

# Advanced Quantum Mechanics

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# Quantum Algorithms

We will study the three historically most important algorithms:

- Simple ones (Deutsch, Deutsch-Jozsa...)
- Grover (search in a data base)
- Shor (factorization)

What is special about quantum algorithms?

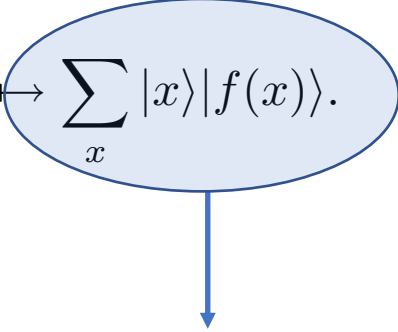
# Quantum Parallelism

Given an input  $x$ , a typical quantum computer computes  $f(x)$  in such a way as

$$U_f : |x\rangle|0\rangle \mapsto |x\rangle|f(x)\rangle, \quad (4.61)$$

where  $U_f$  is a unitary matrix that implements the function  $f$ .

Suppose  $U_f$  acts on the input which is a superposition of many states. Since  $U_f$  is a linear operator, it acts simultaneously on all the vectors that constitute the superposition. Thus the output is also a superposition of all the results;

$$U_f : \sum_x |x\rangle|0\rangle \mapsto \sum_x |x\rangle|f(x)\rangle. \quad (4.62)$$


All values of the function computed at once. Very easy!!  
But... measurements will make the wave function collapse  
giving only one output. No advantage

# Quantum Algorithms

The goal of a quantum algorithm is to operate in such a way that the particular outcome we want to observe has a larger probability to be measured than the other outcomes.



# Deutsch Algorithm

Let  $f : \{0, 1\} \rightarrow \{0, 1\}$  be a binary function. Note that there are only four possible  $f$ , namely

$$\begin{aligned} f_1 : 0 \mapsto 0, 1 \mapsto 0, & \quad f_2 : 0 \mapsto 1, 1 \mapsto 1, \\ f_3 : 0 \mapsto 0, 1 \mapsto 1, & \quad f_4 : 0 \mapsto 1, 1 \mapsto 0. \end{aligned}$$

The first two cases,  $f_1$  and  $f_2$ , are called constant, while the rest,  $f_3$  and  $f_4$ , are balanced. If we only have classical resources, we need to evaluate  $f$  twice to tell if  $f$  is constant or balanced. There is a quantum algorithm, however, with which it is possible to tell if  $f$  is constant or balanced with a single evaluation of  $f$ , as was shown by Deutsch [2].

# Deutsch Algorithm

First we need to turn the classical function  $f(x)$  into a quantum one.

1. Make it reversible.
2. Define it on the computational basis to act like the classical circuit and extend it by linearity.

$$U_f : |x, y\rangle \mapsto |x, y \oplus f(x)\rangle$$

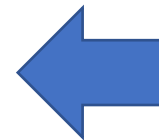
Where  $\oplus$  is addition mod 2.

# Deutsch Algorithm

The algorithm is structured as follows.

1. Start with the state  $|01\rangle$ .
2. Apply an Hadamard on both qubits:  $\frac{1}{2}(|00\rangle - |01\rangle + |10\rangle - |11\rangle)$
3. Apply the operator  $U_f$  implementing the function

$$\begin{aligned} & \frac{1}{2}(|0, f(0)\rangle - |0, 1 \oplus f(0)\rangle + |1, f(1)\rangle - |1, 1 \oplus f(1)\rangle) \\ = & \frac{1}{2}(|0, f(0)\rangle - |0, \neg f(0)\rangle + |1, f(1)\rangle - |1, \neg f(1)\rangle), \end{aligned}$$



Quantum  
parallelism:  
all values  
computed at  
once

# Deutsch Algorithm

## 4. Apply an Hadamard to the first qubit

$$\frac{1}{2\sqrt{2}} [ (|0\rangle + |1\rangle)(|f(0)\rangle - |\neg f(0)\rangle) + (|0\rangle - |1\rangle)(|f(1)\rangle - |\neg f(1)\rangle) ]$$

The wave function reduces to

$$\frac{1}{\sqrt{2}} |0\rangle (|f(0)\rangle - |\neg f(0)\rangle) \quad (5.1)$$

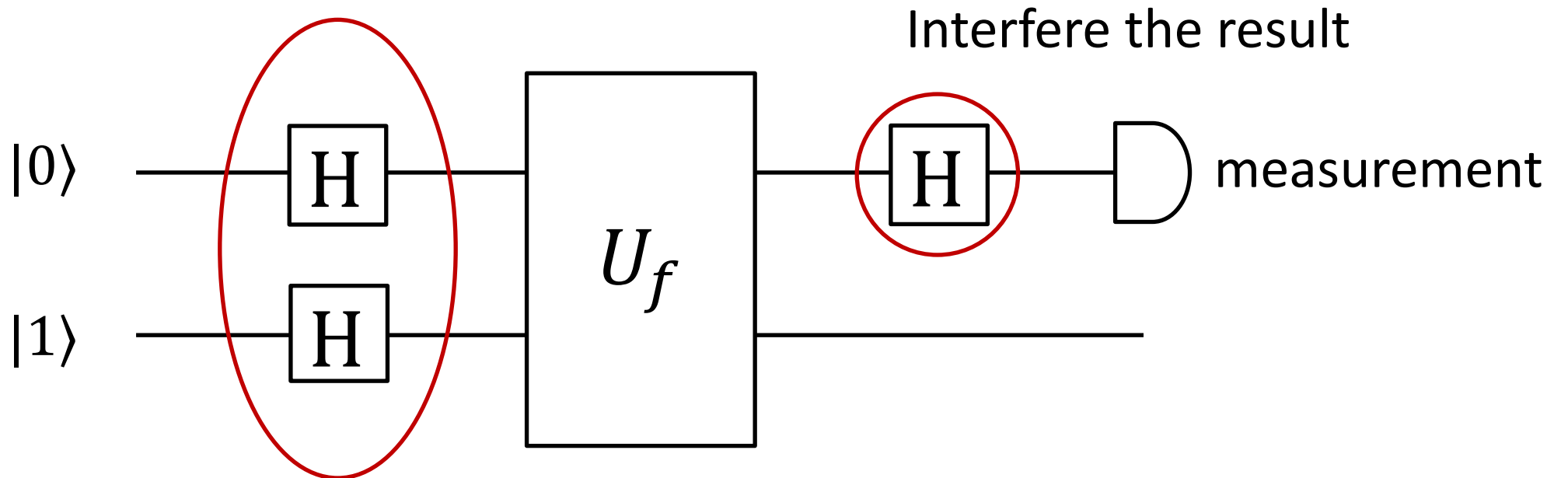
in case  $f$  is constant, for which  $|f(0)\rangle = |f(1)\rangle$ , and

$$\frac{1}{\sqrt{2}} |1\rangle (|f(0)\rangle - |\neg f(0)\rangle) \quad (5.2)$$

if  $f$  is balanced, for which  $|\neg f(0)\rangle = |f(1)\rangle$ . Therefore the measurement of the first qubit tells us whether  $f$  is constant or balanced.

# Deutsch Algorithm

5. Measure the first qubit



Calculate the function on both input values simultaneously

# Deutsch-Jozsa Algorithm

Let us first define the **Deutsch-Jozsa problem**. Suppose there is a binary function

$$f : S_n \equiv \{0, 1, \dots, 2^n - 1\} \rightarrow \{0, 1\}. \quad (5.3)$$

We require that  $f$  be either *constant* or *balanced* as before. When  $f$  is constant, it takes a constant value 0 or 1 irrespective of the input value  $x$ . When it is balanced the value  $f(x)$  for the half of  $x \in S_n$  is 0, while it is 1 for the rest of  $x$ .

It is clear that we need at least  $2^{n-1} + 1$  steps, in the worst case with classical manipulations, to make sure if  $f(x)$  is constant or balanced with 100% confidence. It will be shown below that the number of steps reduces to a single step if we are allowed to use a quantum algorithm.

# Deutsch-Jozsa Algorithm

1. Prepare an  $(n + 1)$ -qubit register in the state  $|\psi_0\rangle = |0\rangle^{\otimes n} \otimes |1\rangle$ . First  $n$  qubits work as input qubits, while the  $(n + 1)$ st qubit serves as a “scratch pad.” Such qubits, which are neither input qubits nor output qubits, but work as a scratch pad to store temporary information are called **ancillas** or **ancillary qubits**.
2. Apply the Walsh-Hadamard transformation to the register. Then we have the state

$$\begin{aligned} |\psi_1\rangle &= U_{\text{H}}^{\otimes n+1} |\psi_0\rangle = \frac{1}{\sqrt{2^n}} (|0\rangle + |1\rangle)^{\otimes n} \otimes \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle) \\ &= \frac{1}{\sqrt{2^n}} \sum_{x=0}^{2^n-1} |x\rangle \otimes \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle). \end{aligned} \tag{5.4}$$

# Deutsch-Jozsa Algorithm

3. Apply  $U_f|x\rangle|c\rangle = |x\rangle|c \oplus f(x)\rangle$

The state changes into

$$\begin{aligned} |\psi_2\rangle &= U_f|\psi_1\rangle \\ &= \frac{1}{\sqrt{2^n}} \sum_{x=0}^{2^n-1} |x\rangle \frac{1}{\sqrt{2}} (|f(x)\rangle - |\neg f(x)\rangle) \\ &= \frac{1}{\sqrt{2^n}} \sum_x (-1)^{f(x)} |x\rangle \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle). \end{aligned} \quad (5.5)$$

Although the gate  $U_f$  is applied once for all, it is applied to *all* the  $n$ -qubit states  $|x\rangle$  simultaneously.



# Deutsch-Jozsa Algorithm

4. The Walsh-Hadamard transformation (4.11) is applied on the first  $n$  qubits next. We obtain

$$|\psi_3\rangle = (W_n \otimes I)|\psi_2\rangle = \frac{1}{\sqrt{2^n}} \sum_{x=0}^{2^n-1} (-1)^{f(x)} U_{\text{H}}^{\otimes n} |x\rangle \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle). \quad (5.6)$$

# On the Hadamard gate

It is instructive to write the action of the one-qubit Hadamard gate in the following form,

$$U_{\text{H}}|x\rangle = \frac{1}{\sqrt{2}}(|0\rangle + (-1)^x|1\rangle) = \frac{1}{\sqrt{2}} \sum_{y \in \{0,1\}} (-1)^{xy} |y\rangle,$$

where  $x \in \{0, 1\}$ , to find the resulting state. The action of the Walsh-Hadamard transformation on  $|x\rangle = |x_{n-1} \dots x_1 x_0\rangle$  yields

$$\begin{aligned} W_n|x\rangle &= (U_{\text{H}}|x_{n-1}\rangle)(U_{\text{H}}|x_{n-2}\rangle) \dots (U_{\text{H}}|x_0\rangle) \\ &= \frac{1}{\sqrt{2^n}} \sum_{y_{n-1}, y_{n-2}, \dots, y_0 \in \{0,1\}} (-1)^{x_{n-1}y_{n-1} + x_{n-2}y_{n-2} + \dots + x_0y_0} \\ &\quad \times |y_{n-1}y_{n-2} \dots y_0\rangle \\ &= \frac{1}{\sqrt{2^n}} \sum_{y=0}^{2^n-1} (-1)^{x \cdot y} |y\rangle, \end{aligned} \tag{5.7}$$

where  $x \cdot y = x_{n-1}y_{n-1} \oplus x_{n-2}y_{n-2} \oplus \dots \oplus x_0y_0$ .

# Deutsch-Jozsa Algorithm

Coming back to step 4:

4. The Walsh-Hadamard transformation (4.11) is applied on the first  $n$  qubits next. We obtain

$$\begin{aligned} |\psi_3\rangle &= (W_n \otimes I)|\psi_2\rangle = \frac{1}{\sqrt{2^n}} \sum_{x=0}^{2^n-1} (-1)^{f(x)} U_{\text{H}}^{\otimes n} |x\rangle \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle). \quad (5.6) \\ &= \frac{1}{2^n} \left( \sum_{x,y=0}^{2^n-1} (-1)^{f(x)} (-1)^{x \cdot y} |y\rangle \right) \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle). \end{aligned}$$

As we will see, this operation will make the different terms interfere in order to read the desired result

# Deutsch-Jozsa Algorithm

5. The first  $n$  qubits are measured. Suppose  $f(x)$  is constant. Then  $|\psi_3\rangle$  is put in the form

$$|\psi_3\rangle = \frac{1}{2^n} \sum_{x,y} (-1)^{x \cdot y} |y\rangle \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle)$$

up to an overall phase. Now let us consider the summation

$$\frac{1}{2^n} \sum_{x=0}^{2^n-1} (-1)^{x \cdot y}$$

with a fixed  $y \in S_n$ . Clearly it vanishes since  $x \cdot y$  is 0 for half of  $x$  and 1 for the other half of  $x$  unless  $y = 0$ . Therefore the summation yields  $\delta_{y0}$ . Now the state reduces to

$$|\psi_3\rangle = |0\rangle^{\otimes n} \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle),$$

and the measurement outcome of the first  $n$  qubits is always  $00\dots 0$ .

Example with 3 qubits.

Take  $y = 110$ . Then

$$x \cdot y = x_2 \oplus x_1$$

x	$x_2 \oplus x_1$
000	0
001	0
010	1
011	1
100	1
101	1
110	0
111	0

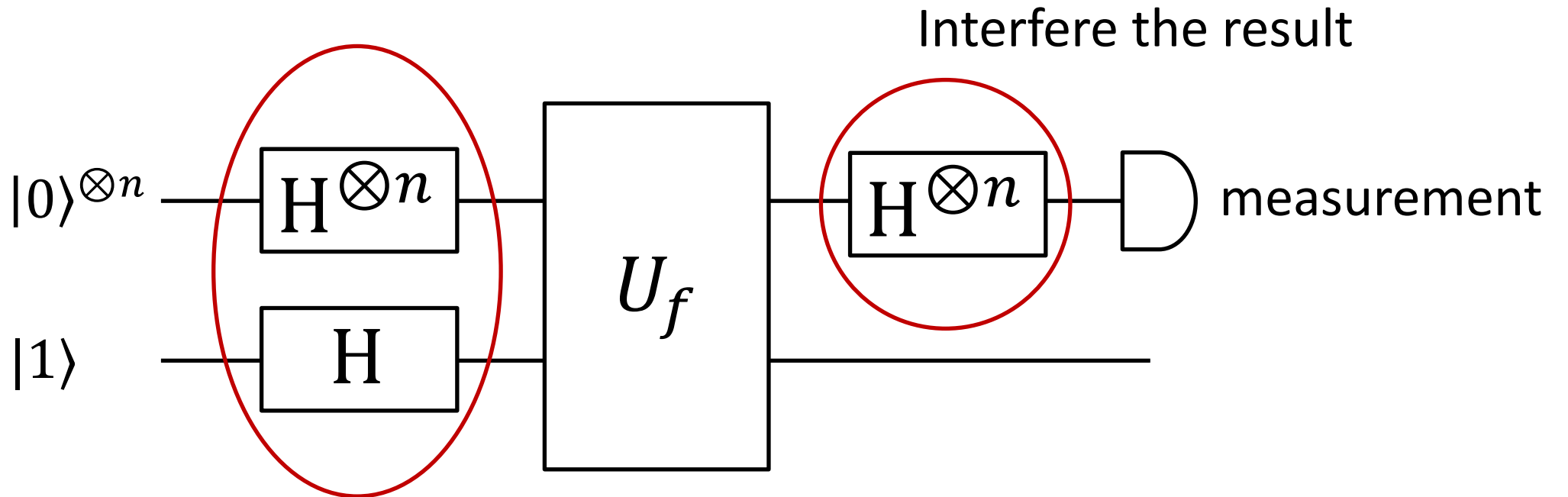
# Deutsch-Jozsa Algorithm

Suppose  $f(x)$  is balanced next. The probability amplitude of  $|y = 0\rangle$  in  $|\psi_3\rangle$  is proportional to

$$\sum_{x=0}^{2^n-1} (-1)^{f(x)} (-1)^{x \cdot 0} = \sum_{x=0}^{2^n-1} (-1)^{f(x)} = 0.$$

Therefore the probability of obtaining measurement outcome  $00 \dots 0$  for the first  $n$  qubits vanishes. In conclusion, the function  $f$  is constant if we obtain  $00 \dots 0$  upon the measurement of the first  $n$  qubits in the state  $|\psi_3\rangle$ , and it is balanced otherwise.

# Deutsch-Jozsa Algorithm



Calculate the function on both input values simultaneously

# Bernstein-Vazirani Algorithm

The **Bernstein-Vazirani algorithm** is a special case of the Deutsch-Jozsa algorithm, in which  $f(x)$  is given by  $f(x) = c \cdot x$ , where  $c = c_{n-1}c_{n-2} \dots c_0$  is an  $n$ -bit binary number [4]. Our aim is to find  $c$  with the smallest number of evaluations of  $f$ . If we apply the Deutsch-Jozsa algorithm with this  $f$ , we obtain

$$|\psi_3\rangle = \frac{1}{2^n} \left[ \sum_{x,y=0}^{2^n-1} (-1)^{c \cdot x} (-1)^{x \cdot y} |y\rangle \right] \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle).$$

Let us fix  $y$  first. If we take  $y = c$ , we obtain

$$\sum_x (-1)^{c \cdot x} (-1)^{x \cdot c} = \sum_x (-1)^{2c \cdot x} = 2^n.$$

# Bernstein-Vazirani Algorithm

If  $y \neq c$ , on the other hand, there will be the same number of  $x$  such that  $c \cdot x = 0$  and  $x$  such that  $c \cdot x = 1$  in the summation over  $x$  and, as a result, the probability amplitude of  $|y \neq c\rangle$  vanishes. By using these results, we end up with

$$|\psi_3\rangle = |c\rangle \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle). \quad (5.9)$$

We are able to tell what  $c$  is by measuring the first  $n$  qubits.



# Exercise

**EXERCISE 5.1** Let us take  $n = 2$  for definiteness. Consider the following cases and find the final wave function  $|\psi_3\rangle$  and evaluate the measurement outcomes and their probabilities for each case.

(1)  $f(x) = 1 \forall x \in S_2$ .

(2)  $f(00) = f(01) = 1, f(10) = f(11) = 0$ .

(3)  $f(00) = 0, f(01) = f(10) = f(11) = 1$ . (This function is neither constant nor balanced.)

**EXERCISE 5.2** Consider the Bernstein-Vazirani algorithm with  $n = 3$  and  $c = 101$ . Work out the quantum circuit depicted in Fig. 5.2 to show that the measurement outcome of the first three qubits is  $c = 101$ .

# Grover's search algorithm

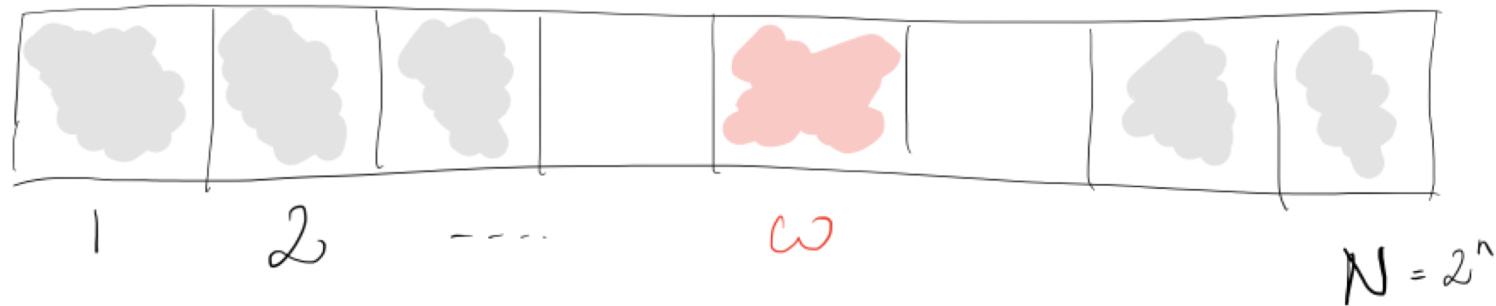
Suppose there is a stack of  $N = 2^n$  files, randomly placed, that are numbered by  $x \in S_n \equiv \{0, 1, \dots, N - 1\}$ . Our task is to find an algorithm which picks out a particular file which satisfies a certain condition.

In mathematical language, this is expressed as follows. Let  $f : S_n \rightarrow \{0, 1\}$  be a function defined by

$$f(x) = \begin{cases} 1 & (x = z) \\ 0 & (x \neq z), \end{cases} \quad (7.1)$$

where  $z$  is the address of the file we are looking for. It is assumed that  $f(x)$  is *instantaneously* calculable, such that this process does not require any computational steps. A function of this sort is often called an oracle as noted in Chapter 5. Thus, the problem is to find  $z$  such that  $f(z) = 1$ , given a function  $f : S_n \rightarrow \{0, 1\}$  which assumes the value 1 only at a single point.

# Grover's search algorithm



Clearly we have to check one file after another in a classical algorithm, and it will take  $O(N)$  steps on average. It is shown below that it takes only  $O(\sqrt{N})$  steps with Grover's algorithm. This is accomplished by *amplifying the amplitude of the vector  $|z\rangle$  while cancelling that of the vectors  $|x\rangle$  ( $x \neq z$ ).*

# Grover's search algorithm

We first need to implement the function  $f(x)$  quantum mechanically. We define  $U_f$  as follows (**oracle**)

$$U_f|x\rangle = (-1)^{f(x)}|x\rangle$$

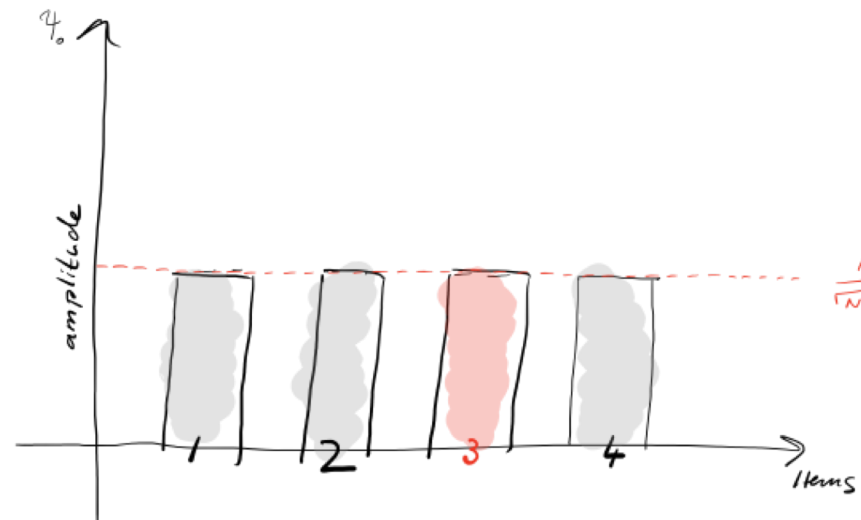
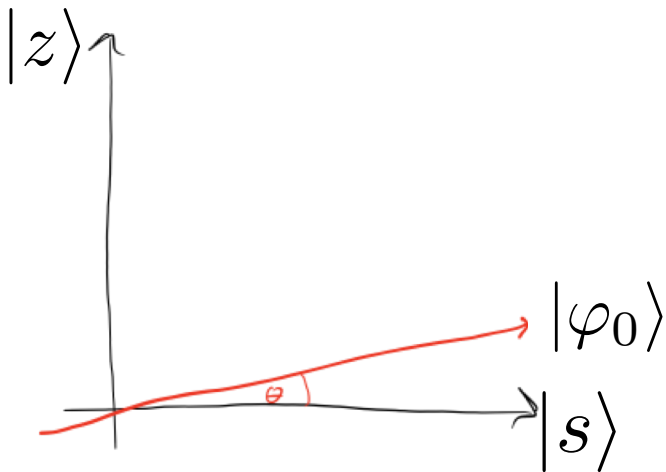
On the computational basis. We see that if  $x$  is an unmarked item, the oracle does nothing to the state. It flips the phase for the marked item. It is easy to see that

$$U_f = I - 2|z\rangle\langle z|$$

# Grover's search algorithm

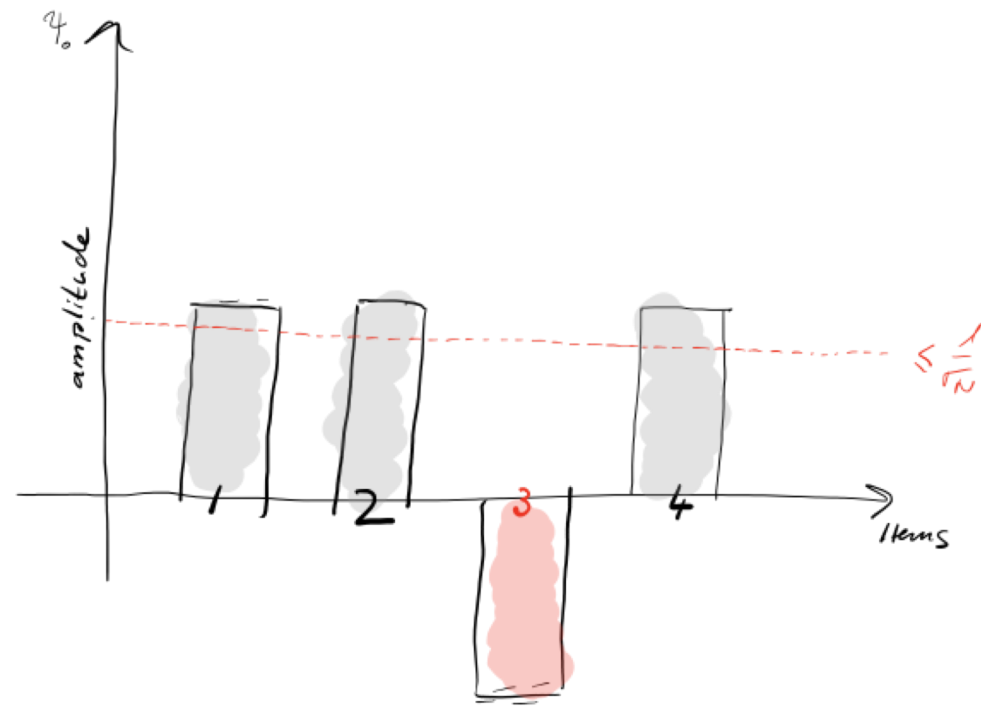
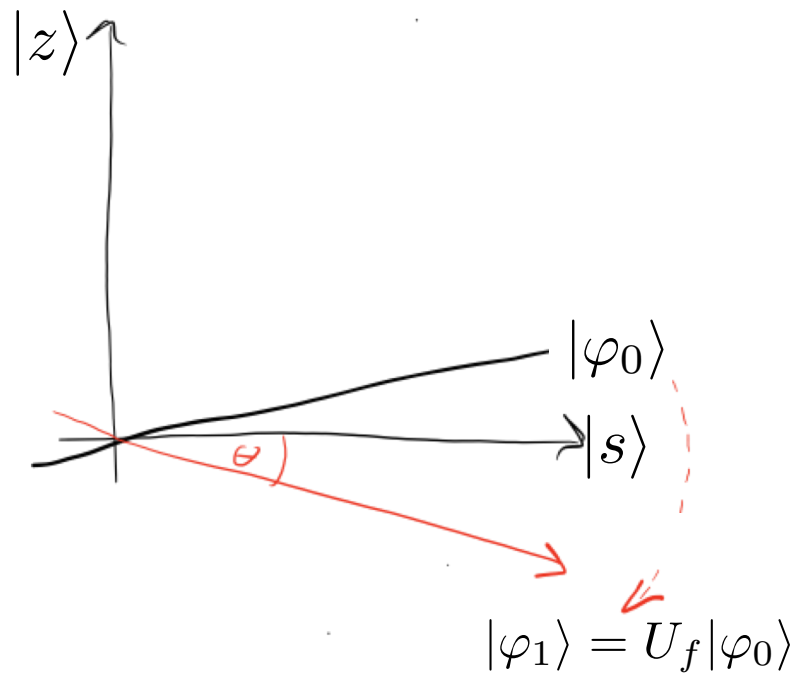
**Step 1:** Create an initially **equal weighted superposition** of all states (this is done with N Hadamard gates):

$$|\varphi_0\rangle = \frac{1}{\sqrt{N}} \sum_{x=0}^{N-1} |x\rangle.$$



# Grover's search algorithm

**Step 2: Apply the oracle  $U_f$ .** Geometrically this corresponds to a reflection of the state  $|z\rangle$  about  $|s\rangle$ . This transformation means that the amplitude in front of the  $|z\rangle$  state becomes negative, which in turn means that the average amplitude has been lowered.



# Grover's search algorithm

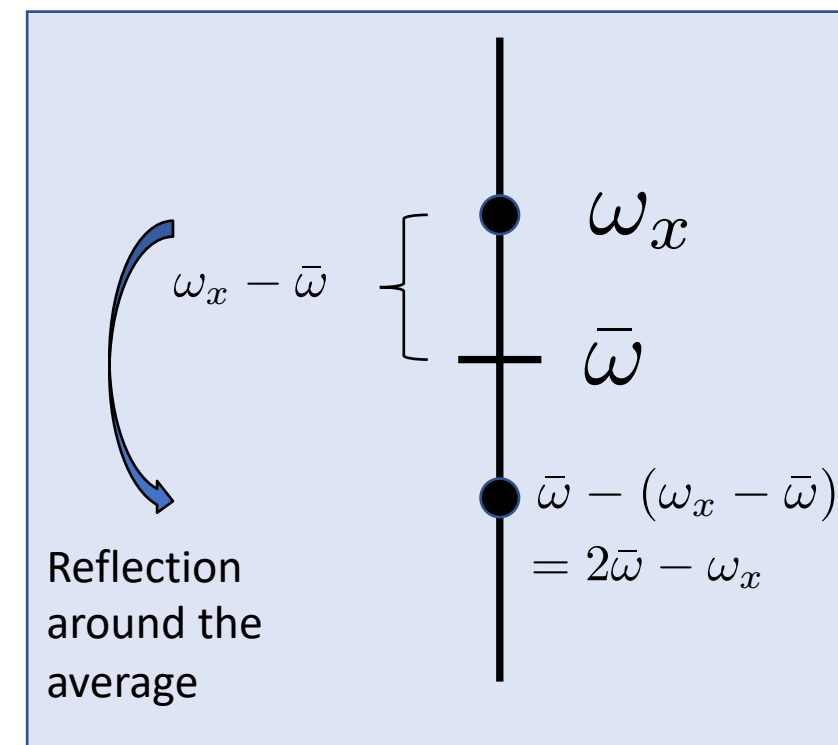
## Step 3: Apply the gate

$$D = -I + 2|\varphi_0\rangle\langle\varphi_0|.$$

The action of the gate is the following

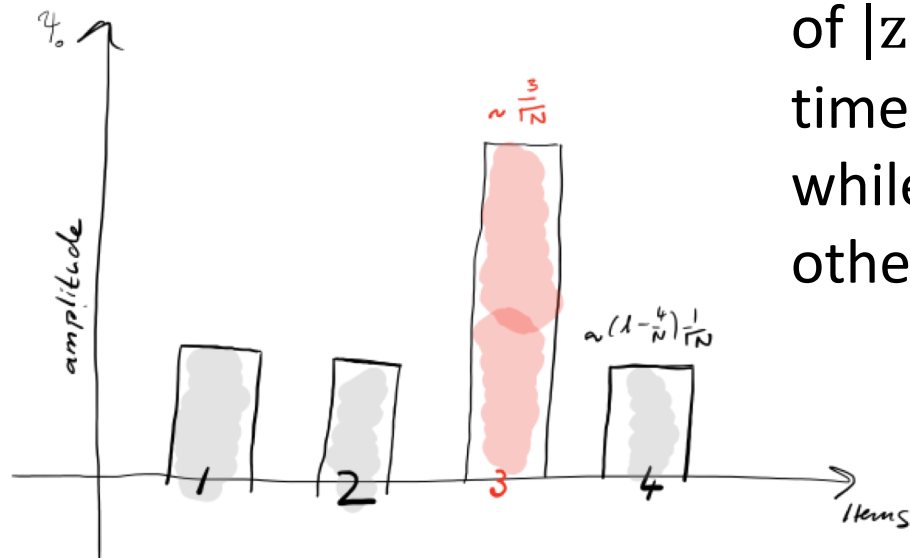
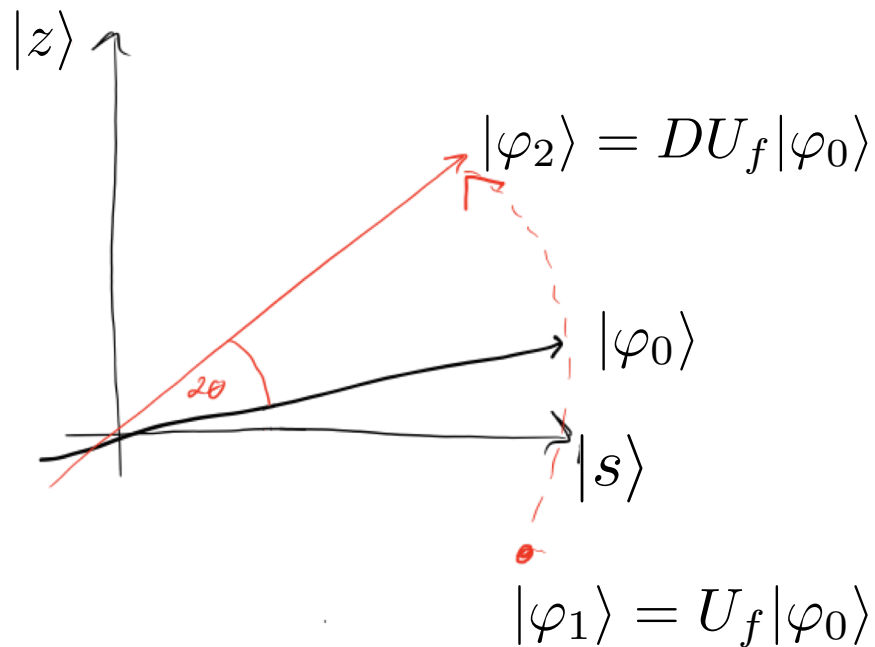
$$\begin{aligned} |\varphi\rangle = \sum_{x=0}^{N-1} \omega_x |x\rangle &\rightarrow D|\varphi\rangle = \left[ \frac{2}{N} \sum_{x,y=0}^{N-1} |x\rangle\langle y| \right] \sum_{z=0}^{N-1} \omega_z |z\rangle - \sum_{x=0}^{N-1} \omega_x |x\rangle \\ &= \frac{2}{N} \left[ \sum_{x=0}^{N-1} |x\rangle \right] \left[ \sum_{y=0}^{N-1} \omega_y \right] - \sum_{x=0}^{N-1} \omega_x |x\rangle = \sum_{x=0}^{N-1} (2\bar{\omega} - \omega_x) |x\rangle \end{aligned}$$

with  $\bar{\omega} = \frac{1}{N} \sum_{x=0}^{N-1} \omega_x$  average



# Grover's search algorithm

In our case we get



Since the average amplitude has been lowered by the first reflection, this transformation boosts the negative amplitude of  $|z\rangle$  to roughly three times its original value, while it decreases the other amplitudes.



# Grover's search algorithm

**Step 3:** go to step 2 and repeat the application of  $U_f$  and  $D$  a sufficient number of times. Let us call  $G_f = D U_f$ .

**PROPOSITION 7.2** Let us write

$$G_f^k |\varphi_0\rangle = a_k |z\rangle + b_k \sum_{x \neq z} |x\rangle \quad (7.17)$$

with the initial condition

$$a_0 = b_0 = \frac{1}{\sqrt{N}}.$$

Then the coefficients  $\{a_k, b_k\}$  for  $k \geq 1$  satisfy the recursion relations

$$a_k = \frac{N-2}{N} a_{k-1} + \frac{2(N-1)}{N} b_{k-1}, \quad (7.18)$$

$$b_k = -\frac{2}{N} a_{k-1} + \frac{N-2}{N} b_{k-1} \quad (7.19)$$

for  $k = 1, 2, \dots$

# Grover's search algorithm

*Proof.* It is easy to see the recursion relations are satisfied for  $k = 1$

Let  $G_f^{k-1}|\varphi_0\rangle = a_{k-1}|z\rangle + b_{k-1}\sum_{x\neq z}|x\rangle$ . Then

$$\begin{aligned} G_f^k|\varphi_0\rangle &= G_f \left( a_{k-1}|z\rangle + b_{k-1}\sum_{x\neq z}|x\rangle \right) \\ &= (-I + 2|\varphi_0\rangle\langle\varphi_0|) \left( -a_{k-1}|z\rangle + b_{k-1}\sum_{x\neq z}|x\rangle \right) \\ &= -b_{k-1}\sum_{x\neq z}|x\rangle + a_{k-1}|z\rangle + \frac{2}{\sqrt{N}}(N-1)b_{k-1}|\varphi_0\rangle - \frac{2a_{k-1}}{\sqrt{N}}|\varphi_0\rangle \\ &= -b_{k-1}\sum_{x\neq z}|x\rangle + a_{k-1}|z\rangle + \frac{2}{N}(N-1)b_{k-1}\sum_x|x\rangle - \frac{2a_{k-1}}{N}\sum_x|x\rangle \\ &= \left[ \frac{N-2}{N}a_{k-1} + \frac{2(N-1)}{N}b_{k-1} \right] |z\rangle + \left[ -\frac{2}{N}a_{k-1} + \frac{N-2}{N}b_{k-1} \right] \sum_{x\neq z}|x\rangle, \end{aligned}$$

and proposition is proved. ■

# Grover's search algorithm

**PROPOSITION 7.3** The solutions of the recursion relations in Proposition 7.2 are explicitly given by

$$a_k = \sin[(2k + 1)\theta], \quad b_k = \frac{1}{\sqrt{N - 1}} \cos[(2k + 1)\theta], \quad (7.20)$$

for  $k = 0, 1, 2, \dots$ , where

$$\sin \theta = \sqrt{\frac{1}{N}}, \quad \cos \theta = \sqrt{1 - \frac{1}{N}}. \quad (7.21)$$

# Grover's search algorithm

*Proof.* Let  $c_k = \sqrt{N-1}b_k$ . The recursion relations (7.18) and (7.19) are written in a matrix form,

$$\begin{pmatrix} a_k \\ c_k \end{pmatrix} = M \begin{pmatrix} a_{k-1} \\ c_{k-1} \end{pmatrix}, \quad M = \begin{pmatrix} (N-2)/N & 2\sqrt{N-1}/N \\ -2\sqrt{N-1}/N & (N-2)/N \end{pmatrix} = \begin{pmatrix} \cos 2\theta & \sin 2\theta \\ -\sin 2\theta & \cos 2\theta \end{pmatrix}.$$

Note that  $M$  is a rotation matrix in  $\mathbb{R}^2$ , and its  $k$ th power is another rotation matrix corresponding to a rotation angle  $2k\theta$ . Thus the above recursion relation is easily solved to yield

$$\begin{pmatrix} a_k \\ c_k \end{pmatrix} = M^k \begin{pmatrix} a_0 \\ c_0 \end{pmatrix} = \begin{pmatrix} \cos 2k\theta & \sin 2k\theta \\ -\sin 2k\theta & \cos 2k\theta \end{pmatrix} \begin{pmatrix} \sin \theta \\ \cos \theta \end{pmatrix} = \begin{pmatrix} \sin[(2k+1)\theta] \\ \cos[(2k+1)\theta] \end{pmatrix}.$$

Replacing  $c_k$  by  $b_k$  proves the proposition. ■

# Grover's search algorithm

We have proved that the application of  $\bar{G}_f$   $k$  times on  $|\varphi_0\rangle$  results in the state

$$G_f^k |\varphi_0\rangle = \sin[(2k+1)\theta] |z\rangle + \frac{1}{\sqrt{N-1}} \cos[(2k+1)\theta] \sum_{x \neq z} |x\rangle. \quad (7.22)$$

Measurement of the state  $U_f^k |\varphi_0\rangle$  yields  $|z\rangle$  with the probability

$$P_{z,k} = \sin^2[(2k+1)\theta]. \quad (7.23)$$

**STEP 4** Our final task is to find the  $k$  that maximizes  $P_{z,k}$ . A rough estimate for the maximizing  $k$  is obtained by putting

$$(2k+1)\theta = \frac{\pi}{2} \rightarrow k = \frac{1}{2} \left( \frac{\pi}{2\theta} - 1 \right). \quad (7.24)$$

# Grover's search algorithm

**PROPOSITION 7.4** Let  $N \gg 1$  and let

$$m = \left\lfloor \frac{\pi}{4\theta} \right\rfloor, \quad (7.25)$$

where  $\lfloor x \rfloor$  stands for the floor of  $x$ . The file we are searching for will be obtained in  $G_f^m |\varphi_0\rangle$  with the probability

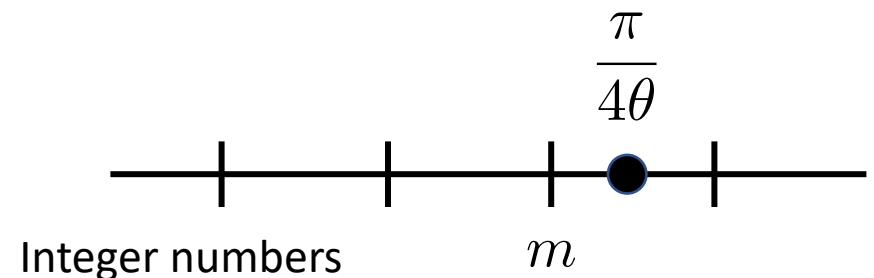
$$P_{z,m} \geq 1 - \frac{1}{N} \quad (7.26)$$

and

$$m = O(\sqrt{N}). \quad (7.27)$$



This is the number of times we repeat the algorithm, which grows with the square root of  $N$



# Grover's search algorithm

*Proof.* Equation (7.25) leads to the inequality  $\pi/4\theta - 1 < m \leq \pi/4\theta$ . Let us define  $\tilde{m}$  by

$$(2\tilde{m} + 1)\theta = \frac{\pi}{2} \rightarrow \tilde{m} = \frac{\pi}{4\theta} - \frac{1}{2}.$$

Observe that  $m$  and  $\tilde{m}$  satisfy

$$|m - \tilde{m}| \leq \frac{1}{2}, \quad (7.28)$$

from which it follows that

$$|(2m + 1)\theta - (2\tilde{m} + 1)\theta| = \left| (2m + 1)\theta - \frac{\pi}{2} \right| \leq \theta. \quad (7.29)$$

# Grover's search algorithm

Considering that  $\theta \sim 1/\sqrt{N}$  is a small number when  $N \gg 1$  and  $\sin x$  is monotonically increasing in the neighborhood of  $x = 0$ , we obtain

$$0 < \sin |(2m + 1)\theta - \pi/2| < \sin \theta$$

or

$$\cos^2[(2m + 1)\theta] \leq \sin^2 \theta = \frac{1}{N}. \quad (7.30)$$

$$= \cos[(2m + 1)\theta]$$

Thus it has been shown that

$$P_{m,z} = \sin^2[(2m + 1)\theta] = 1 - \cos^2[(2m + 1)\theta] \geq 1 - \frac{1}{N}. \quad (7.31)$$

It also follows from  $\theta > \sin \theta = 1/\sqrt{N}$  that

$$m = \left\lfloor \frac{\pi}{4\theta} \right\rfloor \leq \frac{\pi}{4\theta} \leq \frac{\pi}{4} \sqrt{N}. \quad (7.32)$$

■



# Grover's search algorithm

It is important to note that this quantum algorithm takes only  $O(\sqrt{N})$  steps and this is much faster than the classical counterpart which requires  $O(N)$  steps.

Next we will show how how to implement the gates