

Carl-Gustaf Arvid Rossby

Born: December 28, 1898 Stockholm, Sweden

Massachusetts Institute of Technology

University of Chicago

Woods Hole Oceanographic Institution

 $\Omega = 1^{1} \cdot 1 \cdot N f + \dots \cdot 1 = 1 = 1$

Rossby Waves – as seen by Rossby



Platzman 1968



Palmen 1949

The picture is taken from above the South Pole, shows a number of mid latitude cyclones circling Antarctica.



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they form cyclones of cold air.

Rossby waves (hydrodynamics)

Momentum equation in rotating frame (with angular velocity Ω)

$$\frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} = -\frac{1}{\rho} \nabla p - 2\mathbf{\Omega} \times \mathbf{u} + \mathbf{g},$$

where $-2\rho \Omega \times u$ is the Coriolis force.

Ratio of convective and Coriolis terms is called a Rossby number

$$Ro = \frac{U}{L\Omega}$$

When $Ro \ll 1$ then the rotation effects are significant.

Rectangular coordinates:



 $f = 2\Omega \sin \vartheta$ is the Coriolis parameter.

 $\vartheta = 90^{\circ} - \theta$ is the latitude.

These equations can be cast into one equation

$$\frac{\partial}{\partial t} \left[\frac{1}{c^2} \left(\frac{\partial^2}{\partial t^2} + f^2 \right) - \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right) \right] u_y - \frac{\partial f}{\partial y} \frac{\partial u_y}{\partial x} = 0.$$

 $c = \sqrt{gH}$ is the surface gravity speed.

If one neglects the surface elevation $h \approx 0$ or $\frac{h}{H} << 1$

$$\frac{\partial}{\partial t} \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right) u_y + \frac{\partial f}{\partial y} \frac{\partial u_y}{\partial x} = 0.$$

This approximation eliminates surface gravity (or Poincare) waves and induces small change in Rossby wave dispersion relation. At this point came up Rossby with his β – plane approximation

When spatial scales of considered process is less than sphere radius then one can expand the Coriolis parameter at a given latitude as

$$f = f_0 + \beta y,$$

$$\beta = \frac{\partial f}{\partial y} = \frac{2\Omega}{R} \cos \vartheta = const.$$

$$\frac{\partial}{\partial t} \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right) u_y + \beta \frac{\partial u_y}{\partial x} = 0.$$

Fourier analysis of the form $exp(-i\omega t + ik_x x + ik_y y)$ leads to the dispersion relation of Rossby (or planetary) waves

$$\widetilde{\omega} = -\frac{\beta k_x}{k_x^2 + k_y^2}.$$

Rossby waves always propagate in the opposite direction of rotation!

For purely toroidal propagation
$$\widetilde{\omega} = -\frac{1}{k}$$

Phase speed $v_{ph} = \frac{\widetilde{\omega}}{k_x} = -\frac{\beta}{k_x^2 + k_y^2}$.

Long wavelength waves propagate faster!

Group speed
$$\mathbf{v}_{\mathbf{g}} = \left(\frac{\partial \widetilde{\omega}}{\partial k_x}, \frac{\partial \widetilde{\omega}}{\partial k_y}\right) = \left(-\beta \frac{k_y^2 - k_x^2}{\left(k_x^2 + k_y^2\right)^2}, \beta \frac{2k_x k_y}{\left(k_x^2 + k_y^2\right)^2}\right)$$

If there is a constant zonal flow then the phase speed can be written as (Rossby 1939)

$$c = U - \frac{\beta L^2}{4\pi^2},$$

It appears that the waves become stationary when

$$c = U - \frac{\beta L_{s}^{2}}{4\pi^{2}} = 0, \quad L_{s} = 2\pi \sqrt{\frac{U}{\beta}}.$$
$$c = U \left(1 - \frac{L^{2}}{L_{s}^{2}}\right),$$

Long wavelength waves propagate westward and short wavelength waves propagate eastward!

 β =1.6 10⁻¹¹ m⁻¹ s⁻¹ at mid-latitudes.

For the wavelength of 10000 km, one gets the period of 5.6 days.

Phase speed - 20 m/s.

The observed Rossby wave period on the Earth is 4-6 days (Yanai and Maruyama 1966, Wallace 1973, Madden 1979).

The ratio of Rossby wave and Earth rotation periods is around 6!