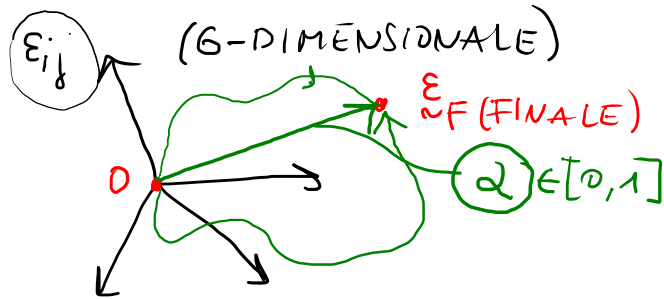


COMPLEMENTI SUGLI ASPETTI ENERGETICI DELL'ELASTICITA'

4/11/22



IPOTESI: PROBL. CONSERVATIVO; $\exists$ ENERGIA (FUNZIONE DI STATO) φ CHE PERMETTE DI AFFERMARE CHE ESSA NON DIPENDA DAL PERCORSO SEGUITO PER ANDARE DA $\epsilon_{ij}=0$ A $\underline{\epsilon}_F$

$$\varphi = \int \underline{\sigma} \cdot d\underline{\epsilon} \Rightarrow \varphi = \frac{1}{2} \underline{\sigma}_F \cdot \underline{\epsilon}_F$$

(φ : POTENZIALE ELASTICO)

$$\varphi = \frac{1}{2} \sigma_{xx} \epsilon_{xx}$$

$$\int_{\text{PERCORSO}} \underline{\sigma} \cdot d\underline{\epsilon} = \varphi(\underline{\epsilon}_F) - \varphi(\underline{\epsilon}_0)$$

$$\int \underline{\sigma} \cdot d\underline{\epsilon} = \int_0^{\underline{\epsilon}_F} d\varphi = \varphi(\underline{\epsilon}_F) - \varphi(\underline{\epsilon}_0)$$

$d\varphi$: DIFFERENZ. ESATTO

$$= \varphi(\underline{\epsilon}_F) \Rightarrow \varphi(\underline{\epsilon}) = \int \underline{\sigma} \cdot d\underline{\epsilon}$$

Se conosco LA FUNZ. $\varphi(\underline{\epsilon})$ ALLORA

$$d\varphi(\underline{\epsilon}) = \underline{\sigma} \cdot d\underline{\epsilon} \quad (\text{DIFF. ESATTO})$$

$$\frac{\partial \varphi}{\partial \epsilon_{xx}} d\epsilon_{xx} + \frac{\partial \varphi}{\partial \epsilon_{yy}} d\epsilon_{yy} + \dots + \frac{\partial \varphi}{\partial \epsilon_{yz}} d\epsilon_{yz} + \dots$$

$\sigma_{xx} d\epsilon_{xx} + \sigma_{yy} d\epsilon_{yy} + \dots + \tau_{yz} d\epsilon_{yz} + \dots$

$$\underline{\varphi(\underline{\varepsilon})} \rightarrow \boxed{\underline{\sigma}_{ij} = \frac{\partial \varphi}{\partial \varepsilon_{ij}}} \quad (*)$$

EL. LIN.
ISOTROPO

ELAST. LIN.

$$\varphi(\underline{\varepsilon}) = \frac{1}{2} \underbrace{\underline{\varepsilon} \cdot \underline{\varepsilon}}_{\underline{\sigma}}$$

VISTO CHE IN TERMINI DI $\underline{\varepsilon}$: $\underline{\sigma} = 2G \underline{\varepsilon} + \lambda (\text{tr } \underline{\varepsilon}) \underline{I}$

$$\Rightarrow \boxed{\varphi(\underline{\varepsilon}) = G \underline{\varepsilon} \cdot \underline{\varepsilon} + \frac{\lambda}{2} (\text{tr } \underline{\varepsilon})^2} \Rightarrow \underline{\sigma} = \frac{\partial \varphi}{\partial \underline{\varepsilon}} = 2G \underline{\varepsilon} + \lambda (\text{tr } \underline{\varepsilon}) \underline{I}$$

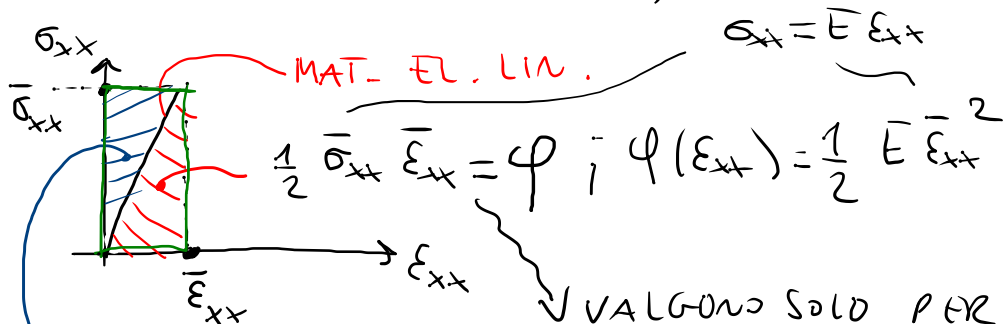
OSS.: 1) L'EQ. (*) VALE PER TUTTI I MATERIALI PER CUI φ È UNA FUNZ. DI STATO. IL CASO DI ELASTICITÀ LINEARE ISOTROPA È SOLO UN CASO PARTICOLARE.

2) I MATERIALI PER CUI $\exists \varphi$ SI CHIAMANO IPERELASTICI ° ELASTICI DI GREEN

3) SE $\varphi \nexists$ (PROBL. NON CONSERVATIVO) $\Rightarrow$ MATERIALI ELASTICI DI CAUCHY.

POTENZ. ELASTICO COMPLEMENTARE (φ^c)

$\sigma_{xx} \leftarrow \square \rightarrow \sigma_{xx}$ (MONOASSIALE)

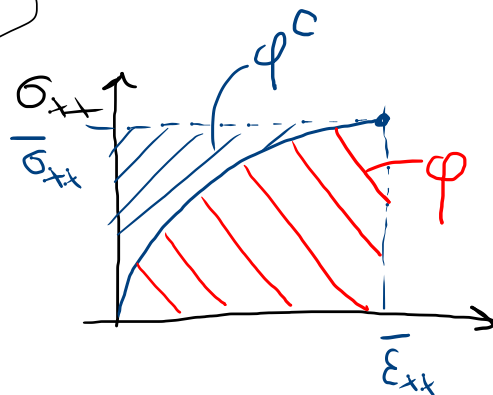


VALGONO SOLO PER MATERIALI

$$\varphi^c = \frac{1}{2} \bar{\sigma}_{xx} \bar{\epsilon}_{xx} \rightarrow \text{ELASTICI } \underline{\text{LINEARI}}$$

$$\varphi + \varphi^c = \bar{\sigma}_{xx} \bar{\epsilon}_{xx}$$

PROPR. CHE VALE
PER TUTTI I MAT.
ELASTICI.



MAT. ELASTICO
NON LINEARE

$$\varphi \neq \varphi^c$$

PROPRIETÀ DI φ^c NEL CASO DI MAT. ELASTICO LINEARE

$$\rightarrow \varphi + \varphi^c = \underset{\sim}{\sigma} \cdot \underset{\sim}{\varepsilon}$$

$$\rightarrow \text{se } \varphi(\underset{\sim}{\varepsilon}) \Rightarrow \varphi^c(\underset{\sim}{\sigma}) \Rightarrow$$

$$\boxed{\varepsilon_{ij} = \frac{\partial \varphi^c}{\partial \sigma_{ij}}}$$

$$\rightarrow \varphi^c(\underset{\sim}{\sigma}) = \frac{1}{2} \underset{\sim}{\sigma} \cdot \underset{\sim}{\varepsilon} = \frac{1}{2} \underset{\sim}{\sigma} \cdot \underset{\sim}{C}^{-1} \underset{\sim}{\sigma}$$

FUNZ. QUADRATICA IN σ_{ij}

NEL CASO ISOTROPICO:

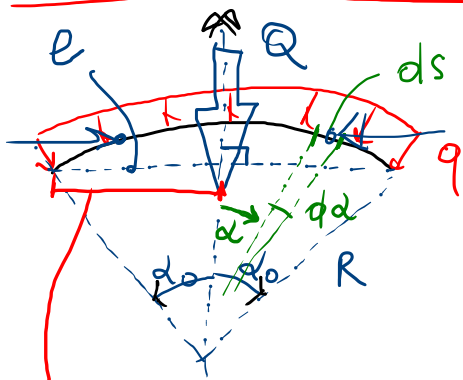
$$\varphi^c(\underset{\sim}{\sigma}) = \frac{1}{2} \left[\frac{1}{E} (\sigma_{xx}^2 + \sigma_{yy}^2 + \sigma_{zz}^2) - \frac{2\nu}{E} (\sigma_{xx}\sigma_{yy} + \sigma_{xx}\sigma_{zz} + \sigma_{yy}\sigma_{zz}) + \frac{2(1+\nu)}{E} (\tau_{xy}^2 + \tau_{xz}^2 + \tau_{yz}^2) \right]$$

$$\varepsilon_{xx} = \frac{\partial \varphi^c}{\partial \sigma_{xx}} = \frac{1}{E} \left[\frac{1}{E} \cancel{2} \sigma_{xx} - \frac{2\nu}{E} (\sigma_{yy} + \sigma_{zz}) \right] = \frac{1}{E} [\sigma_{xx} - \nu (\sigma_{yy} + \sigma_{zz})]$$

GLA' VISTO!
OK!

$$\varepsilon_{xy} = \frac{\partial \varphi^c}{\partial \tau_{xy}} = \frac{1}{2} \left[\cancel{2} \frac{(1+\nu)}{E} 2 \tau_{xy} \right] = \frac{(1+\nu)}{E} \tau_{xy} = \frac{\tau_{xy}}{2G}$$

COMPLEMENTI CARICHI DISTRIBUITI

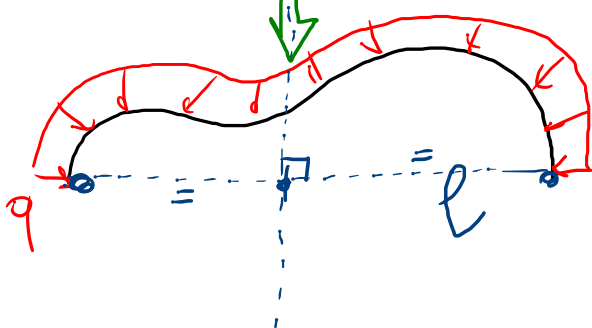


$R \sin \alpha_0$

l : LUNGH. ARCA

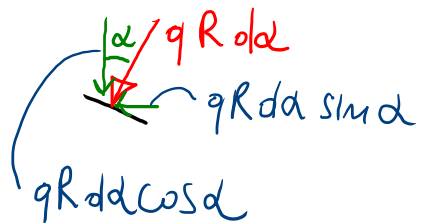
$$Q = ql$$

(SENZA
DIMOSTRAZIONE?)



$$Q = ql$$

$$ds = R d\alpha$$

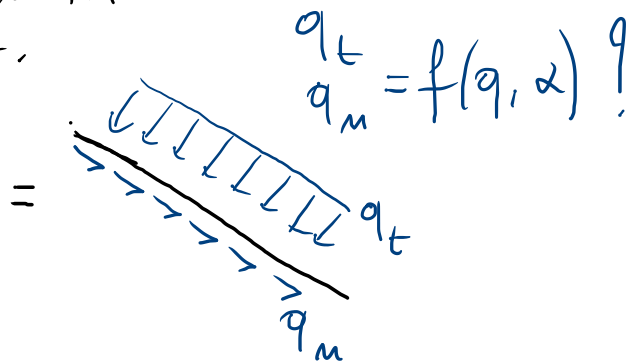
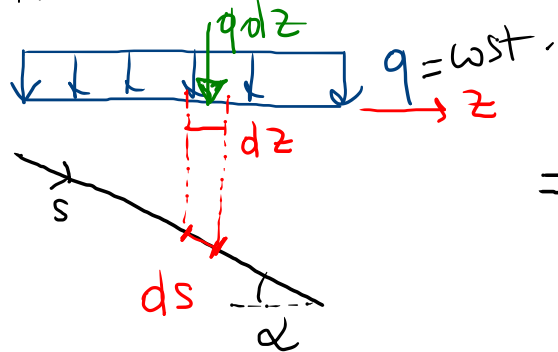


$$Q = \int_{-\alpha_0}^{\alpha_0} qR d\alpha \cos \alpha = qR \int_{-\alpha_0}^{\alpha_0} \cos \alpha d\alpha$$

$$= qR [\sin \alpha]_{-\alpha_0}^{\alpha_0} = qR [\sin \alpha_0 - \sin(-\alpha_0)] + \sin \alpha_0$$

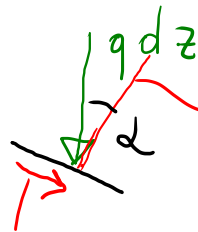
$$Q = 2qR \sin \alpha_0 = ql$$

CARICO SU FALDA INCLINATA



$$q_t = f(q, \alpha)!$$

$$dz = ds \cos \alpha$$



$$q dz \cos \alpha = q_t ds$$

$$q ds \cos^2 \alpha = q_t ds$$

$$q dz \sin \alpha = q_n ds \longrightarrow q ds \cos \alpha \sin \alpha = q_n ds$$

q : NON SI PROIETTA

$q dz$: FORZA $\Rightarrow$

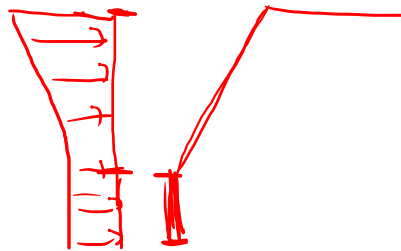
SI PROIETTA

(O SCOMPONE)

$$q_t = q \cos^2 \alpha$$

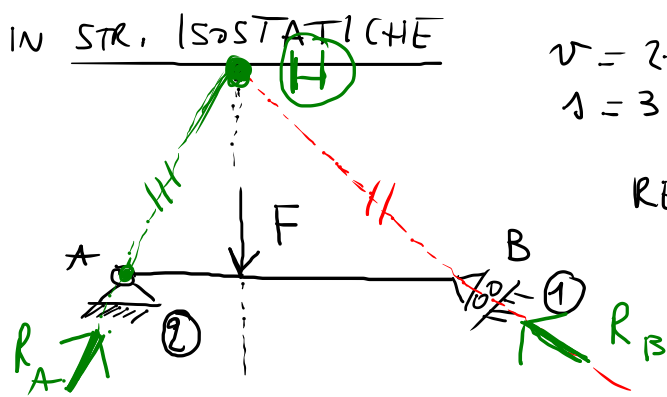
$$q_n = q \cos \alpha \sin \alpha$$

VENTO



SOLUZIONE GRAFICA DELLE STR. ISOSTATICHE

METODO ALTERNATIVO A QUELLO "ANALITICO" PER CALCOLARE LE REAZ. VINCULI
IN STR. ISOSTATICHE



$$\begin{aligned} v &= 2 + 1 = 3 \\ a &= 3, g = 3 \end{aligned} \quad \text{ISO ST.}$$

REAZ. VINCULI?

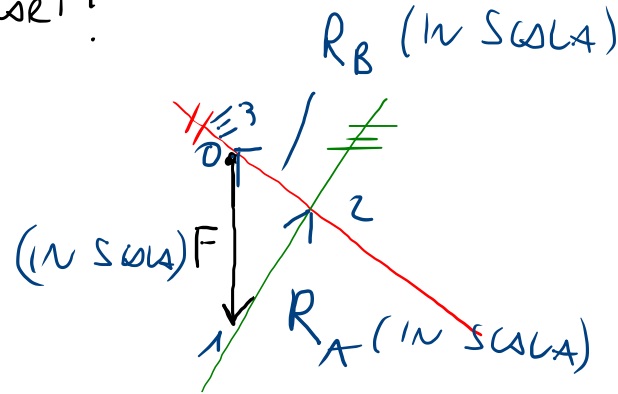
HO 3 FORZE IN GIOCO

R_A, R_B, F

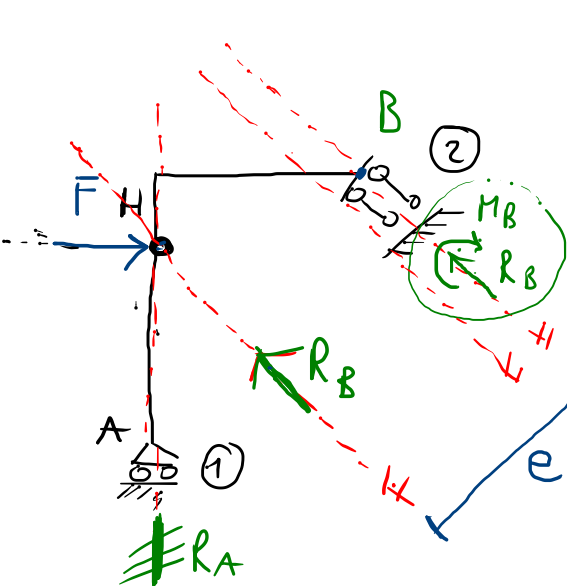
INCOGNITE

PER L'EQUIL. DI 3 FORZE NON PARALLELE OCCORRE

- R_A, R_B, F SI "INCONTRINO" IN UN PUNTO (H)
- CHIUDANO IL TRIANGOLO DELLE FORZE



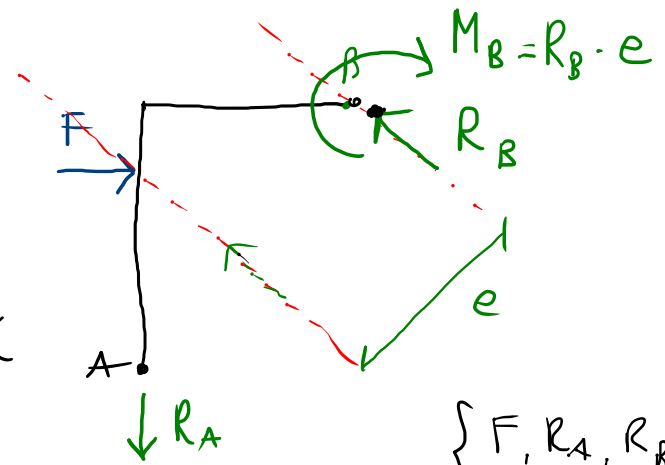
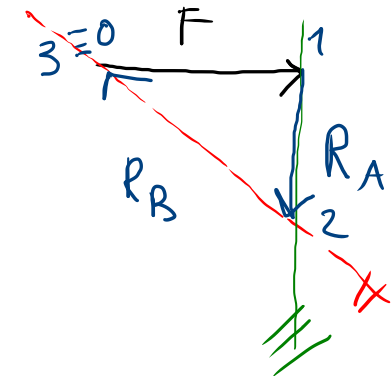
PROBL.
RISOLTO



$v = 3$
 $g = 3$
 $\Delta = 3$
 STR. INST.

REAZ. UNCOUPL. OK
 GRAFIC.

$\{F, R_A, R_B\}$: SIST. NULLO



TRIANGOLO
 CHIUSO

$\{F, R_A, R_B, M_B\}$
 SIST. NULLO