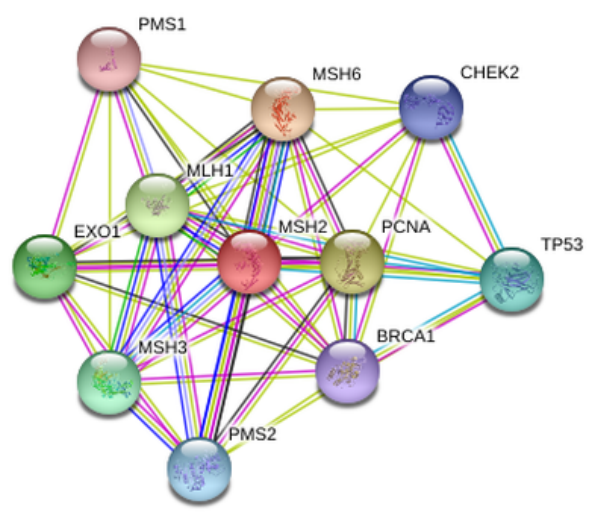
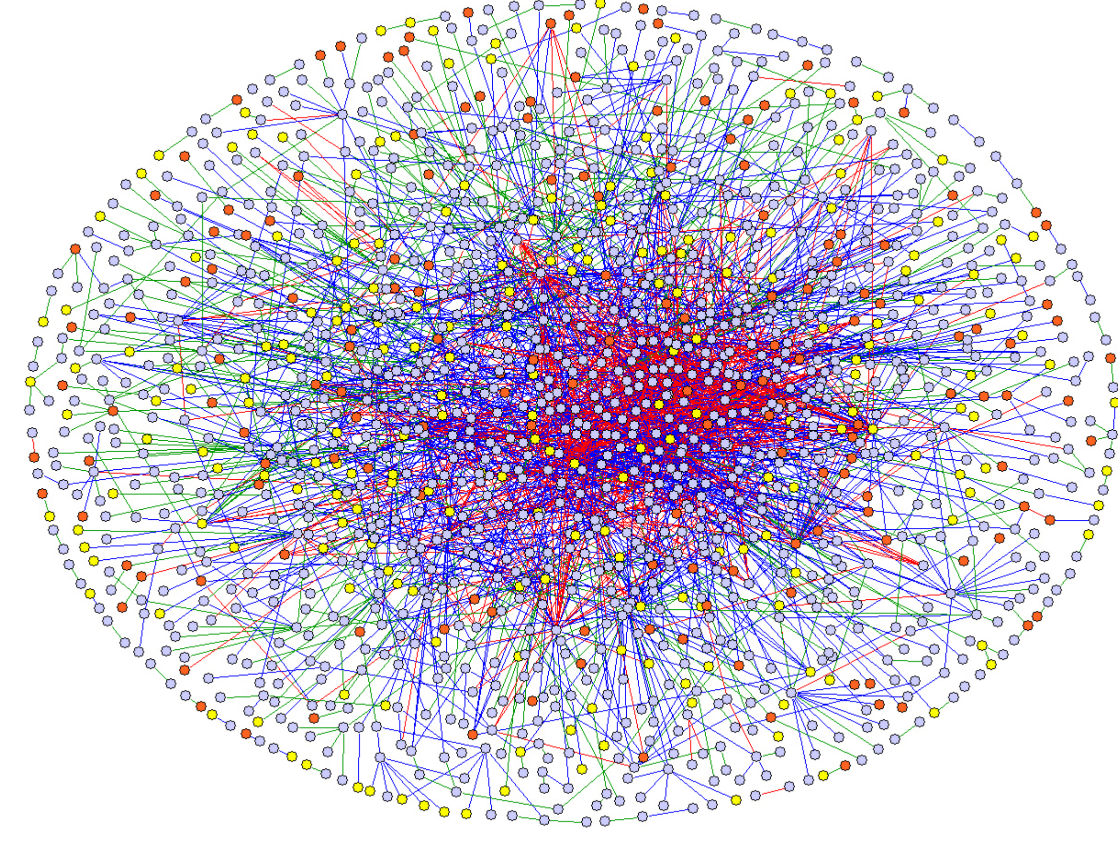
Protein-ligand interactions

BIOMOLECULAR INTERACTIONS

L'espressione e la regolazione delle componenti di una cellula e l'organizzazione dei pathways che sono alla base della sua funzionalità vengono controllati da una complessa rete di interazioni tra biomolecole



…quanto complessa??

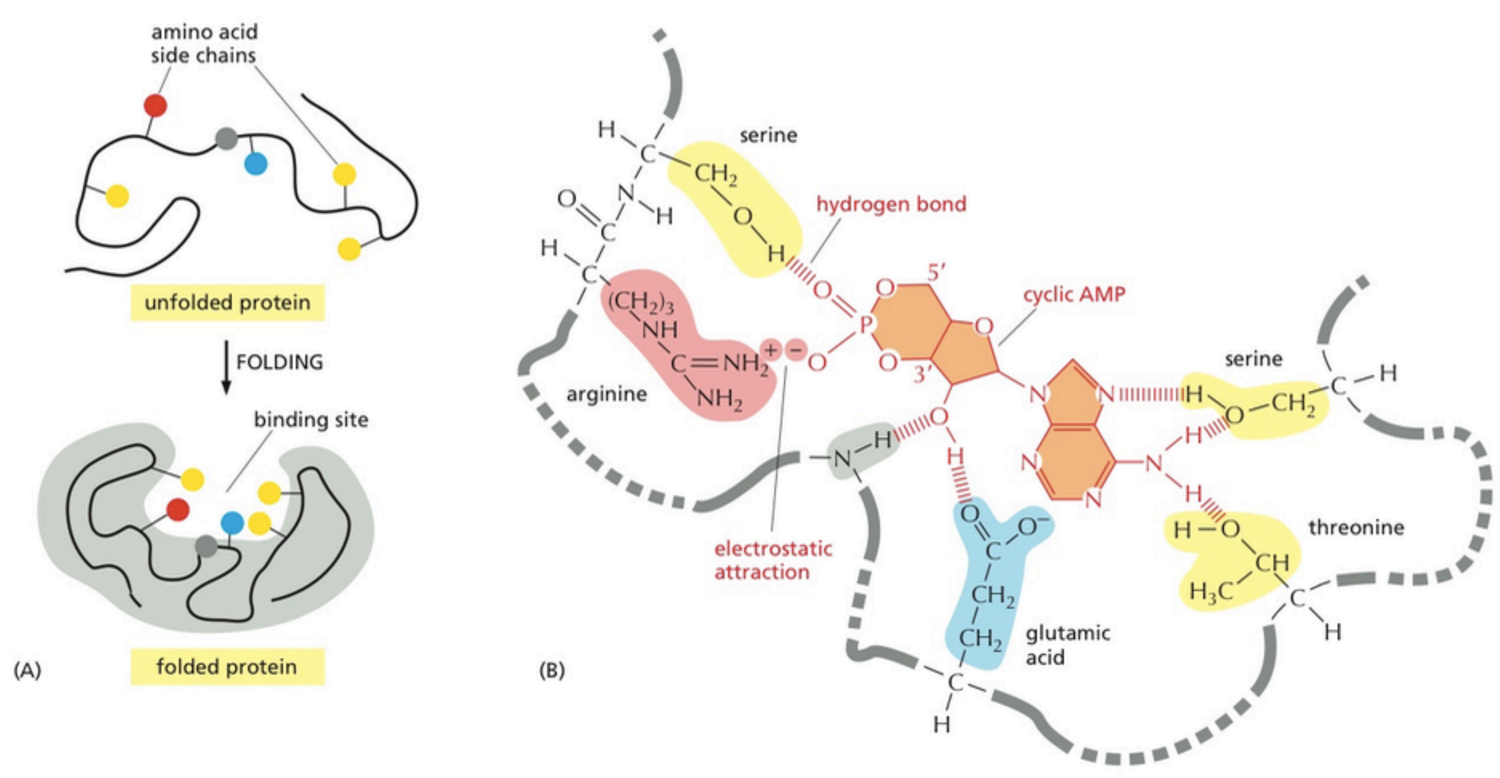


Protein-protein human network

- Sistemi *in vitro*:
  - in soluzione
  - in superficie

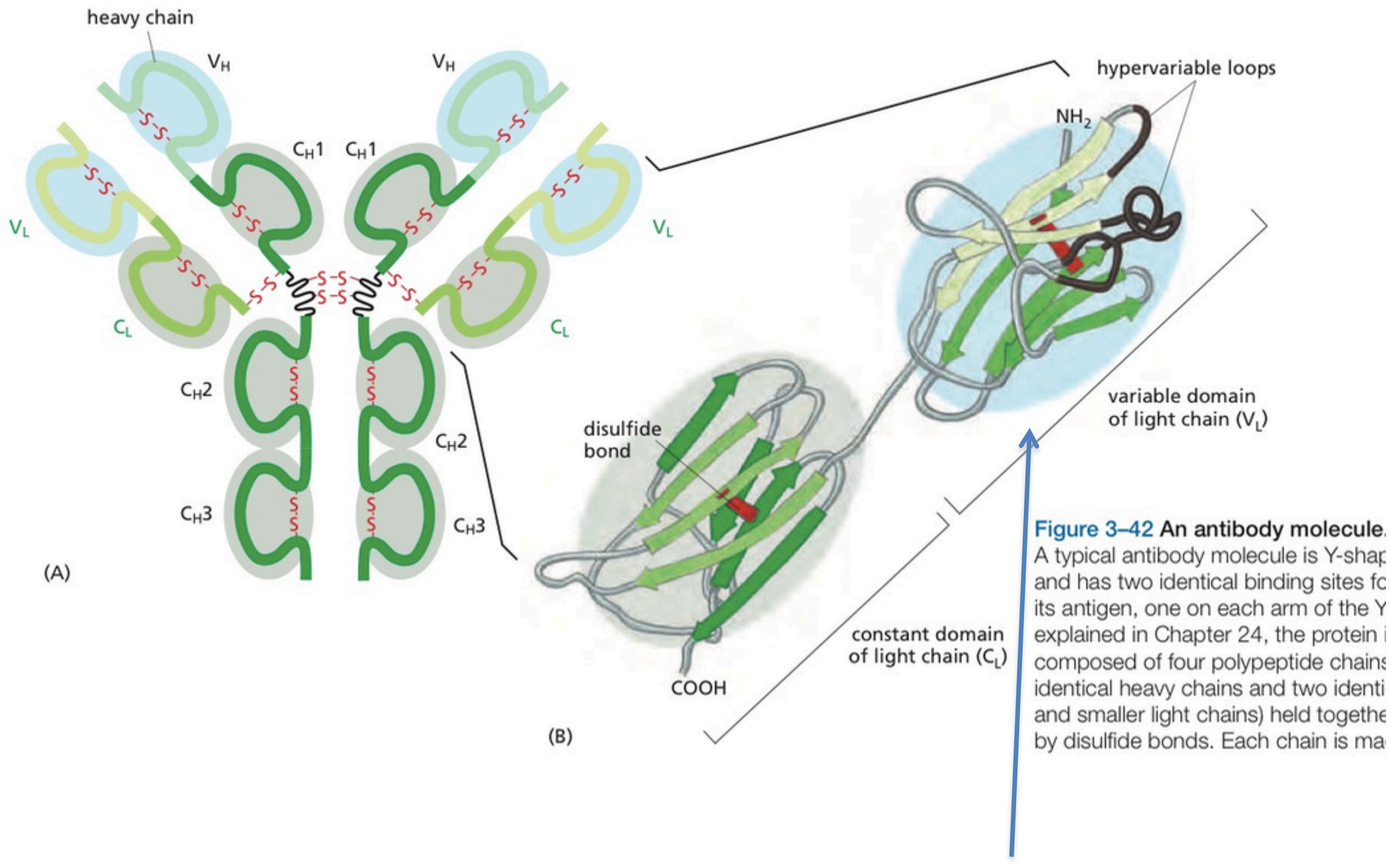
Molecular recognition

A **binding site** is a cavity made by aminoacids from different portions of the chain. Separated regions provide binding sites for different ligands, allowing for protein activity being regulated.



Selectivity depends on the set of non-covalent bonds (additive!!!) and the favorable hybdrophobic interactions that can form simultaneoulsly

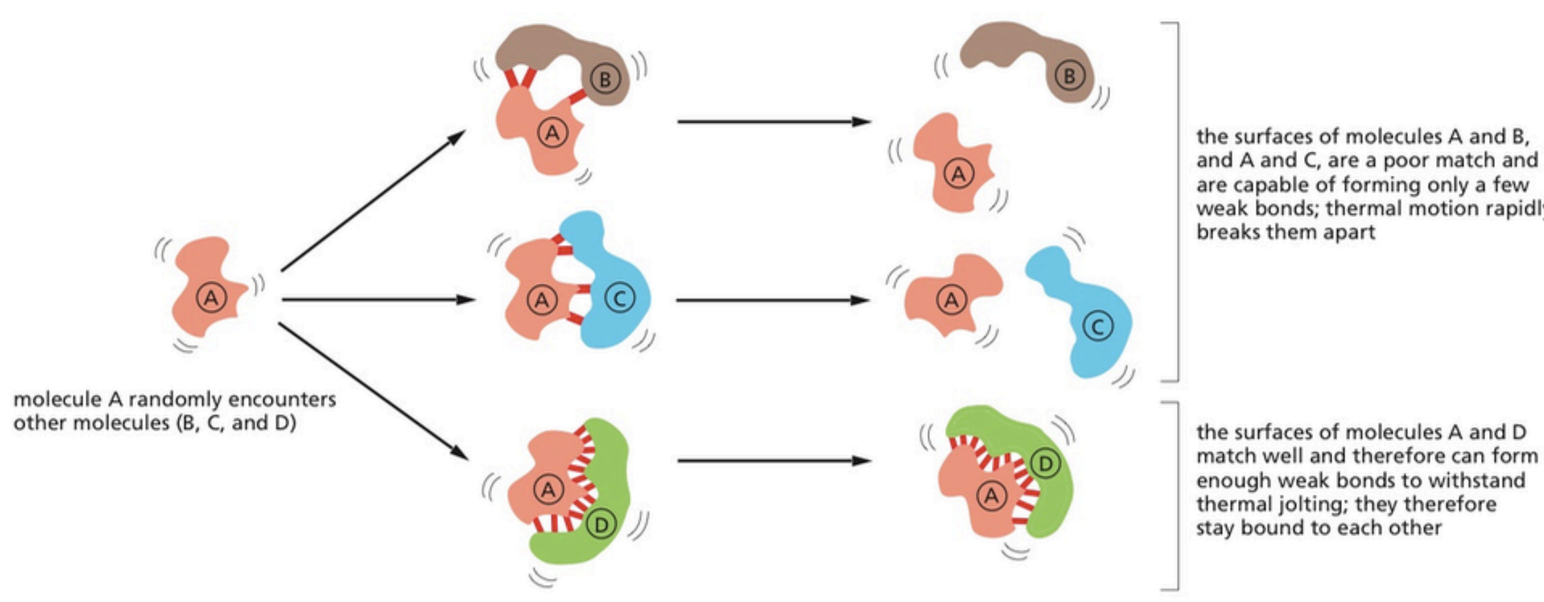
Highest affinity macromolecules: Antibodies



Two equivalent binding sites for the antigen, one per each arm

Molecular recognition

Molecular recognition refers to the process in which **biological macromolecules** interact with each other or with various **small molecules** through **noncovalent interactions** to form a specific **complex**. Characteristics:

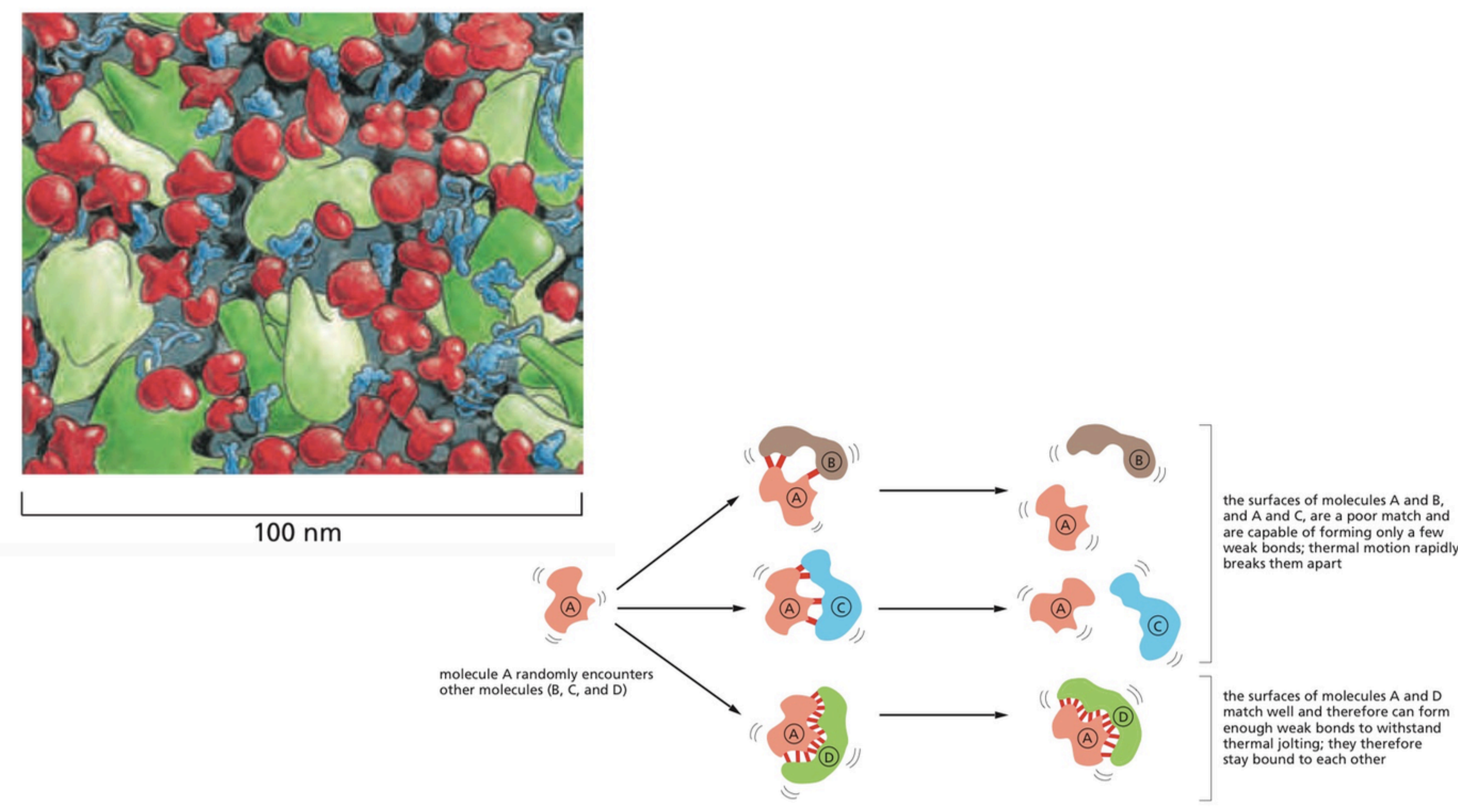


(i) **Specificity**. which distinguishes the highly specific binding partner from less specific partners;

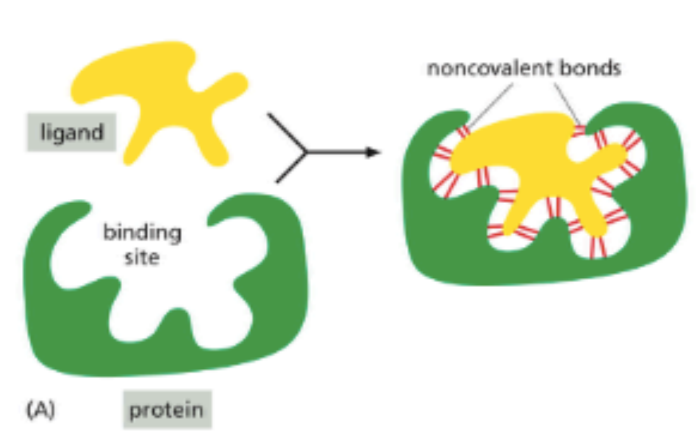
(ii) **Affinity**. high concentration of weakly interacting partners cannot replace the effect of a low concentration of the specific partner interacting with high affinity

Molecular recognition

Molecules in the cell are in a very crowded environment, in continual random thermal movements: rapid “faint” associations and dissociation between molecules are made.



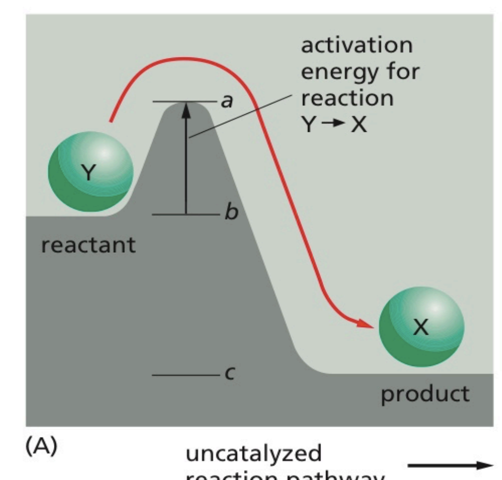
An examples: enzymes

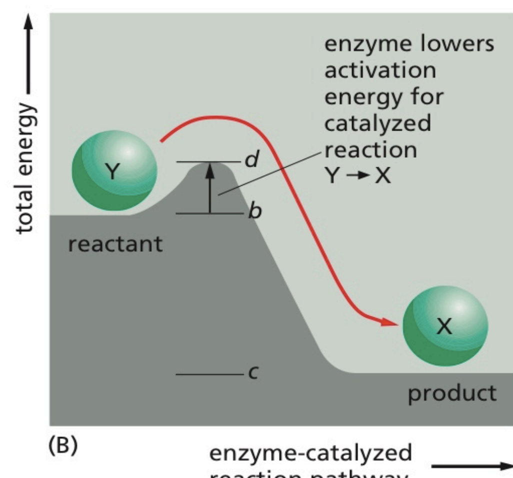


An **actin molecule** needs only to associate with other actin molecules to form a filament.

For **enzymes,** ligand binding is only a necessary first step in their function.
Enzymes cause the chemical transformations that make and break covalent bonds in cells. They bind to one or more ligands, called **substrates**, and convert them into one or more chemically modified products, doing this over and over again with amazing rapidity.

Enzymes speed up reactions, often by a **factor of a million** or more, without themselves being changed—that is, they act as catalysts that permit cells to make or break covalent bonds in a controlled way.
Catalysis of organized sets of chemical reactions by enzymes creates and maintains the cell, making life.





Other examples: enzymes

TABLE 3–1 Some Common Types of Enzymes

Enzyme	Reaction catalyzed
Hydrolases	General term for enzymes that catalyze a hydrolytic cleavage reaction; *nucleases* and *proteases* are more specific names for subclasses of these enzymes
Nucleases	Break down nucleic acids by hydrolyzing bonds between nucleotides. *Endo*– and *exonucleases* cleave nucleic acids *within* and *from the ends of* the polynucleotide chains, respectively
Proteases	Break down proteins by hydrolyzing bonds between amino acids
Synthases	Synthesize molecules in anabolic reactions by condensing two smaller molecules together
Ligases	Join together (ligate) two molecules in an energy-dependent process. DNA ligase, for example, joins two DNA molecules together end-to-end through phosphodiester bonds
Isomerases	Catalyze the rearrangement of bonds within a single molecule
Polymerases	Catalyze polymerization reactions such as the synthesis of DNA and RNA
Kinases	Catalyze the addition of phosphate groups to molecules. Protein kinases are an important group of kinases that attach phosphate groups to proteins
Phosphatases	Catalyze the hydrolytic removal of a phosphate group from a molecule
Oxido-Reductases	General name for enzymes that catalyze reactions in which one molecule is oxidized while the other is reduced. Enzymes of this type are often more specifically named *oxidases*, *reductases*, or *dehydrogenases*
ATPases	Hydrolyze ATP. Many proteins with a wide range of roles have an energy-harnessing ATPase activity as part of their function; for example, motor proteins such as *myosin* and membrane transport proteins such as the *sodium–potassium pump*
GTPases	Hydrolyze GTP. A large family of GTP-binding proteins are GTPases with central roles in the regulation of cell processes

Enzyme names typically end in "-ase," with the exception of some enzymes, such as pepsin, trypsin, thrombin, and lysozyme, that were discovered and named before the convention became generally accepted at the end of the nineteenth century. The common name of an enzyme usually indicates the substrate or product and the nature of the reaction catalyzed. For example, citrate synthase catalyzes the synthesis of citrate by a reaction between acetyl CoA and oxaloacetate.

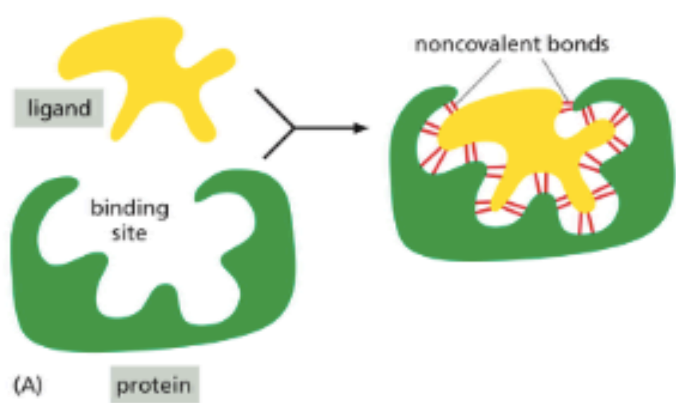
Equilibrium constant

Protein–Ligand Binding Kinetics

describes the process underlying the association between the protein and ligand, particularly focusing on the rate at which these two partners bind to each other

$$P + L \underset{k_{off}}{\overset{k_{on}}{\rightleftharpoons}} PL$$

k_{on} and k_{off} are the kinetic rate constants



Equilibrium constant

Protein–Ligand Binding Kinetics

At the equilibrium, the two reactions balance

$$k_{on}[P][L] = k_{off}[PL]$$

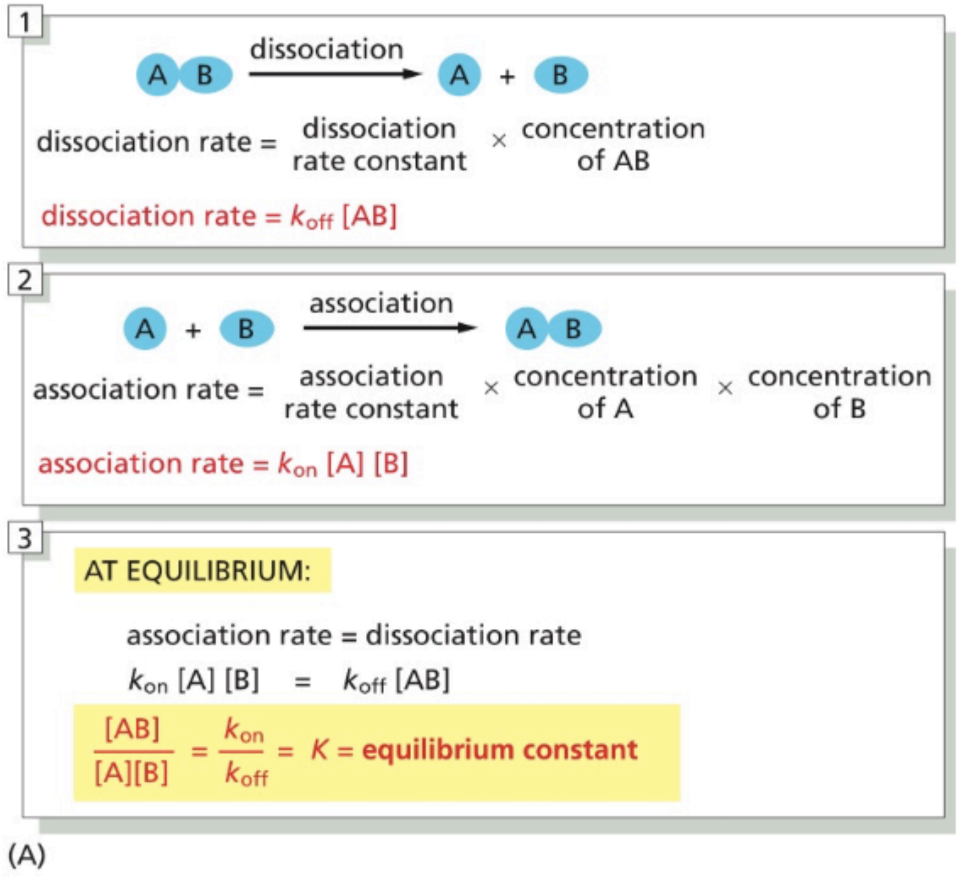
[..] is the **equilibrium concentration**

We define **the binding constant** $\mathbf{K_b}$ (M^{-1}) and the **dissociation constant** $\mathbf{K_d}$ (M) as:

$$K_b = \frac{k_{on}}{k_{off}} = \frac{[PL]}{[P][L]} = \frac{1}{K_d}$$

Therefore, the fast binding rate accompanied by a slow dissociation rate will give a high/low binding/dissociation constant and, hence, a high binding affinity.

Equilibrium constant

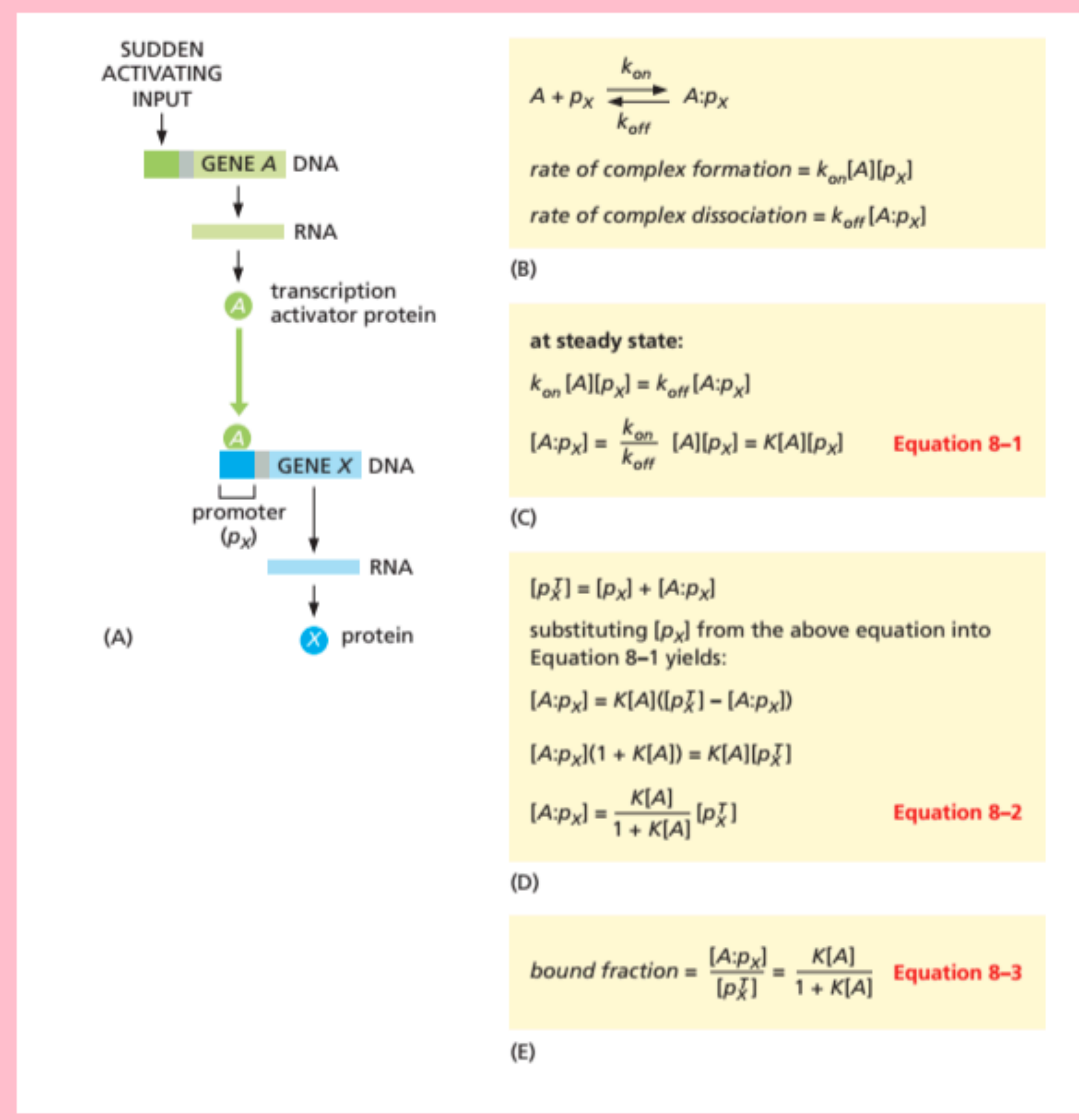


Conveniently we define the EQUILIBRIUM CONSTANT K (also known as association constant or binding constant) as a measure of the strength of the binding. **Half of the binding sites will be occupied by ligand** when the ligand's concentration (in moles/liter) reaches a value that is equal to **1/K**

k_{on} describes how many productive collisions occur per unit time per protein at a given concentration

k_{off} can differ by orders of magnitude (even for different DNA sequences) because it depends on the strength of the noncovalent bonds formed between A and B

Equilibrium constant



k_{on} describes how many productive collisions occur per unit time per protein at a given concentration

k_{off} can differ by orders of magnitude (even for different DNA sequences) because it depends on the strength of the noncovalent bonds formed between A and B

K is sometimes called the association constant, Ka. The larger this constant K, the stronger the interaction between A and p_X

Transient behavior

Let us return to Equation 8–2 (Figure 8–72D), which tells us that when $[A]$ changes, $[A{:}p_X]$ at steady state will also change to a new concentration that we can calculate with precision. However, $[A{:}p_X]$ does not change instantaneously to this value. If we hope to understand the behavior of this system in detail, we must also ask how long it takes $[A{:}p_X]$ to get to its new steady-state value inside the cell. Equation 8–2 cannot answer this question. We need calculus.

The most common strategy for solving this problem is to use ordinary differential equations. The equations that describe biochemical reactions have a simple premise: the rate of change in the concentration of any molecular species X (that is, $d[X]/dt$) is given by the balance of the rate of its appearance with that of its disappearance. For our example, the rate of change in the concentration of the bound promoter complex, $[A{:}p_X]$, is determined by the rates of complex assembly and disassembly. We can incorporate these rates into the differential equation shown in Figure 8–73A (Equation 8–4). When $[A]$ changes, Equation 8–4 can be solved to generate the concentration of $[A{:}p_X]$ as a function of time. Notice that when $k_{on}[A][p_X] = k_{off}[A{:}p_X]$, then $d[A{:}p_X]/dt = 0$ and $[A{:}p_X]$ stops changing. At this point, the system has reached the steady state.

Transient behavior

Calculation of all $[A{:}p_X]$ values as a function of time, using Equation 8–4, allows us to determine the rate at which $[A{:}p_X]$ reaches its steady-state value. Because this value is attained asymptotically, it is often most useful to compare the times needed to get to 50, 90, or 99 percent of this new steady state. The simplest way to determine these values is to solve Equation 8–4 with a method called numerical integration, which involves plugging in values for all of the parameters (k_{on}, k_{off}, etc.) and then using a computer to determine the values of $[A{:}p_X]$ over time, starting from given initial concentrations of $[A]$ and $[p_X]$. For $k_{on} = 0.5 \times 10^7\ \text{sec}^{-1}\ \text{M}^{-1}$, $k_{off} = 0.5 \times 10^{-1}\ \text{sec}^{-1}$ ($K = 10^8\ \text{M}^{-1}$ as above), and $[p_X^T] = 10^{-10}$ M, it takes $[A{:}p_X]$ about 5, 20, and 40 seconds to reach 50, 90, and 99 percent of the new steady-state value following a sudden tenfold change in $[A]$ (Figure 8–73B). Thus, a sudden jump in $[A]$ does not have instantaneous effects, as we might have assumed from looking at the cartoon in Figure 8–72A.

Differential equations therefore allow us to understand the transient dynamics of biochemical reactions. This tool is critical for achieving a deep understanding of cell behavior, in part because it allows us to determine the dependence of the dynamics inside cells on parameters that are specific to the particular molecules involved. For example, if we double the values of both k_{on} and k_{off}, then Equation 8–1 (Figure 8–72C) indicates that the steady-state value of $[A{:}p_X]$ does not change. However, the time it takes to reach 50% of this steady state after a ten-fold

$$\frac{d[A{:}p_X]}{dt} = \textit{rate of complex formation} - \textit{rate of complex dissociation}$$

$$\frac{d[A{:}p_X]}{dt} = k_{on}[A][p_X] - k_{off}[A{:}p_X] \qquad \textbf{Equation 8–4}$$

(A)

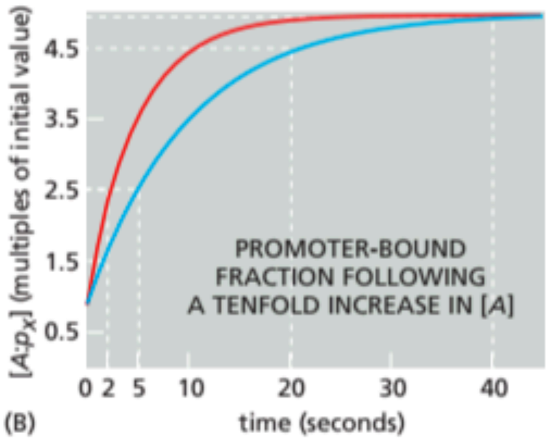


Figure 8–73 Using differential equations to study the dynamics and steady-state behavior of a biological system. (A) Equation 8–4 is an ordinary differential equation for calculating the rate of change in the formation of bound promoter complex in response to a change in other components. (B) Formation of $[A{:}p_X]$ after a tenfold increase in $[A]$, as determined by solving Equation 8–4. In *blue* is the solution corresponding to $k_{on} = 0.5 \times 10^7\ \text{sec}^{-1}\ \text{M}^{-1}$ and $k_{off} = 0.5 \times 10^{-1}\ \text{sec}^{-1}$. In this case, it takes $[A{:}p_X]$ about 5, 20, and 40 seconds to reach 50, 90, and 99 percent of the new steady-state value. For the *red curve*, the k_{on} and k_{off} values are doubled, and the system reaches the same steady state more rapidly.

Transient behavior

$$\text{transcription rate} = \beta \frac{K[A]}{1 + K[A]}$$

$$\text{protein production rate} = \beta \cdot m \frac{K[A]}{1 + K[A]}$$

$$\text{protein degradation rate} = \frac{[X]}{\tau_X}$$

(A)

$$\frac{d[X]}{dt} = \text{protein production rate} - \text{protein degradation rate}$$

$$\frac{d[X]}{dt} = \beta \cdot m \frac{K[A]}{1 + K[A]} - \frac{[X]}{\tau_X} \qquad \text{Equation 8–5}$$

(B)

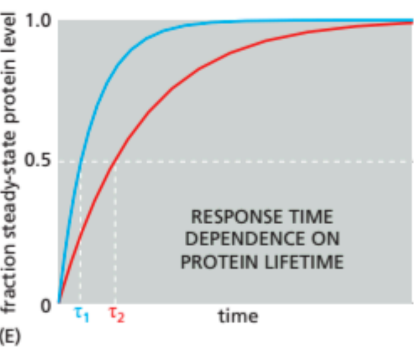
at steady state:

$$[X_{st}] = \beta \cdot m \frac{K[A]}{1 + K[A]} \cdot \tau_X \qquad \text{Equation 8–6}$$

(C)

$$[X](t) = [X_{st}](1 - e^{-\frac{t}{\tau_X}})$$

(D)



(E)

Figure 8–74 Effect of protein lifetime on the timing of the response. (A) Equations for calculation of the rates of gene *X* transcription, protein *X* production, and protein *X* degradation, as explained in the text. (B) Equation 8–5 is an ordinary differential equation for calculating the rate of change in protein *X* in response to changes in other components. (C) When the rate of change in protein *X* is zero (steady state), its concentration can be calculated with Equation 8–6, revealing a direct relationship with protein lifetime (τ). (D) The solution of Equation 8–5 specifies the concentration of protein *X* over time as it approaches its steady-state concentration. (E) Response time depends on protein lifetime. As described in the text, the time that it takes a protein to reach a new steady state is greater when the protein is more stable. Here, the *blue line* corresponds to a protein with a lifetime that is 2.5-fold shorter than the lifetime of the protein in *red*.

change in [*A*] in our example changes from about 5 seconds to 2 seconds (see Figure 8–73B). These insights are not accessible from either cartoons or equilibrium equations. This is an unusually simple example; mathematical descriptions such as differential equations become more indispensible for understanding biological interactions as the number of interactions increases.

Transient behavior

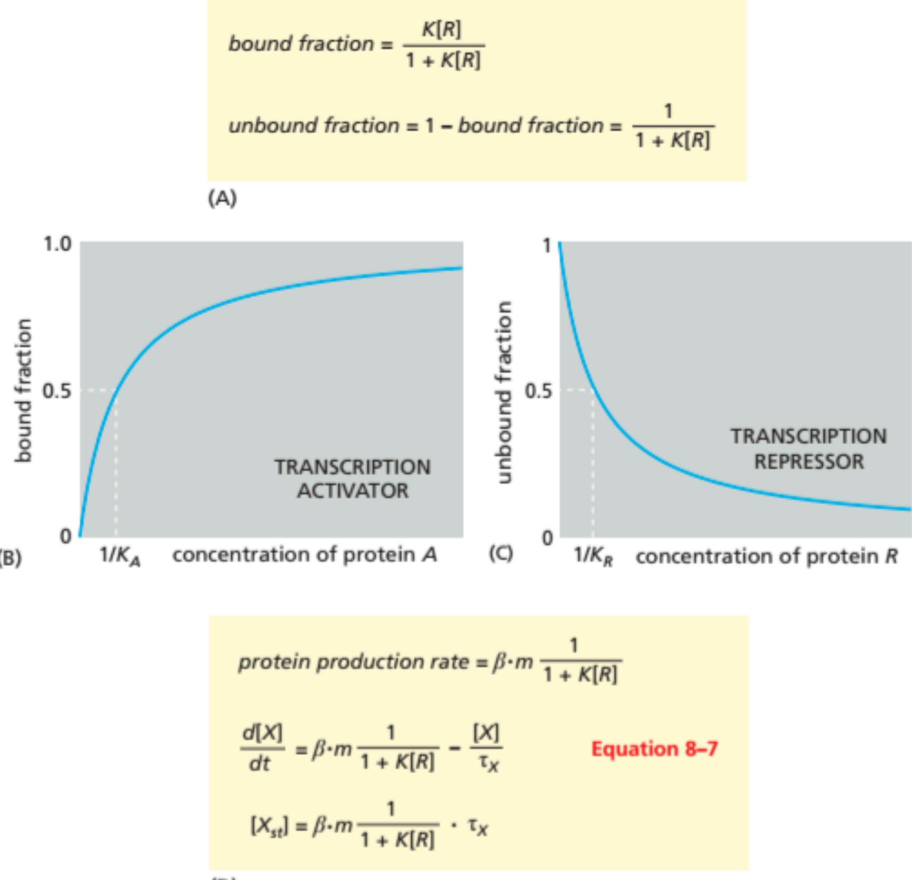
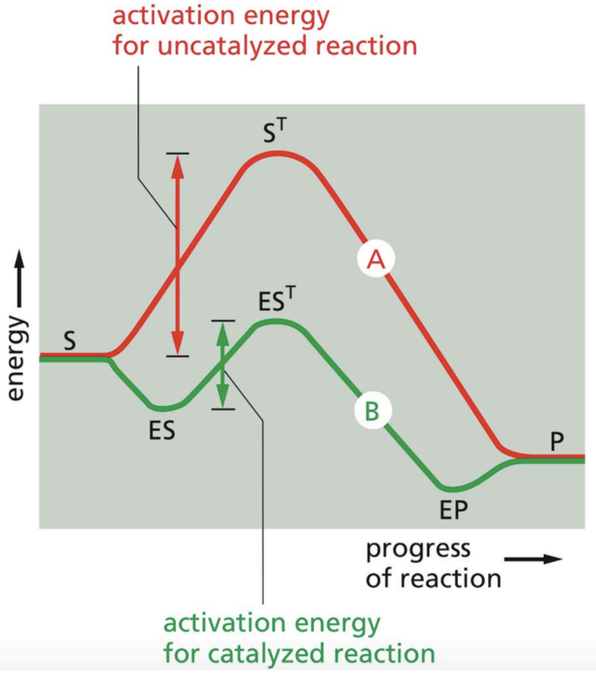
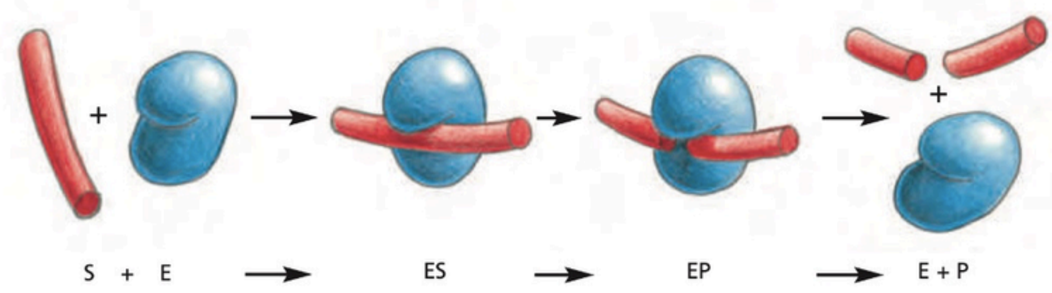


Figure 8–75 How promoter occupancy depends on the binding affinity of a transcription regulator protein. (A) The fraction of a binding site that is occupied by a transcription repressor *R* is determined by an equation that is similar to the one we used for a transcription activator (see Figure 8–72E), except that in the case of a repressor we are interested primarily in the unbound fraction. (B) For a transcription activator *A*, half of the promoters are occupied when $[A] = 1/K_A$. Gene activity is proportional to this bound fraction. (C) For a transcription repressor *R*, gene activity is proportional to the unbound fraction of promoters. As indicated, this fraction is reduced to half of its maximal value when $[R]=1/K_R$. (D) As in the case of the transcription activator *A* (see Figure 8–74), we can derive equations to assess the timing of protein *X* production as a function of repressor concentrations.

Steady state enzyme kinetics





WHY ANALYZE THE KINETICS OF ENZYMES?

Enzymes are the most selective and powerful catalysts known. An understanding of their detailed mechanisms provides a critical tool for the discovery of new drugs, for the large-scale industrial synthesis of useful chemicals, and for appreciating the chemistry of cells and organisms. A detailed study of the rates of the chemical reactions that are catalyzed by a purified enzyme—more specifically how these rates change with changes in conditions such as the concentrations of substrates, products, inhibitors, and regulatory ligands—allows biochemists to figure out exactly how each enzyme works. For example, this is the way that the ATP-producing reactions of glycolysis, shown previously in Figure 2–48, were deciphered—allowing us to appreciate the rationale for this critical enzymatic pathway.

In this Panel, we introduce the important field of enzyme kinetics, which has been indispensable for deriving much of the detailed knowledge that we now have about cell chemistry.

Basic Concepts and Thermodynamic Relationships

Enzyme can process 1000 mol. per second. Meaning they bind a new substrate in a fraction of a milliseconds.
But enzymes and their substrates are present in relatively small numbers in a cell.
How do they find each other so fast?
Rapid binding is possible because the **motions caused by heat energy are enormously fast at the molecular level,** generating:

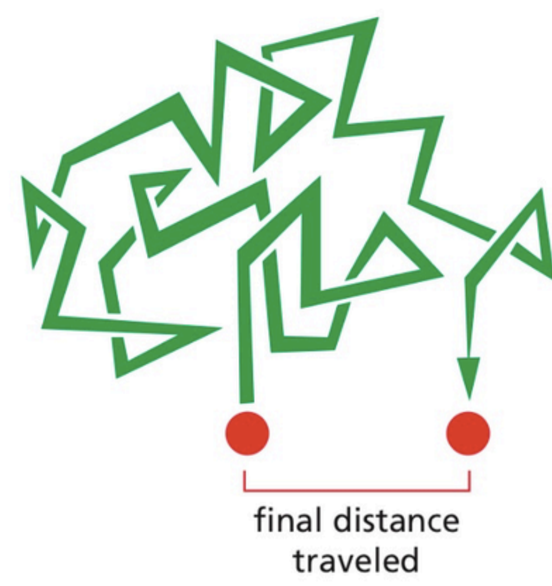
(1) the movement of a molecule from one place to another (translational motion)
(2) the rapid back-and-forth movement of covalently linked atoms with respect to one another (vibrations)
(3) rotations.

All of these motions help to bring the surfaces of interacting molecules together.
The rates of molecular motions can be measured by a variety of spectroscopic techniques.

Alberts, Bruce; Johnson, Alexander; Lewis, Julian; Morgan, David; Raff, Martin; Roberts, Keith; Walter, Peter. Molecular Biology of the Cell (Page 59). Garland Science. Kindle Edition.

Basic Concepts and Thermodynamic Relationships

A large globular protein is constantly tumbling, rotating about its axis about a million times per second.
Molecules are also in constant translational motion, which causes them to explore the space inside the cell very efficiently by wandering through it—a process called ***diffusion.*** A small organic molecule, for example, takes only about one-fifth of a second on average to diffuse a distance of 10 µm, or the whole cell!



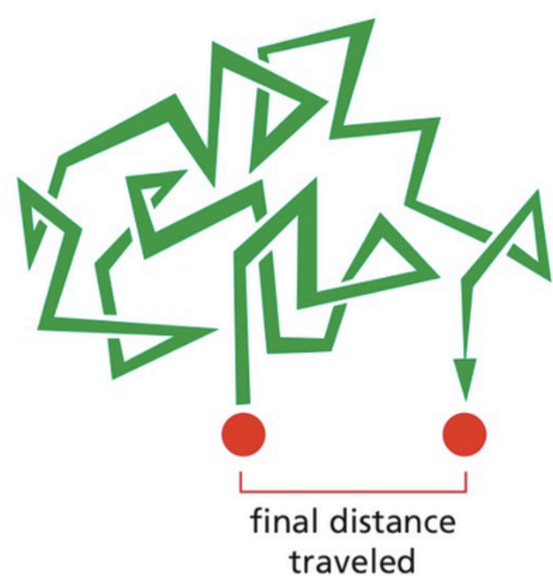
Diffusion with great heat exchange generates a molecule ***random walk***

The **average net distance** that each molecule travels from its starting point is **proportional to the square root of the time** involved.
if it takes a molecule 1 second on average to travel 1 µm, it takes 4 seconds to travel 2 µm, 100 seconds to travel 10 µm, and so on.

Alberts, Bruce; Johnson, Alexander; Lewis, Julian; Morgan, David; Raff, Martin; Roberts, Keith; Walter, Peter. Molecular Biology of the Cell (Page 59). Garland Science. Kindle Edition.

Basic Concepts and Thermodynamic Relationships

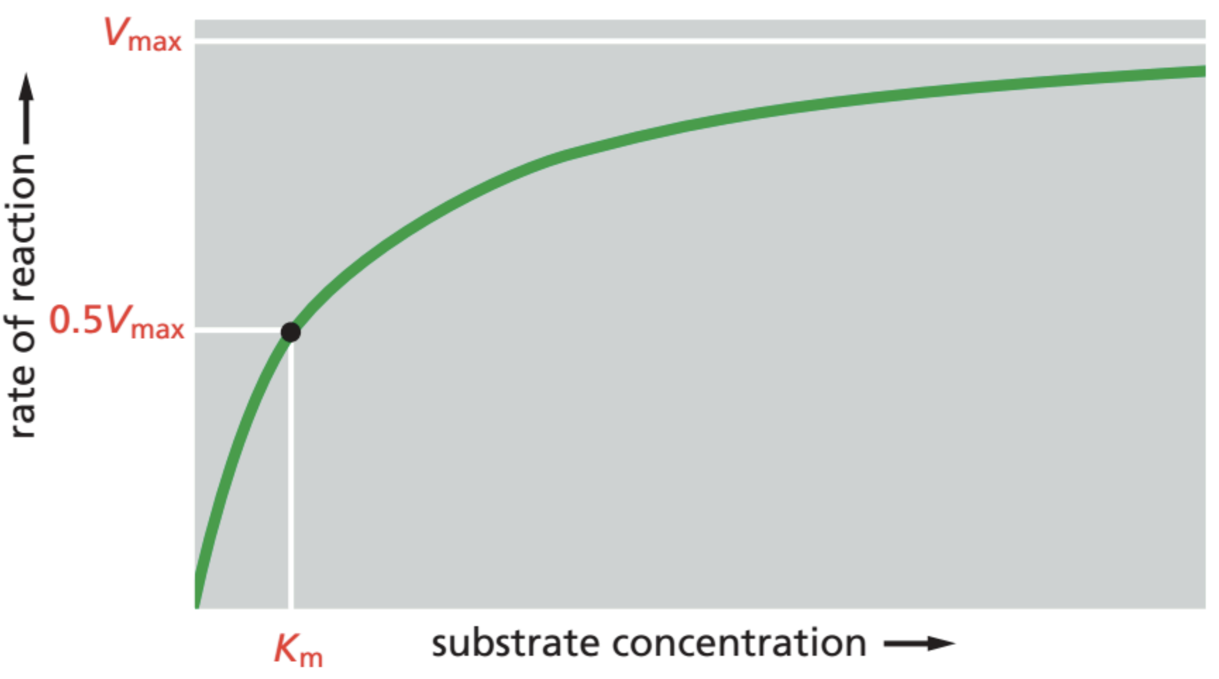
Since enzymes move more slowly than substrates in cells, we can think of them as sitting still. The rate of encounter of each enzyme molecule with its substrate will depend on the concentration of the substrate



Abundant substrates : 0.5 mM
Water: 55.5 M
1 substrate mol per 10^5 water mol!
However, will face 100.000 random collisions per second by the substrate

Alberts, Bruce; Johnson, Alexander; Lewis, Julian; Morgan, David; Raff, Martin; Roberts, Keith; Walter, Peter. Molecular Biology of the Cell (Page 59). Garland Science. Kindle Edition.

Basic Concepts and Thermodynamic Relationships

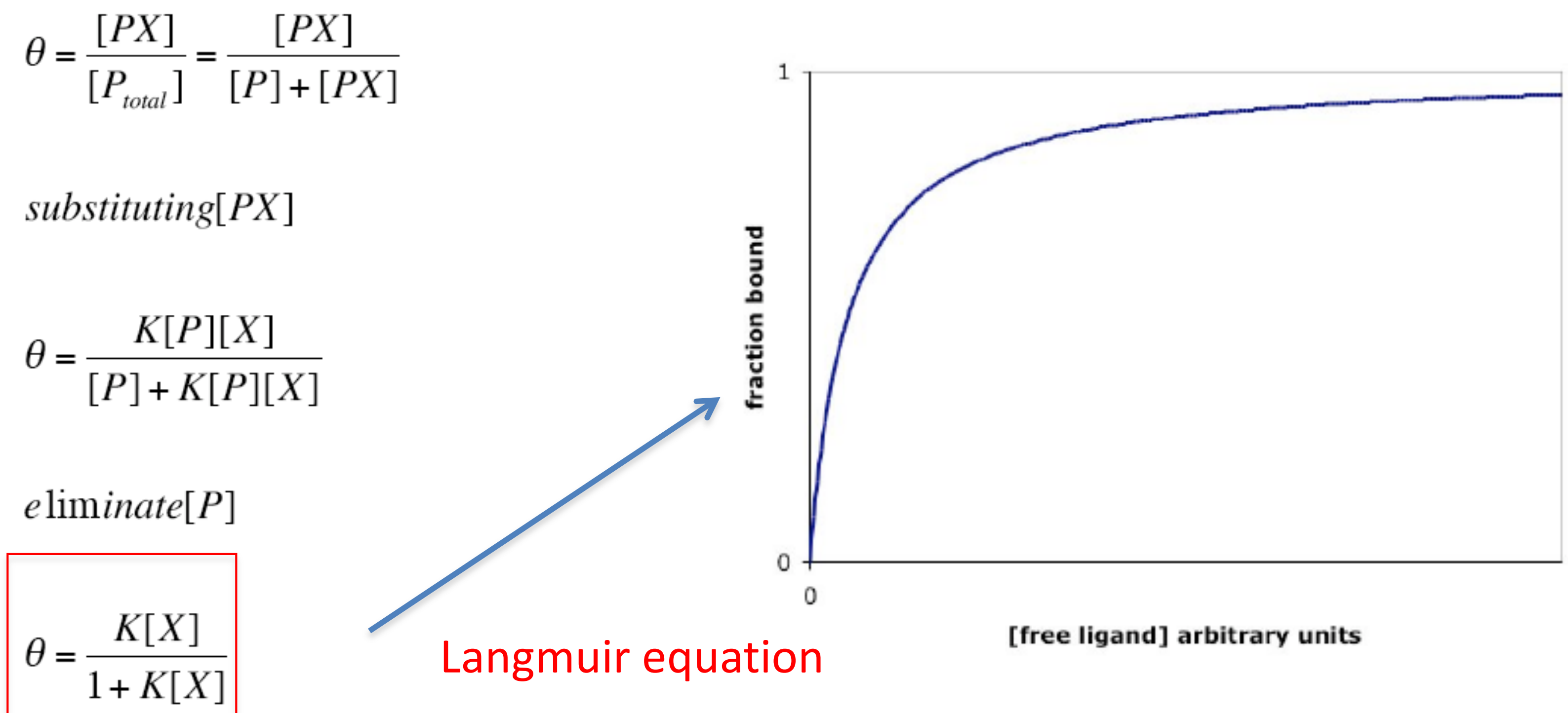


Turnover number = V max/ Enzyme concentration is often about 1000 substrate molecules processed per second

The rate of an enzyme reaction (V) increases as the substrate concentration increases until a maximum value (V max) is reached. At this point all substrate-binding sites on the enzyme molecules are fully occupied, and the rate of reaction is limited by the rate of the catalytic process on the enzyme surface.

For most enzymes, the concentration of substrate at which the reaction rate is half-maximal (**K_m**) is a measure of how tightly the substrate is bound, with a large value of K_m corresponding to weak binding.

In most experiments the fraction of molecules bound to the protein – let's call this fraction q – is the parameter that is easiest to measure. The situation here is similar to that for the protein melting curves. The exact physical nature of the signal does not really matter; all that matters is that the signal is linearly proportional to the fraction of ligands bound.



Brownian motion

Il moto casuale di una piccola particella (con diametro dell'ordine del micron) immersa in un fluido, dovuto ad urti tra la particella e le molecole del fluido stesso, e chiamato moto browniano.

Il **moto browniano** di una particella colloidale `e una sequenza di passi casuali della posizione e dell'orientazione della particella stessa. Tale sequenza `e chiamata **diffusione** ed `e descritta da un'equazione,che permette di comprendere come la posizione della particella evolve nel tempo.

Come si deriva l'equazione di diffusione per particelle all'equilibrio, ovvero in assenza di forze esterne?

Legge di Fick e di legge di conservazione

N particelle browniane, ossia grandi rispetto alle particelle del fluido in cui si muovono, ciascuna avente una posizione $\vec{x}_\alpha$ e una velocita $\vec{v}_\alpha = \frac{d\vec{x}_\alpha}{dt}$ con $1 \leq \alpha \leq$ N.

Definiamo la concentrazione c:

$$c(\vec{x},t) = \sum_{\alpha=1}^{N} \delta(\vec{x} - \vec{x}_\alpha(t))$$

E la densita' di corrente media:

$$\vec{j}(\vec{x},t) = \sum_{\alpha=1}^{N} \delta(\vec{x} - \vec{x}_\alpha(t))\vec{v}_\alpha(t)$$

Legge di Fick e di legge di conservazione

Se il sistema e' isolato, vale la legge di conservazione del numero di particelle, che da' una legge di continuita':

$$\frac{\partial c}{\partial t} = -\vec{\nabla} \cdot \vec{j}$$

In assenza di forze esterne:

$$< c(\vec{x}, t) >= costante$$

In presenza di forze esterne o fluttuazioni statistiche: ci aspettiamo che il sistema tenda a ritornare all'equilibrio.
Assumiamo, empiricamente, che la reazione del sistema avvenga secondo la cosidetta **legge di Fick**, per cui si crea una corrente di particelle nel verso opposto al gradiente iniziale

$$\vec{\nabla} c \neq 0 \qquad \vec{j} \propto -\vec{\nabla} c \rightarrow \vec{j} = -D\vec{\nabla} c \qquad [D] = \frac{m^2}{s}$$

dove **D e' detto coefficiente di diffusione** ed in generale dipende da posizione e tempo

Legge di Fick e di legge di conservazione

Si ottiene: $$\frac{\partial c}{\partial t} = -\vec{\nabla} \cdot \vec{j} = \vec{\nabla} \cdot (D\vec{\nabla} c) = D\vec{\nabla}^2 c$$

Vera se D e' indipendente dalla diffusione.
Nel caso 1-dimensionale:

$$\frac{\partial c}{\partial t} = D\frac{\partial^2 c}{\partial x^2}$$

La soluzione c (x,t) e' una gaussiana. Se inizialmente tutte le particelle sono condensate in un punto, cioe' $\mathrm{c(x,t{=}0) = N\delta(x)}$ si ottiene l'equazione di diffusione:

$$c(x,t) = \frac{1}{(4\pi Dt)^{\frac{1}{2}}} e^{-\frac{x^2}{4Dt}}$$

Legge di Fick e di legge di conservazione

Da: $$< x^2 >= \frac{1}{N}\int x^2 c(x,t)dx = 2Ddt$$

Si ottiene: $$D = \lim_{x\to\infty} \frac{< (x(t) - x(0))^2 >}{2dt}$$

Formula empirica della diffusione, derivate da Einstein.

Equazione di Langevin

Nel quadro della descrizione teorica dei processi diffusivi, nel primo decennio del XX secolo, accanto ad Einstein si inserisce il fisico francese P.Langevin (1872-1946).
Nella descrizione del moto browniano, Langevin parte dalle equazioni del moto.[3]
Si consideri una particella browniana immersa in un fluido, di raggio a, $10^{-9}m < a < 5\times10^{-7}m$ (essa ha quindi dimensioni maggiori rispetto agli atomi del fluido in cui è immersa). Il suo moto è pertanto più lento di quello degli atomi ed è il risultato di collisioni rapide e casuali con essi, dovute a fluttuazioni di densità del fluido.
Considerando il moto unidimensionale, l'equazione di Newton della particella di massa m e velocià $v(t)$ è

$$m\frac{dv(t)}{dt} = F(t) \tag{2.11}$$

dove F(t) è la forza totale agente sulla particella al tempo t, causata dall'interazione di questa con il mezzo circostante. Tale forza è dominata da un termine di attrito viscoso $-\gamma v(t)$ dove $\gamma = 6\pi\eta a$, con a raggio della particella ed η viscosità del fluido. Ci si aspetta anche un contribuito casuale alla forza F(t), denotato con $\xi(t)$, dovuto alle fluttuazioni casuali della densità del fluido.
Dunque le equazioni del moto della particella browniana, dette *equazioni di Langevin* sono

$$\frac{dx(t)}{dt} = v(t) \qquad\qquad \frac{dv(t)}{dt} = -\frac{\gamma}{m}v(t) + \frac{\xi(t)}{m} \tag{2.12}$$

dove la variabile stocastica $\xi(t)$ dà l'effetto di rumore dovuto al fluido.

Equazione di Langevin

dove la variabile stocastica $\xi(t)$ dà l'effetto di rumore dovuto al fluido. Se trascurassimo tale componente stocastica, l'equazione (2.12) diverrebbe

$$\frac{dv(t)}{dt} = -\frac{\gamma}{m}v(t) \tag{2.13}$$

con soluzione

$$v(t) = e^{-\frac{t}{\tau}}v(0) \qquad\qquad \tau = \frac{m}{\gamma} \approx 10^{-3}s \tag{2.14}$$

dove τ è il tempo di rilassamento tipico della velocità della particella browniana.

Secondo l'uguaglianza (2.14), la velocità della particella dovrebbe tendere a zero per tempi lunghi. Ciò non può essere vero in quanto in equiliibrio, per il teorema di equipartizione vale

$$< v^2(t) >_{eq} = \frac{kT}{m} \tag{2.15}$$

mentre da (2.14) risulta

$$< v^2(t) >_{eq} = e^{-\frac{2t}{\tau}} < v^2(0) >_{eq} \to 0 \tag{2.16}$$

Quindi $\xi(t)$ in (2.12) è necessaria per avere il corretto equilibrio.

Il modo browniano è dovuto alle collisioni rapide e casuali tra particelle colloidali e molecole del fluido circostante, causate da fluttuazioni di densità di quest'ultimo. Si suppone che la forza durante tale collisione vari in modo estremamente rapido nel tempo di qualsiasi osservazione; ci si aspetta dunque, per la componente casuale della forza F

$$< \xi(t) >_{\xi} = 0 \tag{2.17}$$

Equazione di Langevin

Si dimostra matematicamente, applicando il teorema del limite centrale a una distribuzione di variabili casuali indipendenti, che la componente casuale della forza ha una distribuzione gaussiana con media zero.[3] A differenza dell'equazione di diffusione (2.7), che è un'equazione deterministica, l'equazione di Langevin è stocastica,in quanto include una forza casuale. Ogni soluzione di tale equazione del moto stocastica è una differente traiettoria casuale.

I processi di diffusione browniana sono caratterizzati, dunque, da due proprietà[4]:

- la crescita lineare nel tempo dello spostamento quadratico medio (MSD da mean-squared displacement)

  $$< \mathbf{r}^2(t) >= 2dDt \tag{3.1}$$

  dove d è la dimensione spaziale e D il coefficiente di diffusione.

- la funzione densità di probabilità di trovare la particella che diffonde, alla posizione $\mathbf{r}$ al tempo t è una gaussiana

  $$P(\mathbf{r}, t) = \frac{1}{(4\pi Dt)^{d/2}} e^{-\frac{\mathbf{r}^2}{4Dt}} \tag{3.2}$$