

$\Delta_{gl}$

$\Delta_1 = -i \partial^\mu A_\mu - i A_\mu \partial^\mu$

$$\begin{aligned}
 & -\frac{1}{2} \text{Tr} \left( (-\partial^2)^{-1} \Delta_1 (-\partial^2)^{-1} \Delta_1 \right) = -\frac{1}{2} \int d^d y \langle y | \text{tr} \left( (-\partial^2)^{-1} \Delta_1 (-\partial^2)^{-1} \Delta_1 \right) | y \rangle = \\
 & = +\frac{1}{2} \int d^d y \langle y | \left[ \text{tr} \left( (-\partial^2)^{-1} \partial^\mu A_\mu^a t_{Ad}^a (-\partial^2)^{-1} \partial^\nu A_\nu^b t_{Ad}^b \right) + \right. \\
 & \quad + \text{Ar} \left( (-\partial^2)^{-1} A_\mu^a t_{Ad}^a \partial^\mu (-\partial^2)^{-1} \partial^\nu A_\nu^b t_{Ad}^b \right) + \\
 & \quad + \text{Ar} \left( (-\partial^2)^{-1} \partial^\mu A_\mu^a t_{Ad}^a (-\partial^2)^{-1} A_\nu^b \partial^\nu t_{Ad}^b \right) + \\
 & \quad \left. + \text{Ar} \left( (-\partial^2)^{-1} A_\mu^a t_{Ad}^a \partial^\mu (-\partial^2)^{-1} A_\nu^b t_{Ad}^b \partial^\nu \right) \right] | y \rangle
 \end{aligned}$$

*Annotations:*  
 - Blue arrows point to the  $\int d^d x$  terms.  
 - Green dotted lines and arrows labeled "cicli e tracce" connect the terms in the trace.  
 - A question mark points to the first term.

Consideriamo il primo termine

$\langle y | k \rangle = e^{i k y}$

1<sup>st</sup> term =  $\frac{1}{2} \int d^d x d^d y \text{Ar} \left[ A_\mu^a(x) t_{Ad}^a A_\nu^b(y) t_{Ad}^b \right]$

$k'$  - completazione

$\langle y | k \rangle = e^{i k y}$

$$\begin{aligned}
 & \langle y | (-\partial^2)^{-1} \partial^\mu | x \rangle \langle x | (-\partial^2)^{-1} \partial^\nu | y \rangle \\
 & \frac{1}{-(i k)^2} i k^\mu e^{i k(y-x)} \quad \frac{1}{-(i k')^2} i k'^\nu e^{-i k'(y-x)} \\
 & - \frac{k^\mu}{k^2} \frac{k'^\nu}{k'^2} e^{i(k-k')(y-x)}
 \end{aligned}$$

2<sup>nd</sup> term =

... ← stessa prima riga del 1<sup>st</sup> term

$$\begin{aligned}
 & \langle y | (-\partial^2)^{-1} | x \rangle \cdot \langle x | \partial^\mu (-\partial^2)^{-1} \partial^\nu | y \rangle \\
 & + \frac{1}{k^2} \cdot \frac{(-) k^\mu k'^\nu}{k'^2} e^{i(k-k')(y-x)}
 \end{aligned}$$

3<sup>rd</sup> ferm

$$-\frac{k^\mu k^\nu}{k^2} \cdot \frac{1}{k^2} e^{i(k-k')(y-x)}$$

4<sup>th</sup> ferm

$$-\frac{k^\nu}{k^2} \frac{k'^\mu}{k'^2} e^{i(k-k')(y-x)}$$



$$1^{st} + 2^{nd} + 3^{rd} + 4^{th} = -\frac{e^{i(k-k')(y-x)}}{k^2 k'^2} (k+k')^\mu (k+k')^\nu$$

Quindi:

$$-\frac{1}{2} \text{Tr} \left( (-\partial^2)^{-1} \Delta_1 (-\partial^2)^{-1} \Delta_1 \right) =$$

$$= -\frac{1}{2} \int d^d x d^d y \int \frac{d^d k}{(2\pi)^d} \frac{d^d k'}{(2\pi)^d} e^{i(k-k')(y-x)} \frac{(k+k')^\mu (k+k')^\nu}{k^2 k'^2} \text{Tr} \left( A_\mu^a(x) t_{ad}^a A_\nu^b(y) t_{da}^b \right)$$

$$\tilde{A}(k) = \int d^d x e^{ikx} A(x) \rightarrow A(x) = \int \frac{d^d p}{(2\pi)^d} e^{-ipx} \tilde{A}(p)$$

$$A(y) = \int \frac{d^d q}{(2\pi)^d} e^{-iqy} \tilde{A}(q)$$



$$-\frac{1}{2} \int \underline{d^d x} \underline{d^d y} \frac{d^d k}{(2\pi)^d} \frac{d^d k'}{(2\pi)^d} \frac{d^d p}{(2\pi)^d} \frac{d^d q}{(2\pi)^d} \frac{e^{i(k-k')y}}{\delta(k-k'-q)} \frac{e^{-i(k-k')x}}{\delta(k-k'+p)} \frac{e^{-ipx}}{e^{-iqy}}$$

$$\cdot \frac{(k+k')^\mu (k+k')^\nu}{k^2 k'^2} \tilde{A}_\mu^a(p) \tilde{A}_\nu^b(q) \text{Ar}(t_{Ad}^a t_{Ad}^b)$$

$$\begin{cases} k-k'-q=0 \\ k-k'+p=0 \end{cases} \Rightarrow \begin{cases} q=-p \\ k'=p+k \end{cases}$$

$$= -\frac{1}{2} \int \frac{d^d p}{(2\pi)^d} \text{Ar}(t_{Ad}^a t_{Ad}^b) \tilde{A}_\mu^a(p) \tilde{A}_\nu^b(-p) \int \frac{d^d k}{(2\pi)^d} \frac{(p+2k)^\mu (p+2k)^\nu}{k^2 (k+p)^2}$$

Riscriviamo anche il primo termine nello sp. dei momenti:

$$\int d^d x A^{\mu a}(x) A_\mu^b(x) \text{Ar}(t_{Ad}^a t_{Ad}^b) \int \frac{d^d k}{(2\pi)^d} \frac{1}{k^2}$$

$$\text{Tr}((- \partial^2) \Delta_L) = \int \underline{d^d x} \int \frac{d^d p}{(2\pi)^d} \frac{d^d q}{(2\pi)^d} \frac{d^d k}{(2\pi)^d} \text{tr}(t_{Ad}^a t_{Ad}^b) \tilde{A}^{\mu a}(p) \tilde{A}_\mu^b(q) \frac{1}{k^2} e^{-ipx} e^{-iqx}$$

$$= \int \frac{d^d p}{(2\pi)^d} \text{tr}(t_{Ad}^a t_{Ad}^b) \tilde{A}_\mu(p) \tilde{A}_\nu(-p) \int \frac{d^d k}{(2\pi)^d} \frac{\delta^{\mu\nu}}{k^2} \leftarrow \text{siamo su sp. Euclideo}$$

$$\text{Tr} \log \Delta_{g^{\mu\nu}} = \int \frac{d^d p}{(2\pi)^d} \text{tr}(t_{Ad}^a t_{Ad}^b) \tilde{A}_\mu(p) \tilde{A}_\nu(-p) \cdot \left\{ \int \frac{d^d k}{(2\pi)^d} \frac{\delta_{\mu\nu}}{k^2} + \right. \\ \left. -\frac{1}{2} \int \frac{d^d k}{(2\pi)^d} \frac{(p+2k)^\mu (p+2k)^\nu}{k^2 (k+p)^2} \right\}$$

Per calcolare  $\{ \dots \}$  usiamo gli integrali di Feynman

$$\int \frac{d^d k}{(2\pi)^d} \frac{1}{(k^2 + a^2)^A} = \frac{\Gamma(A - d/2)}{(a^2)^{A-d/2} (4\pi)^{d/2} \Gamma(A)}$$

$$\Rightarrow \lim_{\epsilon \rightarrow 0} \int \frac{d^d k}{(2\pi)^d} \frac{1}{k^2 + \epsilon^2} = \lim_{\epsilon \rightarrow 0} \frac{\Gamma(1 - d/2)}{\epsilon^{1-d/2} (4\pi)^{d/2}} = 0 \quad \text{for } d > 2$$

$A=1 \quad a=\epsilon$

$$\frac{1}{k^2(k+p)^2} = \int_0^1 d\zeta \frac{1}{[k^2(1-\zeta) + (k^2 + 2kp + p^2)\zeta]^2} = \int_0^1 d\zeta \frac{1}{(k^2 + 2\zeta p \cdot k + \zeta p^2)^2}$$

$$\int \frac{d^d k}{(2\pi)^d} \frac{1}{(k^2 + 2kp + b^2)^A} = \frac{\Gamma(A - d/2)}{(b^2 - p^2)^{A-d/2} (4\pi)^{d/2} \Gamma(A)} \quad \Gamma(n+1) = n!$$

$$\Rightarrow \int \frac{d^d k}{(2\pi)^d} \frac{p^\mu p^\nu}{(k^2 + 2\zeta p \cdot k + \zeta p^2)^2} = \frac{p^\mu p^\nu \Gamma(2 - d/2)}{(\zeta(1-\zeta)p^2)^{2-d/2} (4\pi)^{d/2}}$$

$$\int \frac{d^d k}{(2\pi)^d} \frac{k^\mu}{(k^2 + 2kp + b^2)^A} = -\frac{p^\mu \Gamma(A - d/2)}{(b^2 - p^2)^{A-d/2} (4\pi)^{d/2} \Gamma(A)}$$

$$\Rightarrow \int \frac{d^d k}{(2\pi)^d} \frac{2k^\mu p^\nu + 2k^\nu p^\mu}{(k^2 + 2\zeta p \cdot k + \zeta p^2)^2} = -\frac{4\zeta p^\mu p^\nu \Gamma(2 - d/2)}{(\zeta(1-\zeta)p^2)^{2-d/2} (4\pi)^{d/2}}$$

$$\int \frac{d^d k}{(2\pi)^d} \frac{k^\mu k^\nu}{(k^2 + 2kp + b^2)^A} = \frac{1}{(4\pi)^{d/2} \Gamma(A)} \left[ \frac{p^\mu p^\nu \Gamma(A - d/2)}{(b^2 - p^2)^{A-d/2}} + \frac{1}{2} \delta^{\mu\nu} \frac{\Gamma(A - 1 - d/2)}{(b^2 - p^2)^{A-1-d/2}} \right]$$

$$\Rightarrow \int \frac{d^d k}{(2\pi)^d} \frac{4k^\mu k^\nu}{(k^2 + 2\zeta p \cdot k + \zeta p^2)^2} = \frac{4\zeta^2 p^\mu p^\nu \Gamma(2 - d/2)}{(\zeta(1-\zeta)p^2)^{2-d/2} (4\pi)^{d/2}} + \frac{1}{2} \frac{4\delta^{\mu\nu} \Gamma(1 - d/2)}{(\zeta(1-\zeta)p^2)^{1-d/2} (4\pi)^{d/2}}$$

$$\Rightarrow -\frac{1}{2} \int \frac{d^d k}{(2\pi)^d} \frac{(p+2k)^\mu (p+2k)^\nu}{k^2 (k+p)^2} = -\frac{1}{2} \int_0^1 d\xi \left\{ \frac{(1-2\xi)^2 p^\mu p^\nu \Gamma(2-d/2)}{(\xi(1-\xi)p^2)^{2-d/2} (4\pi)^{d/2}} + \frac{2 \delta^{\mu\nu} \Gamma(1-d/2)}{(\xi(1-\xi)p^2)^{1-d/2} (4\pi)^{d/2}} \right\}$$

$$4\xi^2 - 4\xi + 1 = (2\xi - 1)^2$$

$$d = 2\omega \quad (\omega \rightarrow 2)$$

espandiamo la  
funz. in d attorno  
a d=4

$$\Gamma(1-d/2) = \frac{1}{1-d/2} \Gamma(2-d/2)$$

Parte divergente  
in d → 4:

Parte divergente in d → 4:

$$\int_0^1 d\xi \frac{(1-2\xi)^2 p^\mu p^\nu}{(4\pi)^2} \frac{1}{2-\omega}$$

$$\int_0^1 d\xi \xi(1-\xi) \frac{\delta^{\mu\nu} p^2}{(4\pi)^2} \left( -\frac{2}{2-\omega} \right)$$

$$\int_0^1 d\xi (1-4\xi+4\xi^2) = 1 - 2 + \frac{4}{3} = \frac{1}{3}$$

$$\int_0^1 \xi(1-\xi) d\xi = \frac{1}{2} - \frac{1}{3} = \frac{1}{6}$$

$$\rightsquigarrow \text{Parte div.} = -\frac{1}{2} \frac{1}{48\pi^2} \frac{1}{2-\omega} (p^\mu p^\nu - p^2 \delta^{\mu\nu})$$

⇓

$$\text{Tr log } \Delta_{gh} \approx \int \frac{d^d p}{(2\pi)^d} \text{tr} (t_{Ad}^a t_{Ad}^b) \tilde{A}_{(p)}^{\mu a} \tilde{A}_{(-p)}^{\nu b} \left( -\frac{1}{2} \frac{1}{3} \frac{1}{16\pi^2} \frac{1}{2-\omega} (p^\mu p^\nu - p^2 \delta^{\mu\nu}) \right) + \dots$$

parte divergente

$$= -\frac{1}{2} \frac{1}{3} \frac{1}{(4\pi)^2} \underbrace{\text{tr} (t_{Ad}^a t_{Ad}^b)}_{\substack{C(\text{Adj}) \delta^{ab} \\ C_2(G)}} \underbrace{\int \frac{d^d p}{(2\pi)^d} \tilde{A}_{(p)}^{\mu a} \tilde{A}_{(-p)}^{\nu b} (p^\mu p^\nu - p^2 \delta^{\mu\nu})}_{\downarrow} \frac{1}{2-\omega}$$

Ora vedremo che qta è l'espressione  
di  $g^2 S[A]$  (la parte quadratica in A)  
nello spazio dei momenti.

↓

Parte quadratica di azione S:

$$\begin{aligned}
 & \frac{1}{4g^2} \int (\partial_\mu A_\nu^a - \partial_\nu A_\mu^a) (\partial^\mu A^{\nu a} - \partial^\nu A^{\mu a}) d^d x = \\
 & = \frac{1}{4g^2} \int \frac{d^d x}{(2\pi)^d} \frac{d^d q_1}{(2\pi)^d} \frac{d^d q_2}{(2\pi)^d} (-iq_{1\mu} \tilde{A}_\nu^a(q_1) + iq_{1\nu} \tilde{A}_\mu^a(q_1)) (-iq_2^\mu \tilde{A}^{\nu a}(q_2) + iq_2^\nu \tilde{A}^{\mu a}(q_2)) \\
 & \quad \cdot \underbrace{e^{-iq_1 x - iq_2 x}}_{\rightarrow \delta(q_1 + q_2)} \quad q_1 = -q_2 \equiv p \\
 & = \frac{1}{4g^2} \int \frac{d^d p}{(2\pi)^d} (p_\mu \tilde{A}_\nu^a(p) - p_\nu \tilde{A}_\mu^a(p)) (p^\mu \tilde{A}^{\nu a}(-p) - p^\nu \tilde{A}^{\mu a}(-p)) = \\
 & = -\frac{1}{2g^2} \int \frac{d^d p}{(2\pi)^d} (p^\mu p^\nu - p^2 \delta^{\mu\nu}) \tilde{A}_\mu^a(p) \tilde{A}_\nu^a(-p)
 \end{aligned}$$

### Contributo di $\Delta_{gauge}$

$$\Delta_{gauge}^{\mu\nu} = -D^2 \delta^{\mu\nu} - 2i F^{\alpha\mu\nu} t_{Ad}^a = \underbrace{\Delta_{gh} \delta^{\mu\nu}} - 2i F^{\alpha\mu\nu} t_{Ad}^a$$

prendo traccia su sugli indici di Lorentz, otteniamo un fattore 4. ← op. come nei ghost ma ora abbiamo  $\delta^{\mu\nu}$

$$\begin{aligned}
 \text{Tr log } \Delta_{gauge} &= \text{Tr log} \left( -\partial^2 \delta^{\mu\nu} + \overbrace{(\Delta_1 + \Delta_2) \delta^{\mu\nu} - 2i F^{\alpha\mu\nu} t_{Ad}^a}^{(\dots)} \right) = \\
 &= \text{Tr log} \underbrace{(-\partial^2 \delta^{\mu\nu})}_{\text{cont. trascurabile}} + \text{Tr log} (1 + (-\partial^2)^{-1} (\dots)) \approx \text{termini con } \Delta_1, \Delta_2 \text{ trascurabili} \\
 &\approx \text{Tr} [(-\partial^2)^{-1} (\dots)] - \frac{1}{2} \text{Tr} [(-\partial^2)^{-1} (\dots)]^2 = \text{Tr log } \Delta_{gh} \delta^{\mu\nu} \\
 & \quad \uparrow \text{termini "crist."} \\
 & \quad \text{F}_{\mu\nu} \Delta_1 \delta^{\mu\nu} \text{ sono zero} \\
 & \quad \text{perch\u00e9 } F_{\mu\nu} \delta^{\mu\nu} = 0 \\
 &= 4 \text{ Tr log } \Delta_{gh} + F^{\mu\nu} \text{-terms} \leftarrow \text{quadratici in } F
 \end{aligned}$$

$$\begin{aligned} \text{div } F^{\mu\nu} \text{ terms} &= \overbrace{-\frac{1}{2} (-2i)^2}^{+2} \text{Tr} \left( (-\partial^2)^{-1} F^{\mu\nu} t_{Ad}^a (-\partial^2)^{-1} F_{\nu\mu}^b t_{Ad}^b \right) \\ &= -2 \int d^d y \langle y | (-\partial^2)^{-1} F^{\mu\nu a} (-\partial^2)^{-1} F_{\mu\nu}^b (-\partial^2)^{-1} | y \rangle \text{Tr} (t_{Ad}^a t_{Ad}^b) \\ &= -2 \int d^d y d^d x \langle y | (-\partial^2)^{-1} | x \rangle \langle x | (-\partial^2)^{-1} | y \rangle F^{\mu\nu a}(x) F_{\mu\nu}^b(y) \text{Tr} (t_{Ad}^a t_{Ad}^b) \\ &= -2 \int d^d y d^d x \frac{d^d k}{(2\pi)^d} \frac{d^d k'}{(2\pi)^d} \frac{1}{k^2} \frac{1}{k'^2} F^{\mu\nu a}(x) F_{\mu\nu}^b(y) \text{Tr} (t^a t^b) e^{ik(y-x)} e^{ik'(x-y)} \\ &= -2 \int d^d y d^d x \frac{d^d k}{(2\pi)^d} \frac{d^d k'}{(2\pi)^d} \frac{d^d p}{(2\pi)^d} \frac{d^d q}{(2\pi)^d} \frac{1}{k^2} \frac{1}{k'^2} \text{Tr} (t^a t^b) e^{-ipx} e^{-iqy} e^{i(k-k')y} e^{i(k'-k)x} \end{aligned}$$

trans. Fourier

$$\cdot (-i)^2 (p^\sigma \tilde{A}_\rho^\sigma(p) - p^\rho \tilde{A}_\sigma^\rho(p)) (q_\sigma \tilde{A}_\rho^\sigma(q) - q_\rho \tilde{A}_\sigma^\rho(q))$$

$$\delta(k'-k-p) \quad \delta(k-k'-q) \rightarrow \begin{matrix} q = -p \\ k' = k+p \end{matrix}$$

$$= -\frac{1}{2} \int \frac{d^d p}{(2\pi)^d} \tilde{A}_\mu^a(p) \tilde{A}_\nu^b(-p) \text{Tr} (t^a t^b) \cdot \int \frac{d^d k}{(2\pi)^d} \frac{4 (p^\sigma \delta^{\mu\sigma} - p^\sigma \delta^{\mu\sigma}) (p_\sigma \delta^\nu - p_\sigma \delta^\nu)}{k^2 (k+p)^2}$$

$$= 4 \int \frac{d^d p}{(2\pi)^d} \tilde{A}_\mu^a(p) \tilde{A}_\nu^b(-p) \text{Tr} (t^a t^b) (p^\mu p^\nu - p^2 \delta^{\mu\nu}) \int d\xi \frac{\Gamma(2-d/2)}{(\xi(1-\xi)p^2)^{2-d/2}} (4\pi)^{d/2}$$

$$\sim \text{div.} \quad \frac{4 C(Adj)}{(4\pi)^2} \delta^{ab} \frac{1}{2-\omega} \int \frac{d^d p}{(2\pi)^d} \tilde{A}_\mu^a(p) \tilde{A}_\nu^b(-p) (p^\mu p^\nu - p^2 \delta^{\mu\nu})$$



$$\begin{aligned} \text{Tr log } \Delta_{\text{gauge}} \Big|_{\text{divergent part}} &\approx 4 \text{Tr log } \Delta_{gh} \Big|_{\text{divergent part}} + \\ &+ \frac{4}{(4\pi)^2} C_2(G) \delta^{ab} \frac{1}{2-\omega} \int \frac{d^d p}{(2\pi)^d} \tilde{A}_\mu^a(p) \tilde{A}_\nu^b(-p) (p^\mu p^\nu - p^2 \delta^{\mu\nu}) \end{aligned}$$

$$= -\frac{4^2}{3} \frac{1}{(4\pi)^2} c_2(G) \delta^{ab} \frac{1}{2-\omega} \int \frac{d^d p}{(2\pi)^d} \hat{A}_\mu^a(p) \hat{A}_\nu^b(-p) (p^\mu p^\nu - p^2 \delta^{\mu\nu})$$

$$+ \frac{4}{(4\pi)^2} c_2(G) \delta^{ab} \frac{1}{2-\omega} \int \frac{d^d p}{(2\pi)^d} \hat{A}_\mu^a(p) \hat{A}_\nu^b(-p) (p^\mu p^\nu - p^2 \delta^{\mu\nu})$$

$$= \frac{10}{3} \frac{1}{(4\pi)^2} c_2(G) \delta^{ab} \frac{1}{2-\omega} \int \frac{d^d p}{(2\pi)^d} \hat{A}_\mu^a(p) \hat{A}_\nu^b(-p) (p^\mu p^\nu - p^2 \delta^{\mu\nu})$$