

MECCANICA RAZIONALE

legge circale & Ambiensole
Nazionale

20 Aprile 2021

Dinamica

\rightarrow

$$\frac{d}{dt} \left(\frac{\partial k}{\partial \dot{q}_i} \right) - \frac{\partial k}{\partial q_i} = Q_i$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_i} \right) - \frac{\partial L}{\partial q_i} = 0$$

$$\text{con } L = (k - V)$$

$$\left\{ \begin{array}{l} R^i = \frac{\partial}{\partial \dot{q}_i} P \\ M^e(O) = \end{array} \right.$$

$$M^e(O) = \frac{\partial}{\partial \dot{q}_i} L(O) + \omega \times I^2$$

$$P = \sum_{B \in S} m_B v_B = M v_G$$

$$L(O) = \sum_{B \in S} (x_B - x_0) \times m_B v_B$$

Abbia uno studio sulla $\underline{L}(0)$

per un corpo rigido

K per un rigido

$$\cdot K = \frac{1}{2} M v_G^2 + \frac{1}{2} \underline{\omega} \cdot I_G(\underline{\omega}) \quad \text{generale}$$

$$\rightarrow \cdot K = \frac{1}{2} \underline{\omega} \cdot I_0(\underline{\omega}) \quad \text{o fissa}$$

$$\cdot K = \frac{1}{2} I_r \omega^2 \quad \text{e fissa}$$

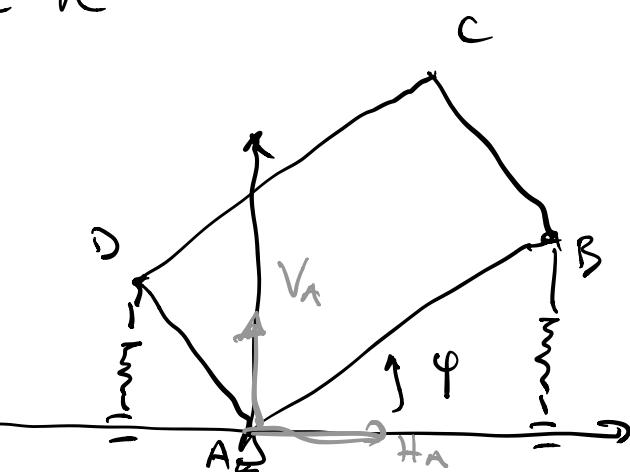
$$\cdot K = \frac{1}{2} M v_G^2 \quad \text{freddo}$$

$\underline{L}(0)$ per un rigido

$$\cdot O \in R \quad \underline{L}(0) = M(\underline{x}_0 - \underline{x}_0) \wedge \underline{v}_0 + I_0(\underline{\omega})$$

$$\cdot O \notin R \quad \underline{L}(0) = \underline{L}(A) + (\underline{x}_A - \underline{x}_0) \wedge M_A$$

$$A \in R$$



Esercizio

$$\overline{AB} = 2l$$

$$\overline{BC} = l$$

la cui rete tangolare
essere

Eq. und mto

$$K = \frac{1}{2} I_{A,3} \dot{\varphi}^2 = \frac{1}{2} \left[\frac{m}{3} (l^2 + a l^2) \dot{\varphi}^2 \right]$$
$$= \frac{1}{2} \left(\frac{5}{3} m l^2 \dot{\varphi}^2 \right)$$
$$\frac{m (a^2 + b^2)}{3}$$

$$V = \frac{c}{2} \bar{y}_B^2 + \frac{c}{2} \bar{y}_D^2 = \frac{c}{2} (2l)^2 \sin^2 \varphi + \frac{c}{2} l^2 \cos^2 \varphi$$
$$= \frac{c}{2} l^2 (3 \sin^2 \varphi + 1)$$

$$L = \frac{1}{2} \left(\frac{5}{3} m l^2 \dot{\varphi} \right) - \frac{c}{2} l^2 (3 \sin^2 \varphi + 1)$$

$$\boxed{\frac{d}{dt} \frac{\partial L}{\partial \dot{\varphi}} - \frac{\partial L}{\partial \varphi} = 0}$$

$$\frac{5}{3} m l^2 \ddot{\varphi} + 3 c l^2 \sin \varphi \cos \varphi = 0$$

condizioni iniziali

$$\begin{cases} \varphi(0) = \varphi_0 = \frac{\pi}{4} \\ \dot{\varphi}(0) = \omega_0 \end{cases}$$

Energie: $\bar{E} = K + V$

$$\begin{cases} = \frac{1}{2} \left(\frac{5}{3} m l^2 \dot{\varphi}^2 \right) + \frac{c}{2} l^2 (3 \sin^2 \varphi + 1) \\ = \frac{1}{2} \left(\frac{5}{3} m l^2 \omega_0^2 \right) + \frac{c}{2} l^2 \left(\frac{3}{2} + 1 \right) \end{cases}$$

Reaktion: Vinculon in A

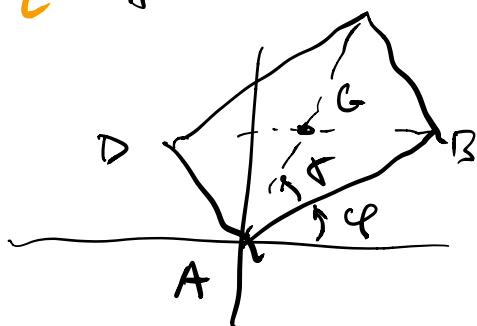
Additions $H_A, V_A \rightarrow ECD$ per la risultante

$$R = H_A e_1 + V_A e_2 - y_B e_2 - y_D e_2$$

$$= \text{P} \quad (\text{da } \ddot{x}_G) \rightarrow \text{in } \ddot{x}_G \leftarrow \text{in } \ddot{y}_G e_2$$

$$H_A = m \ddot{x}_G$$

$$V_A = c \cdot l \sin \varphi + l \cos \varphi + m \ddot{y}_G$$



$$\begin{matrix} x_G \\ y_G \end{matrix} \rightarrow \begin{matrix} \ddot{x}_G \\ \ddot{y}_G \end{matrix}$$

$$\left\{ \begin{array}{l} x_G = \overline{AG} \cos(\varphi + \psi) \\ y_G = \overline{AG} \sin(\varphi + \psi) \end{array} \right.$$

$$\overline{AG} = \frac{1}{2} \sqrt{l^2 + c^2 l^2} = \frac{l}{2} \sqrt{5} \quad \overline{AC} = \sqrt{l^2 + c l^2} = l \sqrt{5}$$

$$\therefore \sin \gamma = \frac{\overline{BC}}{\overline{AC}} = \frac{1}{\sqrt{5}}, \quad \omega \cdot \gamma = \frac{\overline{AB}}{\overline{AC}} = \frac{2}{\sqrt{5}}$$

$$\left\{ \begin{array}{l} \dot{x}_G = - \overline{AG} \sin(\gamma + \varphi(\tau)) \dot{\varphi}(\tau) \\ \dot{y}_G = \overline{AG} \cos(\gamma + \varphi(\tau)) \dot{\varphi}(\tau) \end{array} \right.$$

$$\begin{cases} \ddot{x}_G = -\bar{AG} \sin(\gamma + \varphi(t)) \ddot{\varphi}(t) - \bar{AG} \cos(\gamma + \varphi(t)) \dot{\varphi}(t)^2 \\ \ddot{y}_G = \bar{AG} \cos(\gamma + \varphi(t)) \ddot{\varphi}(t) - \bar{AG} \sin(\gamma + \varphi(t)) \dot{\varphi}(t)^2 \end{cases}$$

Usiamo le di. mōs & conservazione
dell' energia \rightarrow relazioni in funzione
di φ

$$\begin{aligned} \frac{1}{2} \left(\frac{r}{3} \omega t^2 \dot{\varphi}^2 \right) + \frac{c}{2} l^2 (3 \sin^2 \varphi + 1) &= E \\ \rightarrow \left[\dot{\varphi}^2 = \frac{E}{\frac{r}{3} \omega t^2} - \frac{1}{3} \frac{c}{\omega t^2} \left[E - c \frac{l^2}{2} (3 \sin^2 \varphi + 1) \right] \right. \\ &\quad \left. =: f(\varphi) \right] \end{aligned}$$

$$\text{Eq. mōs } \frac{r}{3} \omega t^2 \ddot{\varphi} + 3 \frac{c}{2} l^2 \sin \varphi \cos \varphi \dot{\varphi}^2 = 0$$

$$\boxed{\ddot{\varphi} = - \frac{9}{5} \frac{c}{\omega} \sin \varphi \cos \varphi =: g(\varphi)}$$

$$H_A = -m \bar{AG} \sin(\gamma + \varphi) g(\varphi) - \bar{AG} m \cos(\gamma + \varphi) f(\varphi)$$

$$\begin{aligned} V_A &= 2cl \sin \varphi + cl \cos \varphi + \\ &\quad + m \bar{AC} \left[\cos(\gamma + \varphi) g(\varphi) - \sin(\gamma + \varphi) f(\varphi) \right] \end{aligned}$$

\rightarrow abbiamo H_A, V_A in funzione di φ

Uttiene queste espressioni per fissare
un limite per H_A , V_A .

Potremmo trovare max, ma è
 $\omega_1 \in [-1, 1]$

ad esempio $|\ddot{\varphi}| \leq \frac{7}{5} \frac{c}{m}$ ---

(suggerito)

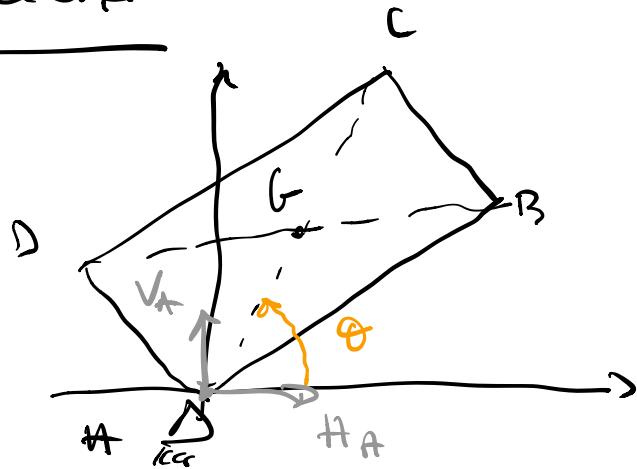
Allo stesso modo $\dot{\varphi}^2 \leq$ valore max

Troviamo

$|H_A| \leq$ esprimere che
non dipende

$|V_A| \leq$ più da φ

Esecuzio



$\downarrow \varphi$

lavoro
rettangolare
presso verticale

1. eq. di moto

2. reazioni in A

$$\text{per } \left[\theta = -\frac{\pi}{2} \right]$$

$$\Theta(0) = 0, \dot{\Theta}(0) = 0$$

$$\Gamma \theta = \varphi + \dot{\varphi}$$

es
prevede

$$Eq. \text{ di moto } \ddot{\theta} = k - V$$

$$k = \frac{1}{2} I_{A,3} \ddot{\theta} = \frac{1}{2} \left(\frac{5}{3} m l^2 \right) \ddot{\theta}$$

$\frac{1}{3} m (a^2 + b^2)$

$$V = mg AG \sin \theta$$

$$\bar{AG} = \frac{1}{2} \sqrt{4l^2 + l^2} = \frac{\sqrt{5}}{2} l$$

$$Allora \quad \ddot{\theta} = k - V$$

$$= \frac{1}{2} \left(\frac{5}{3} m l^2 \right) \ddot{\theta} - mg \frac{\sqrt{5}}{2} l \sin \theta$$

$$\frac{d}{d\theta} \left(\frac{\partial L}{\partial \dot{\theta}} \right) - \frac{\partial L}{\partial \theta} = \frac{5}{3} m l^2 \ddot{\theta} + mg \frac{\sqrt{5}}{2} l \cos \theta = 0$$

2) Reazioni vincolari in A : G.C.D.

$$\underline{R} = H_A \underline{e}_1 + V_A \underline{e}_2 - mg \underline{e}_3 = \underline{P}$$

$$\begin{cases} H_A = m \ddot{x}_G \\ V_A = m \ddot{y}_G + mg \end{cases} \quad \begin{aligned} x_G &= \bar{AG} \cos \theta(\tau) \\ y_G &= \bar{AG} \sin \theta(\tau) \end{aligned}$$

$$\ddot{x}_G = \frac{d}{d\theta} \left(- \bar{AG} \sin \theta(\tau) \dot{\theta}(\tau) \right) =$$

$$= - \bar{AG} \cos \theta(\tau) \ddot{\theta}(\tau) - \bar{AG} \sin \theta(\tau) \dot{\theta}(\tau)$$

$$\ddot{\theta}_G = \frac{d}{dt} \left(\bar{A} \bar{G} \cos \theta(t) \quad \dot{\theta}(t) \right) =$$

$$= - \bar{A} \bar{G} \underbrace{\sin \theta(t)}_{\uparrow} \dot{\theta}(t)^2 + \bar{A} \bar{G} \cos \theta(t) \ddot{\theta}(t)$$

Udiamo le conservazione dell' energia

$$E = K + V = \frac{1}{2} \left(\frac{5}{3} m l^2 \dot{\theta}^2 \right) + mg \frac{l \sqrt{5}}{2} \sin \theta$$

$$= E|_{\theta=0} = 0 + 0 = 0$$

condizioni iniziali $\dot{\theta}(0) = 0$
 $\theta(0) = 0$

$$\hookrightarrow \dot{\theta}^2 = - \frac{g}{l} \frac{3}{\sqrt{5}} \sin \theta$$

Dall' eq. di moto

$$\ddot{\theta} = - \frac{3\sqrt{5}}{10} \frac{g}{l} e^{-\theta}$$

Audiamo a vedere $H_A \in \mathcal{V}_A$ per

$$\theta = -\frac{\pi}{2}$$

$$\dot{\theta}^2 \Big|_{\theta=-\frac{\pi}{2}} = + \frac{3}{\sqrt{5}} \frac{g}{l}$$

$$\dot{\theta} \Big|_{\theta = -\frac{\pi}{2}} = 0$$

$$\ddot{x}_G = 0$$

Per $\theta = -\frac{\pi}{2}$

$$\begin{aligned}\ddot{y}_G &= -\left(\frac{l}{2}\sqrt{3}\right) \Leftrightarrow \left(\frac{3}{4}\frac{\pi}{l}\right) \\ &= \frac{3}{2}\frac{\pi}{l}\end{aligned}$$

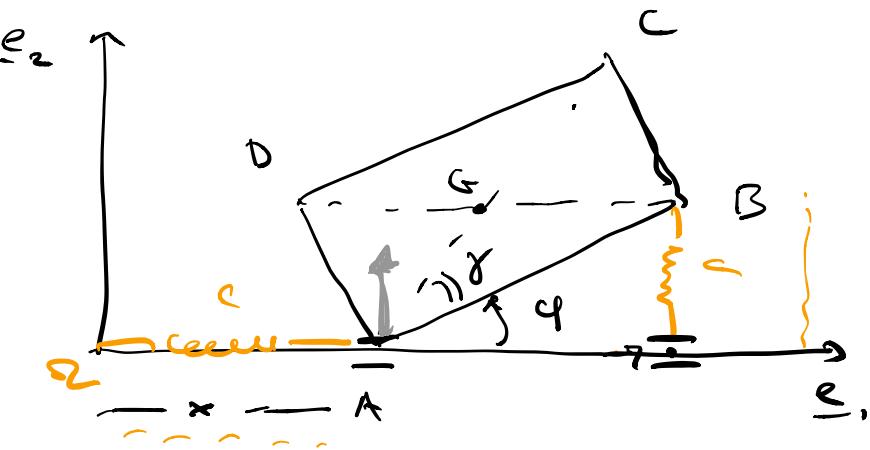
Per finire

$$\left\{ \begin{array}{l} H_A = 0 \\ V_A = mg + m\frac{3}{2}g = \frac{5}{2}mg \end{array} \right.$$

\Leftrightarrow
stabilità
dalla
dimensione (\ddot{y}_G)

Seconda parte

Esercizio



lavoro

rettangolare

molla : C

coordinate

libre

(x, q)

Eq. di moto

$$K = \frac{1}{2} m v_G^2 + \frac{1}{2} I_{G,3} \omega^2$$

$$\omega = \dot{\varphi} \stackrel{e_3}{=} \quad I_{G,3} = \frac{M}{12} (a^2 + b^2) = \frac{M s}{12} l^2$$

$\frac{M}{12} (a^2 + b^2)$

$$\underline{v}_G = \frac{d}{dt} \underline{x}_G$$

$$\underline{x}_G = \left[\underline{x}_{(t)} + \overline{AG} \cos(\gamma + \varphi) \right] e_1$$

↑

$$+ \overline{AG} \sin(\gamma + \varphi) e_2$$

↑

dove $\overline{AG} = l \frac{\sqrt{s}}{2}$, $\sin \gamma = \frac{1}{\sqrt{s}}$, $\cos \gamma = \frac{2}{\sqrt{s}}$

$$\underline{v}_G = \frac{d}{dt} \underline{x}_G = \left[\dot{\underline{x}}_{(t)} - \overline{AG} \sin(\gamma + \varphi) \dot{\varphi} \right] e_1$$

$$+ \overline{AG} \cos(\gamma + \varphi) \dot{\varphi} e_2$$

$$\underline{v}_G^2 = \underline{v}_G \cdot \underline{v}_G = \left(\dot{\underline{x}}_{(t)} - \overline{AG} \sin(\gamma + \varphi) \dot{\varphi} \right)^2$$

+

$$\left(\overline{AG} \cos(\gamma + \varphi) \dot{\varphi} \right)^2$$

$$= \dot{\underline{x}}_{(t)}^2 + \overline{AG}^2 \dot{\varphi}^2 \left(\sin^2(\gamma + \varphi) + \cos^2(\gamma + \varphi) \right)$$

$$= 2 \dot{x}(t) \overline{AG} \dot{\varphi}(t) \sin(\gamma + \varphi)$$

Koeffizienten für die Welle

$$K = \frac{1}{2} m \dot{x}_0^2 + \frac{1}{2} I_{0,0} \omega^2 =$$

$$= \frac{1}{2} m \left(\dot{x}^2 + \frac{\dot{\varphi}^2 \frac{5}{4} l^2 - l \sqrt{5} \dot{x} \dot{\varphi} \sin(\gamma + \varphi)}{+ \dot{\varphi}^2 \frac{5}{12} l^2} \right) \quad \frac{5}{4} + \frac{5}{12} = \frac{5}{3}$$

$$V = \frac{c}{2} x^2 + \frac{c}{2} (2l \sin \varphi)^2 = \frac{c}{2} x^2 + 2cl^2 \sin^2 \varphi$$

$$L = K - V = \frac{1}{2} m \left(\dot{x}^2 + \frac{5}{3} l^2 \dot{\varphi}^2 + - l \sqrt{5} \dot{x} \dot{\varphi} \sin(\gamma + \varphi) \right) - \left(\frac{c}{2} x^2 + 2cl^2 \sin^2 \varphi \right)$$

eq. der Bewegung

$$\Rightarrow \frac{d}{dt} \frac{\partial L}{\partial \dot{x}} - \frac{\partial L}{\partial x} = \frac{\partial L}{\partial \dot{x}} = \frac{\partial}{\partial \dot{x}} \left(\frac{1}{2} m \dot{x}^2 \right) = m \ddot{x}$$

$$= \frac{d}{dt} \left[m \ddot{x} - \frac{1}{2} m l \sqrt{5} \dot{\varphi} \sin(\gamma + \varphi) \right] + C_x$$

$$= \boxed{m \ddot{x} - \frac{1}{2} m l \sqrt{5} \dot{\varphi} \sin(\gamma + \varphi) - \frac{1}{2} m l \sqrt{5} \dot{\varphi}^2 \cos(\gamma + \varphi) + C_x = 0}$$

$$\begin{aligned}
 4) \quad & \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\varphi}} - \frac{\partial L}{\partial \varphi} \right) = 0 \\
 &= \frac{d}{dt} \left[\frac{5}{3} \ddot{\varphi} \omega l^2 - \frac{1}{2} m l \sqrt{5} \dot{x} \sin(\varphi + \psi) \right] \\
 &\quad - \left[-\frac{1}{2} m l \sqrt{5} \dot{x} \dot{\varphi} \cos(\varphi + \psi) - u c l^2 \cos \varphi \sin \varphi \right] \\
 &= \frac{5}{3} \ddot{\varphi} \omega l^2 - \frac{1}{2} m l \sqrt{5} \dot{x} \sin(\varphi + \psi) + \\
 &\quad - \frac{1}{2} m l \sqrt{5} \dot{x} \dot{\varphi} \cos(\varphi + \psi) + \\
 &\quad + \frac{1}{2} m l \sqrt{5} \dot{x} \dot{\varphi} \cos(\varphi + \psi) + u c l^2 \cos \varphi \sin \varphi \\
 &= \underbrace{\frac{5}{3} \omega l^2 \ddot{\varphi}}_{=} - \underbrace{\frac{1}{2} m l \sqrt{5} \dot{x} \sin(\varphi + \psi)}_{=} + u c l^2 \cos \varphi \sin \varphi = 0
 \end{aligned}$$

La Ecuacion resultante es A = \sqrt{A}

eq. cardinale de la $\ddot{\varphi}_2$

$$\begin{aligned}
 V_A - c y_B &= V_A - c 2l \sin \varphi = \omega \tilde{q} G \\
 \left(\frac{R}{L} = \frac{d}{dt} P \text{ tiempo } \ddot{\varphi}_2 \right)
 \end{aligned}$$

$$V_A = \omega l c \sin \varphi + \omega (\bar{AG} \cos(\varphi + \psi) \ddot{\varphi} - \bar{AC} \sin(\varphi + \psi) \ddot{\varphi}_2)$$

$\ddot{\varphi}, \dot{\varphi}^2 \rightarrow$ coordinate libere.

Conservazione energie \rightarrow non ci sono
perdite di energia solo

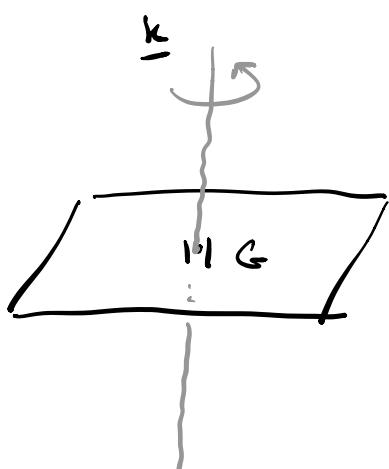
$$\ddot{\varphi}, \ddot{x}, \dot{\varphi}, \dot{x}$$

Per trovare V_A è funzione di $(\dot{x}, \dot{\varphi})$
dobbiamo risolvere le eq. del moto.

Tutto fatto

Rigido con un reale fissa

Energia



Lamine rettangolare



massa fissa

Lamele cilindriche in



Vincoli fissi

$$R = Mg + F_0^2 = P = 0$$

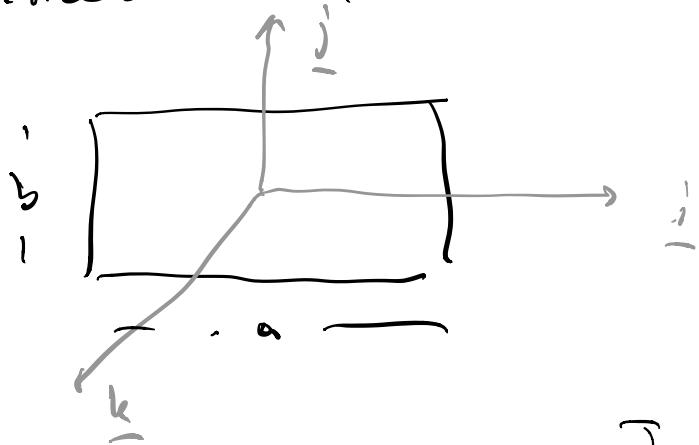
ECD :

perché G
fissa

$$M(G) = \underline{\mu}_G^2 = \frac{d}{dt} \underline{L}(G)$$

La velocità angolare $\underline{\omega} = \omega \underline{k}$

Introduciamo la Fisica



$$\underline{L}(G) = I_G(\underline{\omega})$$

$$= \omega J_3 \underline{k}$$

$$J_3 = \frac{M}{12} (a^2 + b^2)$$

$$\frac{d}{dt} \underline{L}(G) = \dot{\omega} J_3 \underline{k}$$

Ritroviamo E.L.D.

- momenti lungo \underline{k} : $J_3 \dot{\omega} = 0$

$$\rightarrow \omega(t) = \text{costante} = \omega_0$$

rotazione
uniforme

- lungo $\underline{i}, \underline{j}$: eq. dei momenti
 ω di $\underline{F}_G^z = 0$

- ribaltarsi $\underline{F}_G^z = -Mg$

Note : se c'est autre = sol

$$\text{creepus} - \nu \underline{\omega} = -\nu \underline{\omega} \underline{k}$$

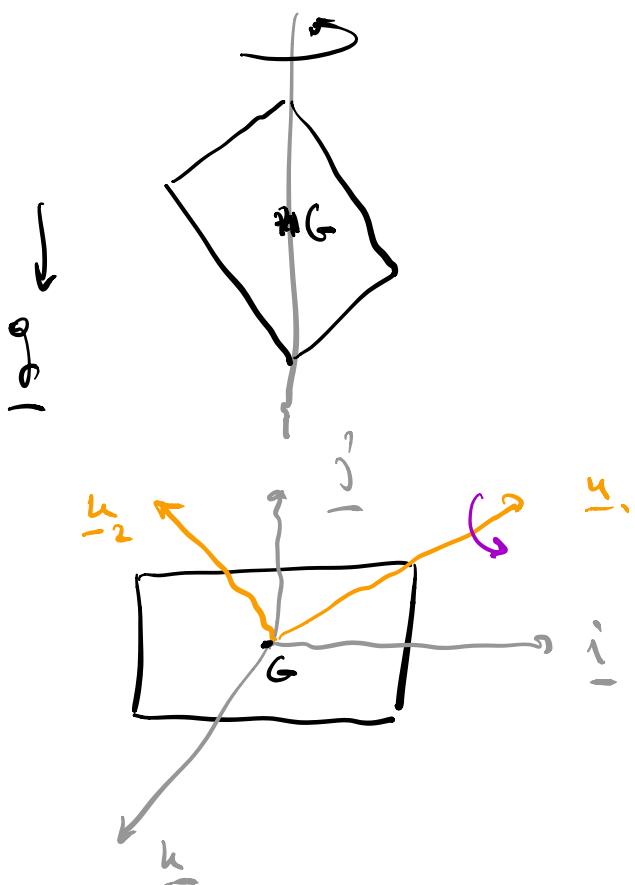
eq. moment: long. \underline{k}

$$J_3 \dot{\omega} = -\nu \omega$$

$$\hookrightarrow \omega(\tau) = \omega_0 \exp\left(-\frac{\nu}{J_3} \tau\right)$$

condition
initial

Fréqu



l'axe d. rotation
est la diagonale

Tense solide

$$S(G; \underline{u}_1, \underline{u}_2, \underline{k})$$

$$I_G = \begin{pmatrix} I_{11} & I_{12} & 0 \\ I_{21} & I_{22} & 0 \\ 0 & 0 & I_{33} \end{pmatrix}$$

$$\underline{\omega} = \omega \underline{u}_1, \quad \underline{u}_1 = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}$$

$$\underline{L}(G) = I_G(\underline{\omega}) = \underline{\omega} \cdot I_G(\underline{u}_.)$$

$$= \underline{\omega} \left(I_{11} \underline{u}_1 + I_{12} \underline{u}_2 \right)$$

Abbrievio: $\underline{\omega}$

$$\frac{d}{dt} \underline{L}(G) = I_G(\dot{\underline{\omega}}) + \underline{\omega} \wedge I_G(\underline{\omega})$$

↑
Perid-

$$= \dot{\underline{\omega}} \cdot I_G(\underline{u}_.) + \underline{\omega}^2 \underline{u}_. \wedge I_G(\underline{u}_.)$$

$$= \dot{\underline{\omega}} \cdot I_{11} \underline{u}_1 + \dot{\underline{\omega}} \cdot I_{12} \underline{u}_2 + \underline{\omega}^2 I_{12} \underline{k}$$

$$ECD: \quad \dot{\underline{M}}(G) = \dot{\underline{\mu}}_G = \frac{d}{dt} \underline{L}(G)$$

$$\left\{ \begin{array}{l} I_{11} \dot{\underline{\omega}} = 0 \quad \text{eq. di moto} \\ \underline{\mu}_G \cdot \underline{u}_2 = I_{12} \dot{\underline{\omega}} \quad \text{momenti di} \\ \underline{\mu}_G \cdot \underline{k} = I_{12} \underline{\omega}^2 \quad \text{reazione} \end{array} \right.$$

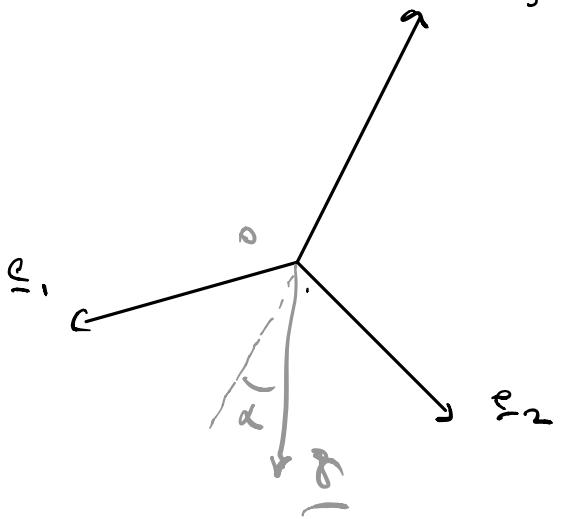
$\left[\begin{array}{l} I_{12} \neq 0 \end{array} \right]$

Ejemplo

Tensión fija

$$\sum (\underline{0}, \underline{\epsilon}_1, \underline{\epsilon}_2, \underline{\epsilon}_3)$$

$$\underline{\epsilon}_3 = \underline{u}_1$$



$$\underline{\epsilon} = g \sin \alpha \underline{\epsilon}_2 - g \cos \alpha \underline{\epsilon}_3$$

ECD

$$\left\{ \begin{array}{l} M_g + \underline{F}^r = \underline{P} \\ (\underline{\epsilon}_0 - \underline{\epsilon}_0) \times M_g + \underline{M}_0^r = \frac{d}{dt} \underline{L}(O) \end{array} \right.$$

↑ ↑

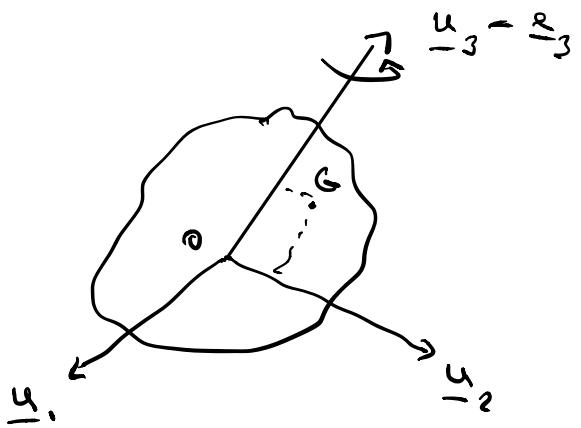
\underline{u} \underline{P}

$$(\underline{\epsilon})_s = R_{(q)}^T (\underline{\epsilon})_\Sigma = \begin{pmatrix} \cos q & \sin q & 0 \\ -\sin q & \cos q & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} 0 \\ \sin \varphi \\ -\cos \varphi \end{pmatrix}$$

$$= g (\sin \varphi \sin q \underline{u}_1 + \sin \varphi \cos q \underline{u}_2 - \cos \varphi \underline{u}_3)$$

Tensión s-fija

$$S(O; \underline{u}_1, \underline{u}_2, \underline{u}_3)$$



$$(\underline{\epsilon}_G - \underline{\epsilon}_0) =$$

$$= \gamma_G \underline{u}_2 + \tau_G \underline{u}_3$$

Allora scriviamo del per

$$\begin{aligned}
 (\underline{x}_3 - \underline{x}_0) \wedge M_g = & (y_G \underline{u}_2 + t_G \underline{u}_3) \wedge \\
 & \wedge (\sin \alpha \sin \varphi \underline{u}_1 + \sin \alpha \cos \varphi \underline{u}_2 - \cos \alpha \underline{u}_3) \\
 = M_g & (-y_G \sin \alpha \sin \varphi \underline{u}_3 - y_G \cos \alpha \underline{u}_1 \\
 & + t_G \sin \alpha \sin \varphi \underline{u}_2 - t_G \sin \alpha \cos \varphi \underline{u}_1)
 \end{aligned}$$

$$\underline{L}(0) = I_0(\omega \underline{u}_3) \quad \omega = \omega \underline{u}_3$$

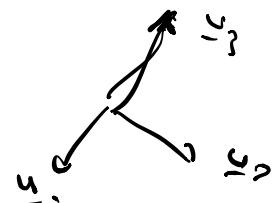
$$\frac{d}{dt} \underline{L}(0) = \dot{\omega} I_0(\underline{u}_3) + \omega^2 \underline{u}_3 \wedge I_0(\underline{u}_3)$$

Per il momento

$$I_0(\underline{u}_3) = \begin{pmatrix} I_{11} & I_{12} & I_{13} \\ I_{12} & I_{22} & I_{23} \\ I_{13} & I_{23} & I_{33} \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} = \begin{pmatrix} I_{13} \\ I_{23} \\ I_{33} \end{pmatrix}$$

$$\underline{u}_3 \wedge I_0(\underline{u}_3) = \underline{u}_3 \wedge (I_{13} \underline{u}_1 + I_{23} \underline{u}_2 + I_{33} \underline{u}_3)$$

$$= -I_{23} \underline{u}_1 + I_{13} \underline{u}_2$$



$$\frac{d}{dt} \underline{L}(0) = \underline{\mu}_1 (I_{13} \dot{\omega} - I_{23} \omega^2) + \underline{\mu}_2 (I_{23} \dot{\omega} + I_{13} \omega^2) + \underline{\omega} \dot{\underline{I}}_{33} \underline{\mu}_3$$

setzt aus Tatsache in die $\omega = \dot{\varphi}$

$$(x_0 - \bar{x}_0) \times Mg + \mu_0^2 = \frac{d}{dt} \underline{L}(0)$$

$$\underline{\mu}_1 \quad \mu_0^2 \underline{\mu}_1 = Mg (y_0 \cos \alpha + r_0 \sin \alpha \cos \varphi) + (I_{13} \ddot{\varphi} - I_{23} \dot{\varphi}^2)$$

$$\underline{\mu}_2 : \mu_0^2 \underline{\mu}_2 = -Mg r_0 \sin \alpha \sin \varphi + (I_{23} \ddot{\varphi} + I_{13} \dot{\varphi}^2)$$

$$\underline{\mu}_3 : I_{33} \ddot{\varphi} = - (Mg y_0 \sin \alpha) \underline{\sin \varphi}$$

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