

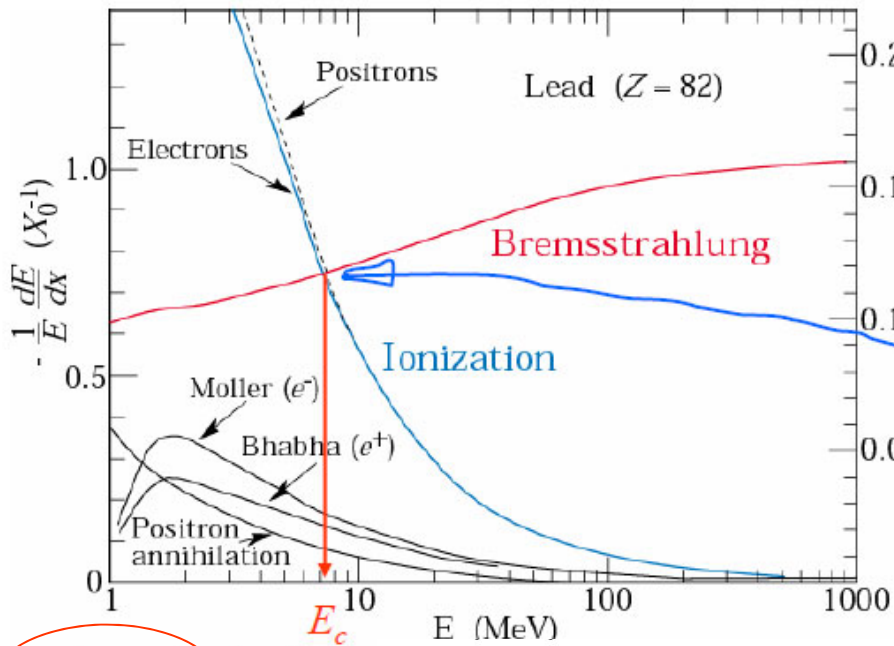
BREMMSTRAHLUNG

Table 2.3. Radiation lengths for various absorbers

Material	[gm/cm ²]	[cm]
Air	36.20	30050
H ₂ O	36.08	36.1
NaI	9.49	2.59
Polystyrene	43.80	42.9
Pb	6.37	0.56
Cu	12.86	1.43
Al	24.01	8.9
Fe	13.84	1.76
BGO	7.98	1.12
BaF ₂	9.91	2.05
Scint.	43.8	42.4

Table 2.2. Critical energies of some materials

Material	Critical energy [MeV]
Pb	9.51 ←
Al	51.0
Fe	27.4
Cu	24.8
Air (STP)	102
Lucite	100
Polystyrene	109
NaI	17.4
Anthracene	105
H ₂ O	92

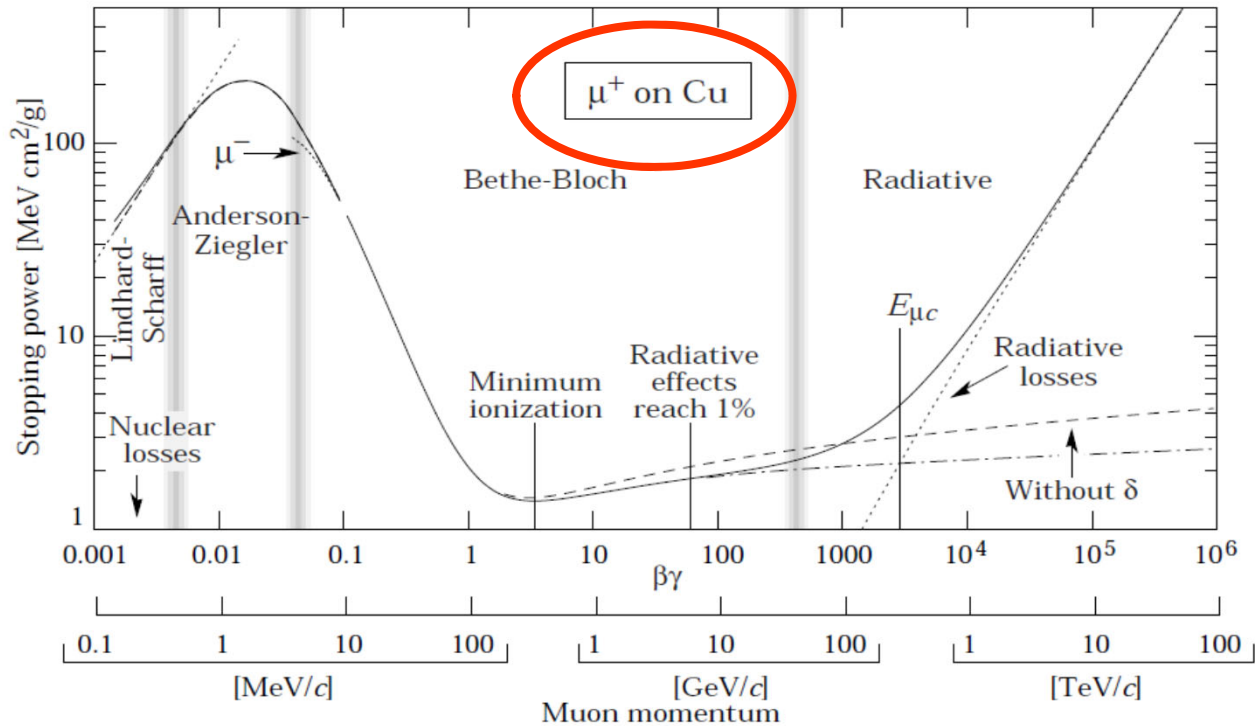


Critical Energy

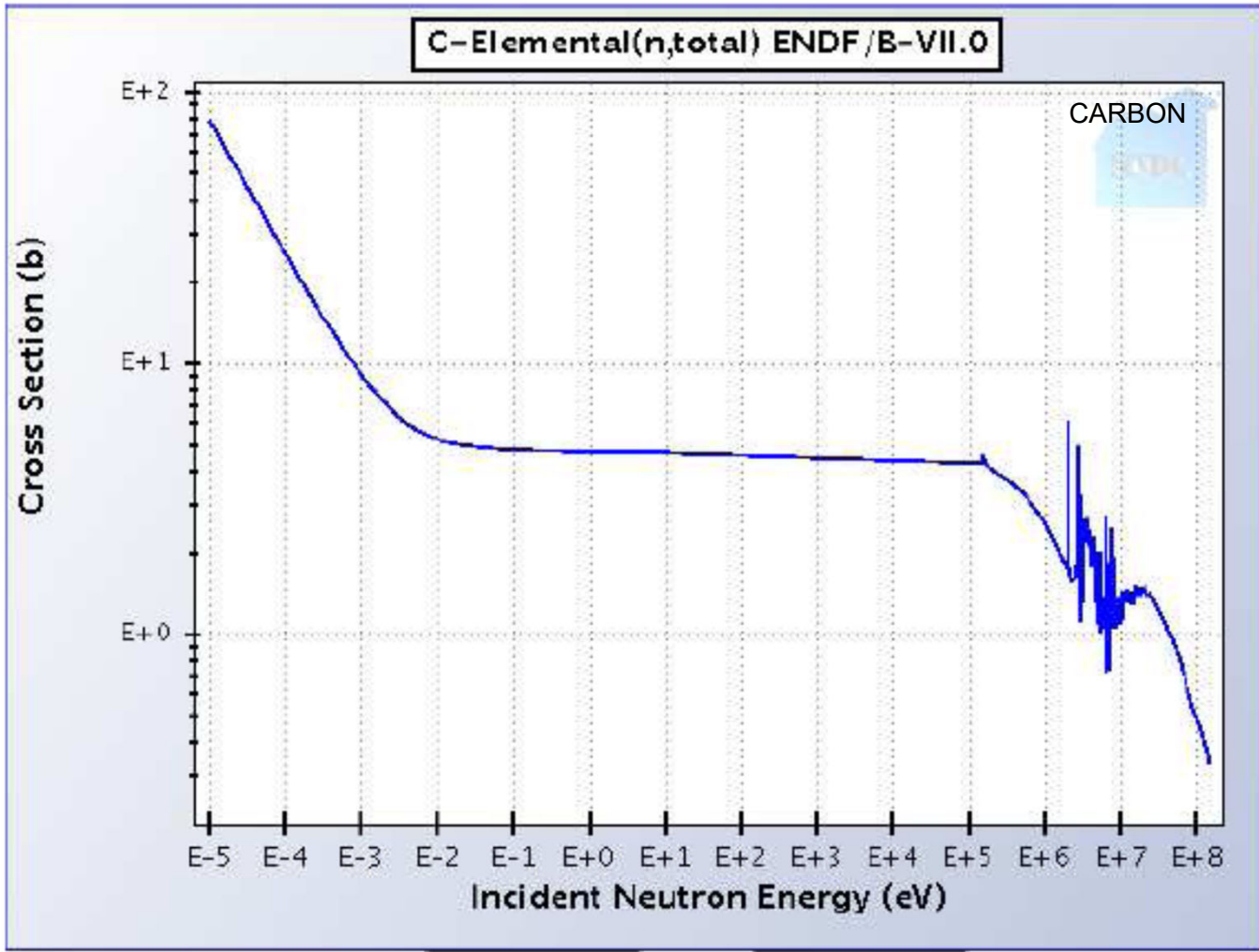
Critical energy E_c
Ionization = Bremsstrahlung

Electron (positron) intera

$E_c \propto Z^{-1}$



NNDCC Databases: | | | | | | |



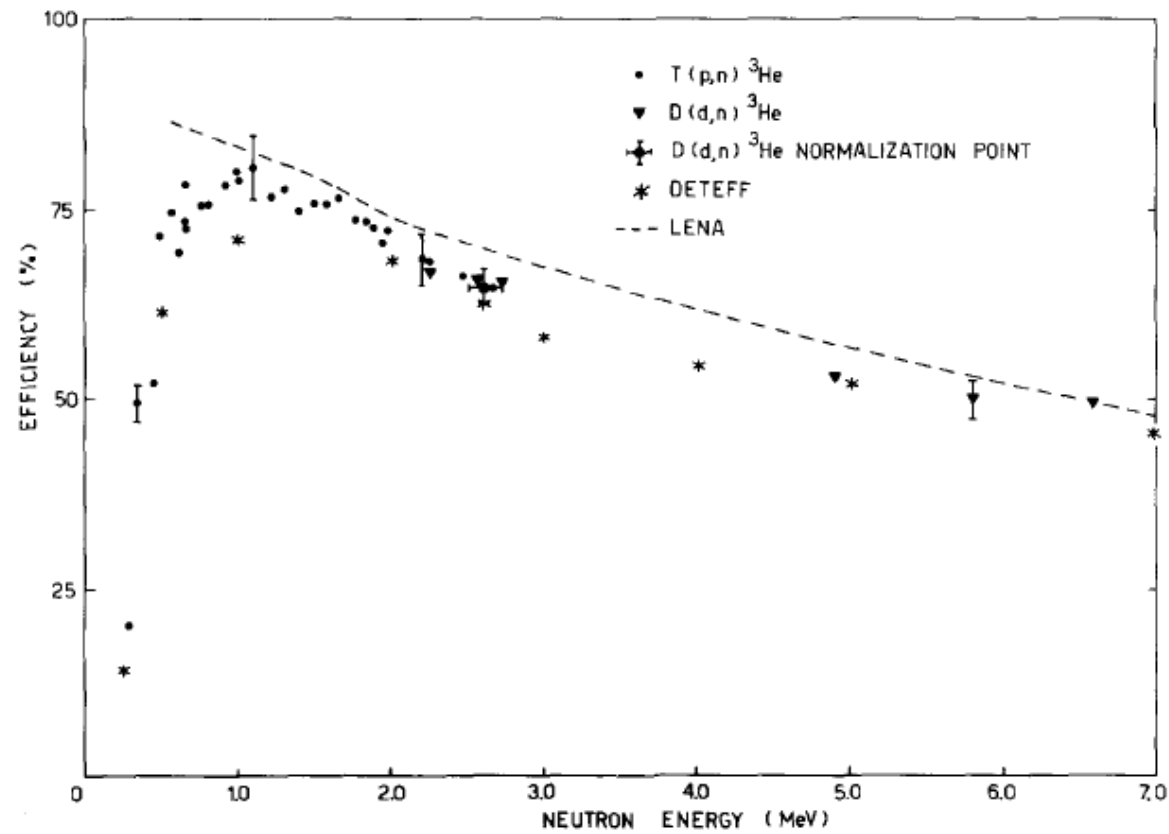


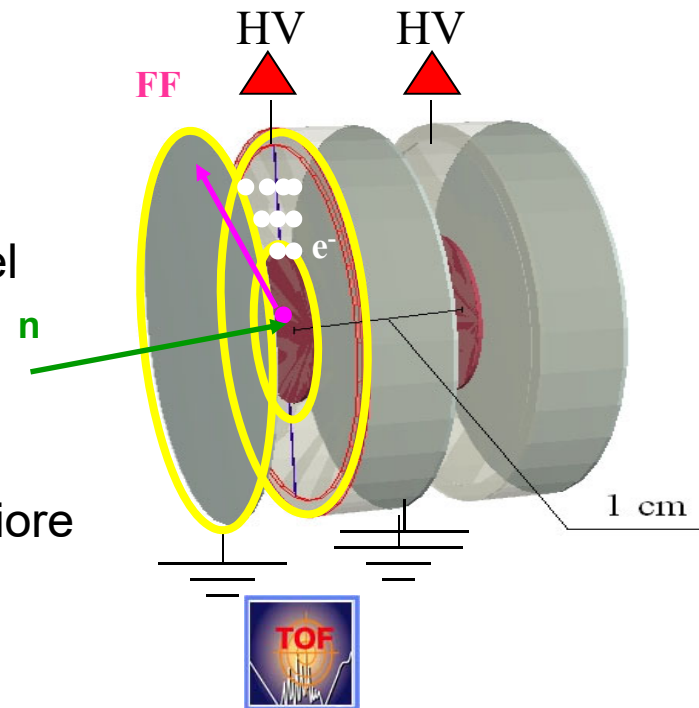
Fig. 1. Experimental neutron detection efficiency of the 10.0 cm \times 10.0 cm NE104 scintillator. DETEFF and LENA refer to the Monte Carlo calculations of refs. 7, 8.



MAIN

Principio di funzionamento

- La camera a ionizzazione è composta da uno stack di celle che si ripetono modularmente
- ogni cella è composta da 3 elettrodi
 - finestra di Al messa a terra
 - supporto di Al, connesso a High Voltage (il target è “dipinto” su entrambi i lati del supporto)
 - finestra di Al messa a terra
- quando un n colpisce il target, può causare fissione. Il FF è emesso nel gas e lo ionizza. Si formano coppie *elettrone-ione* e gli *elettroni* driftano verso l'elettrodo a potenziale maggiore





Pulse shape analysis

- La perdita specifica di energia ($-dE/dx$) di particelle cariche in un mezzo è descritta dalla formula di Bethe:

carica (e) e rest mass (m)
dell' elettrone

densità atomica (N), numero atomico (Z) e
potenziale di ionizzazione (I) del mezzo

$$-\frac{dE}{dx} = \frac{4\pi e^4 z^2}{m_0 v^2} \cdot N \cdot Z \left[\ln \frac{2m_0 v^2}{I} - \ln \left(1 - \frac{v^2}{c^2} \right) - \frac{v^2}{c^2} \right]$$

carica (ze) e velocità (v) della
particella che provoca ionizzazione

- I termini elevati al quadrato incidono maggiormente
- A parità di cammino particelle con numero atomico maggiore (z) perdono più energia nel mezzo.

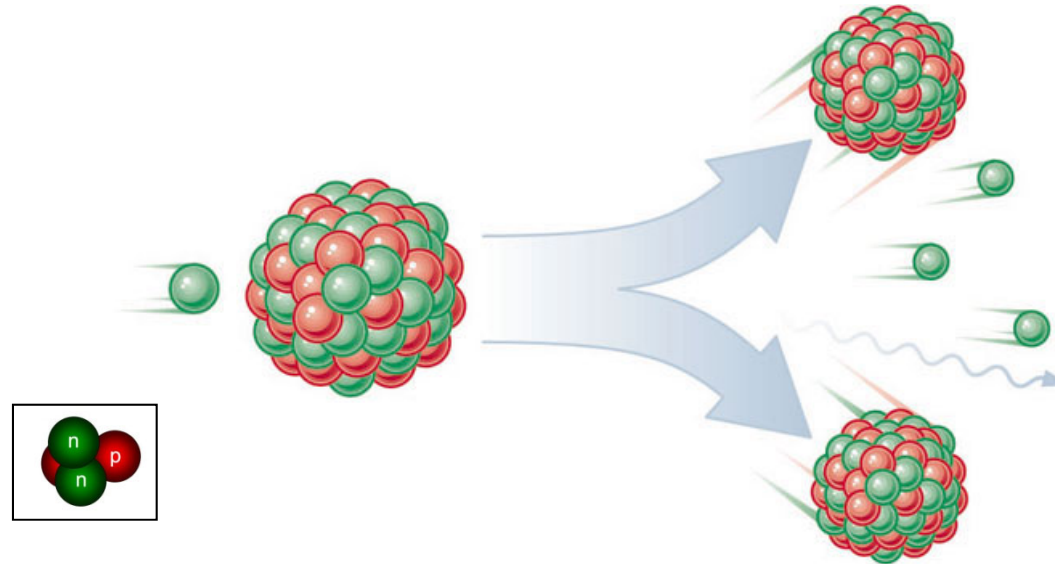




MAIN

Il metodo del rapporto

- La corrente di particelle uscenti è determinata contando il numero di FFs



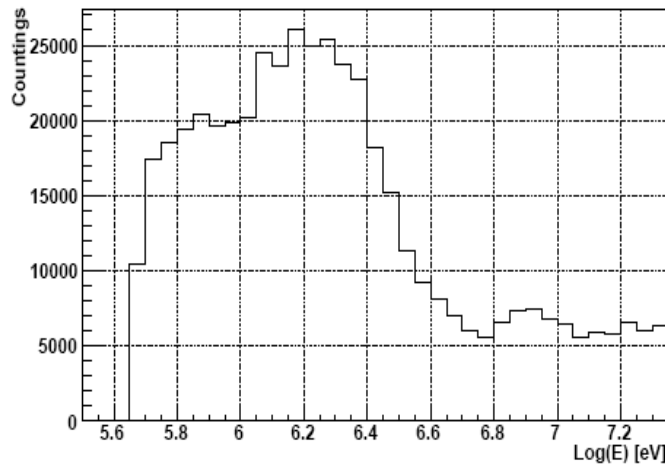
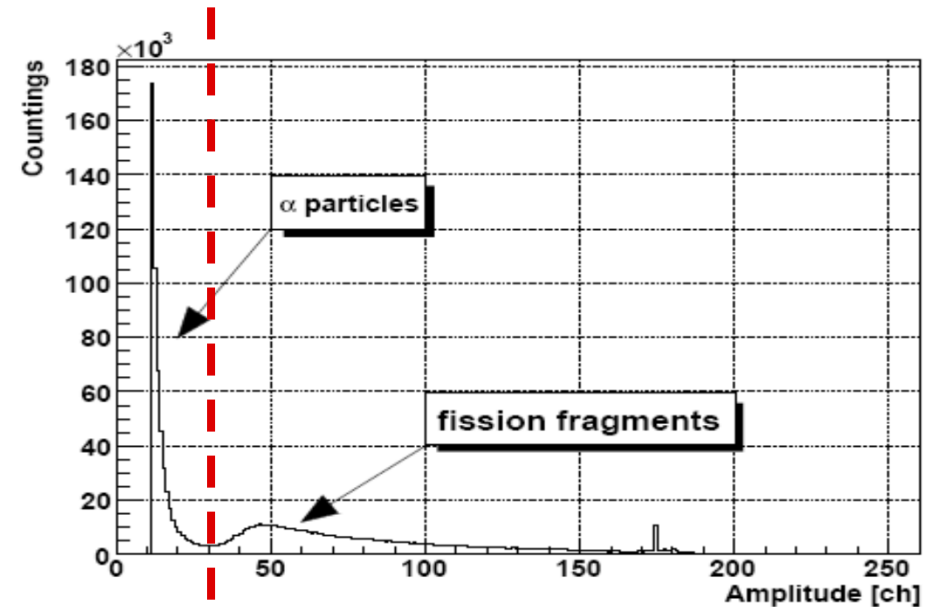
$$R_b = \sigma I_a N \begin{cases} R_b = \text{emission rate} \longrightarrow \text{Determinato contando i FFs} \\ N = \text{nuclei/area} \longrightarrow \text{Noto se nota la geometria del campione} \\ I_a = \text{neutron "beam" intensity?} \longrightarrow \text{Ratio method} \end{cases}$$





Estrazione $\sigma_{(n,f)}(E_n)$

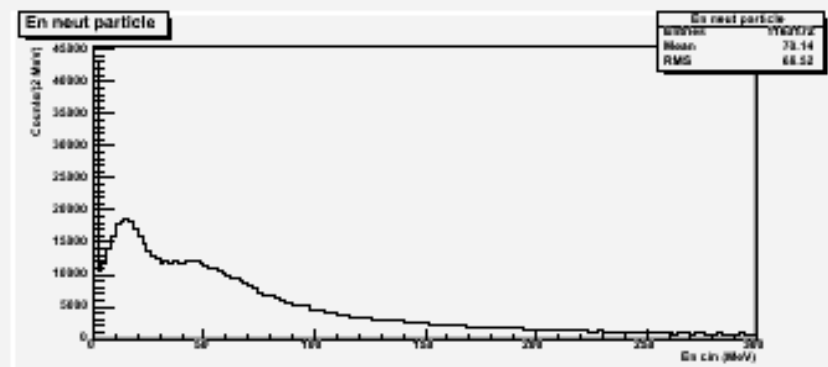
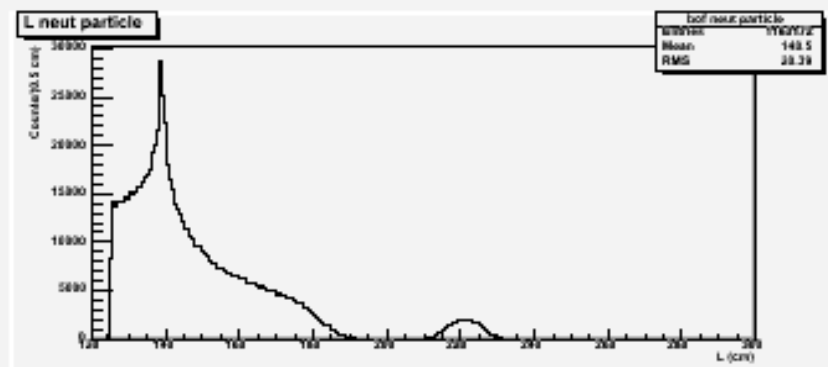
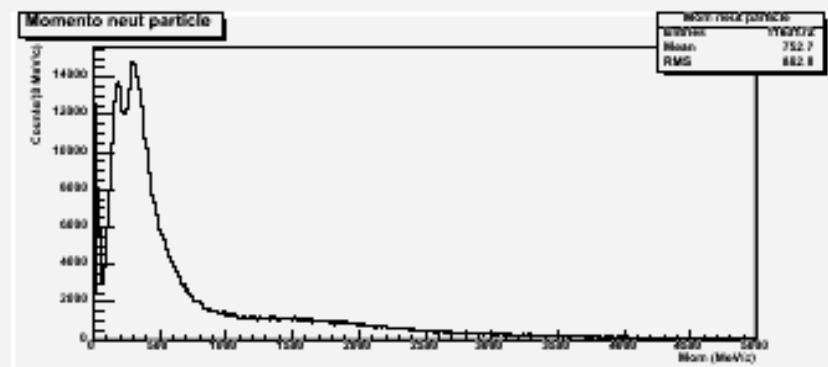
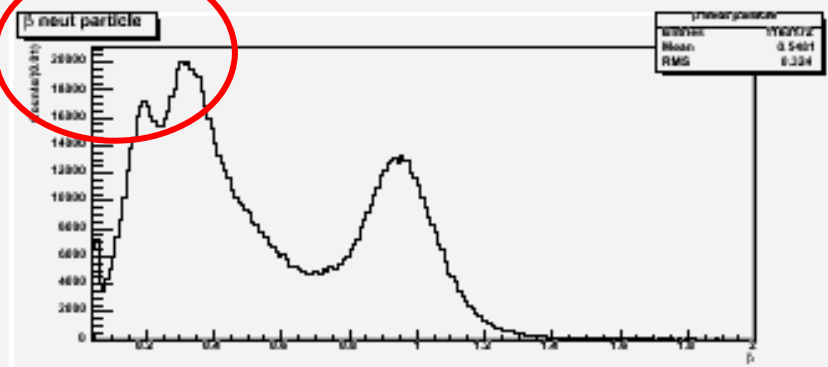
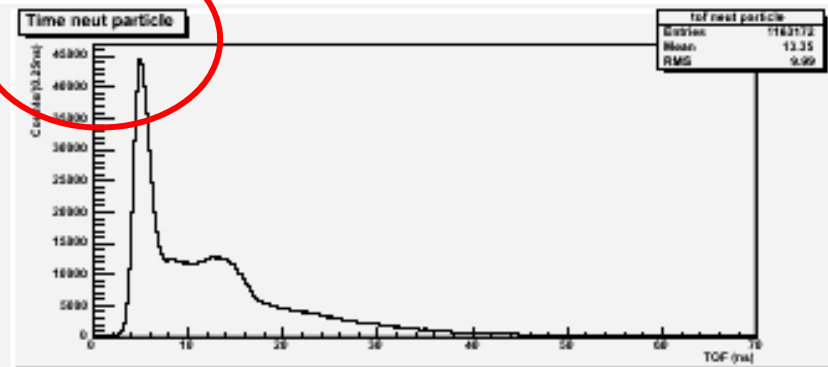
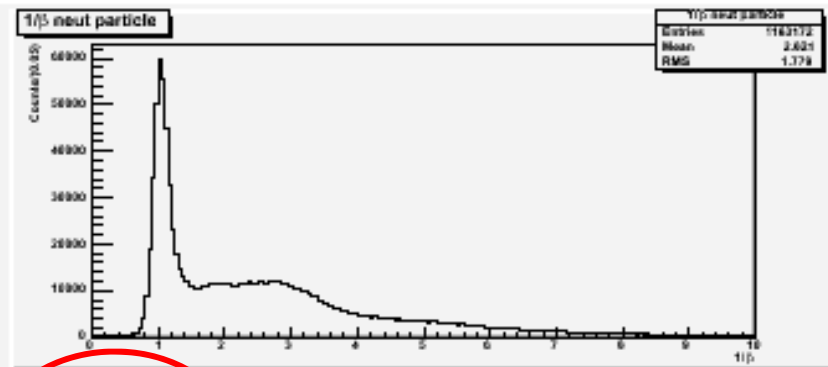
- Come visto prima i FFs possono venire isolati in base all'analisi della forma dell'impulso del segnale indotto dalla fissione



- Si riempie l'istogramma del # di FFs rivelati in corrispondenza di ogni E_n

$$\sigma_{xxx(n,f)} = \sigma_{235ENDF(n,f)} \cdot \frac{Y_{xxx}}{Y_{235}} \cdot \frac{m_{235}}{m_{xxx}} \cdot \frac{A_{xxx}}{A_{235}}$$





Neutron detection efficiency

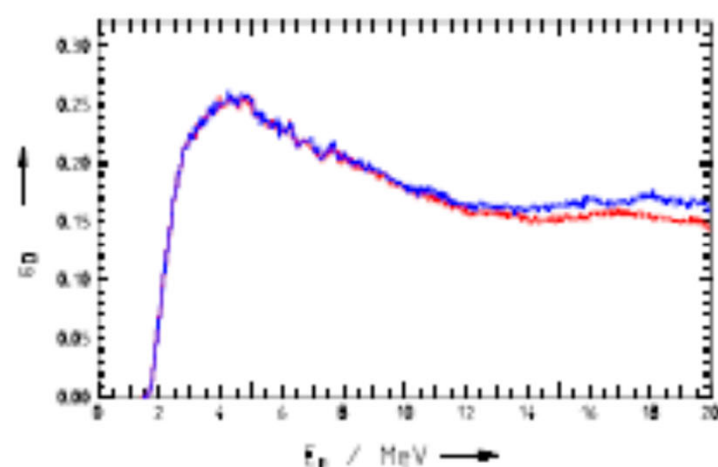
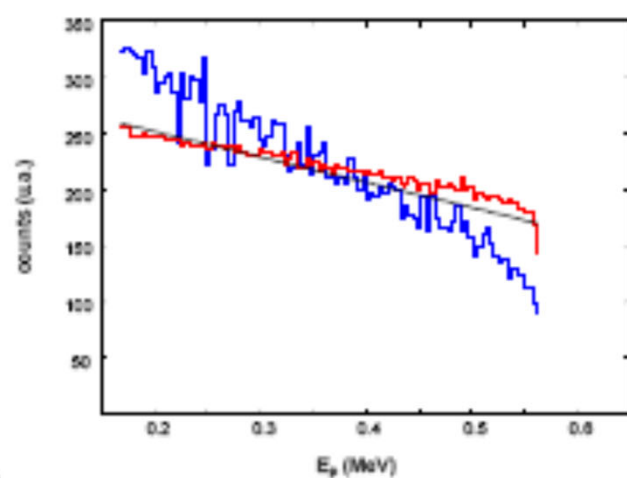


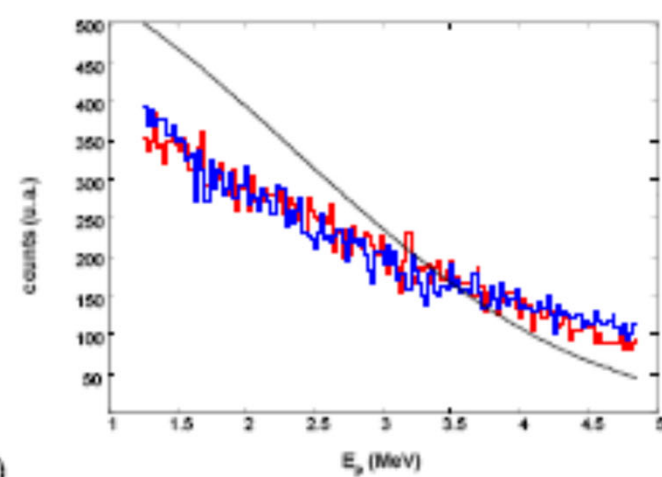
Figure 2: Neutron detection efficiency ϵ_D of an NE213 scintillator, 255 mm in diameter and 51 mm in length, for neutron energies up to 20 MeV and a threshold of about 2 MeV

($E_{\text{thr}} > 0.46$ MeV)

- (a) simulated (blue) with version 7 of the NEFF-code [128]
- (b) simulation corrected for compliance with the experimental response in the low-amplitude region (red).



a)



b)

Figure 1: Proton recoil spectra obtained with monoenergetic neutrons (blue), calculated with MCNPX (red) and with the Snidow algorithm (black): a) SP2-4 with neutrons at 565 keV. b) SP6 with neutrons at 4.5 MeV.