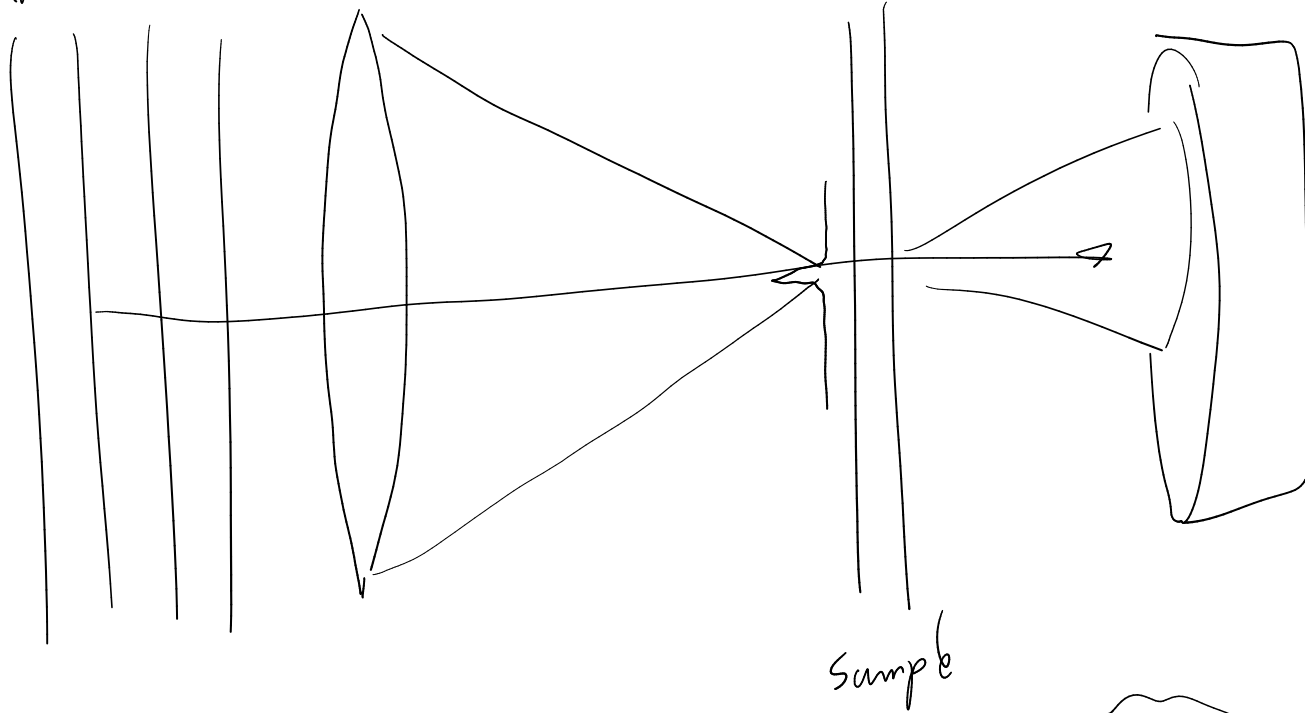
e.g. transmission



# Scanning vs. full field systems

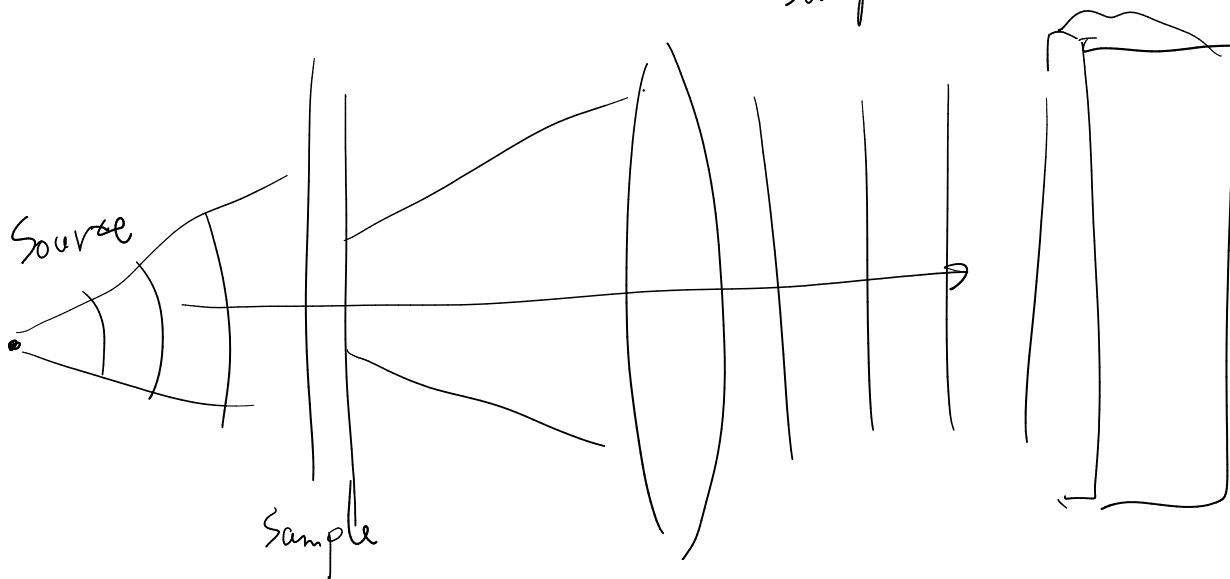
Transmission probe: the reciprocity theorem

plane wave



detector

achievable  
resolution is the  
same



detector