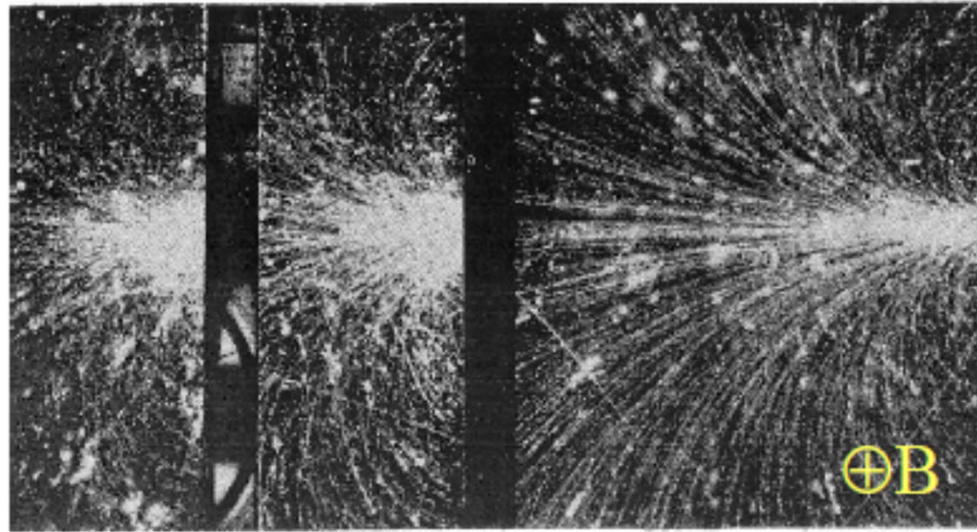
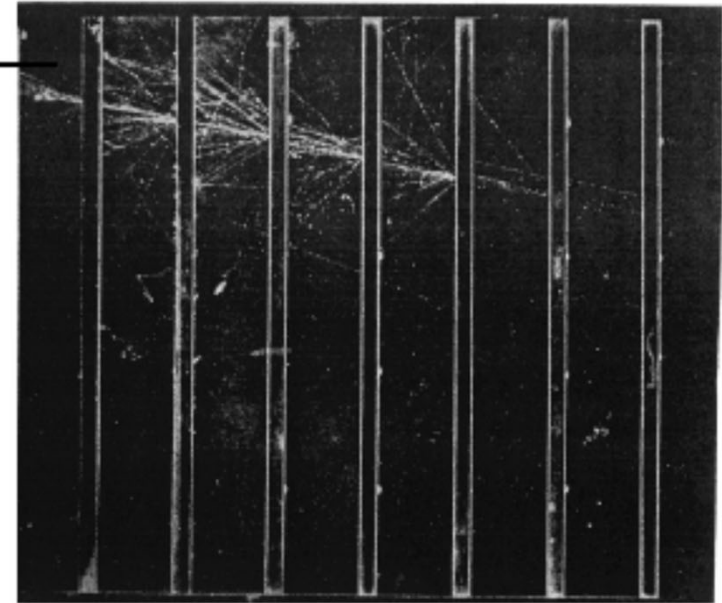


# How a shower looks like

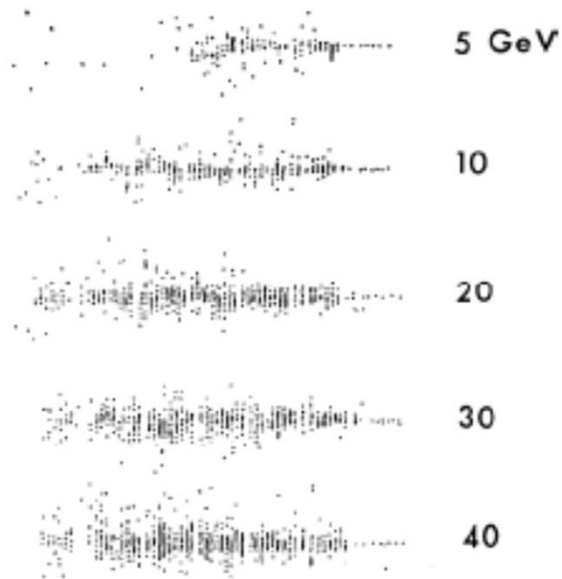


Electron shower in lead. 7500 gauss in cloud chamber. CALTECH

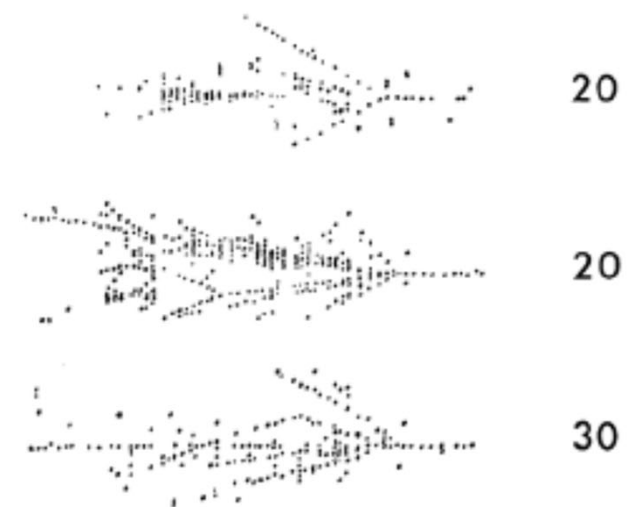


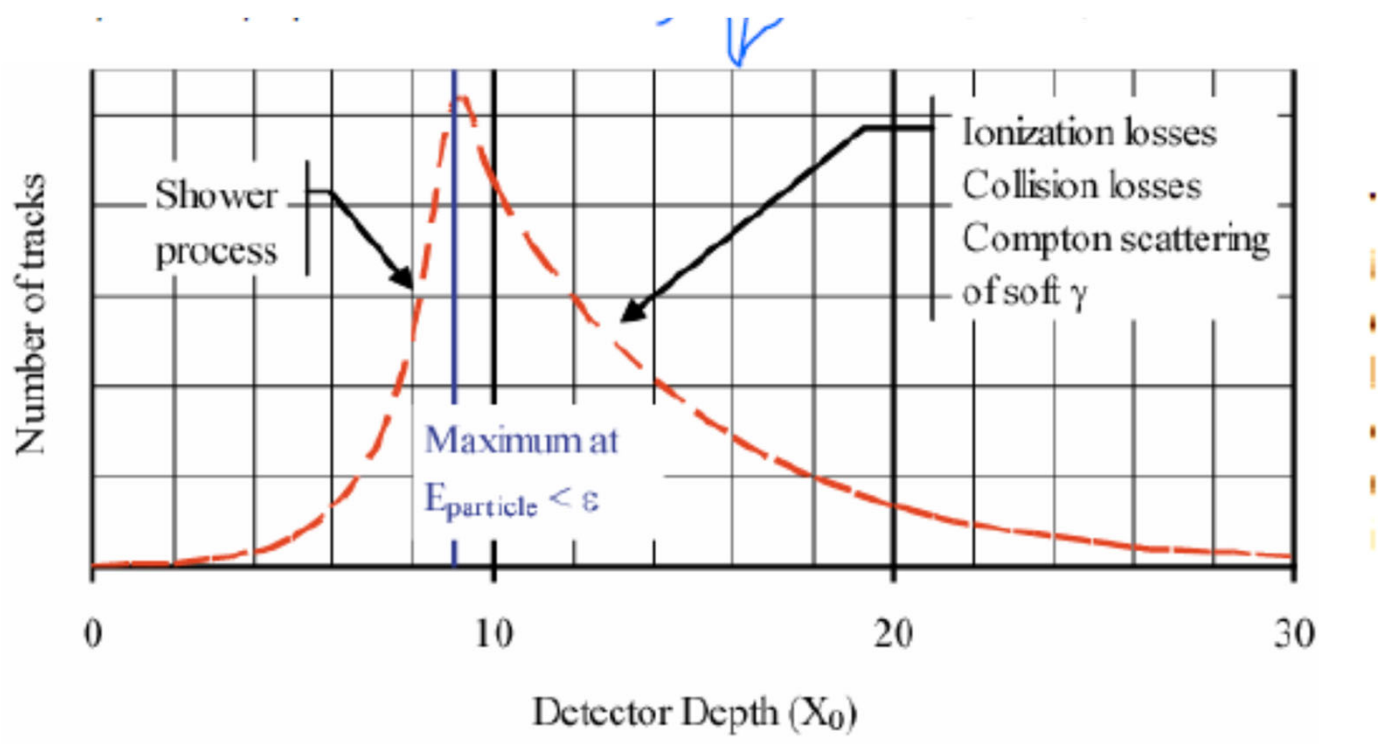
Electron shower in lead. Cloud chamber. W.B. Fretter, UCLA

## Electron showers

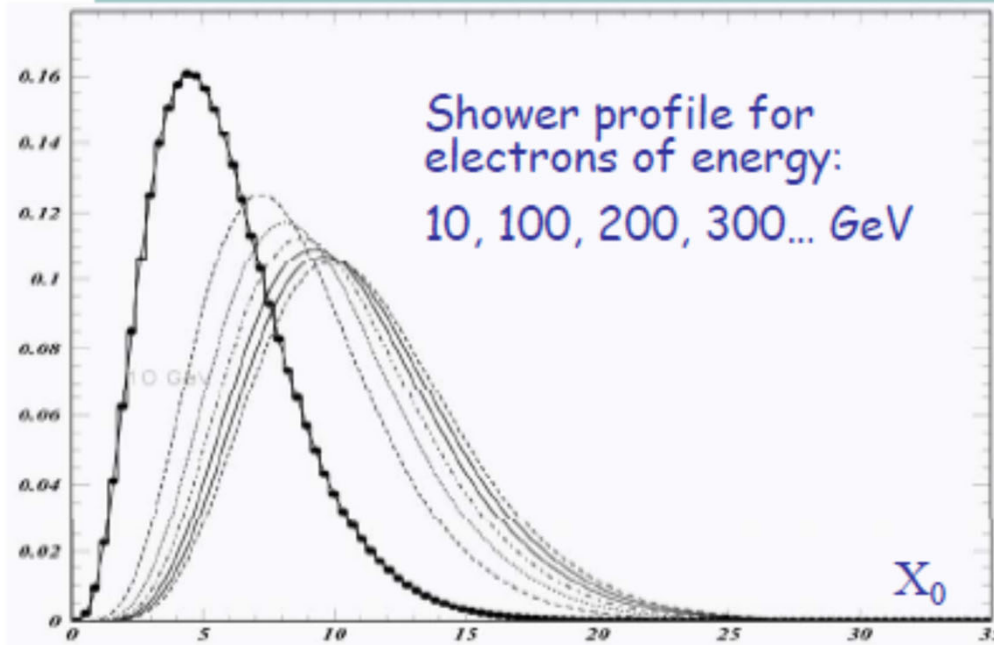


## Hadron showers





# EM showers: longitudinal profile



$$t_{\max} = 1.4 \ln(E_0/E_c)$$

$$N_{\text{tot}} \propto E_0/E_c$$

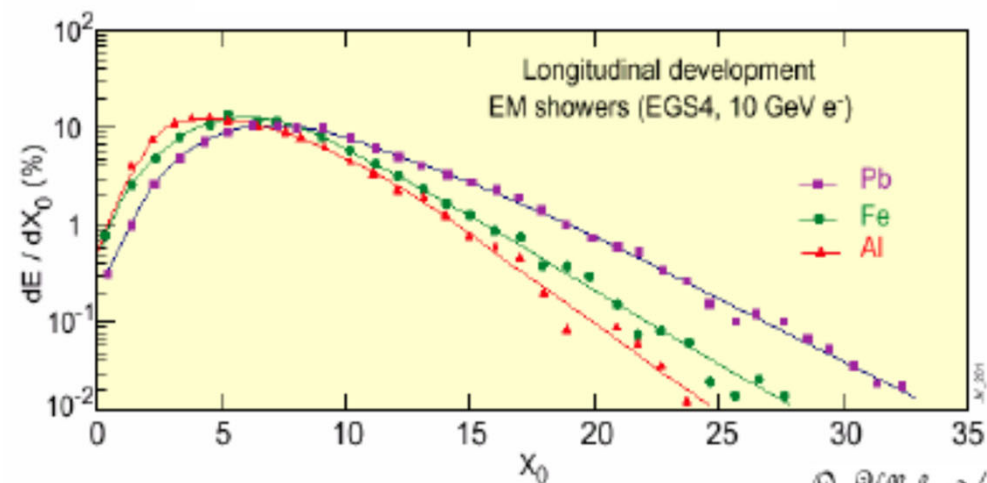
Longitudinal containment:

$$t_{95\%} = t_{\max} + 0.08Z + 9.6$$

$E_c \propto 1/Z$   $\rightarrow$  •shower max  
•shower tail

Shower parametrization

$$\frac{dE}{dt} \propto t^\alpha e^{\beta t}$$



# EM showers: transverse profile

## Transverse shower profile

- Multiple scattering make electrons move away from shower axis
- Photons with energies in the region of minimal absorption can travel far away from shower axis

Molière radius sets transverse shower size, it gives the average lateral deflection of critical energy electrons after traversing  $1X_0$

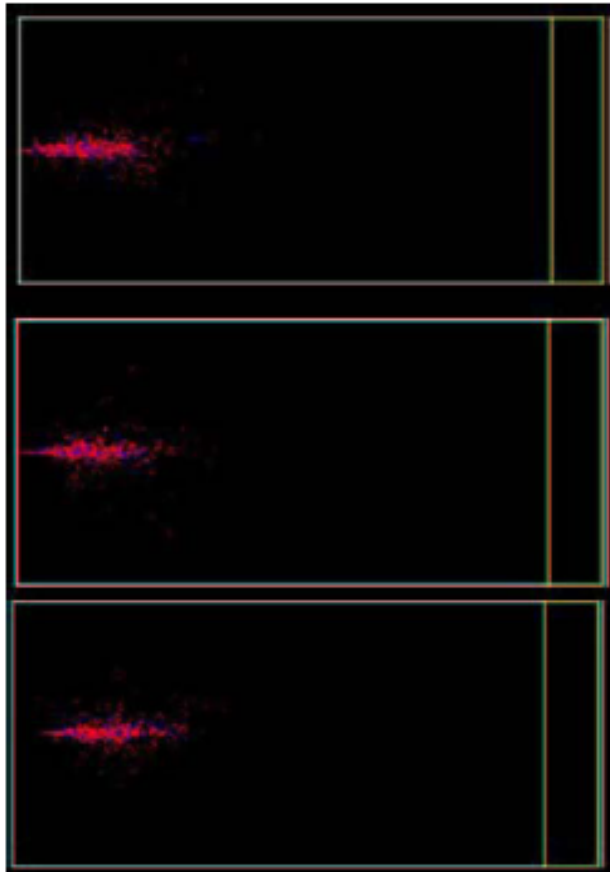
$$R_M = \frac{21\text{MeV}}{E_c} X_0$$

$$R_M \propto \frac{X_0}{E_c} \propto \frac{A}{Z} (Z \gg 1)$$

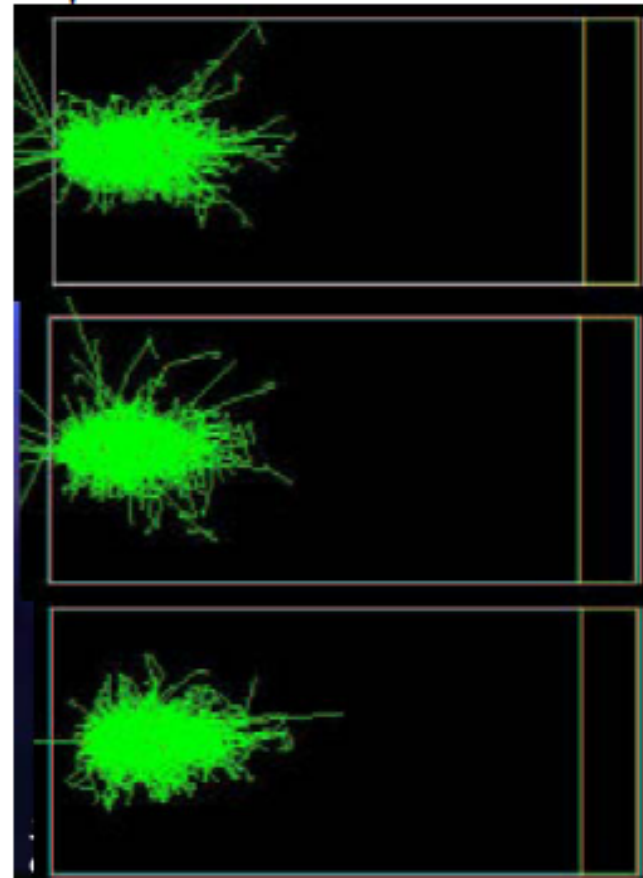
75%  $E_0$  within  $1R_M$ , 95% within  $2R_M$ , 99% within  $3.5R_M$

20 GeV  $\gamma$  in copper (simulation)

charged particles only

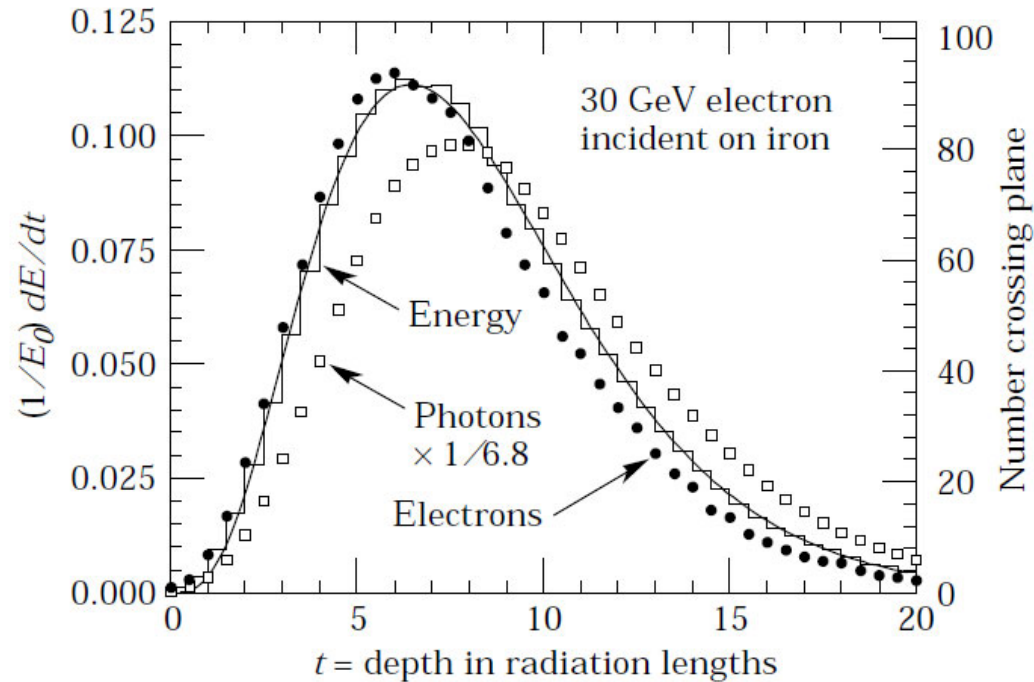


all particles





## Electromagnetic calorimeters.



**Figure 27.18:** An EGS4 simulation of a 30 GeV electron-induced cascade in iron. The histogram shows fractional energy deposition per radiation length, and the curve is a gamma-function fit to the distribution. Circles indicate the number of electrons with total energy greater than 1.5 MeV crossing planes at  $X_0/2$  intervals (scale on right) and the squares the number of photons with  $E \geq 1.5$  MeV crossing the planes (scaled down to have same area as the electron distribution).







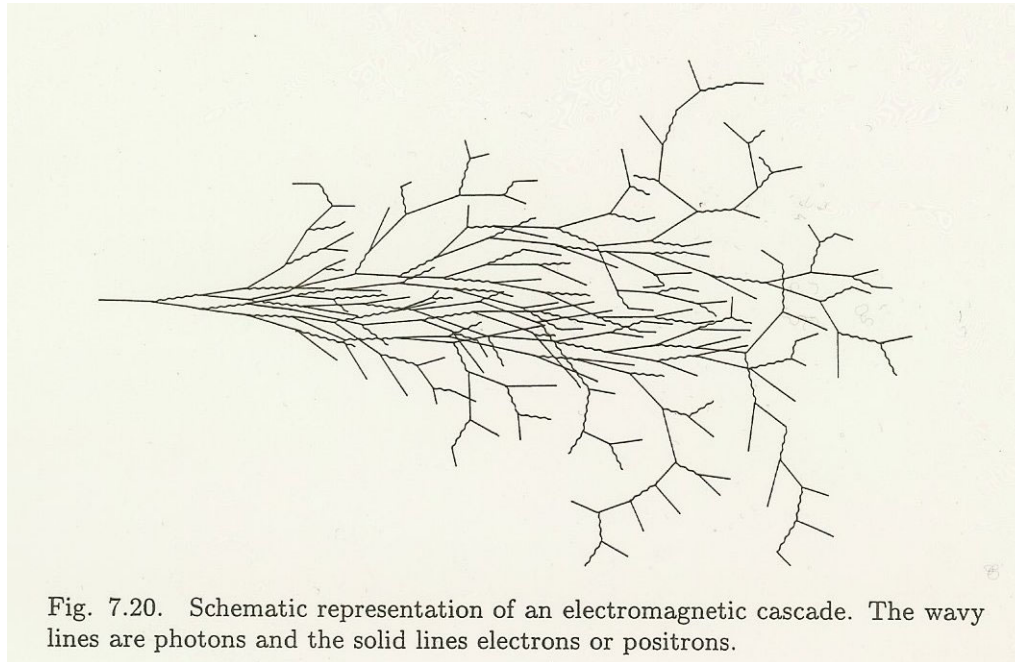
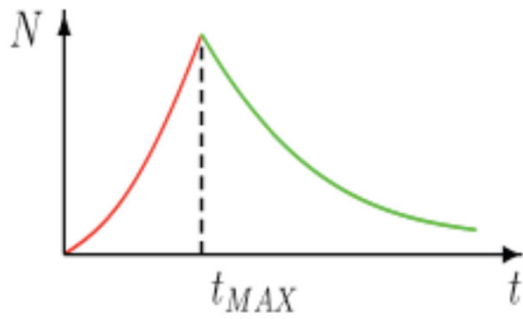


Fig. 7.20. Schematic representation of an electromagnetic cascade. The wavy lines are photons and the solid lines electrons or positrons.



**Figure 11.2** Shower profiles in lead. The number of electrons should be multiplied by a normalization factor of 0.79. (D. Müller, Phys. Rev. D 5: 2677, 1972.)

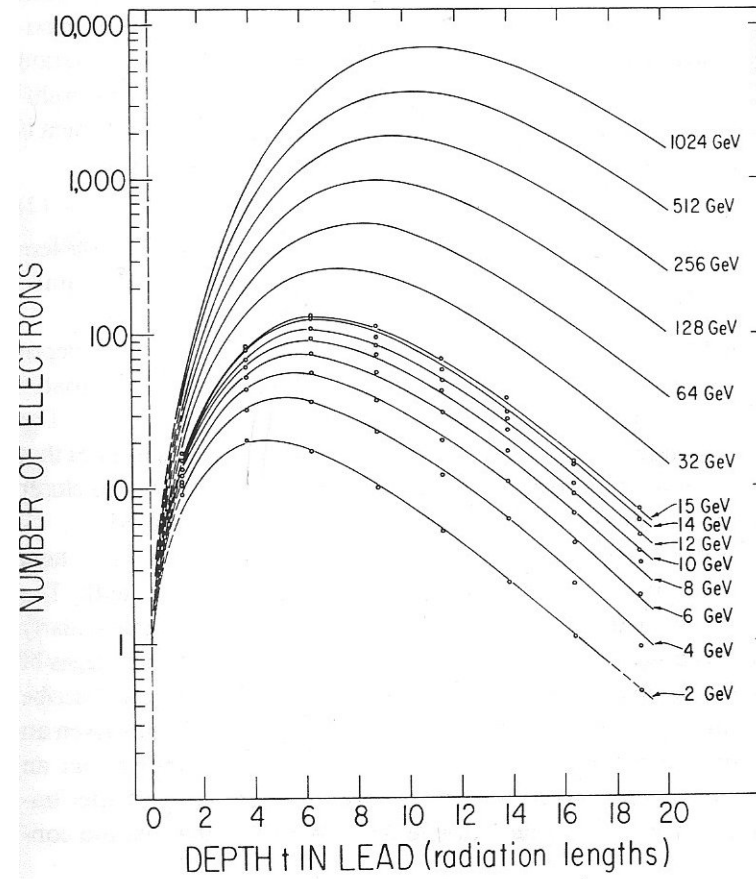
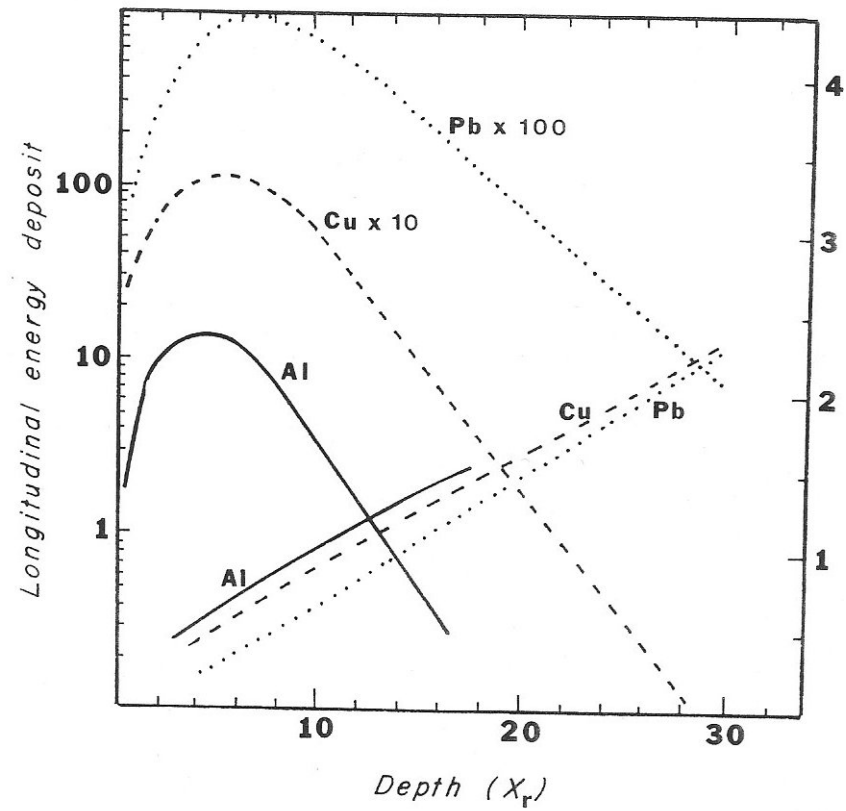


Figure 11.3 Longitudinal development of electromagnetic showers different materials. Right scale shows radii for 90% shower container (C. Fabjan and T. Ludlam, adapted with permission from the Annu Review of Nuclear and Particle Science, Vol. 32, © 1982 by Annu Reviews, Inc.)



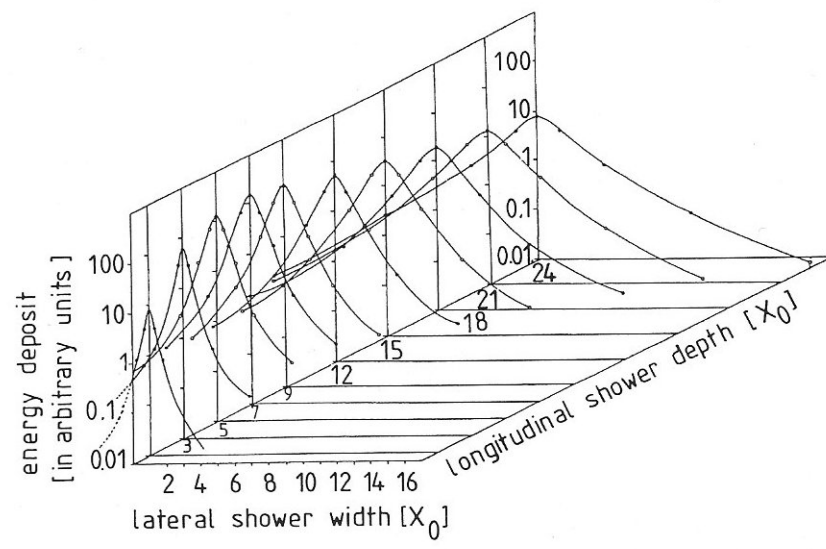
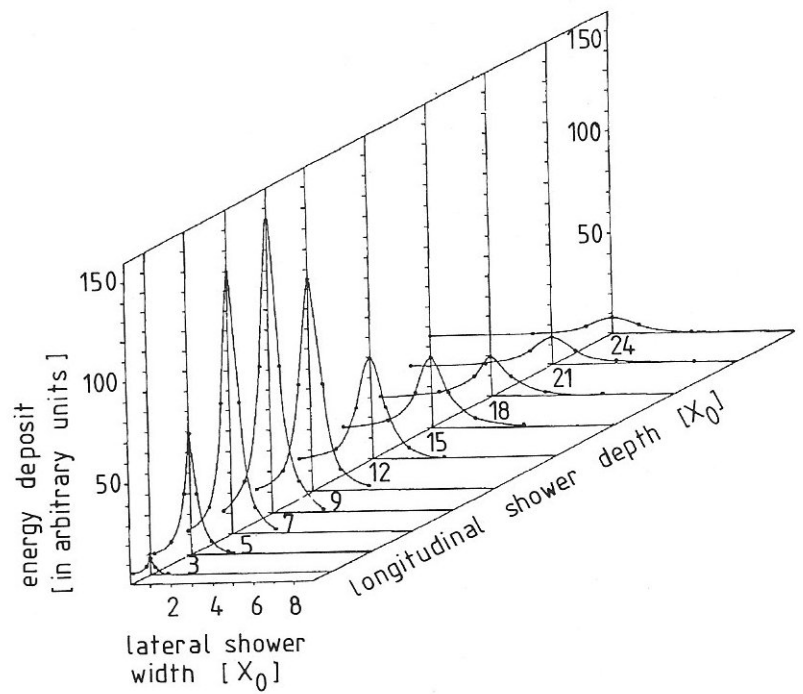
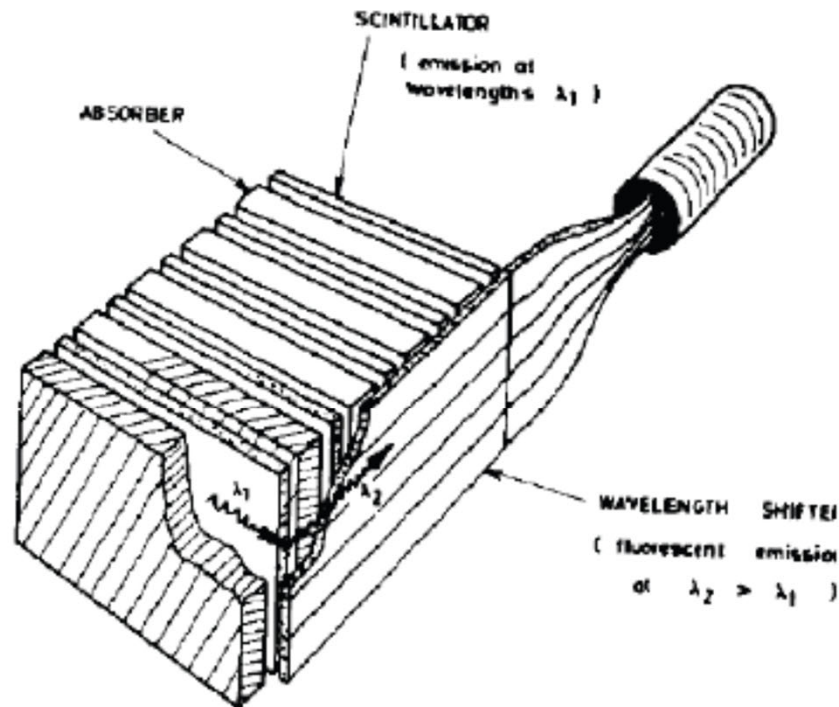


Fig. 7.23. Longitudinal and lateral development of an electron shower (6 GeV) in lead shown with linear and logarithmic scales (based on [504, 505]).



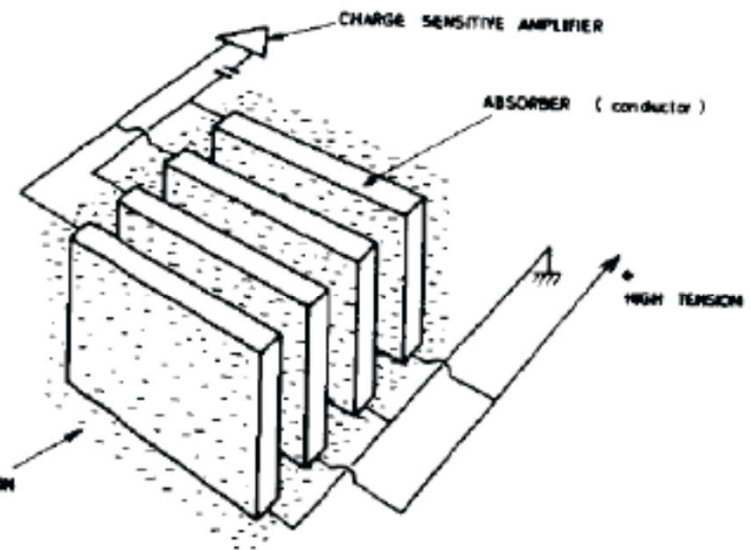
## ■ Basic readout types for sampling calorimeters

□ Metal-scintillator sandwich structure



C. Fabjan and T. Ludlam, 1987

□ Metal-liquid argon ionization chamber

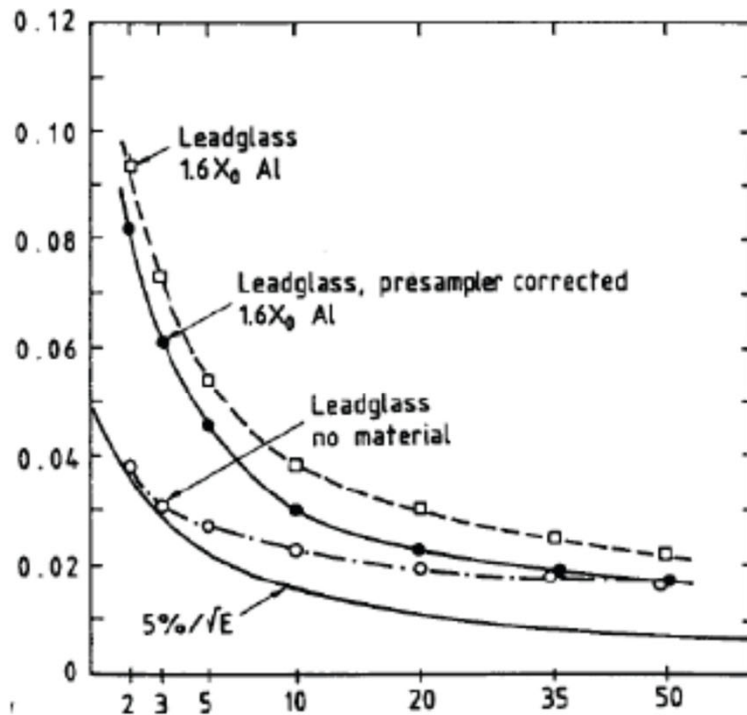


C. Fabjan and T. Ludlam, 1987

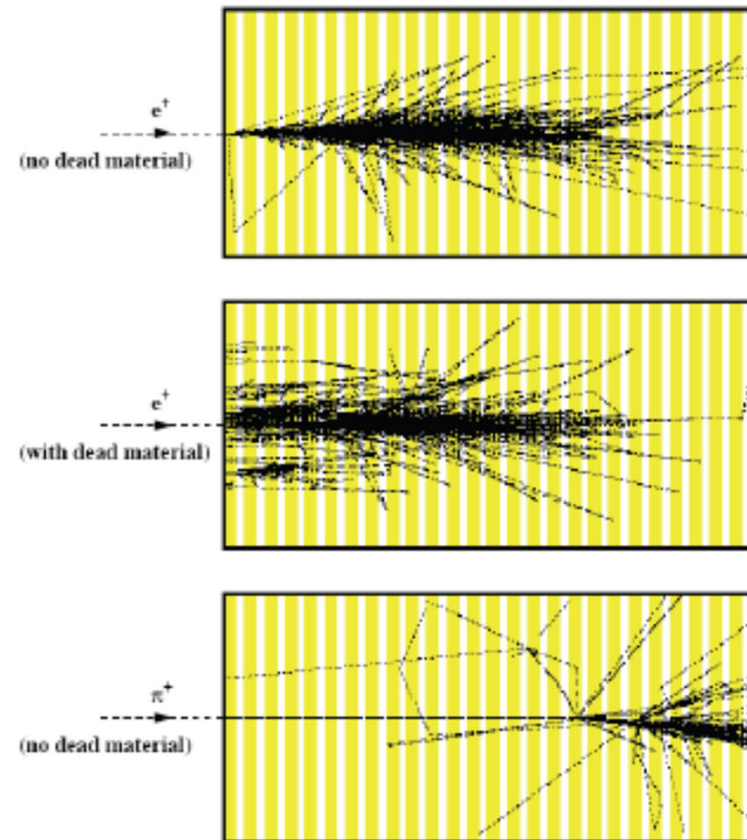


## ■ Energy resolution: Limitations

### □ Dead material effects



OPAL collaboration, C. Beard et al. NIM A 286 (1990) 117.



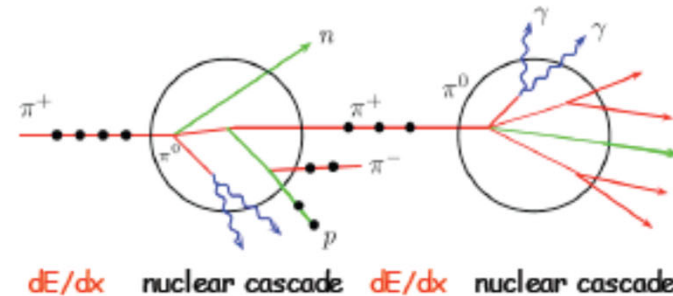


# Hadronic showers

## ■ Hadronic shower development

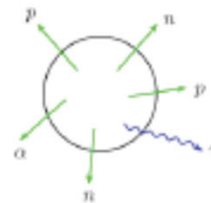
- General comment: Complexity of hadronic and nuclear processes produce multitude of effects that determine the functioning and performance of hadron calorimeters
  - Many channels compete in the development of hadronic showers
  - Larger variations in the deposited and visible energy
  - More complicated to optimize
- Sizeable electromagnetic (e) besides hadronic (h) shower contribution mainly from  $\pi^0$  decay (1/3 of pions)
- **Invisible energy** due to delayed emitted photons in nuclear reactions, soft neutrons and binding energy
- Visible energy smaller for hadronic (h) than for electromagnetic (e) showers: Ratio of response  $e/h > 1$
- Larger intrinsic fluctuations for hadronic than electromagnetic showers
- Improvements: Increase visible energy to get  $e/h=1$ : Compensation (Compensation for the loss of invisible energy)
- Discussed instr. effects for e showers also hold for h showers

**Step 1:** Production of energetic hadrons with a mean free path given by the nuclear interaction length:

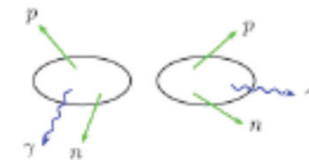


**Step 2:** Hadronic collisions with material nuclei (significant part of the primary part of primary energy is consumed in nuclear processes):

Evaporation



Evaporation followed by evaporation





## Processes contributing to the energy deposition in hadronic cal.

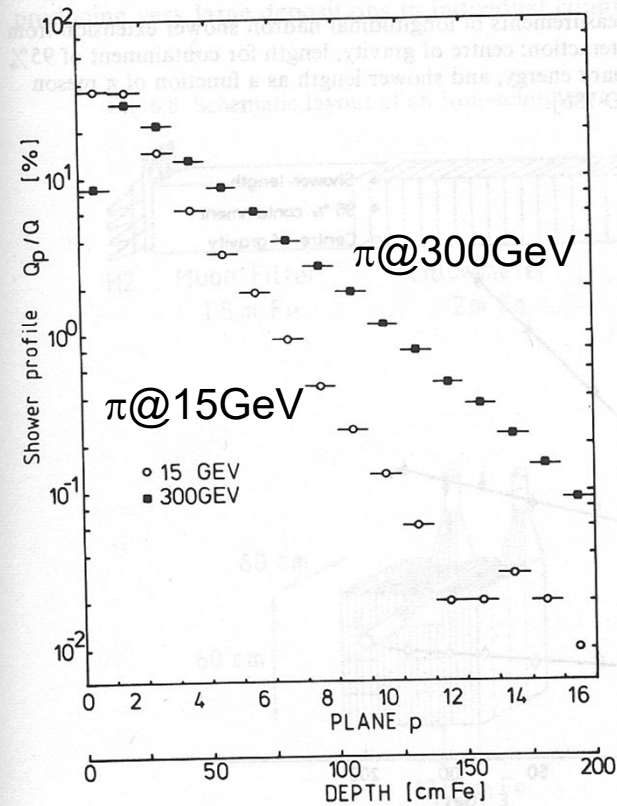
Table 11.2. *Average fractional energy deposition for a 10-GeV proton in an iron/liquid argon calorimeter*

Process	Percent of total
Secondary proton ionization	31.6
Electromagnetic cascade ( $\pi^0$ )	21.0
Nuclear binding energy plus neutrino energy	20.6
Secondary $\pi^\pm$ ionization	8.2
Neutrons with $E > 10$ MeV	4.9
Neutrons with $E < 10$ MeV	3.9
Residual nuclear excitation energy	3.7
$Z > 1$ ionization	2.4
Primary proton ionization	2.3
Other	1.4

Source: T. Gabriel and W. Schmidt, Oak Ridge National Laboratory report, ORNL/TM-5105, 1975.

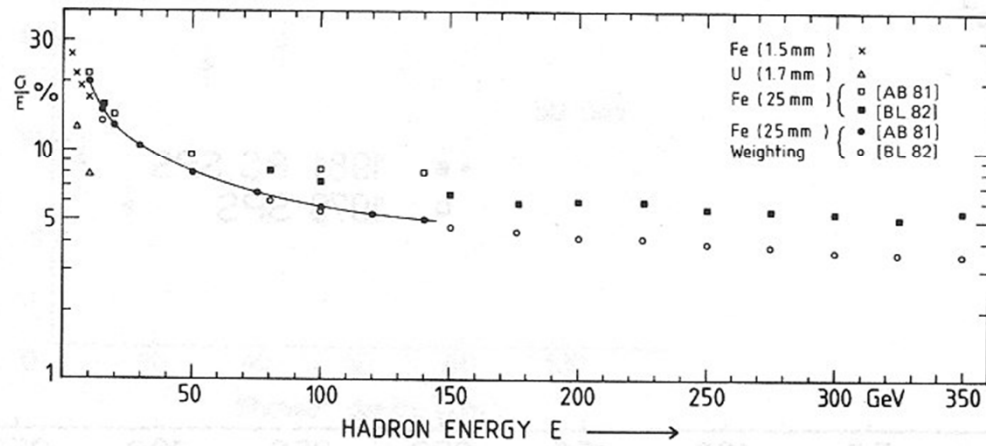
## Longitudinal distribution of energy in hadron calorimeter.

Fig. 6.6. Longitudinal distribution of energy deposited in a hadron calorimeter;  $Q_p$  is the energy deposited in counter  $p$  consisting of five layers with 2.5 cm iron and 0.5 cm scintillator each;  $Q = \sum Q_p$ . Measurements for  $\pi$  mesons of 15 GeV and 300 GeV [BL 82].

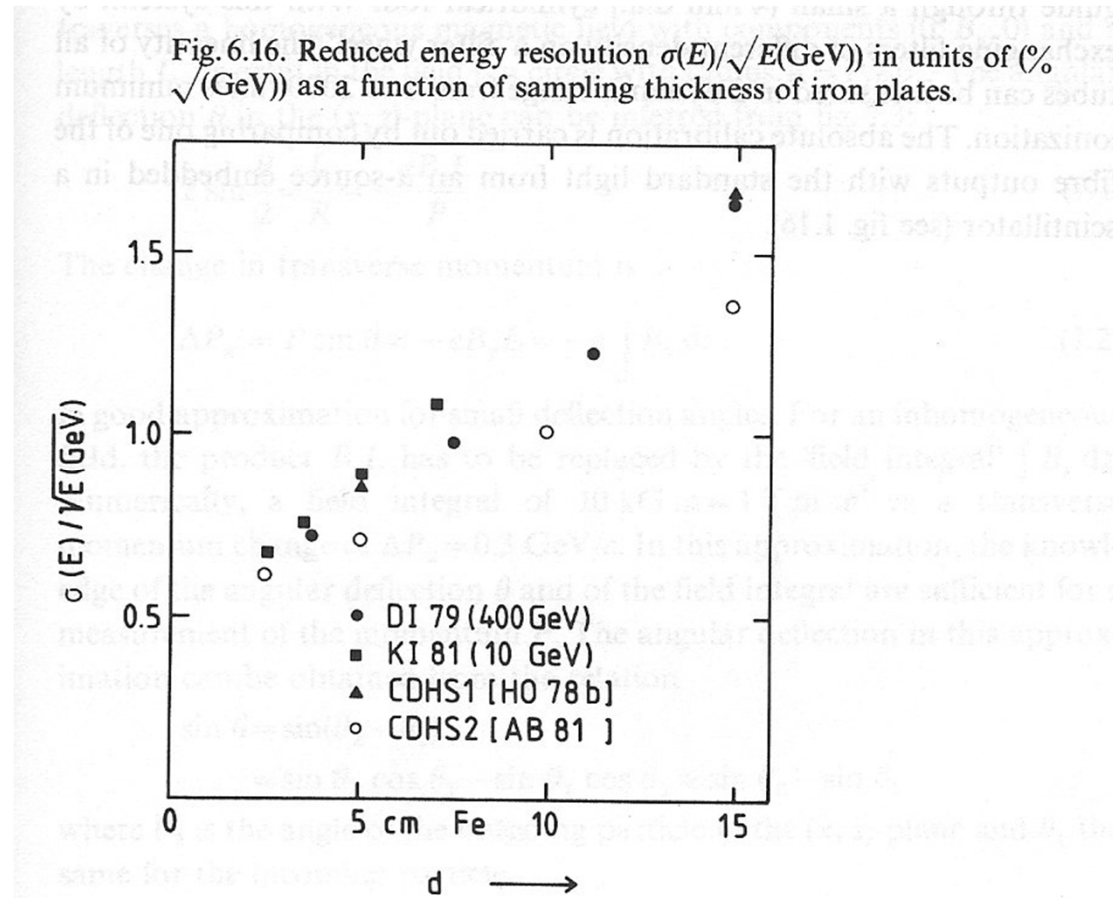


## Energy resolution

Fig. 6.13. Relative energy resolution  $\sigma(E)/E$  for different calorimeters; Fe (1.5 mm) and U (1.7 mm) from [FA 77]; Fe (25 mm) from [AB 81, BL 82].



## Dependence of resolution on sampling characteristics



# Technologies

## Electromagnetic calorimeters

### ◆ Crystals

(current calorimeters) <b>CsI (Tl)</b>	$2.3\%/^4\sqrt{E} \oplus 1.9\%$	<i>BaBar</i>
	$1.5\%/^4\sqrt{E} \oplus 1.2\%$	<i>BELLE</i>
<b>PbWO<sub>4</sub></b>	$2.8\%/ \sqrt{E} \oplus 0.6\%$	<i>CMS</i>
	$3.3\%/ \sqrt{E}$ (low noise term)	<i>ALICE</i>

### ◆ *LAr/Pb* (accordion)

$10\%/ \sqrt{E} \oplus 0.7\%$       *ATLAS*

### ◆ *Scint./Pb* (shashlik)

$10\%/ \sqrt{E} \oplus 1\%$       *LHCb*

# Technologies

## Hadron Calorimeters

◆ <i>Scint. / Brass</i> ( <i>WLS readout</i> )	$\sim 100\% \sqrt{E} \oplus 4.5\%$	<i>CMS</i>
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◆ <i>LAr / Brass</i>	$\sim 60\% \sqrt{E} \oplus 3\%$	<i>ATLAS (end-cap)</i>
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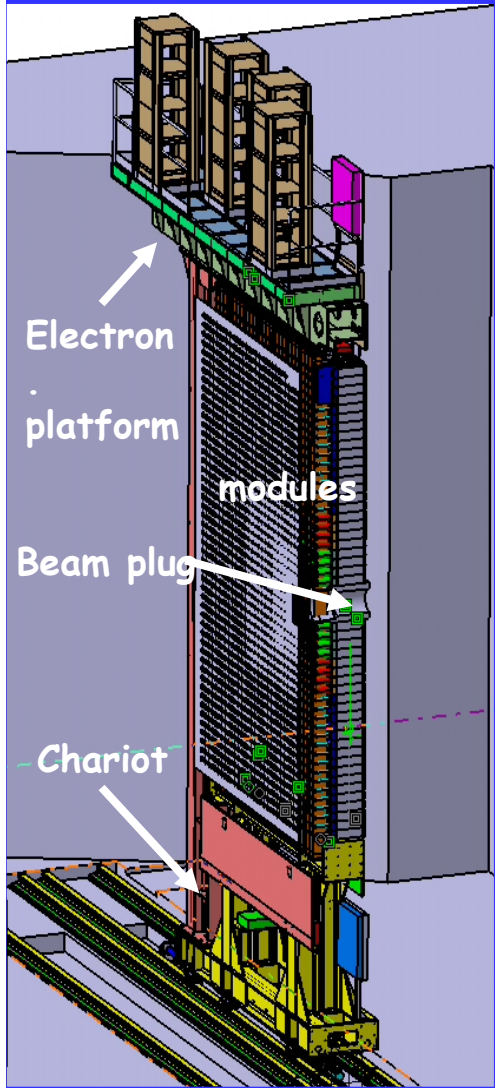
◆ <i>Scint / Fe (WLS readout)</i> ( <i>tiles oriented parallel to the beam</i> )	$\sim 50\% \sqrt{E} \oplus 3\%$	<i>ATLAS (barrel)</i>
---	---------------------------------	-----------------------

◆ <i>Scint / Fe (WLS readout)</i> ( <i>similar to ATLAS tile calorimeter,</i> <i>but planar geometry, 5.4 <math>\lambda</math> depth</i> )	$\sim 70\% \sqrt{E} \oplus 10\%$	<i>LHCb</i>
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*not compensated calorimeters*  
*optimization of the jet energy resolution important !*



Two halves on chariots  
and electronics platform  
on top

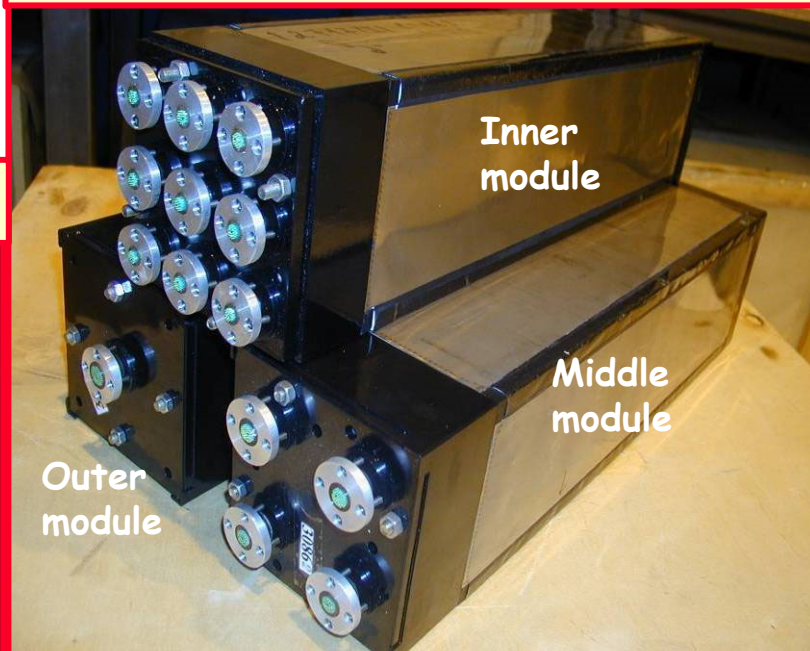


Overview  
of LHCb ECAL

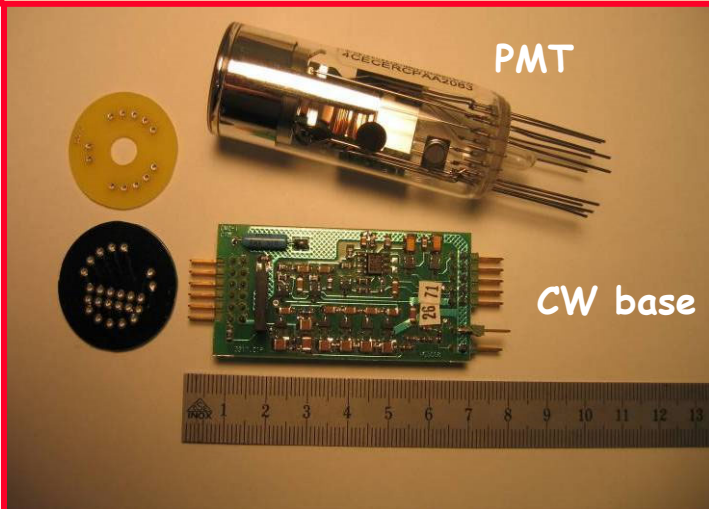
Fibres with loops



3312 shashlik modules with



PMT and CW base

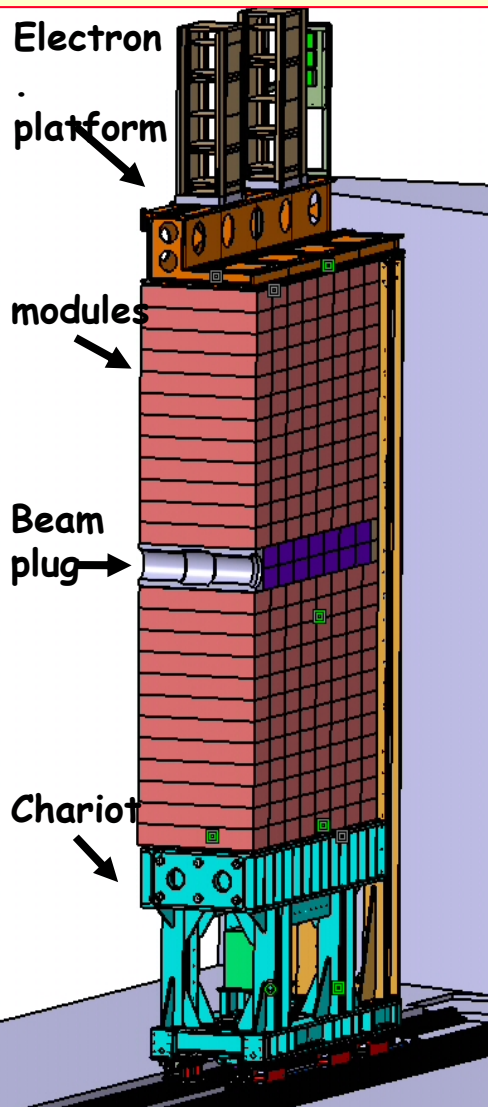


Scintillators, lead-

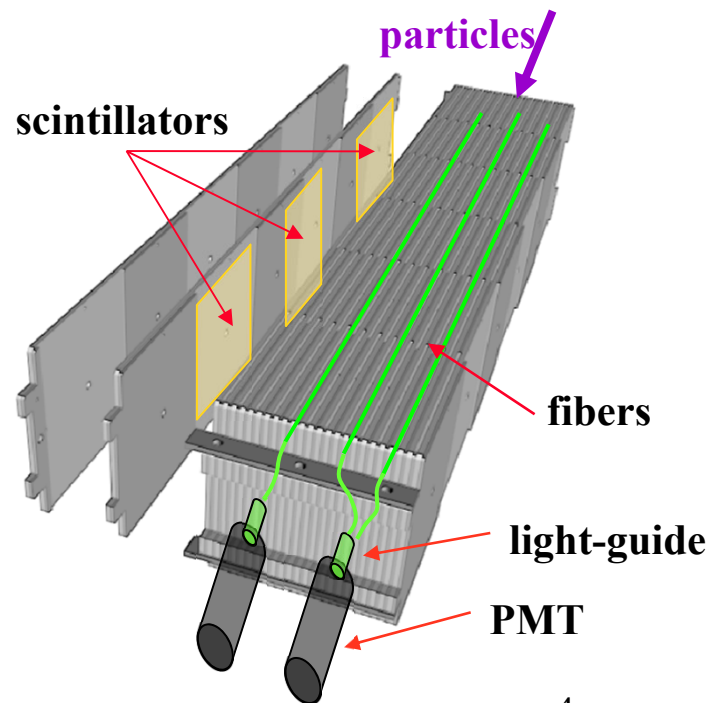
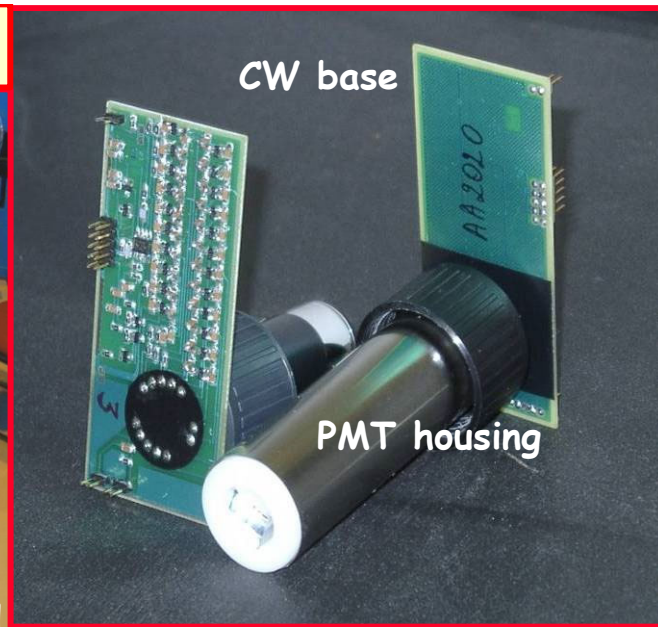
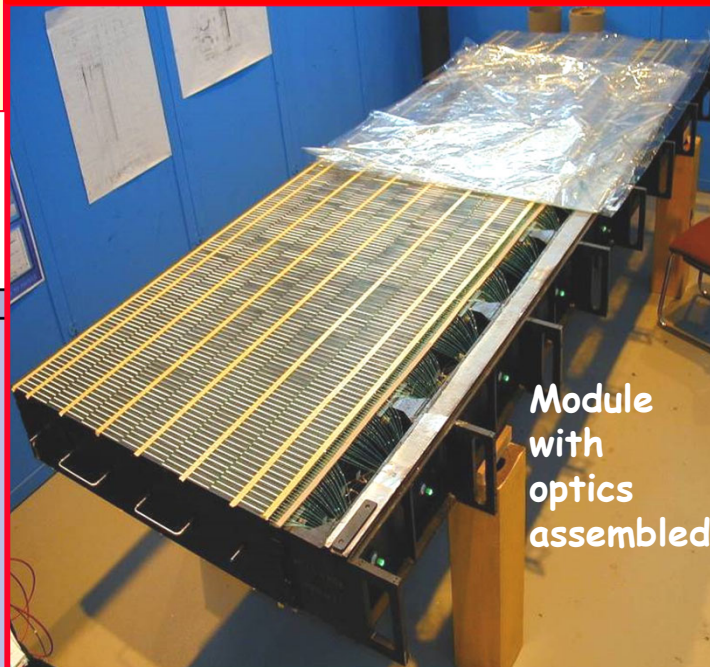




# Overview of LHCb HCAL



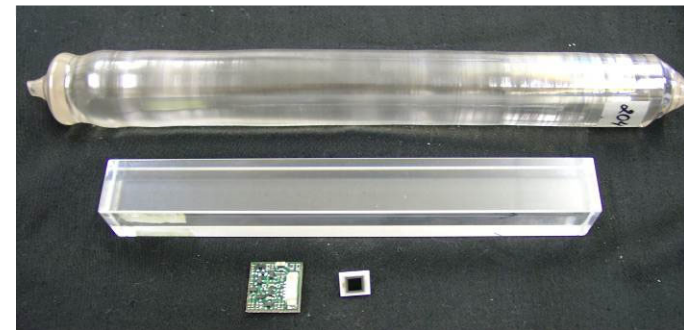
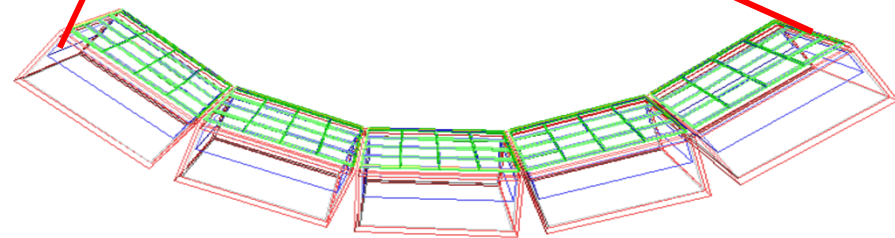
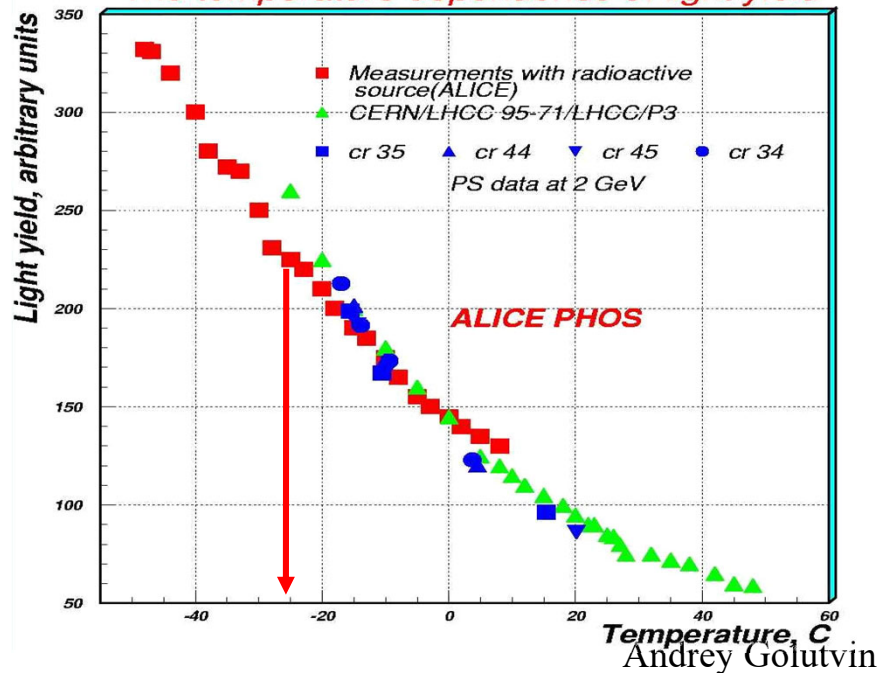
## 52 modules with longitudinal tiles



# ALICE PHOS electromagnetic calorimeter

- 17920 PWO crystals
- distance to IP: 4.6m
- coverage in pseudo-rapidity:  
 $|\Delta\eta| < 0.12$
- coverage in azimuthal angle:  
 $\Delta\Phi < 100^\circ$
- crystal size:  $22 \times 22 \times 180 \text{ mm}^3$
- Depth :  $20X_0$
- photo readout: APD + CSP
- operating temperature:  $-25^\circ\text{C}$

The temperature dependence of light yield

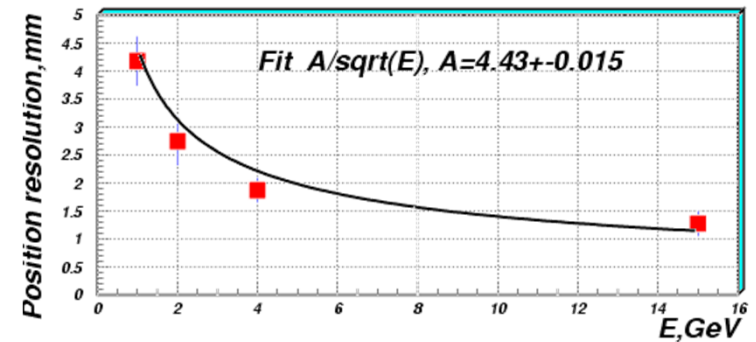
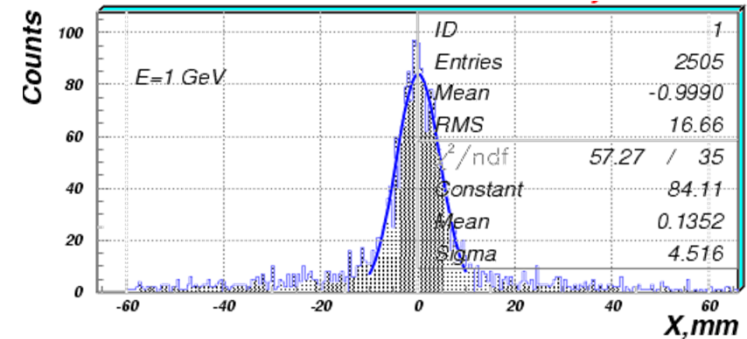
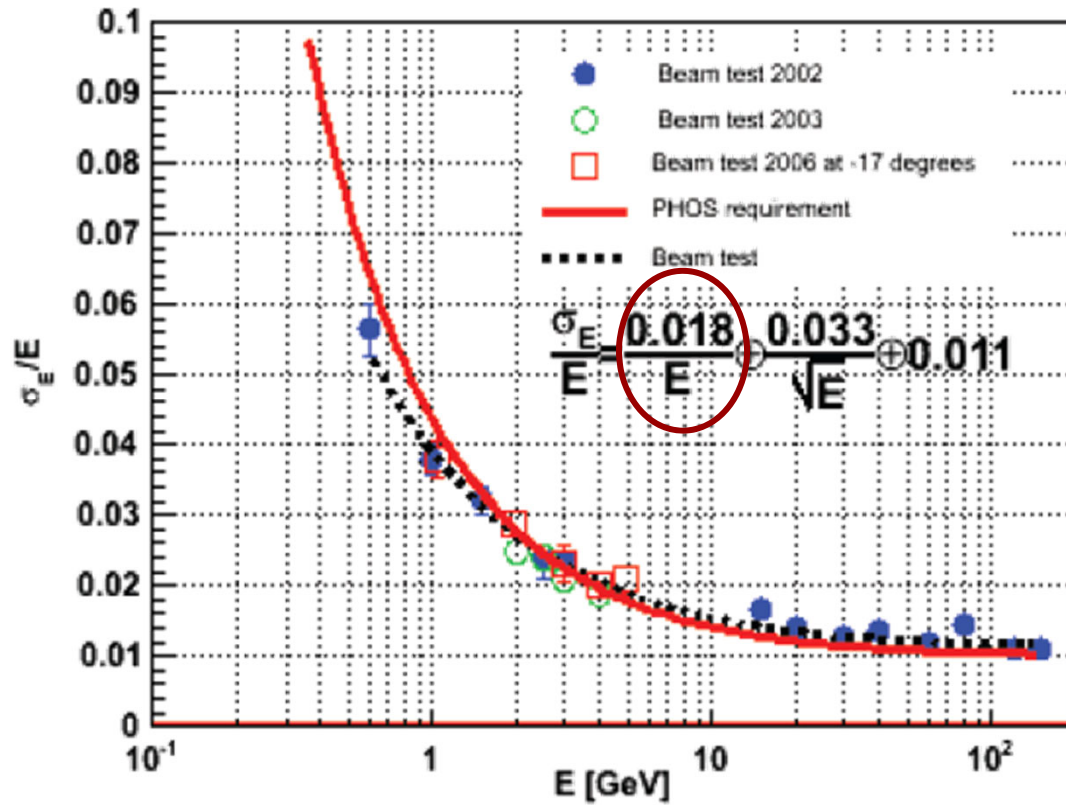


VCI 2007



# PHOS: energy and position resolution

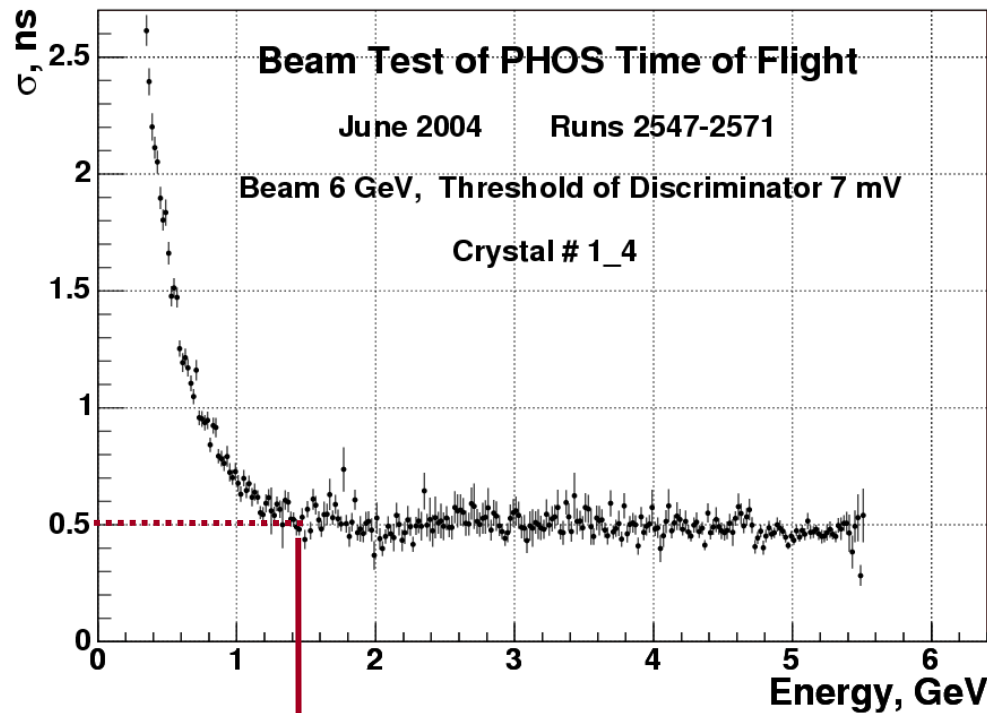
(Daicu Zhou at QM2006)



$\langle \sigma_x \rangle \sim 2.7 \text{ mm @ } 2 \text{ GeV}$

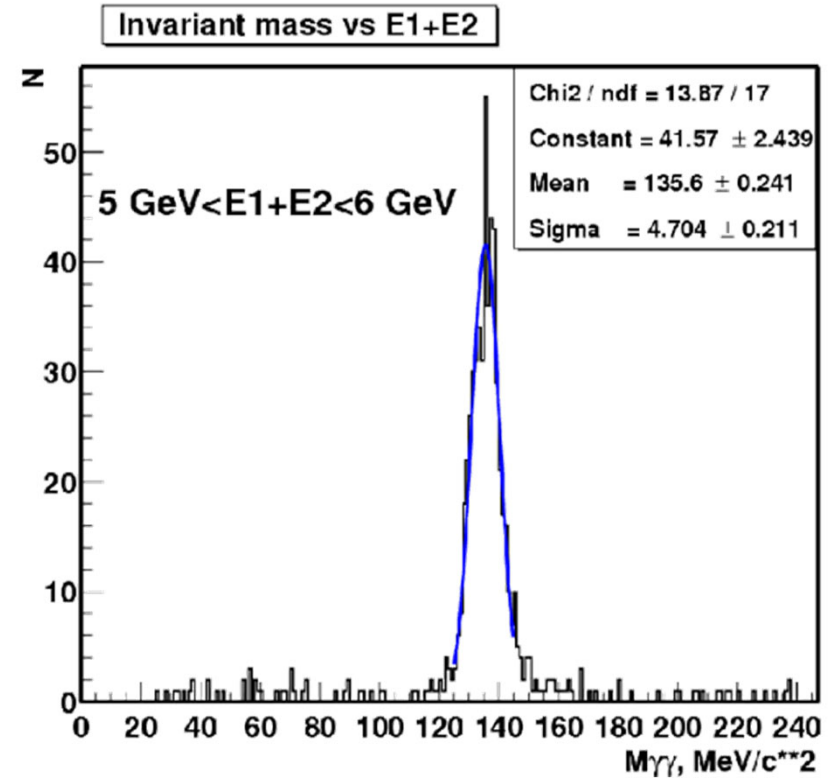
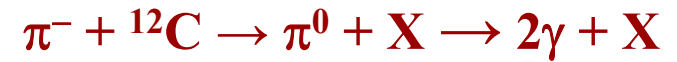
$\langle \sigma(E)/E \rangle \sim 3\%$ ,  $\sigma(\sigma(E)/E) \sim 0.1\%$  @ 2 GeV

Timing resolution measurement with the electron beam. Standard start-stop method with an external trigger



$\sigma \sim 0.5 \text{ ns at } E > 1.5 \text{ GeV}$

## Invariant mass spect



$\sigma(\pi^0) = 4.7 \text{ MeV}/c^2$

$\pi^0$  is hard to see in ion-ion collisions

→ potential problem for intercalibration with data