Advanced Quantum Mechanics

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Quantum Algorithms

We will study the three historically most important algorithms:

- Simple ones (Deutsch, Deutsch-Jozsa...)
- Grover (search in a data base)
- Shor (factorization)

What is special about quantum algorithms?

Quantum Parallelism

Given an input x, a typical quantum computer computes f(x) in such a way as

$$U_f: |x\rangle|0\rangle \mapsto |x\rangle|f(x)\rangle, \qquad (4.61)$$

where U_f is a unitary matrix that implements the function f.

Suppose U_f acts on the input which is a superposition of many states. Since U_f is a linear operator, it acts simultaneously on all the vectors that constitute the superposition. Thus the output is also a superposition of all the results;

$$U_f: \sum_{x} |x\rangle |0\rangle \mapsto \sum_{x} |x\rangle |f(x)\rangle.$$
(4.62)

All values of the function computed at once. Very easy!! But... measurements will make the wave function collapse giving only one output. No advantage

Quantum Algorithms

The goal of a quantum algorithm is to operate in such a way that the particular outcome we want to observe has a larger probability to be measured than the other outcomes.

Let $f:\{0,1\}\to \{0,1\}$ be a binary function. Note that there are only four possible f, namely

$$f_1: 0 \mapsto 0, \ 1 \mapsto 0, \quad f_2: 0 \mapsto 1, \ 1 \mapsto 1,$$

$$f_3: 0 \mapsto 0, \ 1 \mapsto 1, \quad f_4: 0 \mapsto 1, \ 1 \mapsto 0.$$

The first two cases, f_1 and f_2 , are called <u>constant</u>, while the rest, f_3 and f_4 , are <u>balanced</u>. If we only have classical resources, we need to evaluate f twice to tell if f is constant or balanced. There is a quantum algorithm, however, with which it is possible to tell if f is constant or balanced with a single evaluation of f, as was shown by Deutsch [2].

First we need to turn the classical function f(x) into a quantum one.

- 1. Make it reversible.
- 2. Define it on the computational basis to act like the classical circuit and extend it by linearity.

$$U_f : |x, y\rangle \mapsto |x, y \oplus f(x)\rangle$$

Where \oplus is addition mod 2.

The algorithm is structured as follows.

- 1. Start with the state $|01\rangle$.
- 2. Apply an Hadamard on both qubits: $\frac{1}{2}(|00\rangle |01\rangle + |10\rangle |11\rangle)$
- 3. Apply the operator U_f implementing the function

$$\begin{aligned} &\frac{1}{2}(|0, f(0)\rangle - |0, 1 \oplus f(0)\rangle + |1, f(1)\rangle - |1, 1 \oplus f(1)\rangle) \\ &= \frac{1}{2}(|0, f(0)\rangle - |0, \neg f(0)\rangle + |1, f(1)\rangle - |1, \neg f(1)\rangle), \end{aligned}$$

Quantum parallelism: all values computed at once

4. Apply an Hadamard to the first qubit

$$\frac{1}{2\sqrt{2}} \left[(|0\rangle + |1\rangle) (|f(0)\rangle - |\neg f(0)\rangle) + (|0\rangle - |1\rangle) (|f(1)\rangle - |\neg f(1)\rangle) \right]$$

The wave function reduces to

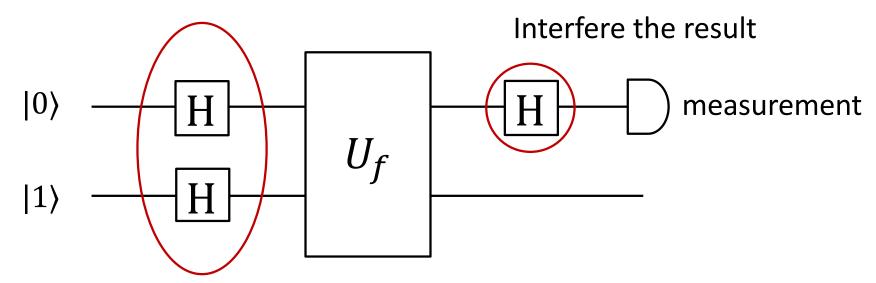
$$\frac{1}{\sqrt{2}}|0\rangle(|f(0)\rangle - |\neg f(0)\rangle) \tag{5.1}$$

in case f is constant, for which $|f(0)\rangle = |f(1)\rangle$, and

$$\frac{1}{\sqrt{2}}|1\rangle(|f(0)\rangle - |\neg f(0)\rangle) \tag{5.2}$$

if f is balanced, for which $|\neg f(0)\rangle = |f(1)\rangle$. Therefore the measurement of the first qubit tells us whether f is constant or balanced.

5. Measure the first qubit



Calculate the function on both input values simultaneously

Let us first define the **Deutsch-Jozsa problem**. Suppose there is a binary function

$$f: S_n \equiv \{0, 1, \dots, 2^n - 1\} \to \{0, 1\}.$$
 (5.3)

We require that f be either *constant* or *balanced* as before. When f is constant, it takes a constant value 0 or 1 irrespetive of the input value x. When it is balanaced the value f(x) for the half of $x \in S_n$ is 0, while it is 1 for the rest of x.

It is clear that we need at least $2^{n-1} + 1$ steps, in the worst case with classical manipulations, to make sure if f(x) is constant or balanced with 100% confidence. It will be shown below that the number of steps reduces to a single step if we are allowed to use a quantum algorithm.

- 1. Prepare an (n + 1)-qubit register in the state $|\psi_0\rangle = |0\rangle^{\otimes n} \otimes |1\rangle$. First n qubits work as input qubits, while the (n + 1)st qubit serves as a "scratch pad." Such qubits, which are neither input qubits nor output qubits, but work as a scratch pad to store temporary information are called **ancillas** or **ancillary qubits**.
- 2. Apply the Walsh-Hadamard transforamtion to the register. Then we have the state

$$|\psi_1\rangle = U_{\mathrm{H}}^{\otimes n+1} |\psi_0\rangle = \frac{1}{\sqrt{2^n}} (|0\rangle + |1\rangle)^{\otimes n} \otimes \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle)$$
$$= \frac{1}{\sqrt{2^n}} \sum_{x=0}^{2^n-1} |x\rangle \otimes \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle).$$
(5.4)

3. Apply $U_f|x\rangle|c\rangle = |x\rangle|c\oplus f(x)\rangle$

The state changes into

$$\psi_{2} \rangle = U_{f} |\psi_{1}\rangle$$

$$= \frac{1}{\sqrt{2^{n}}} \sum_{x=0}^{2^{n}-1} |x\rangle \frac{1}{\sqrt{2}} (|f(x)\rangle - |\neg f(x)\rangle)$$

$$= \frac{1}{\sqrt{2^{n}}} \sum_{x} (-1)^{f(x)} |x\rangle \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle).$$
(5.5)

Although the gate U_f is applied once for all, it is applied to all the *n*-qubit states $|x\rangle$ simultaneously.

4. The Walsh-Hadamard transformation (4.11) is applied on the first n qubits next. We obtain

$$|\psi_3\rangle = (W_n \otimes I)|\psi_2\rangle = \frac{1}{\sqrt{2^n}} \sum_{x=0}^{2^n - 1} (-1)^{f(x)} U_{\mathrm{H}}^{\otimes n} |x\rangle \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle). \quad (5.6)$$

On the Hadamard gate

It is instructive to write the action of the one-qubit Hadamard gate in the following form,

$$U_{\rm H}|x\rangle = \frac{1}{\sqrt{2}}(|0\rangle + (-1)^x|1\rangle) = \frac{1}{\sqrt{2}}\sum_{y\in\{0,1\}}(-1)^{xy}|y\rangle,$$

where $x \in \{0, 1\}$, to find the resulting state. The action of the Walsh-Hadamard transformation on $|x\rangle = |x_{n-1} \dots x_1 x_0\rangle$ yields

$$W_{n}|x\rangle = (U_{\rm H}|x_{n-1}\rangle)(U_{\rm H}|x_{n-2}\rangle)\dots(U_{\rm H}|x_{0}\rangle)$$

$$= \frac{1}{\sqrt{2^{n}}}\sum_{y_{n-1},y_{n-2},\dots,y_{0}\in\{0,1\}} (-1)^{x_{n-1}y_{n-1}+x_{n-2}y_{n-2}+\dots+x_{0}y_{0}}$$

$$\times |y_{n-1}y_{n-2}\dots y_{0}\rangle$$

$$= \frac{1}{\sqrt{2^{n}}}\sum_{y=0}^{2^{n}-1} (-1)^{x \cdot y}|y\rangle, \qquad (5.7)$$

where $x \cdot y = x_{n-1}y_{n-1} \oplus x_{n-2}y_{n-2} \oplus \ldots \oplus x_0y_0$.

Coming back to step 4:

4. The Walsh-Hadamard transformation (4.11) is applied on the first n qubits next. We obtain

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$$= \frac{1}{2^{n}} \left(\sum_{x,y=0}^{2^{n}-1} (-1)^{f(x)} (-1)^{x \cdot y} |y\rangle \right) \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle).$$

As we will see, this operation will make the different terms interfere in order to read the desired result

5. The first *n* qubits are measured. Suppose f(x) is constant. Then $|\psi_3\rangle$ is put in the form

$$|\psi_3\rangle = \frac{1}{2^n} \sum_{x,y} (-1)^{x \cdot y} |y\rangle \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle)$$

up to an overall phase. Now let us consider the summation

$$\frac{1}{2^n} \sum_{x=0}^{2^n - 1} (-1)^{x \cdot y}$$

with a fixed $y \in S_n$. Clearly it vanishes since $x \cdot y$ is 0 for half of x and 1 for the other half of x unless y = 0. Therefore the summation yields δ_{y0} . Now the state reduces to

$$|\psi_3\rangle = |0\rangle^{\otimes n} \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle),$$

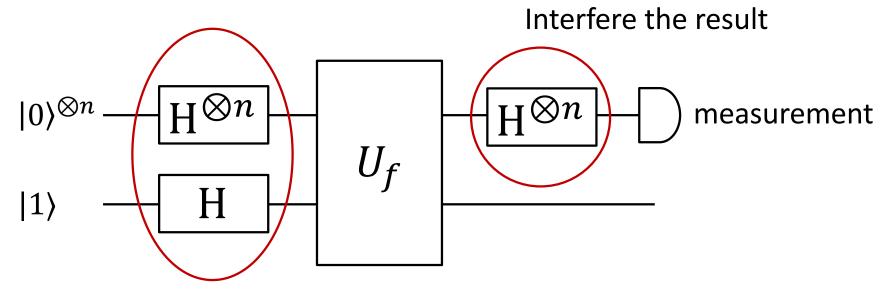
and the measurement outcome of the first n qubits is always 00...0.

Example with 3 qubits.		
Take y = 110. Then		
$\mathbf{x} \cdot \mathbf{y} = \mathbf{x}_2 \oplus \mathbf{x}_1$		
	1	
X	$x_2 \oplus x_1$	
000	0	
001	0	
010	1	
011	1	
100	1	
101	1	
110	0	
111	0	

Suppose f(x) is balanced next. The probability amplitude of $|y = 0\rangle$ in $|\psi_3\rangle$ is proportional to

$$\sum_{x=0}^{2^{n}-1} (-1)^{f(x)} (-1)^{x \cdot 0} = \sum_{x=0}^{2^{n}-1} (-1)^{f(x)} = 0.$$

Therefore the probability of obtaining measurement outcome 00...0 for the first *n* qubits vanishes. In conclusion, the function *f* is constant if we obtain 00...0 upon the meaurement of the first *n* qubits in the state $|\psi_3\rangle$, and it is balanced otherwise.



Calculate the function on both input values simultaneously

Bernstein-Vazirani Algorithm

The **Bernstein-Vazirani algorithm** is a special case of the Deutsch-Jozsa algorithm, in which f(x) is given by $f(x) = c \cdot x$, where $c = c_{n-1}c_{n-2} \dots c_0$ is an *n*-bit binary number [4]. Our aim is to find *c* with the smallest number of evaluations of *f*. If we apply the Deutsch-Jozsa algorithm with this *f*, we obtain

$$|\psi_3\rangle = \frac{1}{2^n} \left[\sum_{x,y=0}^{2^n - 1} (-1)^{c \cdot x} (-1)^{x \cdot y} |y\rangle \right] \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle).$$

Let us fix y first. If we take y = c, we obtain

$$\sum_{x} (-1)^{c \cdot x} (-1)^{x \cdot c} = \sum_{x} (-1)^{2c \cdot x} = 2^{n}.$$

Bernstein-Vazirani Algorithm

If $y \neq c$, on the other hand, there will be the same number of x such that $c \cdot x = 0$ and x such that $c \cdot x = 1$ in the summation over x and, as a result, the probability amplitude of $|y \neq c\rangle$ vanishes. By using these results, we end up with

$$|\psi_3\rangle = |c\rangle \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle). \tag{5.9}$$

We are able to tell what c is by measuring the first n qubits.

Exercise

EXERCISE 5.1 Let us take n = 2 for definiteness. Consider the following cases and find the final wave function $|\psi_3\rangle$ and evaluate the measurement outcomes and their probabilities for each case.

(1) $f(x) = 1 \ \forall x \in S_2$. (2) f(00) = f(01) = 1, f(10) = f(11) = 0. (3) f(00) = 0, f(01) = f(10) = f(11) = 1. (This function is neither constant nor balanced.)

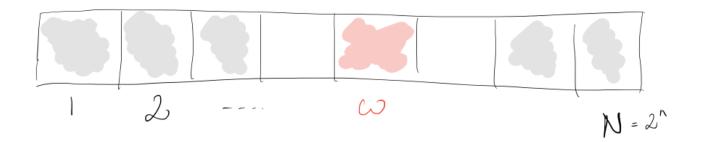
EXERCISE 5.2 Consider the Bernstein-Vazirani algorithm with n = 3 and c = 101. Work out the quantum circuit depicted in Fig. 5.2 to show that the measurement outcome of the first three qubits is c = 101.

Suppose there is a stack of $N = 2^n$ files, randomly placed, that are numbered by $x \in S_n \equiv \{0, 1, \dots, N-1\}$. Our task is to find an algorithm which picks out a particular file which satisfies a certain condition.

In mathematical language, this is expressed as follows. Let $f: S_n \to \{0, 1\}$ be a function defined by

$$f(x) = \begin{cases} 1 \ (x = z) \\ 0 \ (x \neq z), \end{cases}$$
(7.1)

where z is the address of the file we are looking for. It is assumed that f(x) is *instantaneously* calculable, such that this process does not require any computational steps. A function of this sort is often called an oracle as noted in Chapter 5. Thus, the problem is to find z such that f(z) = 1, given a function $f: S_n \to \{0, 1\}$ which assumes the value 1 only at a single point.



Clearly we have to check one file after another in a classical algorithm, and it will take O(N) steps on average. It is shown below that it takes only $O(\sqrt{N})$ steps with Grover's algorithm. This is accomplished by *amplifying* the amplitude of the vector $|z\rangle$ while cancelling that of the vectors $|x\rangle$ $(x \neq z)$.

We first needs to implement the function f(x) quantum mechanically. We define U_f as follows (**oracle**)

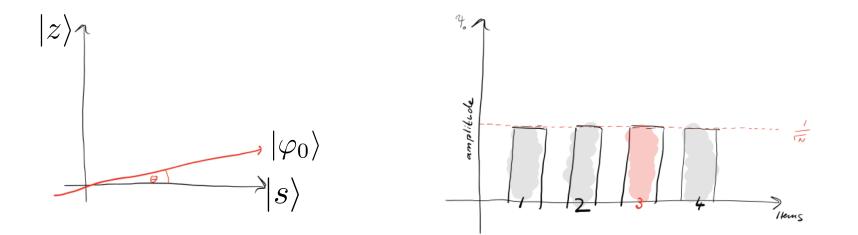
$$U_f|x\rangle = (-1)^{f(x)}|x\rangle$$

On the computational basis. We see that if x is an unmarked item, the oracle does nothing to the state. It flips the phase for the marked item. It is easy to see that

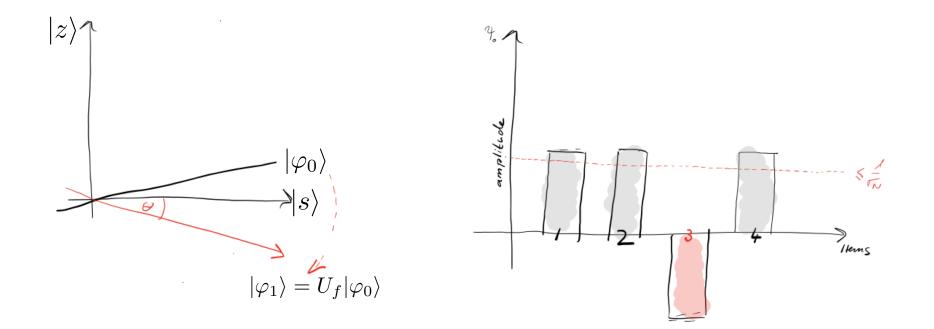
$$U_f = I - 2|z\rangle\langle z|$$

Step 1: Create an initially **equal weighted superposition** of all states (this is done with N Hadamard gates):

$$|\varphi_0\rangle = \frac{1}{\sqrt{N}} \sum_{x=0}^{N-1} |x\rangle.$$



Step 2: Apply the oracle U_f. Geometrically this corresponds to a reflection of the state $|z\rangle$ about $|s\rangle$. This transformation means that the amplitude in front of the $|z\rangle$ state becomes negative, which in turn means that the average amplitude has been lowered.



Grover's search algorithm Step 3: Apply the gate

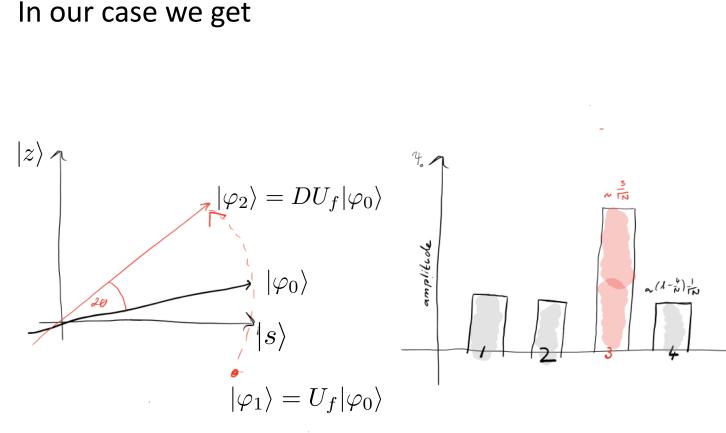
$$D = -I + 2|\varphi_0\rangle\langle\varphi_0|.$$

The action of the gate is the following

$$\begin{bmatrix} \omega_x - \bar{\omega} \\ \bar{\omega}$$

$$\begin{aligned} |\varphi\rangle &= \sum_{x=0}^{N-1} \omega_x |x\rangle \ \to \ D|\varphi\rangle = \left[\frac{2}{N} \sum_{x,y=0}^{N-1} |x\rangle \langle y| \right] \sum_{z=0}^{N-1} \omega_z |z\rangle - \sum_{x=0}^{N-1} \omega_x |x\rangle \\ &= \frac{2}{N} \left[\sum_{x=0}^{N-1} |x\rangle \right] \left[\sum_{y=0}^{N-1} \omega_y \right] - \sum_{x=0}^{N-1} \omega_x |x\rangle = \sum_{x=0}^{N-1} (2\bar{\omega} - \omega_x) |x\rangle \end{aligned}$$

with
$$\bar{\omega} = \frac{1}{N} \sum_{x=0}^{N-1} \omega_x$$
 average



Since the average amplitude has been lowered by the first reflection, this transformation boosts the negative amplitude of $|z\rangle$ to roughly three times its original value, while it decreases the other amplitudes.

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Step 3: go to step 2 an repeat the application of U_f and D a sufficient number of times. Let us call $G_f = D U_f$.

PROPOSITION 7.2 Let us write

$$G_f^k |\varphi_0\rangle = a_k |z\rangle + b_k \sum_{x \neq z} |x\rangle \tag{7.17}$$

with the initial condition

$$a_0 = b_0 = \frac{1}{\sqrt{N}}.$$

Then the coefficients $\{a_k, b_k\}$ for $k \ge 1$ satisfy the recursion relations

$$a_k = \frac{N-2}{N}a_{k-1} + \frac{2(N-1)}{N}b_{k-1},$$
(7.18)

$$b_k = -\frac{2}{N}a_{k-1} + \frac{N-2}{N}b_{k-1} \tag{7.19}$$

for k = 1, 2, ...

Proof. It is easy to see the recursion relations are satified for k = 1Let $G_f^{k-1}|\varphi_0\rangle = a_{k-1}|z\rangle + b_{k-1}\sum_{x\neq z}|x\rangle$. Then $G_f^k |\varphi_0\rangle = G_f \left(a_{k-1} |z\rangle + b_{k-1} \sum_{i} |x\rangle \right)$ $= \left(-I + 2|\varphi_0\rangle\langle\varphi_0|\right) \left(-a_{k-1}|z\rangle + b_{k-1}\sum_{x\neq z}|x\rangle\right)$ $= -b_{k-1} \sum_{n \neq n} |x\rangle + a_{k-1} |z\rangle + \frac{2}{\sqrt{N}} (N-1)b_{k-1} |\varphi_0\rangle - \frac{2a_{k-1}}{\sqrt{N}} |\varphi_0\rangle$ $= -b_{k-1}\sum_{x \neq z} |x\rangle + a_{k-1}|z\rangle + \frac{2}{N}(N-1)b_{k-1}\sum_{x \neq z} |x\rangle - \frac{2a_{k-1}}{N}\sum_{x \neq z} |x\rangle$ $= \left[\frac{N-2}{N}a_{k-1} + \frac{2(N-1)}{N}b_{k-1}\right]|z\rangle + \left[-\frac{2}{N}a_{k-1} + \frac{N-2}{N}b_{k-1}\right]\sum_{k}|x\rangle,$

and proposition is proved.

PROPOSITION 7.3 The solutions of the recursion relations in Proposition 7.2 are explicitly given by

$$a_k = \sin[(2k+1)\theta], \quad b_k = \frac{1}{\sqrt{N-1}}\cos[(2k+1)\theta], \quad (7.20)$$

for k = 0, 1, 2, ..., where

$$\sin\theta = \sqrt{\frac{1}{N}}, \quad \cos\theta = \sqrt{1 - \frac{1}{N}}.$$
(7.21)

Proof. Let $c_k = \sqrt{N-1}b_k$. The recursion relations (7.18) and (7.19) are written in a matrix form,

$$\begin{pmatrix} a_k \\ c_k \end{pmatrix} = M \begin{pmatrix} a_{k-1} \\ c_{k-1} \end{pmatrix}, \ M = \begin{pmatrix} (N-2)/N & 2\sqrt{N-1}/N \\ -2\sqrt{N-1}/N & (N-2)/N \end{pmatrix} = \begin{pmatrix} \cos 2\theta & \sin 2\theta \\ -\sin 2\theta & \cos 2\theta \end{pmatrix}$$

Note that M is a rotation matrix in \mathbb{R}^2 , and its kth power is another rotation matrix corresponding to a rotation angle $2k\theta$. Thus the above recursion relation is easily solved to yield

$$\begin{pmatrix} a_k \\ c_k \end{pmatrix} = M^k \begin{pmatrix} a_0 \\ c_0 \end{pmatrix} = \begin{pmatrix} \cos 2k\theta & \sin 2k\theta \\ -\sin 2k\theta & \cos 2k\theta \end{pmatrix} \begin{pmatrix} \sin \theta \\ \cos \theta \end{pmatrix} = \begin{pmatrix} \sin[(2k+1)\theta] \\ \cos[(2k+1)\theta] \end{pmatrix}.$$

Replacing c_k by b_k proves the proposition.

We have proved that the application of $G_f\ k$ times on $|\varphi_0\rangle$ results in the state

$$G_f^k |\varphi_0\rangle = \sin[(2k+1)\theta]|z\rangle + \frac{1}{\sqrt{N-1}}\cos[(2k+1)\theta]\sum_{x\neq z} |x\rangle.$$
 (7.22)

Measurement of the state $U_f^k |\varphi_0\rangle$ yields $|z\rangle$ with the probability

$$P_{z,k} = \sin^2[(2k+1)\theta]. \tag{7.23}$$

STEP 4 Our final task is to find the k that maximizes $P_{z,k}$. A rough estimate for the maximizing k is obtained by putting

$$(2k+1)\theta = \frac{\pi}{2} \to k = \frac{1}{2}\left(\frac{\pi}{2\theta} - 1\right).$$
 (7.24)

PROPOSITION 7.4 Let $N \gg 1$ and let

$$m = \left\lfloor \frac{\pi}{4\theta} \right\rfloor,\tag{7.25}$$

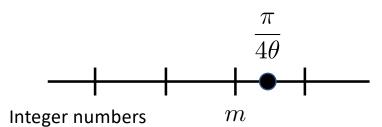
where $\lfloor x \rfloor$ stands for the floor of x. The file we are searching for will be obtained in $G_f^m |\varphi_0\rangle$ with the probability

$$P_{z,m} \ge 1 - \frac{1}{N} \tag{7.26}$$

and



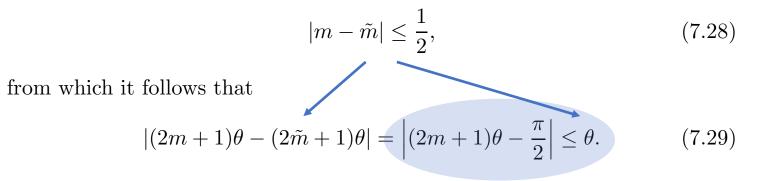
This is the number of times we repeat the algorithm, which grows with the square root of N



Proof. Equation (7.25) leads to the inequality $\pi/4\theta - 1 < m \le \pi/4\theta$. Let us define \tilde{m} by

$$(2\tilde{m}+1)\theta = \frac{\pi}{2} \to \tilde{m} = \frac{\pi}{4\theta} - \frac{1}{2}.$$

Observe that m and \tilde{m} satisfy



Considering that $\theta \sim 1/\sqrt{N}$ is a small number when $N \gg 1$ and $\sin x$ is monotonically increasing in the neighborhood of x = 0, we obtain

$$0 < \sin |(2m+1)\theta - \pi/2| < \sin \theta$$

$$\Rightarrow = \cos[(2m+1)\theta]$$

$$\cos^{2}[(2m+1)\theta] \le \sin^{2}\theta = \frac{1}{N}.$$
(7.30)

Thus it has been shown that

or

$$P_{m,z} = \sin^2[(2m+1)\theta] = 1 - \cos^2[(2m+1)\theta] \ge 1 - \frac{1}{N}.$$
 (7.31)

It also follows from $\theta > \sin \theta = 1/\sqrt{N}$ that

$$m = \left\lfloor \frac{\pi}{4\theta} \right\rfloor \le \frac{\pi}{4\theta} \le \frac{\pi}{4}\sqrt{N}.$$
(7.32)

It is important to note that this quantum algorithm takes only $O(\sqrt{N})$ steps and this is much faster than the classical counterpart which requires O(N) steps.

Next we will show how how to implement the gates