

Δ_{gh}

$$\Delta_{gh} = -D^2 = -\partial^2 + \Delta_1 + \Delta_2$$

$$[A^\mu, [A_\mu, \cdot]] = A^{\mu\alpha} A_\mu^\beta \cdot$$

$$\Delta_1 = -i \partial^\mu A_\mu - i A_\mu \partial^\mu = -i \{ \partial^\mu, A_\mu \} = -i \{ \partial^\mu, A_\mu^a \} T_{Adj}^a \cdot T_{Adj}^a T_{Adj}^b$$

$$\text{Tr log } \Delta_{gh} = \text{Tr log } \underbrace{(-\partial^2 + \Delta_1 + \Delta_2)}_{(-\partial^2)(1 + (-\partial^2)^{-1}(\Delta_1 + \Delta_2))} = \text{Tr log } (-\partial^2) + \text{Tr log } (1 + (-\partial^2)^{-1}(\Delta_1 + \Delta_2))$$

$$= \text{Tr log } (-\partial^2) + \text{Tr} [(-\partial^2)^{-1}(\Delta_1 + \Delta_2)] - \frac{1}{2} \text{Tr} [((-\partial^2)^{-1}(\Delta_1 + \Delta_2))^2] + \dots$$

termini cost
che trascuriamo

lineare in A_μ
e $\text{tr } t^a = 0$

termini di
contenuto potenze di
 A^μ superiori a 2
(Infiniti vengono cancellati
da parti cubica e
quartica di Str)

$|x\rangle \quad |k\rangle$

$$f(x) = \langle x | f \rangle$$

$$\int d^d x |x\rangle \langle x| = 1$$

$$\int \frac{d^d k}{(2\pi)^d} |k\rangle \langle k| = 1$$

$$\langle x | k \rangle = e^{ikx}$$

$$\rightarrow \langle x | (-\partial^2)^{-1} | y \rangle = \frac{1}{(2\pi)^d} \int d^d k \frac{1}{k^2} e^{ik(x-y)}$$

$$\text{Tr } M = \sum_i e_i \cdot \text{Me}_i$$

$$\text{Tr} [(-\partial^2)^{-1} \Delta_2] = \text{Tr} [(-\partial^2)^{-1} A^{\mu\alpha} t_{Ad}^a A_\mu^b t_{Ad}^b] =$$

traccia sulle funz. e sullo sp. di rep. Adj

$$= \int d^d x \langle x | \text{tr} [(-\partial^2)^{-1} A^{\mu\alpha} t_{Ad}^a A_\mu^b t_{Ad}^b] | x \rangle =$$

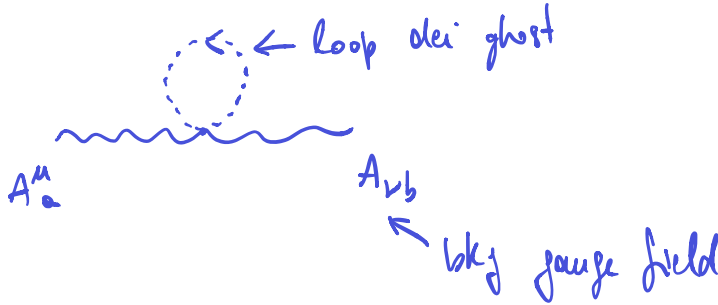
traccia su sp. rep. Adj

$$= \int d^d x d^d y \langle x | (-\partial^2)^{-1} | y \rangle \langle y | A^{\mu\alpha} A_\mu^b | x \rangle \text{tr} (t_{Ad}^a t_{Ad}^b)$$

$$= \int d^d x d^d y \frac{d^d k}{(2\pi)^d} \frac{1}{k^2} e^{ik(x-y)} \text{Tr}(t_{Ad}^a t_{Ad}^b) A^{\mu a}(y) A_\mu^b(x) \delta(y-x)$$

$$= \int d^d x A^{\mu a}(x) A_\mu^b(x) \text{Tr}(t_{Ad}^a t_{Ad}^b) \int \frac{d^d k}{(2\pi)^d} \frac{1}{k^2} \leftarrow \text{DIVERGENZA in } d=4$$

termine in $\text{Sc}_{eff}(A)$ che contribuisce alle funz. a 2pt. (tree-level)



$$\text{Tr}(t_{Ad}^a t_{Ad}^b) = C(\text{Adj}) \delta^{ab} = c_2(G) \delta^{ab}$$

$$\Delta_1 = -i \partial^\mu A_\mu - i A_\mu \partial^\mu$$

$$\begin{aligned} -\frac{1}{2} \text{Tr}((- \partial^2)^{-1} \Delta_1 (- \partial^2)^{-1} \Delta_1) &= -\frac{1}{2} \int d^d y \langle y | \text{tr}((- \partial^2)^{-1} \Delta_1 (- \partial^2)^{-1} \Delta_1) | y \rangle = \\ &= +\frac{1}{2} \int d^d y \langle y | \left[\text{Tr}((- \partial^2)^{-1} \partial^\mu A_\mu^a t_{Ad}^a (- \partial^2)^{-1} \partial^\nu A_\nu^b t_{Ad}^b) + \right. \\ &\quad + \text{Tr}((- \partial^2)^{-1} A_\mu^a t_{Ad}^a \partial^\mu (- \partial^2)^{-1} \partial^\nu A_\nu^b t_{Ad}^b) + \\ &\quad + \text{Tr}((- \partial^2)^{-1} \partial^\mu A_\mu^a t_{Ad}^a (- \partial^2)^{-1} A_\nu^b \partial^\nu t_{Ad}^b) + \\ &\quad \left. + \text{Tr}((- \partial^2)^{-1} A_\mu^a t_{Ad}^a \partial^\mu (- \partial^2)^{-1} A_\nu^b t_{Ad}^b \partial^\nu) \right] | y \rangle \end{aligned}$$

Consideriamo il primo termine

$$\langle y | h \rangle = e^{iky}$$

$$1^{st} \text{ term} = \frac{1}{2} \int d^d x d^d y \text{Tr} \left[A_\mu^a(x) t_{Ad}^a A_\nu^b(y) t_{Ad}^b \right]$$

$\int \frac{d^d k}{(2\pi)^d} |k| < k |$ k' - complete

$$((-\partial^2)^{-1}) \uparrow i \frac{k^\mu}{k^2}$$

$$\begin{aligned} & \langle y | (-\partial^2)^{-1} \partial^\mu |x\rangle \langle x | (-\partial^2)^{-1} \partial^\nu |y\rangle \\ & \frac{1}{-(ik)^2} ik^\mu e^{ik(y-x)} \quad \frac{1}{-(ik')^2} ik'^\nu e^{-ik'(y-x)} \\ & - \frac{k^\mu}{k^2} \frac{k'^\nu}{k'^2} e^{i(k-k')(y-x)} \end{aligned}$$

2nd term =

$$\begin{aligned} & \dots \\ & \langle y | (-\partial^2)^{-1} |x\rangle \cdot \langle x | \partial^\mu (-\partial^2)^{-1} \partial^\nu |y\rangle \\ & - \frac{1}{k^2} \cdot \frac{k^\mu k'^\nu}{k'^2} e^{i(k-k')(y-x)} \end{aligned}$$

3rd term

$$- \frac{k^\mu k^\nu}{k^2} \cdot \frac{1}{k'^2} e^{i(k-k')(y-x)}$$

4th term

$$\begin{aligned} & - \frac{k^\nu}{k^2} \frac{k'^\mu}{k'^2} e^{i(k-k')(y-x)} \\ & \underbrace{\hspace{10em}} \\ & - \frac{e^{i(k-k')(y-x)}}{k^2 k'^2} (k+k')^\mu (k+k')^\nu \end{aligned}$$

$$- \frac{1}{2} \int d^d x d^d y \int \frac{d^d k}{(2\pi)^d} \int \frac{d^d k'}{(2\pi)^d} e^{i(k-k')(y-x)} \frac{(k+k')^\mu (k+k')^\nu}{k^2 k'^2} \text{Tr} (A_\mu^a(x) t_{Ad}^a A_\nu^b(y) t_{Ad}^b)$$

$$\tilde{A}(k) = \int d^d x e^{ikx} A(x) \rightarrow A(x) = \int \frac{d^d p}{(2\pi)^d} e^{-ipx} \tilde{A}(p)$$

$$-\frac{1}{2} \int \underline{d^d x} \underline{d^d y} \frac{d^d k}{(2\pi)^d} \frac{d^d k'}{(2\pi)^d} \frac{d^d p}{(2\pi)^d} \frac{d^d q}{(2\pi)^d} \frac{e^{i(k-k')y}}{\delta(k-k'-q)} \frac{e^{-i(k-k')x}}{\delta(k-k'+p)} \frac{e^{-ipx}}{e^{-iqy}}$$

$$\cdot \frac{(k+k')^\mu (k+k')^\nu}{k^2 k'^2} \tilde{A}_\mu^a(p) \tilde{A}_\nu^b(q) \text{Ar}(t_{Ad}^a t_{Ad}^b)$$

$$\begin{cases} k-k'-q=0 \\ k-k'+p=0 \end{cases} \Rightarrow \begin{cases} q=-p \\ k'=p+k \end{cases}$$

$$= -\frac{1}{2} \int \frac{d^d p}{(2\pi)^d} \text{Ar}(t_{Ad}^a t_{Ad}^b) \tilde{A}_\mu^a(p) \tilde{A}_\nu^b(-p) \int \frac{d^d k}{(2\pi)^d} \frac{(p+2k)^\mu (p+2k)^\nu}{k^2 (k+p)^2}$$

Riscriviamo anche il primo termine nello sp. dei momenti:

$$\text{Tr}((- \partial^2) \Delta_L) = \int \underline{d^d x} \int \frac{d^d p}{(2\pi)^d} \int \frac{d^d q}{(2\pi)^d} \frac{d^d k}{(2\pi)^d} \text{tr}(t_{Ad}^a t_{Ad}^b) \tilde{A}_\mu^a(p) \tilde{A}_\nu^b(q) \frac{1}{k^2} e^{-ipx} e^{-iqx}$$

$$= \int \frac{d^d p}{(2\pi)^d} \text{tr}(t_{Ad}^a t_{Ad}^b) \tilde{A}_\mu^a(p) \tilde{A}_\nu^b(-p) \int \frac{d^d k}{(2\pi)^d} \frac{\delta^{\mu\nu}}{k^2} \leftarrow \text{siamo su sp. Euclideo}$$

$$\text{Tr} \log \Delta_{g^{\mu\nu}} = \int \frac{d^d p}{(2\pi)^d} \text{tr}(t_{Ad}^a t_{Ad}^b) \tilde{A}_\mu^a(p) \tilde{A}_\nu^b(-p) \cdot \left\{ \int \frac{d^d k}{(2\pi)^d} \frac{\delta^{\mu\nu}}{k^2} + \right. \\ \left. -\frac{1}{2} \int \frac{d^d k}{(2\pi)^d} \frac{(p+2k)^\mu (p+2k)^\nu}{k^2 (k+p)^2} \right.$$

Per calcolare $\{ \dots \}$ usiamo gli integrali di Feynman

$$\int \frac{d^d k}{(2\pi)^d} \frac{1}{(k^2 + a^2)^A} = \frac{\Gamma(A - d/2)}{(a^2)^{A-d/2} (4\pi)^{d/2} \Gamma(A)}$$

$$\Rightarrow \lim_{\epsilon \rightarrow 0} \int \frac{d^d k}{(2\pi)^d} \frac{1}{k^2 + \epsilon^2} = \lim_{\epsilon \rightarrow 0} \frac{\Gamma(1 - d/2)}{\epsilon^{1-d/2} (4\pi)^{d/2}} = 0 \quad \text{for } d > 2$$

$A=1 \quad a=\epsilon$

$$\frac{1}{k^2(k+p)^2} = \int_0^1 d\zeta \frac{1}{[k^2(1-\zeta) + (k^2 + 2kp + p^2)\zeta]^2} = \int_0^1 d\zeta \frac{1}{(k^2 + 2\zeta p \cdot k + \zeta p^2)^2}$$

$$\int \frac{d^d k}{(2\pi)^d} \frac{1}{(k^2 + 2kp + b^2)^A} = \frac{\Gamma(A - d/2)}{(b^2 - p^2)^{A-d/2} (4\pi)^{d/2} \Gamma(A)} \quad \Gamma(n+1) = n!$$

$$\Rightarrow \int \frac{d^d k}{(2\pi)^d} \frac{p^\mu p^\nu}{(k^2 + 2\zeta p \cdot k + \zeta p^2)^2} = \frac{p^\mu p^\nu \Gamma(2 - d/2)}{(\zeta(1-\zeta)p^2)^{2-d/2} (4\pi)^{d/2}}$$

$$\int \frac{d^d k}{(2\pi)^d} \frac{k^\mu}{(k^2 + 2kp + b^2)^A} = -\frac{p^\mu \Gamma(A - d/2)}{(b^2 - p^2)^{A-d/2} (4\pi)^{d/2} \Gamma(A)}$$

$$\Rightarrow \int \frac{d^d k}{(2\pi)^d} \frac{2k^\mu p^\nu + 2k^\nu p^\mu}{(k^2 + 2\zeta p \cdot k + \zeta p^2)^2} = -\frac{4\zeta p^\mu p^\nu \Gamma(2 - d/2)}{(\zeta(1-\zeta)p^2)^{2-d/2} (4\pi)^{d/2}}$$

$$\int \frac{d^d k}{(2\pi)^d} \frac{k^\mu k^\nu}{(k^2 + 2kp + b^2)^A} = \frac{1}{(4\pi)^{d/2} \Gamma(A)} \left[\frac{p^\mu p^\nu \Gamma(A - d/2)}{(b^2 - p^2)^{A-d/2}} + \frac{1}{2} \delta^{\mu\nu} \frac{\Gamma(A - 1 - d/2)}{(b^2 - p^2)^{A-1-d/2}} \right]$$

$$\Rightarrow \int \frac{d^d k}{(2\pi)^d} \frac{4k^\mu k^\nu}{(k^2 + 2\zeta p \cdot k + \zeta p^2)^2} = \frac{4\zeta^2 p^\mu p^\nu \Gamma(2 - d/2)}{(\zeta(1-\zeta)p^2)^{2-d/2} (4\pi)^{d/2}} + \frac{1}{2} \frac{4\delta^{\mu\nu} \Gamma(1 - d/2)}{(\zeta(1-\zeta)p^2)^{1-d/2} (4\pi)^{d/2}}$$

$$\Rightarrow -\frac{1}{2} \int \frac{d^d k}{(2\pi)^d} \frac{(p+2k)^\mu (p+2k)^\nu}{k^2 (k+p)^2} = -\frac{1}{2} \int_0^1 d\xi \left\{ \frac{(1-2\xi)^2 p^\mu p^\nu \Gamma(2-d/2)}{(\xi(1-\xi)p^2)^{2-d/2} (4\pi)^{d/2}} + \frac{2 \delta^{\mu\nu} \Gamma(1-d/2)}{(\xi(1-\xi)p^2)^{1-d/2} (4\pi)^{d/2}} \right\}$$

$$4\xi^2 - 4\xi + 1 = (2\xi - 1)^2$$

$$d = 2\omega \quad (\omega \rightarrow 2)$$

espandiamo la
funct. in d attorno
a d=4

$$\Gamma(1-d/2) = \frac{1}{1-d/2} \Gamma(2-d/2)$$

Parte divergente
in d → 4:

Parte divergente in d → 4:

$$\int_0^1 d\xi \frac{(1-2\xi)^2 p^\mu p^\nu}{(4\pi)^2} \frac{1}{2-\omega}$$

$$\int_0^1 d\xi \xi(1-\xi) \frac{\delta^{\mu\nu} p^2}{(4\pi)^2} \left(-\frac{2}{2-\omega} \right)$$

$$\int_0^1 d\xi (1-4\xi+4\xi^2) = 1 - 2 + \frac{4}{3} = \frac{1}{3}$$

$$\int_0^1 \xi(1-\xi) d\xi = \frac{1}{2} - \frac{1}{3} = \frac{1}{6}$$

$$\rightsquigarrow \text{Parte div.} = -\frac{1}{2} \frac{1}{48\pi^2} \frac{1}{2-\omega} (p^\mu p^\nu - p^2 \delta^{\mu\nu})$$

⇓

$$\text{Tr log } \Delta_{gh} \approx \int \frac{d^d p}{(2\pi)^d} \text{tr} (t_{Ad}^a t_{Ad}^b) \tilde{A}_{(p)}^{\mu a} \tilde{A}_{(-p)}^{\nu b} \left(-\frac{1}{2} \frac{1}{3} \frac{1}{16\pi^2} \frac{1}{2-\omega} (p^\mu p^\nu - p^2 \delta^{\mu\nu}) \right) + \dots$$

$$= -\frac{1}{2} \frac{1}{3} \frac{1}{(4\pi)^2} \underbrace{\text{tr} (t_{Ad}^a t_{Ad}^b)}_{\substack{C(\text{Adj}) \delta^{ab} \\ C_2(G)}} \underbrace{\int \frac{d^d p}{(2\pi)^d} \tilde{A}_{(p)}^{\mu a} \tilde{A}_{(-p)}^{\nu b} (p^\mu p^\nu - p^2 \delta^{\mu\nu})}_{g^2 S[A]} \frac{1}{2-\omega}$$

Parte quadratica di azione S:

$$\frac{1}{4g^2} \int (\partial_\mu A_\nu^a - \partial_\nu A_\mu^a) (\partial^\mu A^{\nu a} - \partial^\nu A^{\mu a}) d^d x =$$

$$= \frac{1}{4g^2} \int \frac{d^d x}{(2\pi)^d} \frac{d^d q_1}{(2\pi)^d} \frac{d^d q_2}{(2\pi)^d} (-iq_{1\mu} \tilde{A}_\nu^a(q_1) + iq_{1\nu} \tilde{A}_\mu^a(q_1)) (-iq_2^\mu \tilde{A}^{\nu a}(q_2) + iq_2^\nu \tilde{A}^{\mu a}(q_2)) \cdot e^{-iq_1 x - iq_2 x} \rightarrow \delta(q_1 + q_2) \quad q_1 = -q_2 \equiv p$$

$$= \frac{1}{4g^2} \int \frac{d^d p}{(2\pi)^d} (p_\mu \tilde{A}_\nu^a(p) - p_\nu \tilde{A}_\mu^a(p)) (p^\mu \tilde{A}^{\nu a}(-p) - p^\nu \tilde{A}^{\mu a}(-p)) =$$

$$= -\frac{1}{2g^2} \int \frac{d^d p}{(2\pi)^d} (p^\mu p^\nu - p^2 \delta^{\mu\nu}) \tilde{A}_\mu^a(p) \tilde{A}_\nu^a(-p)$$