

11 Novembre

Lemma X spazio di B , $B: X \times X \rightarrow X$

Sia $\alpha < \frac{1}{4|B|}$. Allora $\forall x_0 \in D_X(0, \alpha)$

$\exists!$ $\overline{x} \in D_X(0, 2\alpha)$ t.c. $x = x_0 + B(x, x)$

$$u = e^{\nu t \Delta} u_0 + B(u, u) \quad *$$

$$\begin{cases} \partial_t u - \nu \Delta u = Q(u, u) \\ u|_{t=0} = u_0 \end{cases}$$

$$\begin{cases} \partial_t B(u, v) - \nu \Delta B(u, v) = Q(u, v) \\ B(u, v)|_{t=0} = 0 \end{cases}$$

$$X = L^4([0, T], \dot{H}^{\frac{d-1}{2}}(\mathbb{R}^d, \mathbb{R}^d))$$

Teor Sia $u_0 \in H^{\frac{d}{2}-1}(\mathbb{R}^d, \mathbb{R}^d)$. Allora $\exists T > 0$

ed una soluz di $u = e^{v t \Delta} u_0 + B(u, u)$ in $L^4([0, T], H^{\frac{d-1}{2}})$. La soluzione è unica.

Soddisfer

$$u \in C^0([0, T], H^{\frac{d}{2}-1}), \quad \nabla u \in L^2([0, T], H^{\frac{d}{2}-1})$$

Sia T_{u_0} il tempo di vita.

1) $\exists c > 0 \quad t \leq \dots \quad \|u_0\|_{H^{\frac{d}{2}-1}} \leq c v \Rightarrow T_{u_0} = \infty$

2) Se $T_{u_0} < +\infty$ allora

$$\int_0^{T_{u_0}} \|u(t)\|_{H^{\frac{d-1}{2}}}^4 dt = +\infty$$

3) Se $T_{u_0} < +\infty$ allora

$$\int_0^{T_{u_0}} \|\nabla u\|_{H^{\frac{d}{2}-1}}^2 dt = +\infty$$

Infine se u e v sono due soluzioni allora

$$\begin{aligned} \|u(t) - v(t)\|_{H^{\frac{d}{2}-1}}^2 + v \int_0^t \|\nabla(u-v)\|_{H^{\frac{d}{2}-1}}^2 dt' &\leq \\ &\leq \|u_0 - v_0\|_{H^{\frac{d}{2}-1}}^2 + C v^{-3} \int_0^t (\|u\|_{H^{\frac{d-1}{2}}}^4 + \|v\|_{H^{\frac{d-1}{2}}}^4) dt' \end{aligned}$$

dove $C = C_d$.

Esser viziato

$$u \rightarrow u_\lambda(t, x) = \lambda u(\lambda^2 t, \lambda x)$$

u e u_λ sono insieme soluzioni di NS.

Lo scaling $f(x) \rightarrow f_\lambda(x) = \lambda f(\lambda x)$
preserva la norma di $\dot{H}^{\frac{d}{2}-1}(\mathbb{R}^d)$.

Si dice allora che $\dot{H}^{\frac{d}{2}-1}(\mathbb{R}^d)$ è
critico per NS. Non esiste una funzione

$$T(\cdot) : [0, +\infty) \rightarrow (0, +\infty] \quad t \in \tau.$$

$$T_{u_0} \geq T(|u_0|_{\dot{H}^{\frac{d}{2}-1}})$$

Sostituiamo più avanti $\dot{H}^{\frac{d}{2}-1}(\mathbb{R}^d)$ con
 $L^d(\mathbb{R}^d)$. Cioè sostituiamo $\dot{H}^{\frac{1}{2}}(\mathbb{R}^3)$
con $L^3(\mathbb{R}^3)$

$$\dot{H}^{\frac{1}{2}}(\mathbb{R}^3) \subset L^3(\mathbb{R}^3) \subset \dots$$

$$f \rightarrow f_\lambda(x) = \lambda f(\lambda x)$$

se $u_0 \in H^1$ $\exists T = T(u_0|_{H^1}) \quad t \in \tau.$
 $u \in L^\infty([0, \tau], H^1)$

Questo non è vero in $H^{\frac{d}{2}-2}$ perché

se u è definito $[0, T_{u_0})$ con $T_{u_0} < \infty$

con dato iniziale $u_0(x)$, allora

$$u_\lambda(t, x) = \lambda u(\lambda^2 t, \lambda x) \quad \text{questo vale per } \lambda^2 t < T_{u_0}$$

$$\text{in } \left[0, \frac{T_{u_0}}{\lambda^2}\right]$$

$$T_{u_0 \lambda} = \frac{T_{u_0}}{\lambda^2}$$

$$\|u_0\|_{H^{\frac{d}{2}-2}} = \|u_0 \lambda\|_{H^{\frac{d}{2}-2}}$$

Osservazioni Escobar, Seregin e Sverack
che implicano che se $T_{u_0} < \infty$ allora

$$\|u\|_{L^\infty([0, T_{u_0}), H^{\frac{1}{2}})} = +\infty \quad \square$$

$$\|\nabla u\|_{L^2([0, T_{u_0}), H^{\frac{d-1}{2}})} = +\infty \quad \times$$

$$\|u\|_{L^k([0, T_{u_0}), H^{\frac{d-1}{2}})} = +\infty \quad \times$$

$$\partial_t u - \nu \Delta u = Q(u, u)$$

$$u \in L^k([0, T], \dot{H}^{\frac{d-1}{2}})$$

$$u \in C^0([0, T], H^1)$$

$$Q(u, v) \in L^2([0, T], H^{\frac{d-1}{2}})$$

Lernzettel

$$|Q(u, v)|_{\dot{H}^{\frac{d-2}{2}}} \leq C \|u\|_{\dot{H}^{\frac{d-1}{2}}} \|v\|_{\dot{H}^{\frac{d-1}{2}}}$$

Dim

$$|Q(u, v)|_{\dot{H}^{-\frac{1}{2}}} \leq \|\nabla u \cdot v\|_{\dot{H}^{-\frac{1}{2}}} + \|u \nabla v\|_{\dot{H}^{-\frac{1}{2}}} \leq$$

$$\dot{H}^{\frac{1}{2}}(\mathbb{R}^3) \subset L^3(\mathbb{R}^3)$$

$$\frac{1}{3} = \frac{1}{2} - \frac{\frac{1}{2}}{3} = \frac{1}{2} - \frac{1}{6}$$

$$L^{\frac{3}{2}}(\mathbb{R}^3) \subset \dot{H}^{-\frac{1}{2}}(\mathbb{R}^3)$$

$$\lesssim \|\nabla u \cdot v\|_{L^{\frac{3}{2}}} + \|u \nabla v\|_{L^{\frac{3}{2}}}$$

$$\frac{2}{3} = \frac{1}{2} + \frac{1}{6}$$

$$\leq \|\nabla u\|_{L^2} \|v\|_{L^6} + \|u\|_{L^6} \|\nabla v\|_{L^2}$$

$$\dot{H}^1(\mathbb{R}^3) \hookrightarrow L^6(\mathbb{R}^3)$$

$$\frac{1}{6} = \frac{1}{2} - \frac{1}{3}$$

$$\lesssim \|\nabla u\|_{L^2} \|\nabla v\|_{L^2}$$

$$\frac{d-1}{2} = 1 \quad \text{per } d=3.$$

$$Q(u, v) \in L^2([0, T], \dot{H}^{\frac{d-2}{2}})$$

Lemma

$$\|Q(u, v)\|_{L^2([0, T], \dot{H}^{\frac{d-2}{2}})} \leq C \|u\|_{L^4([0, T], \dot{H}^{\frac{d-1}{2}})} \|v\|_{L^4([0, T], \dot{H}^{\frac{d-1}{2}})}$$

Dim

$$\begin{aligned} \| |Q(u, v)|_{\dot{H}^{\frac{d-2}{2}}} \|_{L^2([0, T])} &\lesssim \| |u|_{\dot{H}^{\frac{d-1}{2}}} |v|_{\dot{H}^{\frac{d-1}{2}}} \|_{L^2([0, T])} \\ &\leq \|u\|_{L^4([0, T], \dot{H}^{\frac{d-1}{2}})} \|v\|_{L^4([0, T], \dot{H}^{\frac{d-1}{2}})}. \end{aligned}$$

$$(\partial_t - \delta \Delta) B(u, v) = Q(u, v)$$

$$B(u, v)|_{t=0} \approx 0$$

$$B(u, v) \in \underbrace{C^0([0, T], \dot{H}^{\frac{d-1}{2}})}_{p=\infty} \cap \underbrace{L^2([0, T], \dot{H}^{\frac{d-1}{2}})}_{p=2}$$

$$L^p([0, T], \dot{H}^{\frac{d-1}{2} + \frac{2}{p}})$$

$$L^4([0, T], \dot{H}^{\frac{d-1}{2}})$$

$$\|B(u, v)\|_{L^4([0, T], \dot{H}^{\frac{d-1}{2}})} \leq \frac{C'}{\gamma^{\frac{3}{4}}} \|Q(u, v)\|_{L^2([0, T], \dot{H}^{\frac{d-2}{2}})}$$

$$\leq \frac{C}{\gamma^{\frac{3}{4}}} \|u\|_{L^4([0, T], \dot{H}^{\frac{d-1}{2}})} \|v\|_{L^4([0, T], \dot{H}^{\frac{d-1}{2}})}$$

$$B: X_T \times X_T \rightarrow X_T \quad X_T = \downarrow$$

$$|B| \leq \frac{C}{\gamma^{\frac{3}{4}}}$$

$$u = e^{\gamma t \Delta} u_0 + B(u, u)$$

$$\text{Se } \| e^{\gamma t \Delta} u_0 \|_{L^4([0, T], \dot{H}^{\frac{d-1}{2}})} < \frac{\gamma^{\frac{3}{4}}}{4C} \leq \frac{1}{4|B|}$$

allora \exists ~~una~~ soluzione $u \in L^4([0, T], \dot{H}^{\frac{d-1}{2}})$

con norma $\|u\|_{L^4([0, T], \dot{H}^{\frac{d-1}{2}})} < 2 \| e^{\gamma t \Delta} u_0 \|_{L^4([0, T], \dot{H}^{\frac{d-1}{2}})} < \frac{\gamma^{\frac{3}{4}}}{2C}$

1) $\|u_0\|_{\dot{H}^{\frac{d}{2}-1}}$ piccolo

$$\| e^{\gamma t \Delta} u_0 \|_{L^4([0, T], \dot{H}^{\frac{d-1}{2}})} \leq \| e^{\gamma t \Delta} u_0 \|_{L^4(\mathbb{R}_+, \dot{H}^{\frac{d-1}{2}})} \leq \gamma^{-\frac{1}{4}} \|u_0\|_{\dot{H}^{\frac{d}{2}-1}} < \frac{\gamma^{\frac{3}{4}}}{4C}, \text{ cioè } u$$

$$\|u_0\|_{\dot{H}^{\frac{d}{2}-1}} < \frac{\gamma}{4C}$$

allora $\| e^{\gamma t \Delta} u_0 \|_{L^4(\mathbb{R}_+, \dot{H}^{\frac{d-1}{2}})} < \frac{\gamma^{\frac{3}{4}}}{4C}$

$\Rightarrow u \in L^4([0, T], \dot{H}^{\frac{d-1}{2}}) \quad \forall T$
 $\boxed{u \in L^4(\mathbb{R}_+, \dot{H}^{\frac{d-1}{2}})}$

2) u_0 generico $\in \dot{H}^{\frac{d-1}{2}}$

$$\left(|e^{\gamma \Delta t} u_0|_{L^4([0, T], \dot{H}^{\frac{d-1}{2}})} \right) \xrightarrow{T \rightarrow 0} 0$$

e prendi $\exists T$ sufficientemente piccolo t.c.

$$\frac{\sqrt{\frac{3}{4}}}{\epsilon C} \Rightarrow u \in L^4([0, T], \dot{H}^{\frac{d-1}{2}})$$

$$u_0 = P_\rho u_0 + (1 - P_\rho) u_0$$

$$P_\rho u_0 = \chi_{[\rho, \infty)}(|z|) \hat{u}_0(z)$$

$\exists \rho$ sufficientemente grande t.c.

$$\| (1 - P_\rho) u_0 \|_{\dot{H}^{\frac{d-1}{2}}} < \frac{\nu}{8C}$$

$$\left(\int_{|z| > \rho} |z|^{d-2} |\hat{u}_0(z)|^2 \right)^{\frac{1}{2}} < \frac{\nu}{8C}$$

$$\| e^{\gamma \Delta t} u_0 \|_{L^4([0, T], \dot{H}^{\frac{d-1}{2}})} \leq \| e^{\gamma \Delta t} P_\rho u_0 \|_{L^4([0, T], \dot{H}^{\frac{d-1}{2}})} + \| e^{\gamma \Delta t} (1 - P_\rho) u_0 \|_{L^4([0, T], \dot{H}^{\frac{d-1}{2}})}$$

$$\| e^{\gamma \Delta t} \chi_{[\rho, \infty)}(|z|) u_0 \|_{L^4([0, T], \dot{H}^{\frac{d-1}{2}})} < \frac{\nu \sqrt{\frac{3}{4}}}{8C}$$

$$\leq \| e^{\gamma \Delta t} \chi_{[\rho, \infty)}(|z|) \sqrt{\rho} \frac{(|z|)^{\frac{d-1}{2}}}{\sqrt{\rho}} u_0 \|_{L^4([0, T], \dot{H}^{\frac{d-1}{2}})}$$

$$\leq \sqrt{\rho} \| e^{\gamma \Delta t} \chi_{[\rho, \infty)}(|z|) u_0 \|_{L^4([0, T], \dot{H}^{\frac{d-1}{2}})}$$

$$\leq \sqrt{\rho} T^{\frac{1}{4}} \| e^{\gamma \Delta t} u_0 \|_{L^\infty(\mathbb{R}_+, \dot{H}^{\frac{d-1}{2}})} \leq \sqrt{\rho} T^{\frac{1}{4}} \| u_0 \|_{\dot{H}^{\frac{d-1}{2}}}$$

$$\left[e^{\nu t \Delta} \chi_{[0, T]}(\sqrt{-\Delta}) u_0 \right]_{L^4(\mathbb{R}^d), \dot{H}^{\frac{d-1}{2}}} \leq \left(\nu^{\frac{1}{2}} T^{\frac{1}{4}} \|u_0\|_{\dot{H}^{\frac{d}{2}-1}} \right)^{\frac{3}{4}} \frac{\nu^{\frac{3}{4}}}{8C}$$

$$\left[e^{\nu t \Delta} \chi_{[T, \infty)}(\sqrt{-\Delta}) u_0 \right]_{L^4(\mathbb{R}^d), \dot{H}^{\frac{d-1}{2}}} \ll \frac{\nu^{\frac{3}{4}}}{8C}$$

Coni si vuole

$$\left[e^{\nu t \Delta} u_0 \right]_{L^4(\mathbb{R}^d), \dot{H}^{\frac{d-1}{2}}} \ll \frac{\nu^{\frac{3}{4}}}{8C}$$

$$\Rightarrow u \in L^4(\mathbb{R}^d), \dot{H}^{\frac{d-1}{2}}$$

$$T < \left(\frac{\nu^{\frac{3}{4}}}{8 \nu^{\frac{1}{2}} C \|u_0\|_{\dot{H}^{\frac{d}{2}-1}}} \right)^4$$

$$\left\{ \begin{array}{l} (\partial_t - \nu \Delta) u = Q(u, u) \\ u|_{t=0} = u_0 \end{array} \right.$$

$$u \in C^0([0, T], \dot{H}^{\frac{d}{2}-1})$$

$$\nabla u \in L^2([0, T], \dot{H}^{\frac{d}{2}-1})$$

Show u, v given due whyon e is $w = u - v$.

$$\begin{cases} \partial_t w - \nu \Delta w = Q(u, u) - Q(v, v) = Q(\overset{w}{u-v}, u+v) \\ w(0) = u_0 - v_0 \end{cases}$$

$$\Delta w = |w(t)|_{\dot{H}^{\frac{d-1}{2}}}^2 + 2\nu \int_0^t |\nabla w|_{\dot{H}^{\frac{d-1}{2}}}^2 dt' =$$

$$= |w(0)|_{\dot{H}^{\frac{d-1}{2}}}^2 + 2 \int_0^t \langle Q(w, u+v), w \rangle_{\dot{H}^{\frac{d-1}{2}}} dt'$$

$$u_0 \left(L^4(\mathbb{R}^d, \dot{H}^{\frac{d-1}{2}}) \right) \langle Q(u, b), c \rangle_{\dot{H}^{\frac{d-1}{2}}} \leq C \|u\|_{\dot{H}^{\frac{d-1}{2}}} \|b\|_{\dot{H}^{\frac{d-1}{2}}} \|c\|_{\dot{H}^{\frac{d-1}{2}}}$$

$$\langle Q(u, b), c \rangle_{\dot{H}^{\frac{d-1}{2}}} \leq \|Q(u, b)\|_{\dot{H}^{\frac{d-2}{2}}} \|c\|_{\dot{H}^{\frac{d-1}{2}}}$$

$$\leq C \|u\|_{\dot{H}^{\frac{d-2}{2}}} \|b\|_{\dot{H}^{\frac{d-1}{2}}} \|c\|_{\dot{H}^{\frac{d-1}{2}}}$$

$$\Delta_w = |w(t)|^2_{\dot{H}^{\frac{d-1}{2}}} + 2\nu \int_0^t |\nabla w|^2_{\dot{H}^{\frac{d-1}{2}}} dt' =$$

$$= |w_0|^2_{\dot{H}^{\frac{d-1}{2}}} + 2 \int_0^t \langle \mathcal{Q}(w, u+v), w \rangle_{\dot{H}^{\frac{d-1}{2}}} dt'$$

$$\Delta_w \leq |w_0|^2_{\dot{H}^{\frac{d-1}{2}}} + 2C \int_0^t |w|_{\dot{H}^{\frac{d-1}{2}}} |\nabla w|_{\dot{H}^{\frac{d-1}{2}}} N dt'$$

$$N = |u|_{\dot{H}^{\frac{d-1}{2}}} + |v|_{\dot{H}^{\frac{d-1}{2}}}$$

$$|w|_{\dot{H}^{\frac{d-1}{2}}} \leq |w|_{\dot{H}^{\frac{d-1}{2}}}^{\frac{1}{2}} |\nabla w|_{\dot{H}^{\frac{d-1}{2}}}^{\frac{1}{2}}$$

$$\Delta_w \leq |w_0|^2_{\dot{H}^{\frac{d-1}{2}}} + 2C \int_0^t \left(|w|_{\dot{H}^{\frac{d-1}{2}}}^{\frac{1}{2}} N \right) \left(|\nabla w|_{\dot{H}^{\frac{d-1}{2}}}^{\frac{1}{2}} \right)$$

$$ab \leq \frac{a^4}{4} + \frac{3}{4} b^{\frac{4}{3}}$$

$$\Delta_w \leq |w_0|^2_{\dot{H}^{\frac{d-1}{2}}} + C_\nu \int_0^t |w|_{\dot{H}^{\frac{d-1}{2}}}^2 N^4 + \nu \int_0^t |w|_{\dot{H}^{\frac{d-1}{2}}}^2 dt$$

$$|w(t)|_{\dot{H}^{\frac{d-1}{2}}}^2 + \nu \int_0^t |\nabla w|_{\dot{H}^{\frac{d-1}{2}}}^2 dt \leq$$

$X(t)$

$$X(t) \leq |w_0|^2_{\dot{H}^{\frac{d-1}{2}}} + C_\nu \int_0^t N^4 \left(|w|_{\dot{H}^{\frac{d-1}{2}}}^2 \right) dt' \leq X$$

$$X(t) \leq |w_0|^2_{\dot{H}^{\frac{d-1}{2}}} + C_\nu \int_0^t N^4(t') X(t') dt'$$

$$X(t') \leq |w_0|^2_{\dot{H}^{\frac{d-1}{2}}} \left(e^{C_\nu \int_0^t N^4(t'') dt''} \right)$$

$$N^4 \sim |u|_{\dot{H}^{\frac{d-1}{2}}}^4 + |v|_{\dot{H}^{\frac{d-1}{2}}}^4$$