

Master Degree Programme in Physics – UNITS Physics of the Earth and of the Environment

LINEAR SYSTEMS

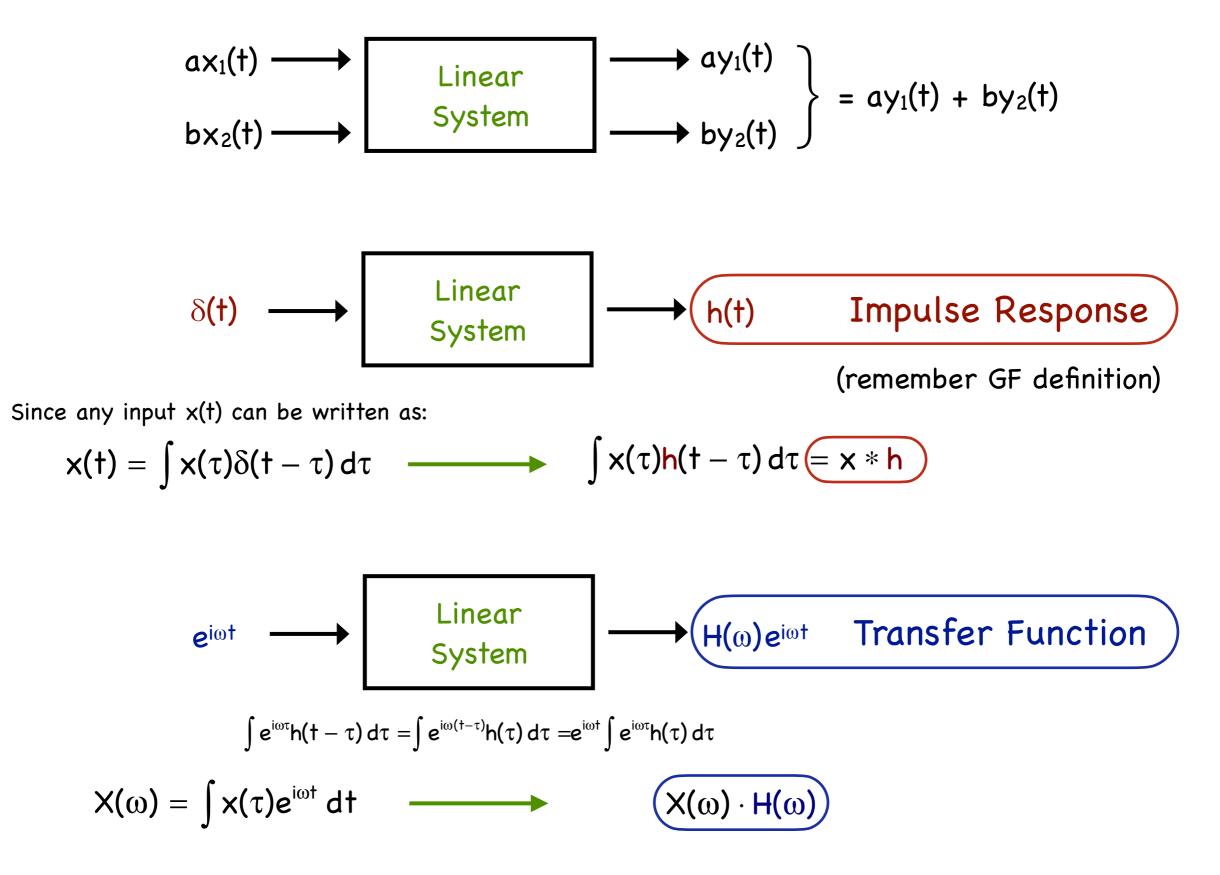
FABIO ROMANELLI

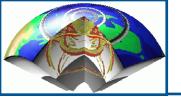
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Linear Systems

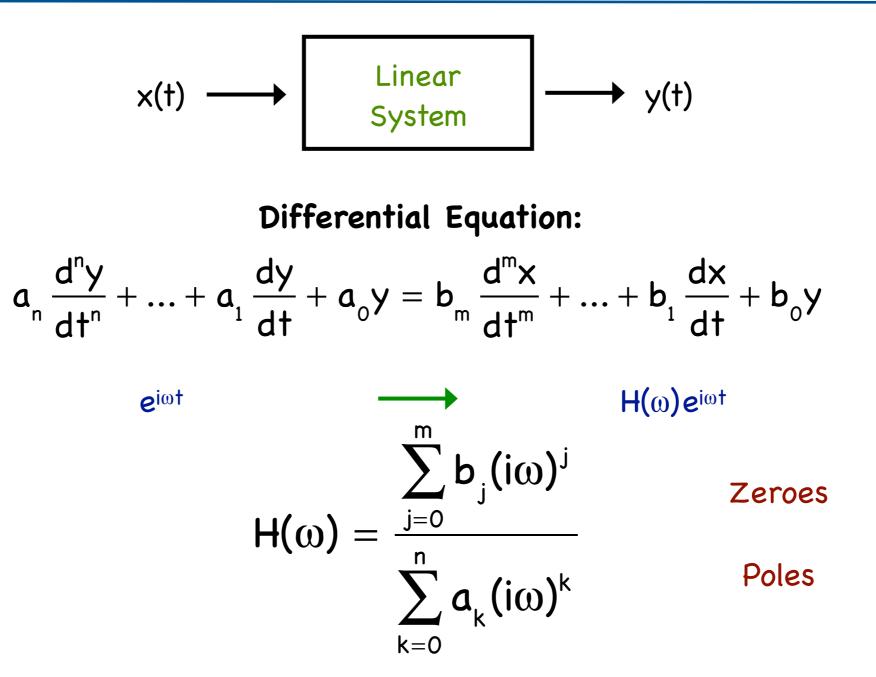






Poles and Zeros





Transfer function of the seismometer:

$$H(\omega) = \frac{Y(\omega)}{U(\omega)} = \frac{-(i\omega)^2}{(i\omega)^2 + 2\gamma(i\omega) + (\omega_0)^2}$$







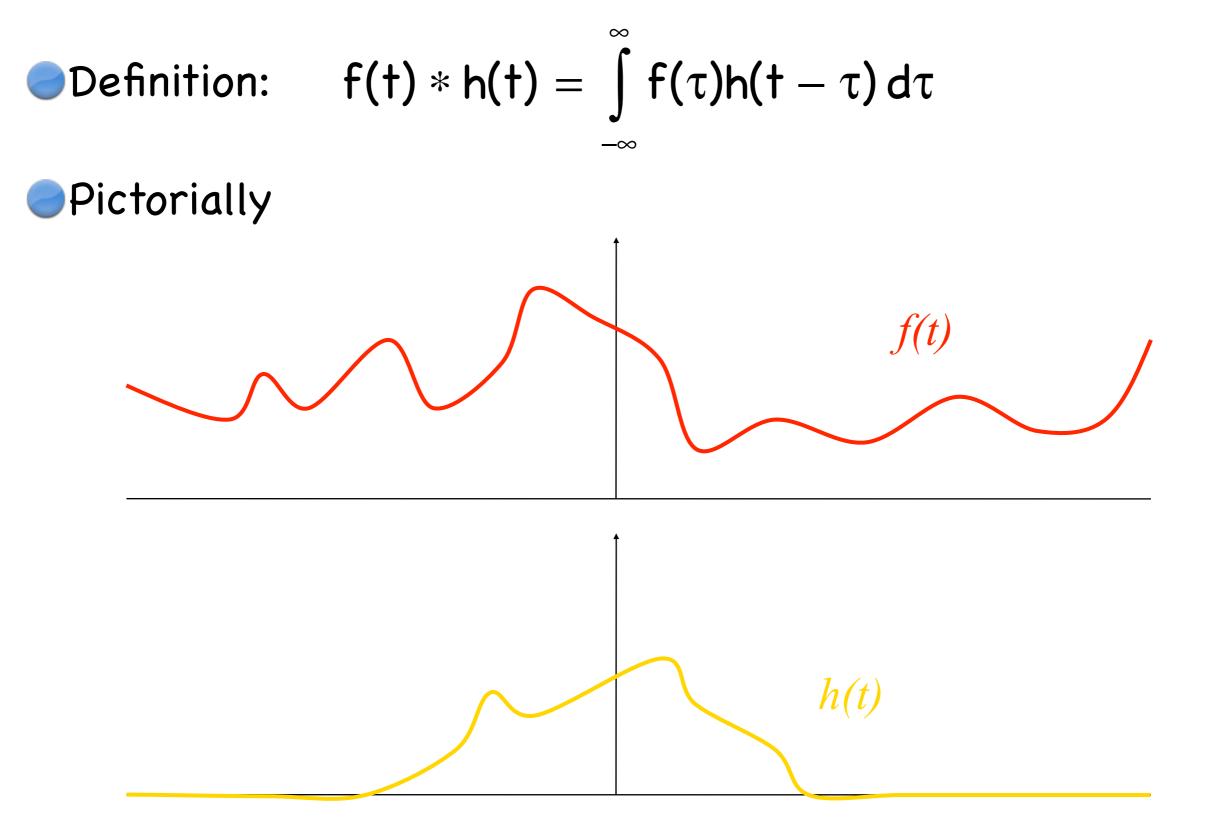


 ∞ $f(t) * h(t) = \int f(\tau)h(t-\tau) d\tau$ $-\infty$



Convolution

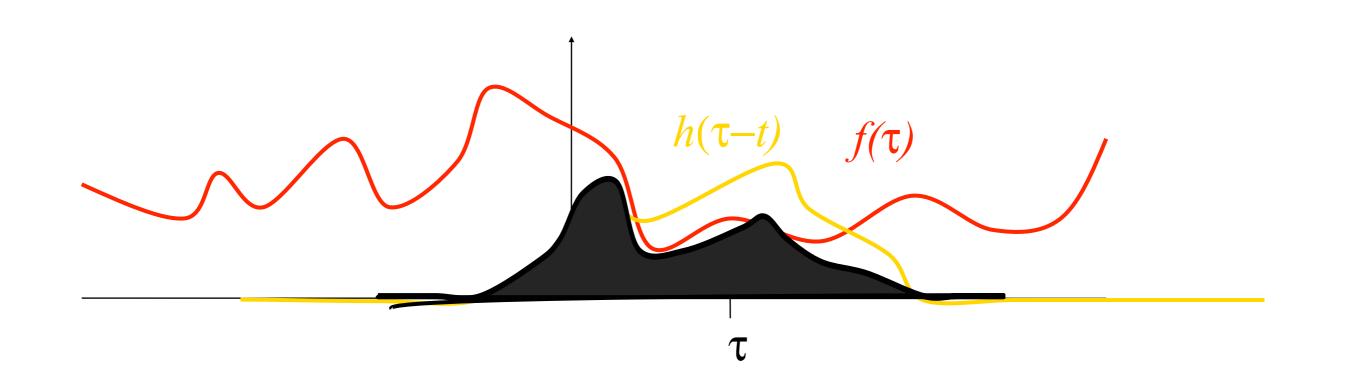


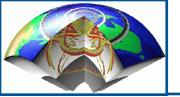










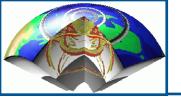






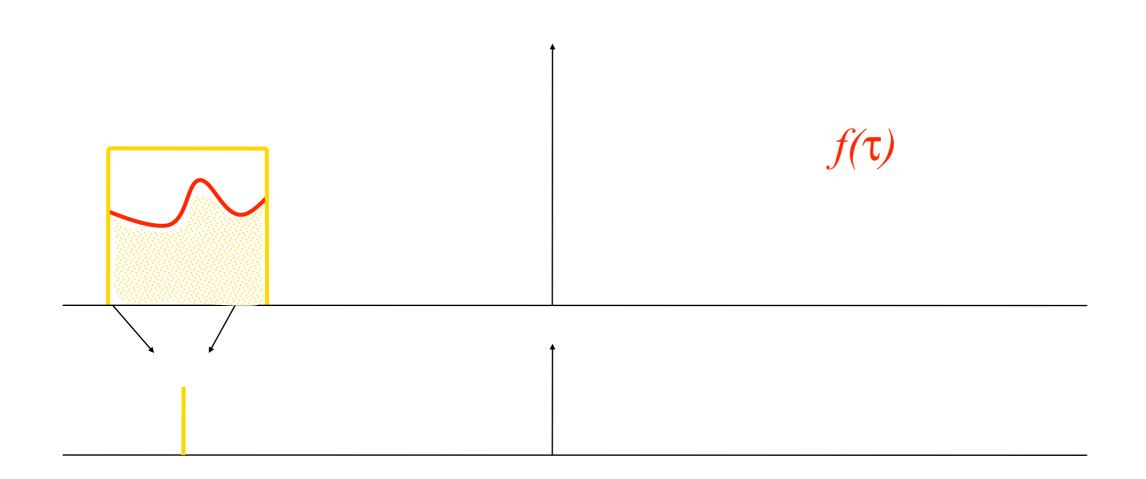
Consider the boxcar function (box filter):

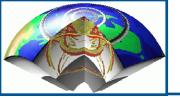
$$h(t) = \begin{cases} 0 & t < -\frac{1}{2} \\ 1 & -\frac{1}{2} \le t \le \frac{1}{2} \\ 0 & t > \frac{1}{2} \end{cases}$$





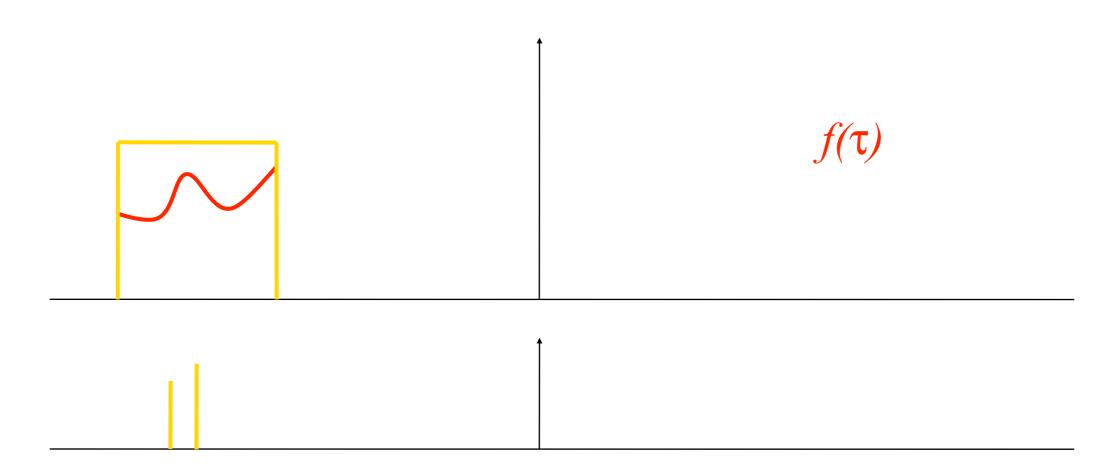


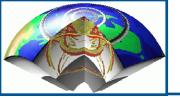








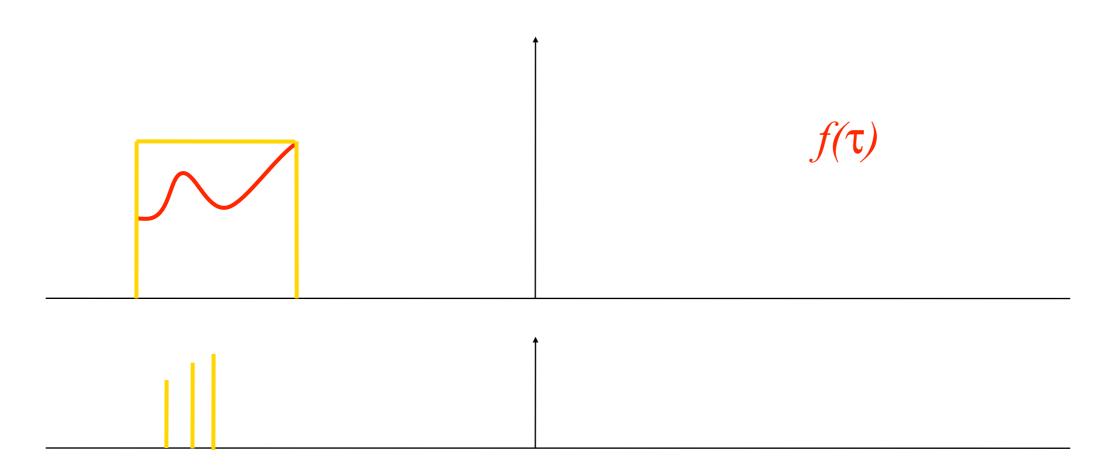


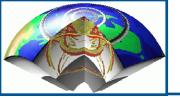






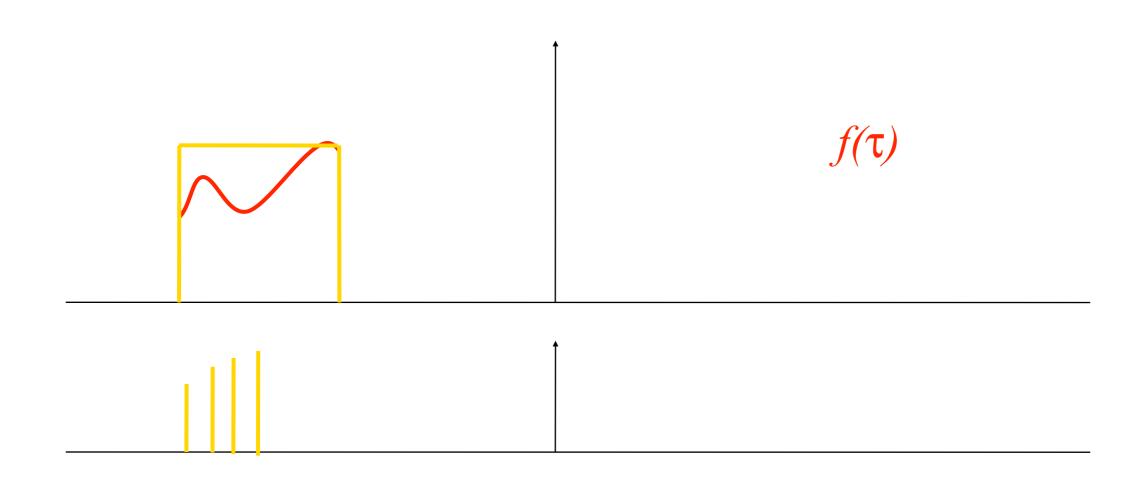
\bigcirc This function windows our function f(t).

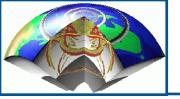






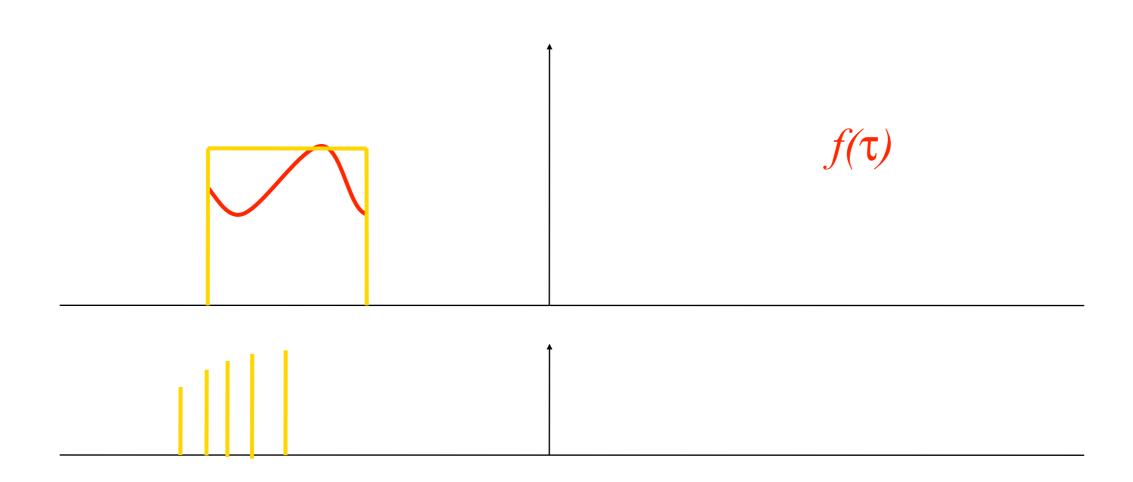


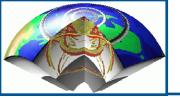






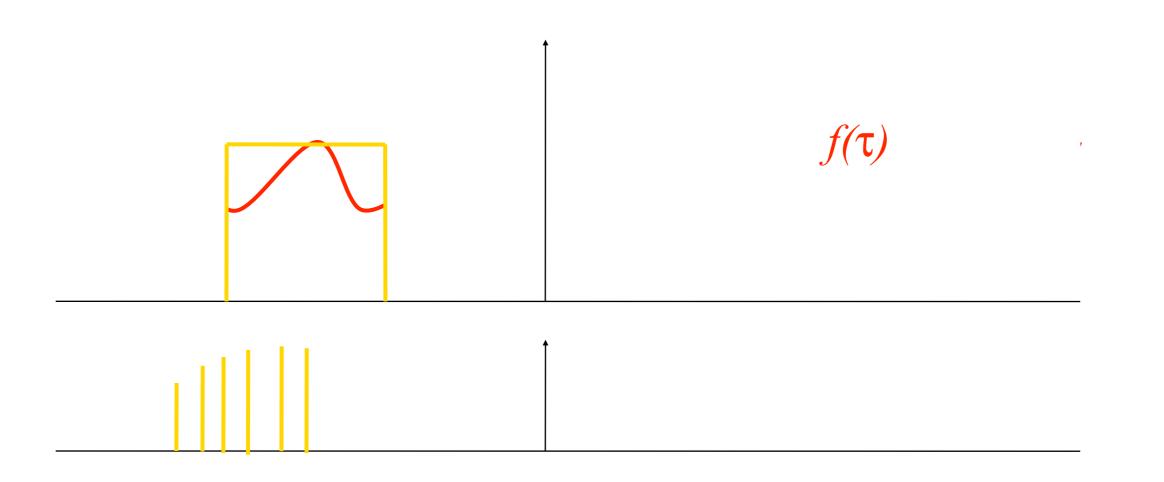


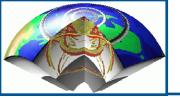






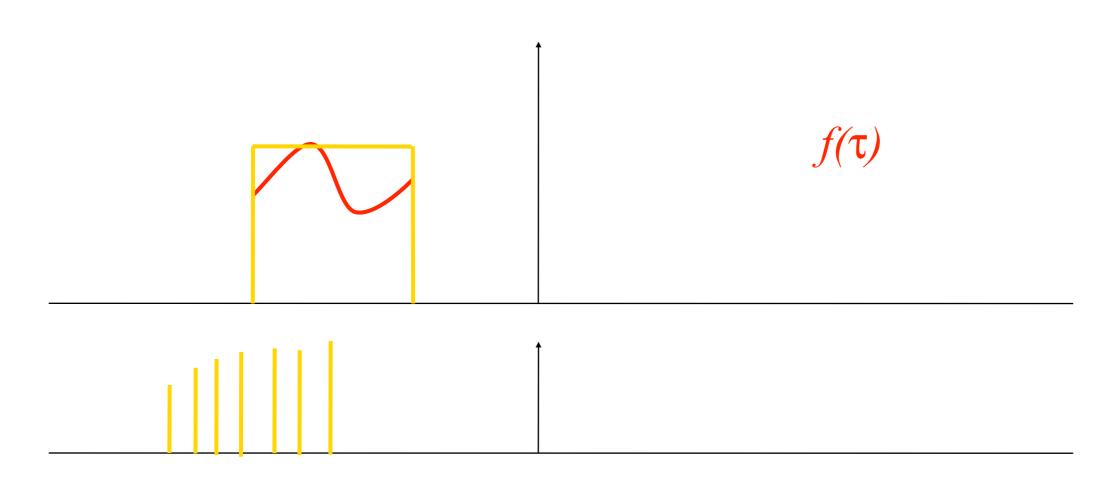


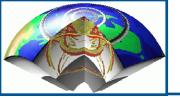






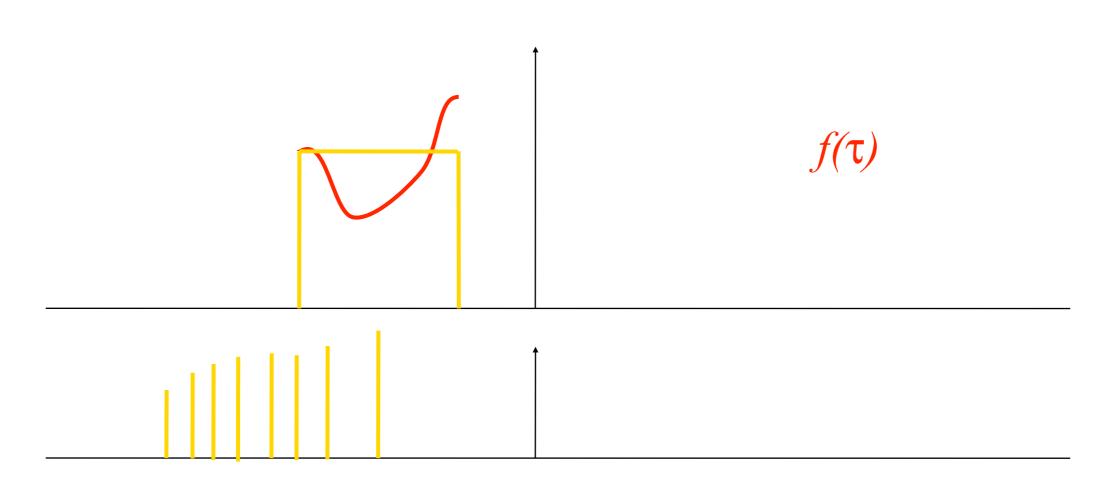


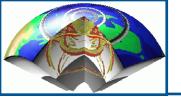






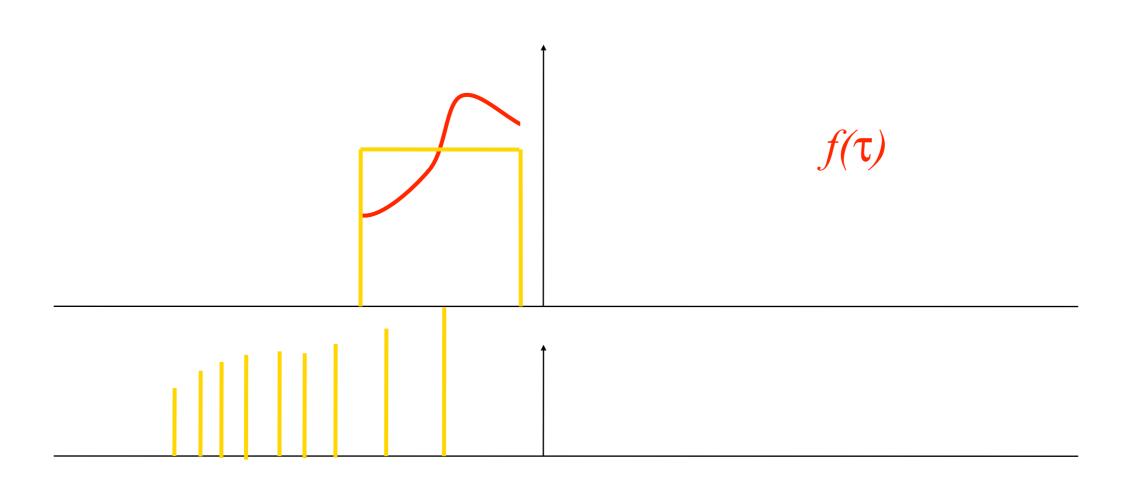


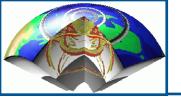






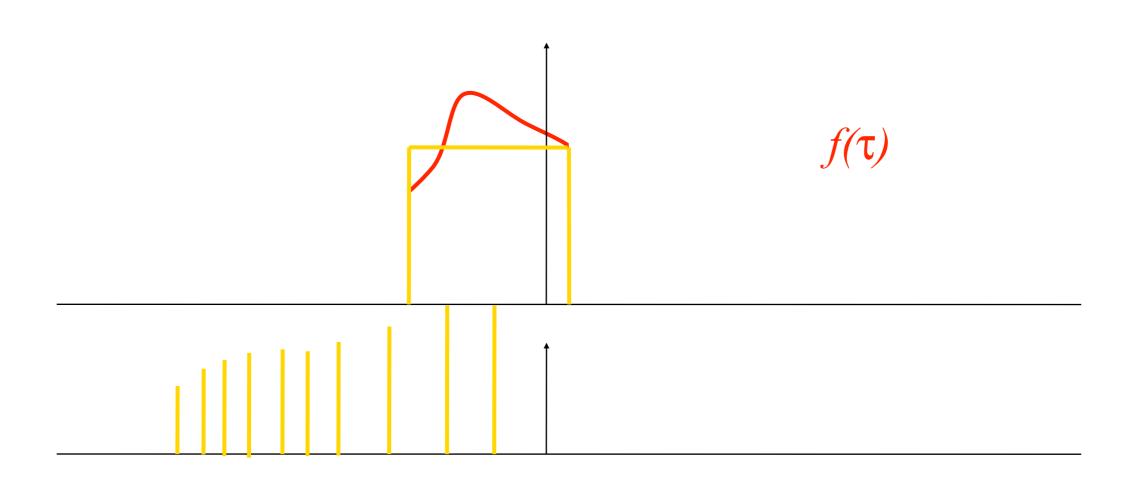


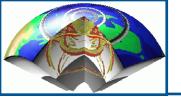






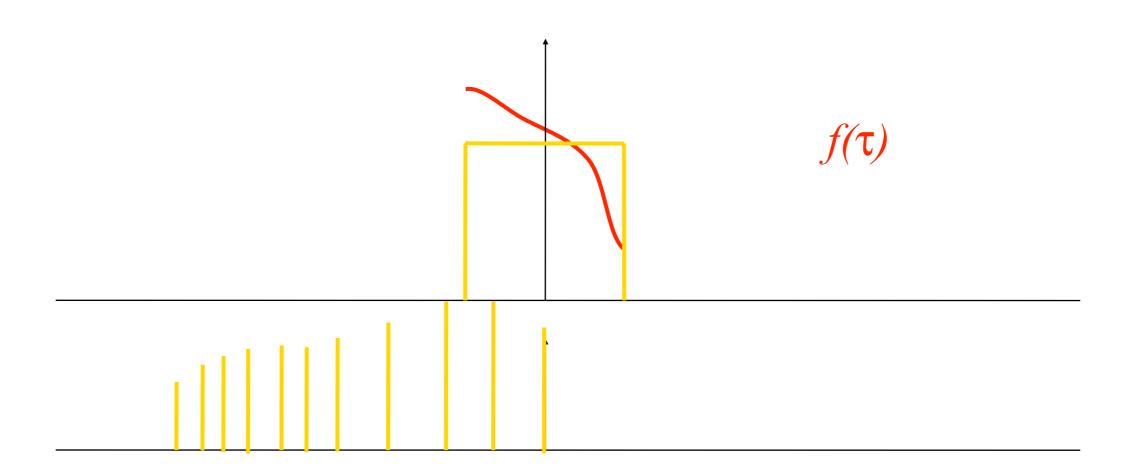


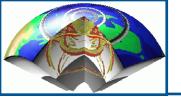






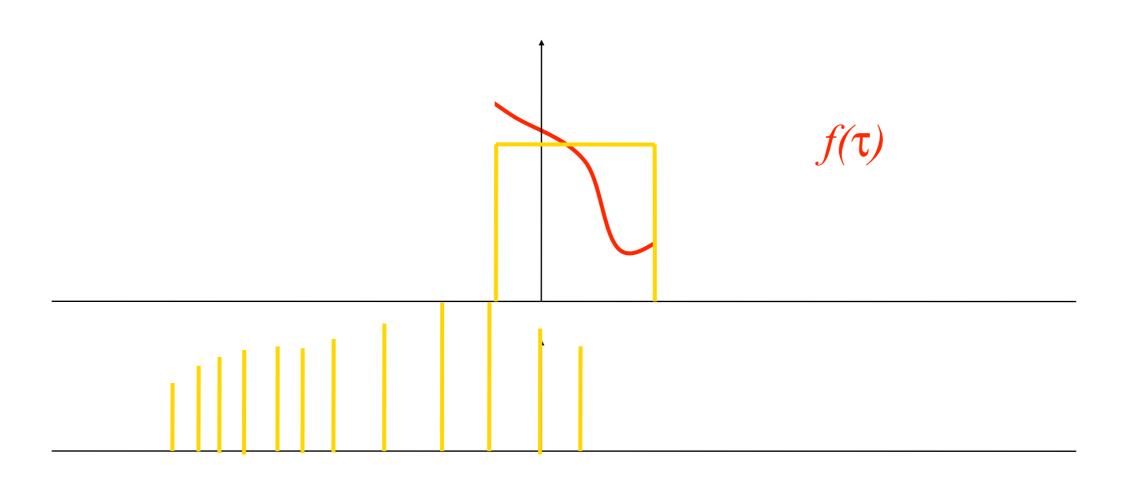


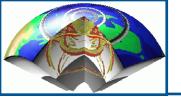






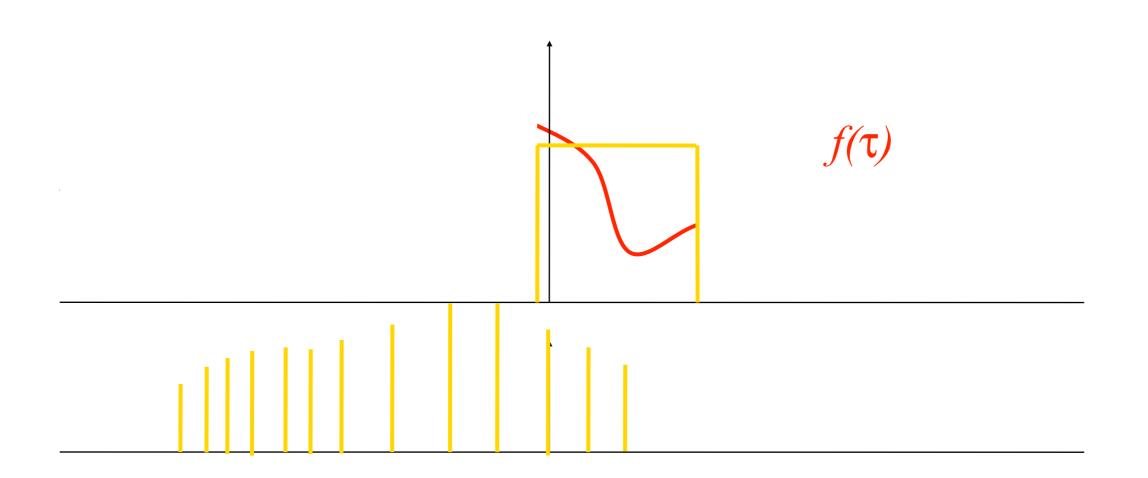


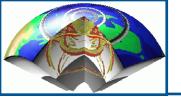






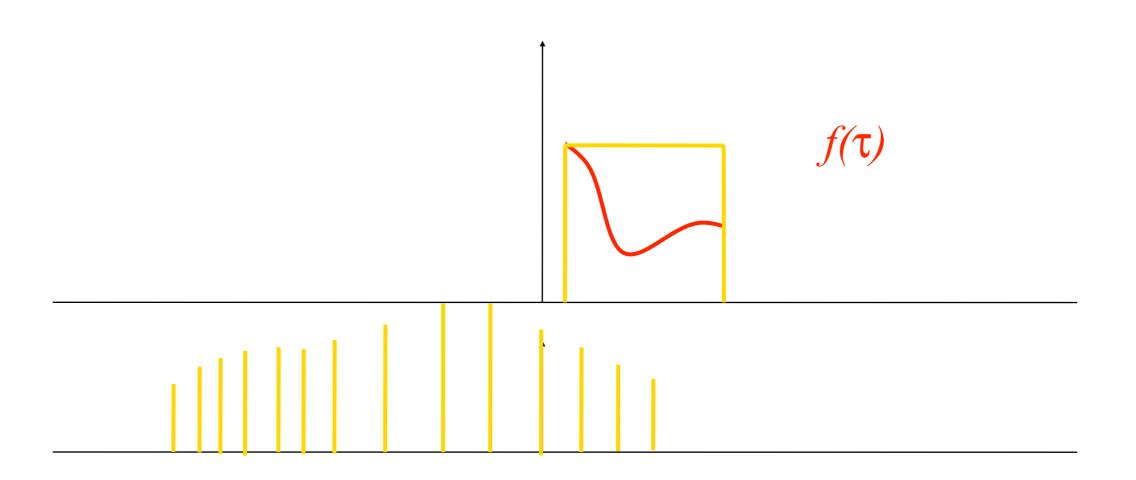


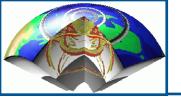






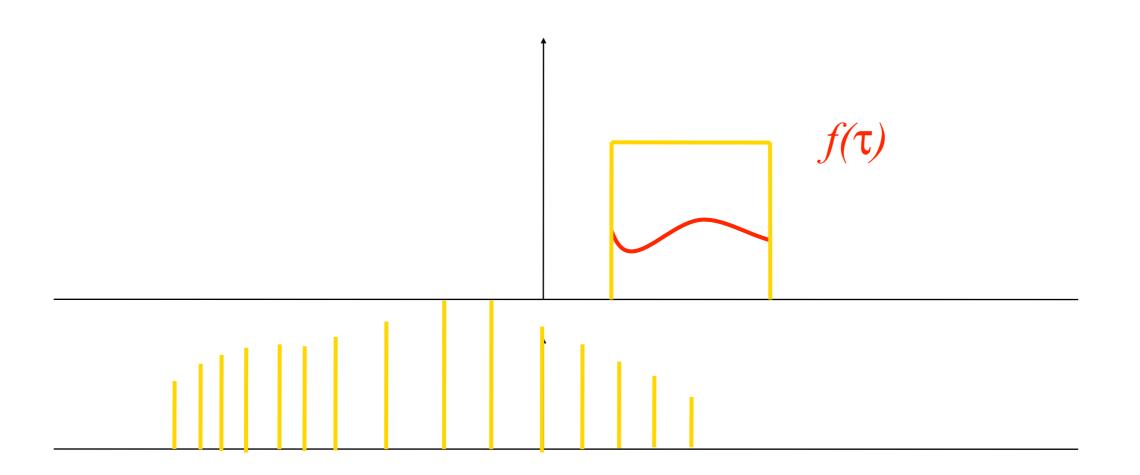


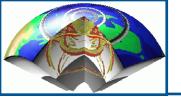






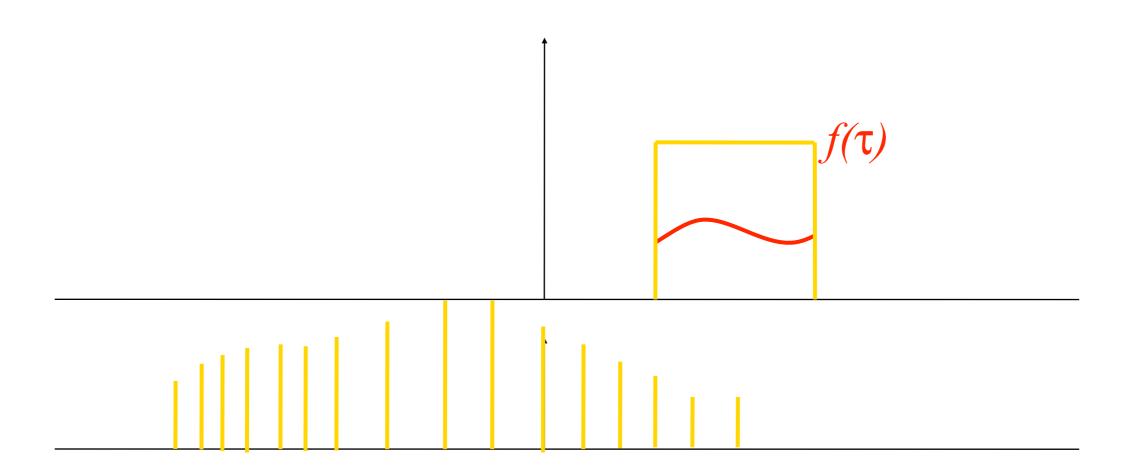


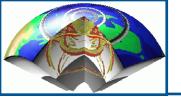






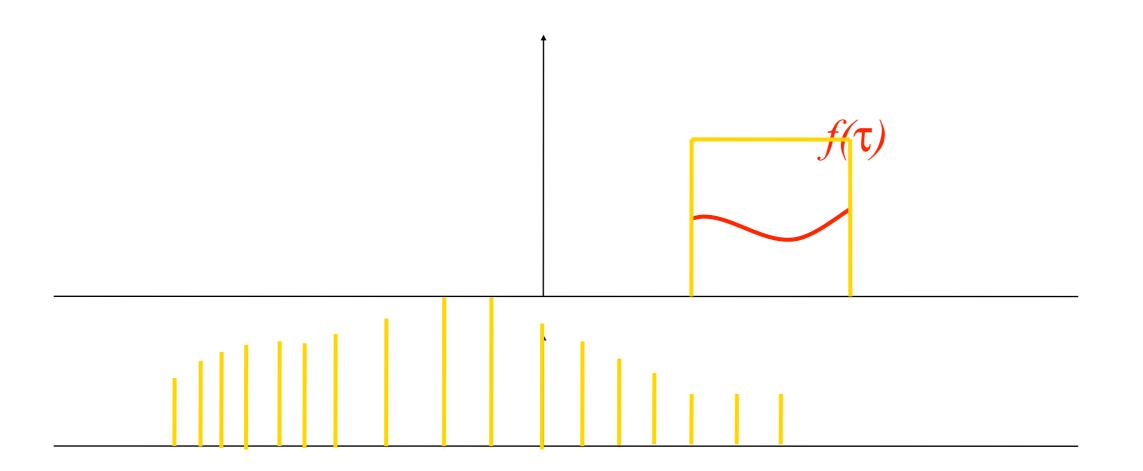


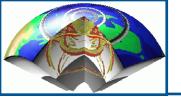






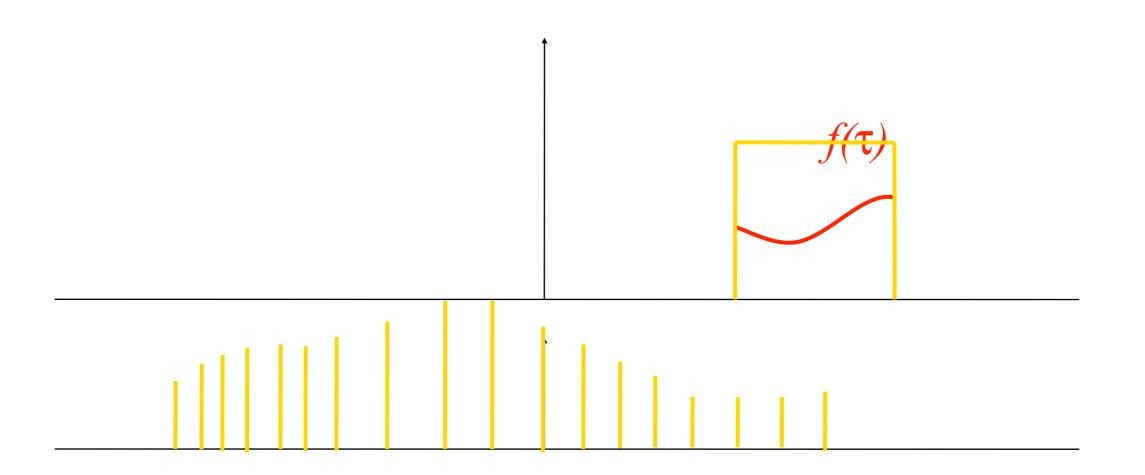


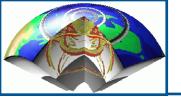






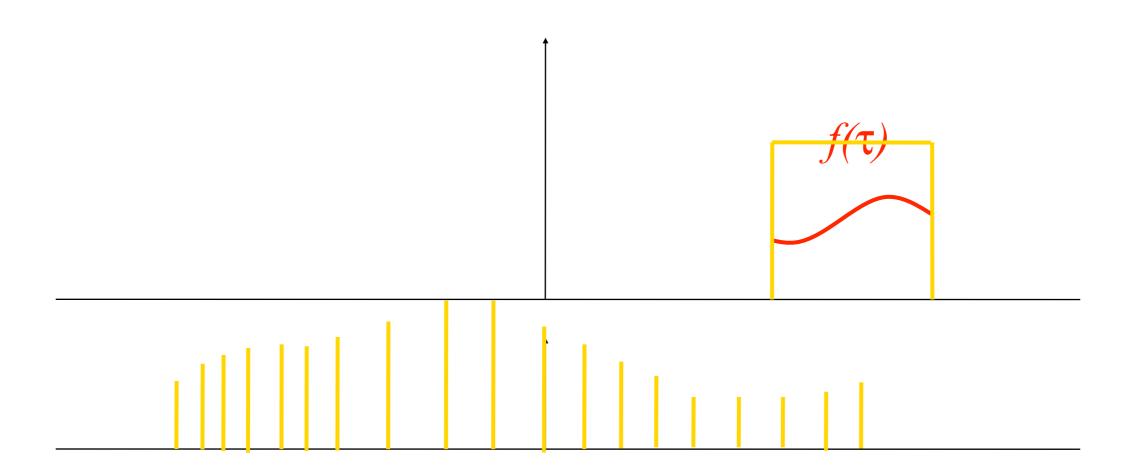


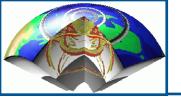






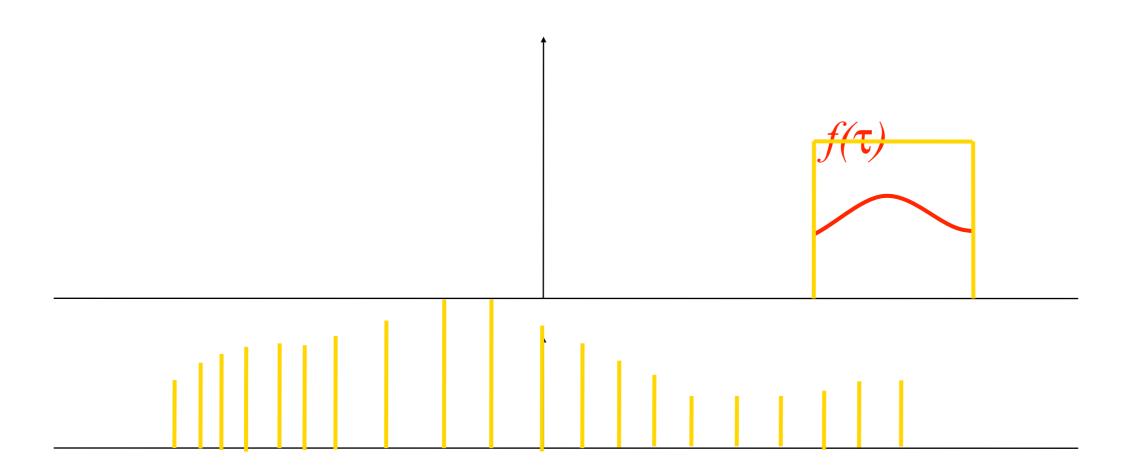


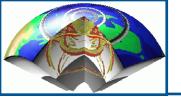






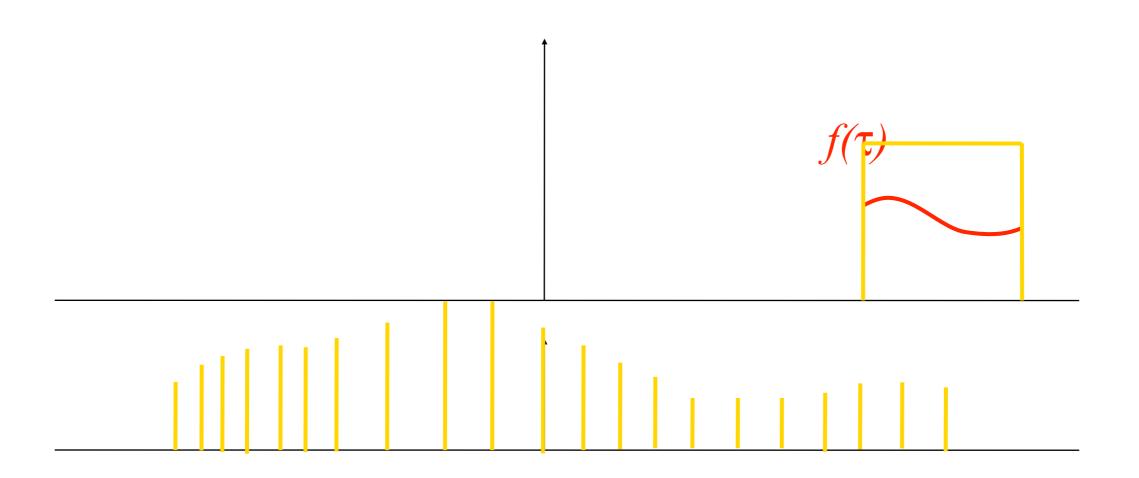


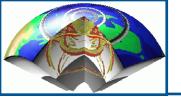






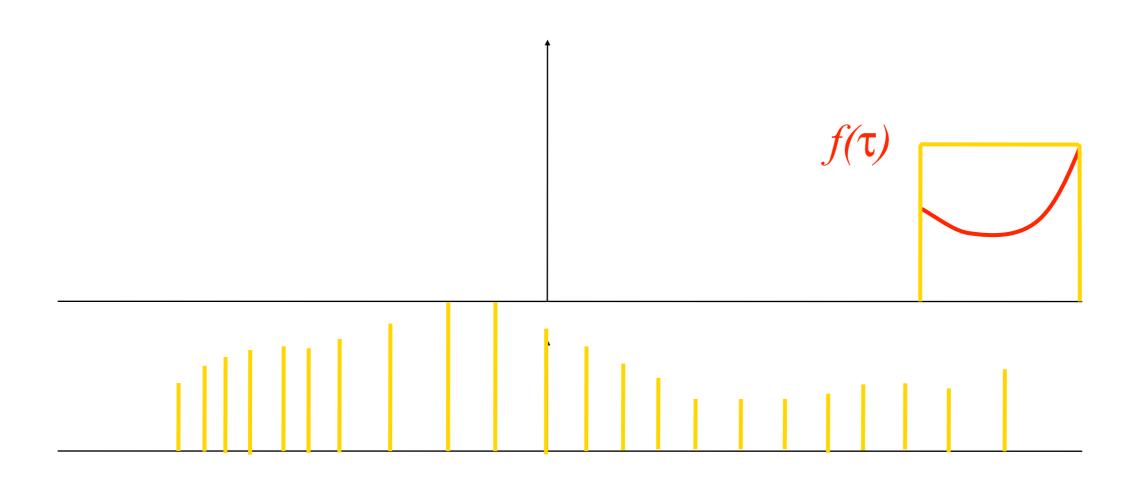


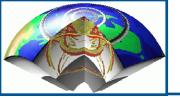








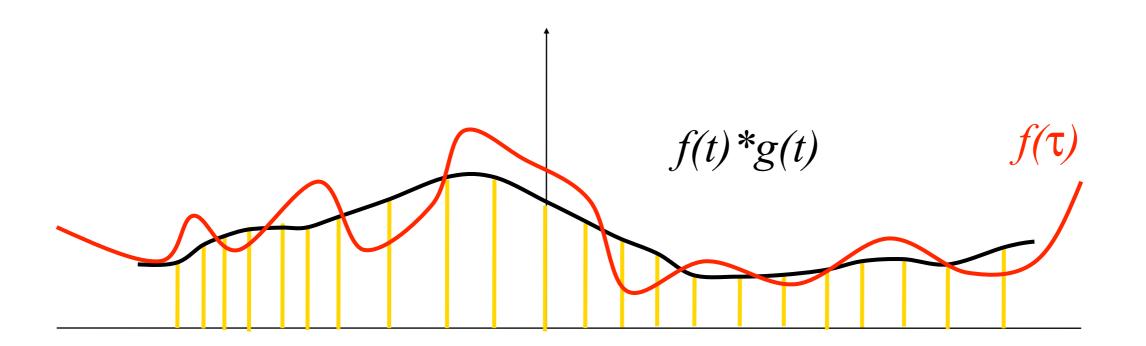


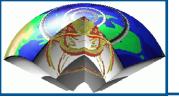


Convolution



This particular convolution smooths out some of the high frequencies in f(t).



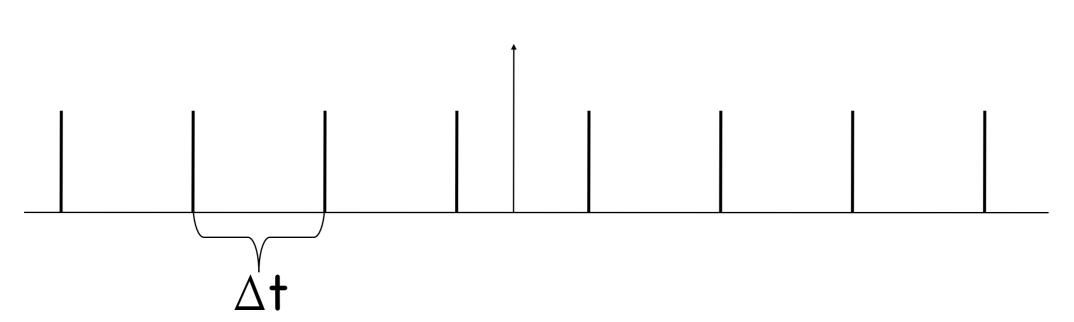


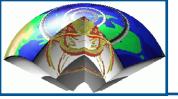


A Sampling Function or Impulse Train is defined by:

$$S_{T}(t) = \sum_{k=-\infty}^{\infty} \delta(t - k\Delta t)$$

where Δt is the sample spacing.



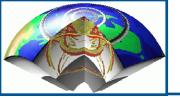




The Fourier Transform of the Sampling Function is itself a sampling function.

The sample spacing is the inverse.

 $S_{\Delta^{\dagger}}(\dagger) \Leftrightarrow S_{\frac{1}{\Delta^{\dagger}}}(\omega)$

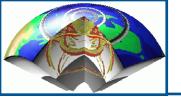




The convolution theorem states that convolution in the spatial domain is equivalent to multiplication in the frequency domain, and viceversa.

$f(t) * g(t) \Leftrightarrow F(\omega) \cdot G(\omega)$

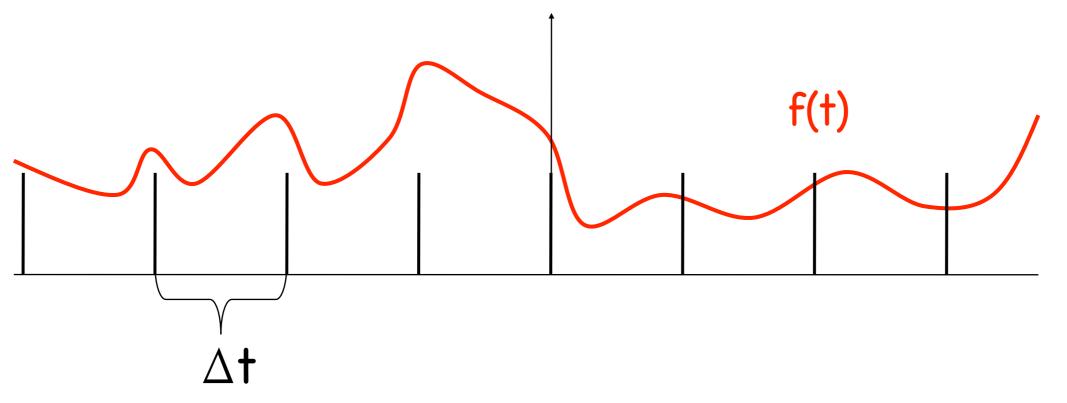
 $\mathsf{f}(\mathsf{t}) \cdot \mathsf{g}(\mathsf{t}) \Leftrightarrow \mathsf{F}(\omega) \ast \mathsf{G}(\omega)$

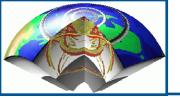




This powerful theorem can illustrate the problems with our point sampling and provide guidance on avoiding aliasing.

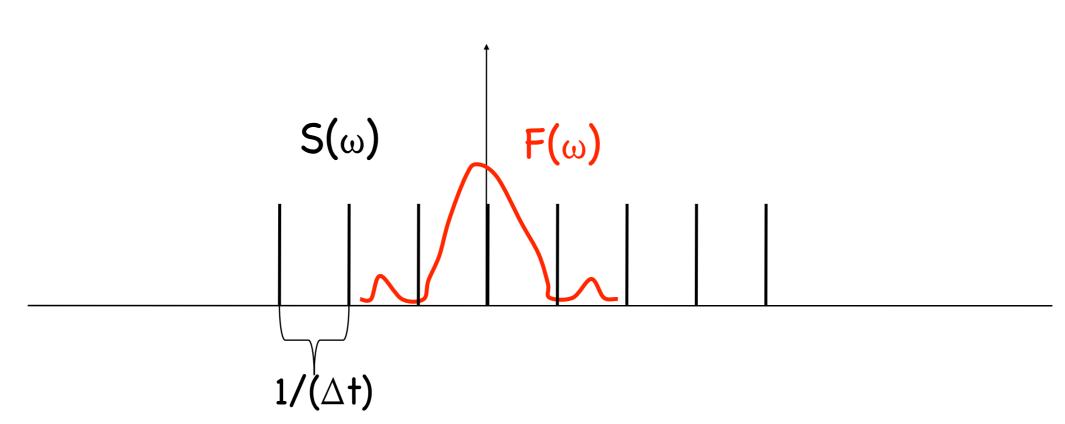
Consider: $f(t) \cdot S_{\Delta t}(t)$

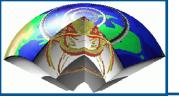






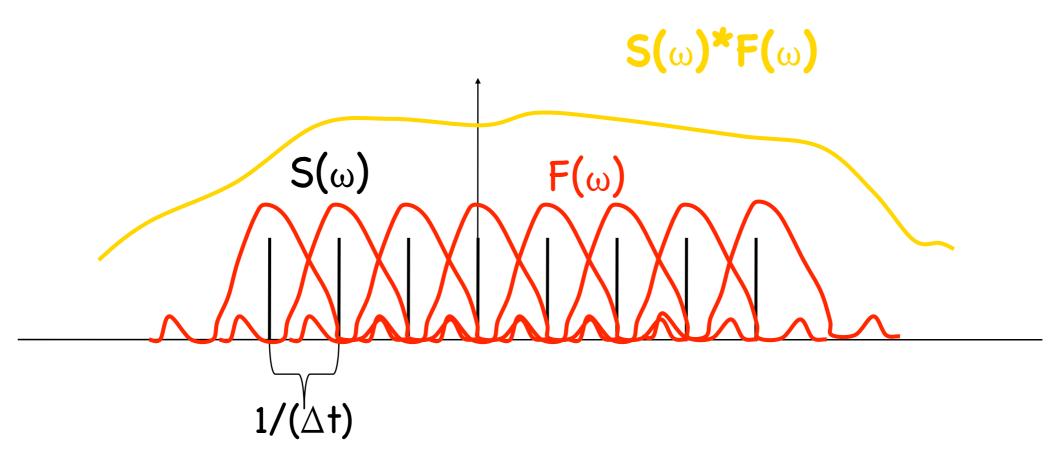
What does this look like in the Fourier domain?

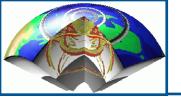






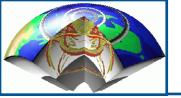
In Fourier domain we would convolve





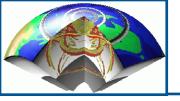


- What this says, is that any frequencies greater than a certain amount will appear intermixed with other frequencies.
- In particular, the higher frequencies for the copy at $1/\Delta t$ intermix with the low frequencies centered at the origin.





- Note, that the sampling process introduces frequencies out to infinity.
- We have also lost the function f(t), and now have only the discrete samples.
- This brings us to our next powerful theory.





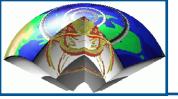
The Shannon Sampling Theorem:

A band-limited signal f(t), with a cutoff frequency of λ , that is sampled with a sampling spacing of Δt may be perfectly reconstructed from the discrete values f[n Δt] by convolution with the sinc(t) function, provided the Nyquist limit: $\lambda < 1/(2\Delta t)$

Why is this?

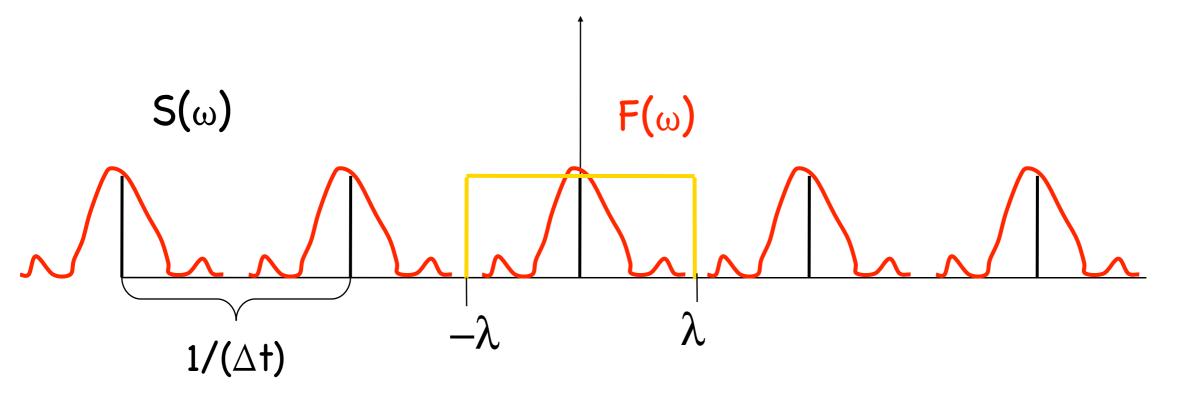
The Nyquist limit will ensure that the copies of $F(\omega)$ do not overlap in the frequency domain.

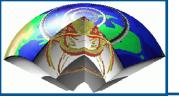
We can completely reconstruct or determine f(t) from $F(\omega)$ using the Inverse Fourier Transform.





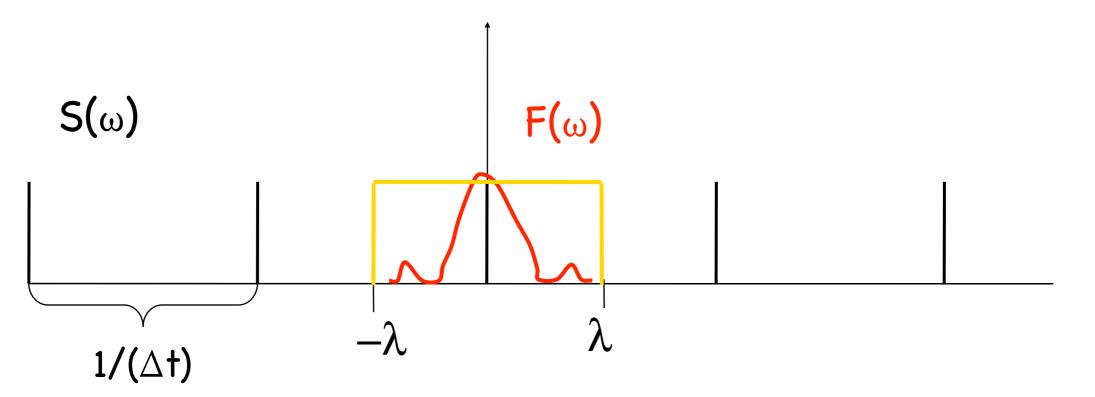
- In order to do this, we need to remove all of the shifted copies of $F(\omega)$ first.
- This is done by simply multiplying $F(\omega)$ by a box function of width 2λ .

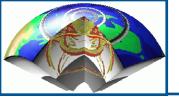






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So, given $f[n\Delta t]$ and an assumption that f(t) does not have frequencies greater than $1/(2\Delta t)$, we can write the formula:

 $f[nT] = f(t) \cdot S_{\Delta t}(t) \Leftrightarrow F(\omega) * S_{\Delta t}(\omega)$

 $\mathsf{F}(\omega) = (\mathsf{F}(\omega) * \mathsf{S}_{\Delta \dagger}(\omega)) \cdot \mathsf{Box}_{1/(2\Delta \dagger)}(\omega)$

therefore,

 $f(t) = f[n\Delta t] * sinc(t)$

http://www.thefouriertransform.com/pairs/box.php

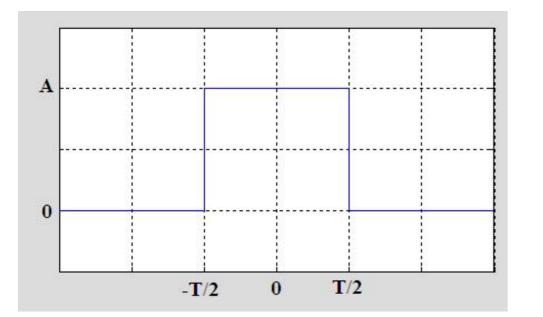
http://195.134.76.37/applets/AppletNyquist/Appl_Nyquist2.html

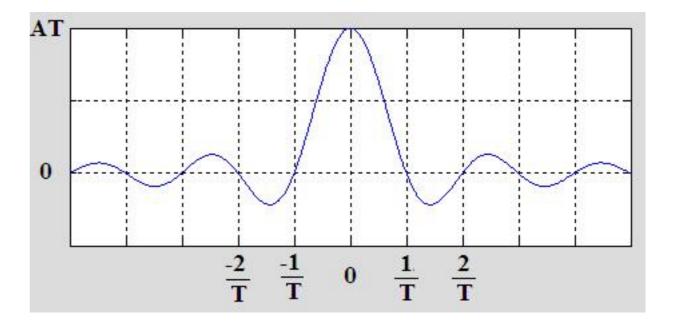






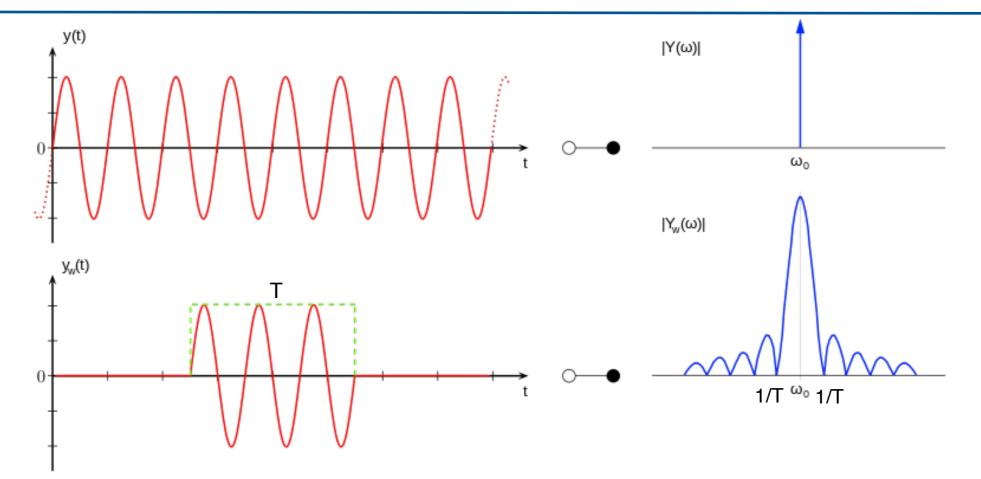
$$\int_{-\infty}^{+\infty} B_{T}(t) e^{-i\omega t} dt = \int_{-T/2}^{+T/2} e^{-i\omega t} dt = \frac{\sin(\pi fT)}{\pi fT}$$
$$\approx \operatorname{sinc}(\pi fT) = \operatorname{sinc}(\omega T / 2)$$





Spectral leakage





 $\Delta f \Delta t \geq \frac{1}{2\sqrt{-N}}$

Resolving power in frequency domain is related to maximum duration in time domain:

$$\Delta f \geq \frac{1}{2\pi T} \left(= \frac{1}{2\pi N \Delta t} \right)$$

Resolving power in time domain decides maximum resolvable frequency:

$$\Delta f \ge \frac{1}{2\Delta t}$$

https://www.youtube.com/watch?v=MBnnXbOM5S4