

“Complementi di Fisica”

Lecture 15

Livio Lanceri
Università di Trieste

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In these lectures

- **Contents**
 - Some results from Quantum Mechanics:
 - hydrogen atom
 - angular momentum and spin
 - identical particles: bosons and fermions
 - Pauli exclusion principle for fermions
 - Some consequences
 - Periodic table of the elements
 - “nearly-free electron gas” in a crystal: filling of available states
- **Reference textbooks**
 - Griffiths
 - Bernstein
 - Taylor-Zafiratos-Dubson

Some QM results

(3-d) Hydrogen atom
angular momentum, spin

Systems with many particles:
fermions and Pauli principle

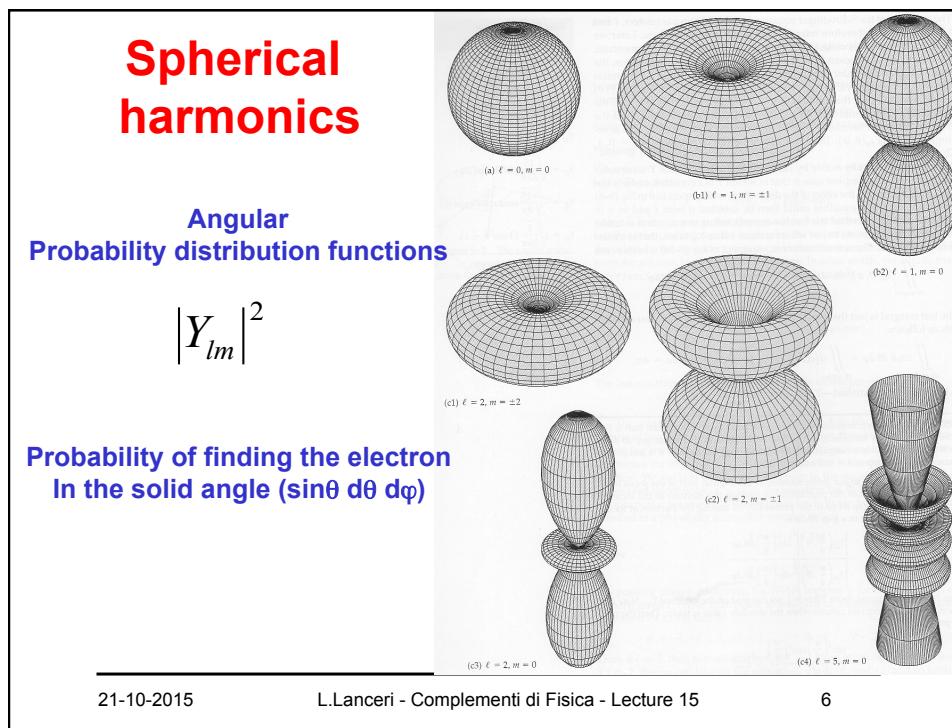
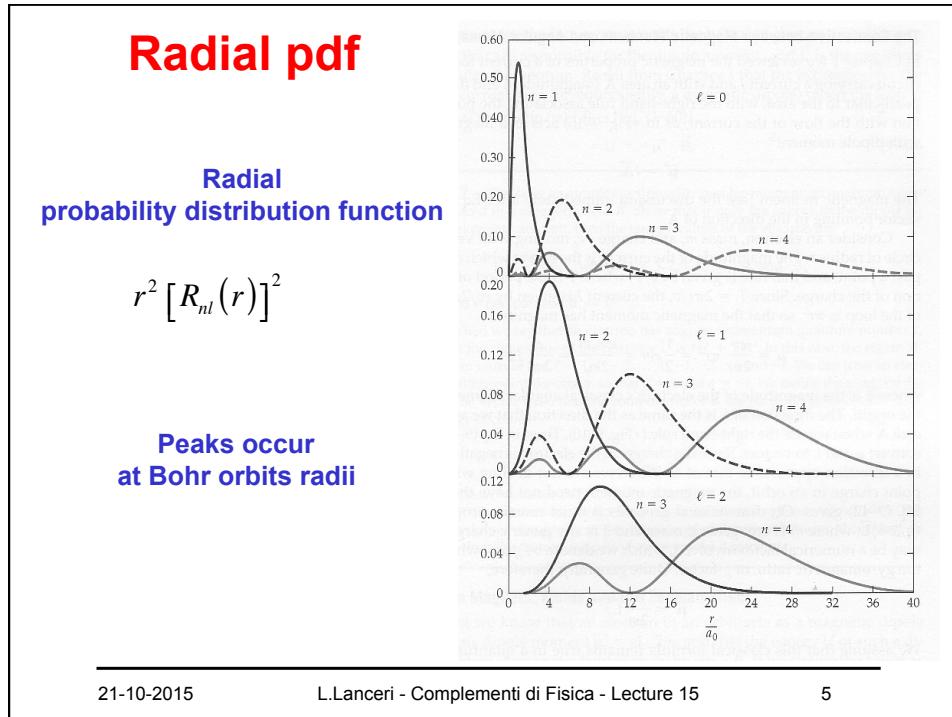
*(These subjects are discussed in more detail
in the textbooks (Bernstein, Griffiths, ...))*

Hydrogen atom

- “simple”: time-independent Schrödinger equation for the electron:
 - central “Coulomb” potential energy $V(r) = q^2/(4\pi\epsilon_0 r)$
 - spherical coordinates (r, θ, ϕ)
 - Separation of variables (3)
 - 3 integer quantum numbers identify each solution
 $\psi_{n,l,m}(r, \theta, \phi) = R_{nl}(r) Y_l^m(\theta, \phi)$

$n = 1, 2, 3 \dots$... principal quantum number	Energy (= Bohr !)
$l = 0, 1, 2, \dots n - 1$... azimuthal quantum number	Angular
$m = -l \text{ to } l$... magnetic orbital quantum number	momentum

$$\hat{H}\psi_{nlm} = E_n\psi_{nlm} \quad \hat{L}^2\psi_{nlm} = \hbar^2 l(l+1)\psi_{nlm} \quad \hat{L}_z\psi_{nlm} = \hbar m\psi_{nlm}$$



Angular momentum

- **Also angular momentum is quantized !**
 - One can only measure simultaneously the magnitude square and one component (the components don't commute !)
 - **Cartesian and spherical coordinates:**
$$\vec{L} = \vec{r} \times \vec{p} \quad \hat{L}_x = y\hat{p}_z - z\hat{p}_y \quad \hat{L}_y = z\hat{p}_x - x\hat{p}_z \quad \hat{L}_z = x\hat{p}_y - y\hat{p}_x$$

$$\hat{L}^2 = \hat{L}_x^2 + \hat{L}_y^2 + \hat{L}_z^2 = -\hbar^2 \left(\frac{\partial^2}{\partial \theta^2} + \cot \theta \frac{\partial}{\partial \theta} + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \right)$$

$$\hat{L}_z = -i\hbar \frac{\partial}{\partial \phi}$$
 - **Eigenvalues and eigenfunctions: spherical harmonics are eigenfunctions of the angular momentum operators**
- $$\hat{L}^2 Y_{lm}(\theta, \phi) = \hbar^2 l(l+1) Y_{lm}(\theta, \phi) \quad l = 1, 2, 3, \dots$$
- $$\hat{L}_z Y_{lm}(\theta, \phi) = \hbar m Y_{lm}(\theta, \phi) \quad -l \leq m \text{ integer} \leq +l$$

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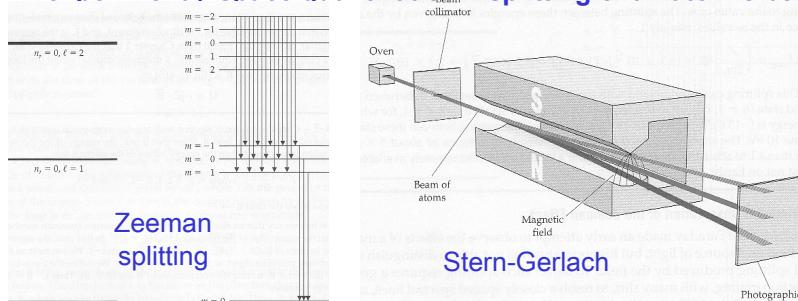
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Magnetic effects

- On dimensional grounds, for a charged particle with angular momentum we expect a magnetic moment and a contribution to potential energy when interacting with an external \vec{B} field:

$$\vec{\mu} = g \frac{q}{2m} \vec{L} \quad U = -\vec{\mu} \cdot \vec{B}$$

“Zeeman effect” (splitting of degenerate levels) and Stern-Gerlach experiment (“space quantization”: splitting of an atomic beam)



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Spin

- Elementary particles carry also an “intrinsic” angular momentum (“spin” S) besides the “orbital” angular momentum (L)
 - The eigenstates are not the spherical harmonics: not functions of θ, ϕ at all!
 - The quantum numbers s, m can be half-integer
 - The magnitude s is specific and fixed for each elementary particle, and is called “spin”
 - Electrons have spin $s = \frac{1}{2}$, with two possible eigenstates: “up” and “down”

$$\hat{S}^2 |sm\rangle = \hbar^2 s(s+1) |sm\rangle \quad s = 0, \frac{1}{2}, 1, \frac{3}{2}, \dots; \quad m = -s, -s+1, \dots, s$$

$$\hat{S}_z |sm\rangle = \hbar m |sm\rangle$$
 electrons: eigenstates and eigenvalues:

$$s = 1/2 \quad \chi_+ = \begin{pmatrix} 1 \\ 2 \\ 2 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \text{ eigenvalue } +\frac{\hbar}{2}$$

$$\chi_- = \begin{pmatrix} 1 \\ 2 \\ -1 \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}, \text{ eigenvalue } -\frac{\hbar}{2}$$

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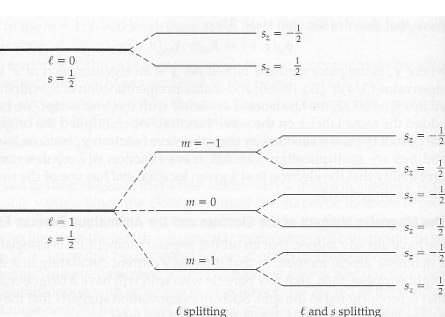
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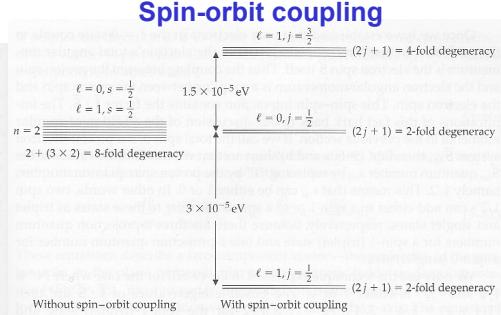
Spin: observable effects

- For example:
 - “Anomalous Zeeman effect”: further level splitting in strong B fields
 - “Fine Structure” level splitting due to “spin-orbit coupling”

Anomalous Zeeman effect



Spin-orbit coupling



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This is not the end...

- **Hydrogen has been a very interesting laboratory:**
 - Orders of magnitude of different effects, treated as “perturbations”, in terms of the a-dimensional “fine structure constant” α , expressing the strength of the electromagnetic coupling:

$$\alpha \equiv \frac{e^2}{4\pi\epsilon_0\hbar c} \cong \frac{1}{137.036}$$

Table 6.1: Hierarchy of corrections to the Bohr energies of hydrogen.

Bohr energies:	of order	$\alpha^2 mc^2$	
Fine structure:	of order	$\alpha^4 mc^2$	← Relativity, spin-orbit
Lamb shift:	of order	$\alpha^5 mc^2$	← Coulomb field quantization
Hyperfine splitting:	of order	$(m/m_p)\alpha^4 mc^2$	← Electron-proton magnetic moments

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Many-particle systems (just a hint...)

- Identical particles
- Bosons and fermions
- Pauli Principle
- Periodic table

Identical particles

- **Many-particle systems? Let's start with two:**

- Wave function, probability distribution, hamiltonian; S.equation

$$\Psi(\vec{r}_1, \vec{r}_2, t) \quad |\Psi(\vec{r}_1, \vec{r}_2, t)|^2 d\vec{r}_1 d\vec{r}_2$$

$$i\hbar \frac{\partial \Psi}{\partial t} = \hat{H}\Psi \quad \hat{H} = -\frac{\hbar^2}{2m_1} \nabla_1^2 - \frac{\hbar^2}{2m_2} \nabla_2^2 + V(\vec{r}_1, \vec{r}_2, t)$$

- For time-independent potentials: time-indep. S.eq. and stationary states

$$\Psi(\vec{r}_1, \vec{r}_2, t) = \psi(\vec{r}_1, \vec{r}_2) e^{-iEt/\hbar}$$

$$-\frac{\hbar^2}{2m_1} \nabla_1^2 \psi - \frac{\hbar^2}{2m_2} \nabla_2^2 \psi + V(\vec{r}_1, \vec{r}_2) \psi = E\psi$$

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Bosons and fermions

- For distinguishable particles (for instance, an electron and a positron):
 - particle 1 is in the (one-particle) state $\psi_a(\mathbf{r}_1)$
 - particle 2 in state $\psi_b(\mathbf{r}_2)$

$$\psi(\vec{r}_1, \vec{r}_2, t) = \psi_a(\vec{r}_1) \psi_b(\vec{r}_2)$$

- But: identical particles (for instance, two electrons) are truly indistinguishable in quantum mechanics:

- There are two possible ways to construct the wave-function:

+ "symmetric": bosons

$$\psi_{\pm}(\vec{r}_1, \vec{r}_2) = A [\psi_a(\vec{r}_1) \psi_b(\vec{r}_2) \pm \psi_b(\vec{r}_1) \psi_a(\vec{r}_2)]$$

- "anti-symmetric": fermions

- All particles with integer spin are bosons

- All particles with half-integer spin are fermions

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Fermions and Pauli principle

- Connection between spin and “statistics” (or wave-function exchange symmetry)
 - can be proven in relativistic QM
 - must be taken as an axiom in non-relativistic QM

- Pauli exclusion principle:
 - Two fermions (anti-symmetric w.f.) cannot occupy the same state! Indeed:

$$\psi_a = \psi_b \Rightarrow \psi_-(\vec{r}_1, \vec{r}_2) = A[\psi_a(\vec{r}_1)\psi_a(\vec{r}_2) - \psi_a(\vec{r}_1)\psi_a(\vec{r}_2)] = 0$$

- It can be shown that:
 - The exchange operator P is a “compatible observable” commuting with H \Rightarrow one can find solutions that are either symmetric or antisymmetric
 - For identical particles, the wave function is required to be symmetric (for bosons) or anti-symmetric (for fermions)

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Pauli Principle: consequences for electrons

- For electrons the total wave-function (including spin) must be anti-symmetric, and they cannot occupy the same state (two per level allowed, with opposite spin).
- The anti-symmetry requirement allows some wave-function configurations, prohibits others: equivalent to an “exchange force”
- Filling of available levels by electrons in a box (neglecting interactions among electrons!): Fermi level= highest energy level occupied at T = 0K (see exercises)
- “degeneracy pressure”: even neglecting electric interactions between electrons, the Pauli principle implies that “the closest that two electrons can get to each other is roughly a half of the DeBroglie wavelength corresponding to the Fermi energy” (see exercises)

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Pauli principle: Periodic table of elements

- Multi-electron atoms are treated by approximate methods:
 - wave functions are modified (and called “orbitals”), but:
 - they are labeled by the same quantum numbers n, l, m, and:
 - Orbitals are filled by electrons following the Pauli exclusion principle: two electrons cannot have the same quantum numbers (state)

Table A.3 Energy States and the Electronic Configuration in Elements 1–14. Atoms are assumed to be in the ground state.

<i>Quantum Numbers</i>	n	1	1	2	2	2	2	2	2	2	3	3	3	3	3	3
	l	0	0	0	0	1	1	1	1	1	0	0	1	1	1	1
	m	0	0	0	0	-1	-1	0	0	1	1	0	0	-1	-1	0
	s	$\frac{1}{2}$	$-\frac{1}{2}$	$\frac{1}{2}$												
	State	1s	1s	2s	2s	2p	2p	2p	2p	2p	3s	3s	3p	3p	3p	3p

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Periodic table of the elements

Atomic Number	Element	Filled States										Electronic Configuration		
		1s	1s	2s	2s	2p	2p	2p	2p	2p	3s	3p	3p	3p
1	H	↑												1s
2	He	↑	↓											1s ²
3	Li	↑	↓	↑										1s ² 2s
4	Be	↑	↓	↑	↓									1s ² 2s ²
5	B	↑	↓	↑	↓	↑								1s ² 2s ² 2p
6	C	↑	↓	↑	↓	↑	↓	↑						1s ² 2s ² 2p ²
7	N	↑	↓	↑	↓	↑	↓	↑	↓	↑				1s ² 2s ² 2p ³
8	O	↑	↓	↑	↓	↑	↓	↑	↓	↑				1s ² 2s ² 2p ⁴
9	F	↑	↓	↑	↓	↑	↓	↑	↓	↑	↑			1s ² 2s ² 2p ⁵
10	Ne	↑	↓	↑	↓	↑	↓	↑	↓	↑	↓	↑		1s ² 2s ² 2p ⁶
11	Na	↑	↓	↑	↓	↑	↓	↑	↓	↑	↓	↑		1s ² 2s ² 2p ⁶ 3s
12	Mg	↑	↓	↑	↓	↑	↓	↑	↓	↑	↓	↑		1s ² 2s ² 2p ⁶ 3s ²
13	Al	↑	↓	↑	↓	↑	↓	↑	↓	↑	↓	↑		1s ² 2s ² 2p ⁶ 3s ² 3p
14	Si	↑	↓	↑	↓	↑	↓	↑	↓	↑	↓	↑		1s ² 2s ² 2p ⁶ 3s ² 3p ²

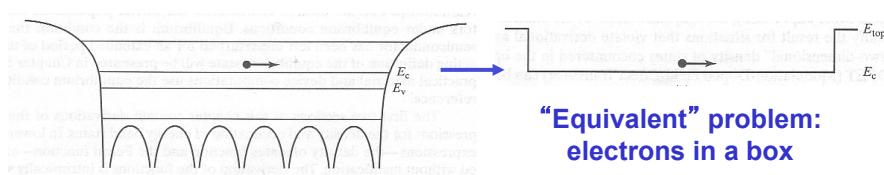
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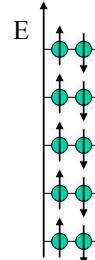
Electrons in a solid ...

- Individual electrons



**“Equivalent” problem:
electrons in a box**

- Interactions among electrons can be neglected to 1st approx.
 - “screening” by ions (L.Landau: “nearly-free electrons”)
- Pauli exclusion principle
 - Obeyed by electrons filling up the available states
- Fermi-Dirac probability distribution
 - Occupation probability for the available states



Lecture 15 - summary

- Some results from Quantum Mechanics
 - hydrogen atom from Schrodinger equation
 - angular momentum and spin
 - identical particles: bosons and fermions
 - Pauli exclusion principle for fermions
- Some consequences
 - Periodic table of the elements
 - “nearly-free electron gas” in a crystal: filling of available states