

(c)

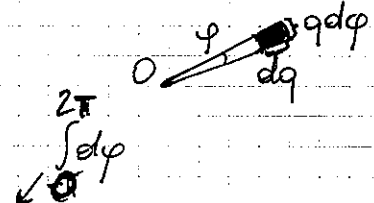
$$g_m(\vec{r}) = \int \frac{1}{4\pi^3 \sqrt{\partial_x \partial_y \partial_z}} d\vec{q} \delta[\vec{r} - \vec{r}_0 - q_x^2 + q_y^2 + q_z^2]$$

Uso coord. cilindriche nel piano (x,y): {q_x, q_y} → {q, φ}

$$d\vec{q} = dq_x dq_y dq_z$$

$$= q dq d\varphi dq_z$$

$$\text{con } q^2 = q_x^2 + q_y^2$$



$$\rightarrow g_m(\vec{r}) = \frac{1}{4\pi^3 \sqrt{\partial_x \partial_y \partial_z}} \cdot 2\pi \int_0^{2\pi} \int q dq dq_z \delta[\vec{r} - \vec{r}_0 - q^2 + q_z^2]$$

$$= \frac{1}{2\pi^2 \sqrt{\partial_x \partial_y \partial_z}} \int_0^{\infty} q dq \left\{ \int_{-\infty}^{+\infty} dq_z \delta[\vec{r} - \vec{r}_0 - q^2 + q_z^2] \right\}$$

Ver. la (*) [formula] di Leibniz:

$$\frac{d}{dq_z} [\vec{r} - \vec{r}_0 - q^2 + q_z^2] = 2q_z$$

$$\vec{r} - \vec{r}_0 - q^2 + q_z^2 = 0 \text{ per } \begin{cases} q_z = \pm \sqrt{q^2 + \vec{r}_0 - \vec{r}} & \text{se } \vec{r} - \vec{r}_0 < q^2 \\ \text{NO ZERI} & \text{per } \vec{r} - \vec{r}_0 > q^2 \end{cases}$$

Allora

viene da $\sum_{z=1}^{\infty} z$, che sono $\pm \dots$

$$g_m(\vec{r}) = \frac{1}{2\pi^2 \sqrt{\partial_x \partial_y \partial_z}} \int_0^{\infty} q dq \begin{cases} 2 \cdot \frac{\Theta(q^2 + \vec{r}_0 - \vec{r})}{2\sqrt{q^2 + \vec{r}_0 - \vec{r}}} & \text{per } q^2 + \vec{r}_0 - \vec{r} > 0 \\ \phi & \text{per } q^2 + \vec{r}_0 - \vec{r} < 0 \end{cases}$$

f. scalino $\Theta(x) = \begin{cases} 1 & x \geq 0 \\ 0 & x < 0 \end{cases}$

K piccolo \leftrightarrow q piccolo, allora considero \int_0^a , $a = q_{\text{max}}$.

Pongo:

$$x \equiv q^2 + \vec{r}_0 - \vec{r}$$

$$2q dq = dx \rightarrow \begin{cases} q=0 \rightarrow x = \vec{r}_0 - \vec{r} \\ q=a \rightarrow x = a^2 + \vec{r}_0 - \vec{r} \end{cases}$$

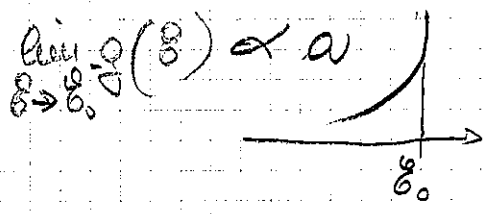
$$g_m(\vec{r}) = \frac{1}{4\pi^2 \sqrt{\partial_x \partial_y \partial_z}} \int_{\vec{r}_0 - \vec{r}}^{a^2 + \vec{r}_0 - \vec{r}} \frac{\Theta(x) dx}{\sqrt{x}}$$

Discutiamo il valore dell' \int a seconda dei limiti di integ.

$\sigma_0 - \sigma > 0$

i limiti dell' \int sono $> 0 \Rightarrow \Theta(x) \equiv 1$ in quel \Rightarrow range

$g_n(\sigma) \propto \int_{\sigma_0 - \sigma}^{a^2 + \sigma_0 - \sigma} \frac{dx}{2\sqrt{x}} = \sqrt{x} \Big|_{\sigma_0 - \sigma}^{a^2 + \sigma_0 - \sigma}$
 $\propto \sqrt{a^2 + \sigma_0 - \sigma} - \sqrt{\sigma_0 - \sigma}$



$\frac{dg_n(\sigma)}{d\sigma} \propto \frac{-1}{\sqrt{a^2 + \sigma_0 - \sigma}} + \frac{1}{\sqrt{\sigma_0 - \sigma}}$
 $\sigma \approx \sigma_0 \Rightarrow \frac{1}{\sqrt{\sigma_0 - \sigma}}$

Passabile rispetto all' $\rightarrow \infty$

$\Rightarrow \frac{dg_n(\sigma)}{d\sigma} \sim (\sigma_0 - \sigma)^{-1/2}$ per $\sigma < \sigma_0$

$\sigma_0 - \sigma < 0$

il limite inf. dell' \int $\bar{\sigma} < 0$, allora l' \int si riduce a:

$g_n(\sigma) \propto \int_0^{a^2 + \sigma_0 - \sigma} \frac{dx}{2\sqrt{x}} = \sqrt{x} \Big|_0^{a^2 + \sigma_0 - \sigma}$
 $\propto \sqrt{a^2 + \sigma_0 - \sigma}$

$\frac{dg_n(\sigma)}{d\sigma} \propto \frac{-1}{\sqrt{a^2 + \sigma_0 - \sigma}} \approx \frac{1}{a}$ costante per $\sigma \approx \sigma_0$

