# Advanced Quantum Mechanics

Angelo Bassi

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#### Quantum Algorithms

We will study the three historically most important algorithms:

- Simple ones (Deutsch, Deutsch-Jozsa...)
- Grover (search in a data base)
- Shor (factorization)

What is special about quantum algorithms?

#### Quantum Parallelism

Più in generale

$$U_s | \vec{x}, y \rangle = | \vec{x} \rangle | y \oplus f_s(\vec{x}) \rangle,$$

Given an input x, a typical quantum computer computes f(x) in such a way as

$$U_f: |x\rangle|0\rangle \mapsto |x\rangle|f(x)\rangle,$$
 (4.61)

where  $U_f$  is a unitary matrix that implements the function f.

Suppose  $U_f$  acts on the input which is a superposition of many states. Since  $U_f$  is a linear operator, it acts simultaneously on all the vectors that constitute the superposition. Thus the output is also a superposition of all the results;

$$U_f: \sum_{x} |x\rangle |0\rangle \longleftrightarrow \sum_{x} |x\rangle |f(x)\rangle. \tag{4.62}$$

All values of the function computed at once. Very easy!! But... measurements will make the wave function collapse giving only one output. No advantage

#### Quantum Algorithms

The goal of a quantum algorithm is to operate in such a way that the particular outcome we want to observe has a larger probability to be measured than the other outcomes.

Let  $f: \{0,1\} \to \{0,1\}$  be a binary function. Note that there are only four possible f, namely

$$f_1: 0 \mapsto 0, \ 1 \mapsto 0, \quad f_2: 0 \mapsto 1, \ 1 \mapsto 1,$$
  
 $f_3: 0 \mapsto 0, \ 1 \mapsto 1, \quad f_4: 0 \mapsto 1, \ 1 \mapsto 0.$ 

The first two cases,  $f_1$  and  $f_2$ , are called <u>constant</u>, while the rest,  $f_3$  and  $f_4$ , are <u>balanced</u>. If we only have classical resources, we need to evaluate f twice to tell if f is constant or balanced. There is a quantum algorithm, however, with which it is possible to tell if f is constant or balanced with a single evaluation of f, as was shown by Deutsch [2].

First we need to turn the classical function f(x) into a quantum one.

- 1. Make it reversible.
- 2. Define it on the computational basis to act like the classical circuit and extend it by linearity.

$$U_f:|x,y\rangle\mapsto|x,y\oplus f(x)\rangle$$

Where  $\oplus$  is addition mod 2.

The algorithm is structured as follows.

- 1. Start with the state  $|01\rangle$ .
- 2. Apply an Hadamard on both qubits:  $\frac{1}{2}(|00\rangle |01\rangle + |10\rangle |11\rangle)$
- 3. Apply the operator U<sub>f</sub> implementing the function

$$\frac{1}{2}(|0, f(0)\rangle - |0, 1 \oplus f(0)\rangle + |1, f(1)\rangle - |1, 1 \oplus f(1)\rangle)$$

$$= \frac{1}{2}(|0, f(0)\rangle - |0, \neg f(0)\rangle + |1, f(1)\rangle - |1, \neg f(1)\rangle),$$



Quantum parallelism: all values computed at once

#### 4. Apply an Hadamard to the first qubit

$$\frac{1}{2\sqrt{2}}\left[\left(|0\rangle+|1\rangle\right)\left(|f(0)\rangle-|\neg f(0)\rangle\right)+\left(|0\rangle-|1\rangle\right)\left(|f(1)\rangle-|\neg f(1)\rangle\right)\right]$$

The wave function reduces to

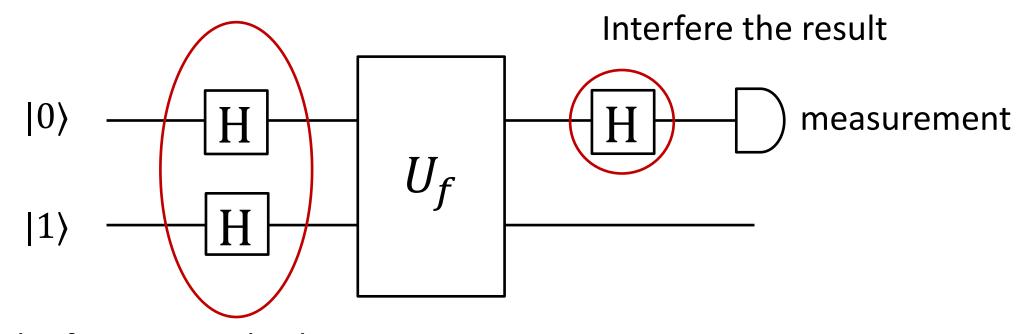
$$\frac{1}{\sqrt{2}}|0\rangle(|f(0)\rangle - |\neg f(0)\rangle) \tag{5.1}$$

in case f is constant, for which  $|f(0)\rangle = |f(1)\rangle$ , and

$$\frac{1}{\sqrt{2}}|1\rangle(|f(0)\rangle - |\neg f(0)\rangle) \tag{5.2}$$

if f is balanced, for which  $|\neg f(0)\rangle = |f(1)\rangle$ . Therefore the measurement of the first qubit tells us whether f is constant or balanced.

5. Measure the first qubit



Calculate the function on both input values simultaneously

Let us first define the **Deutsch-Jozsa problem**. Suppose there is a binary function

$$f: S_n \equiv \{0, 1, \dots, 2^n - 1\} \to \{0, 1\}.$$
 (5.3)

We require that f be either constant or balanced as before. When f is constant, it takes a constant value 0 or 1 irrespetive of the input value x. When it is balanaced the value f(x) for the half of  $x \in S_n$  is 0, while it is 1 for the rest of x.

It is clear that we need at least  $2^{n-1} + 1$  steps, in the worst case with classical manipulations, to make sure if f(x) is constant or balanced with 100% confidence. It will be shown below that the number of steps reduces to a single step if we are allowed to use a quantum algorithm.

- 1. Prepare an (n+1)-qubit register in the state  $|\psi_0\rangle = |0\rangle^{\otimes n} \otimes |1\rangle$ . First n qubits work as input qubits, while the (n+1)st qubit serves as a "scratch pad." Such qubits, which are neither input qubits nor output qubits, but work as a scratch pad to store temporary information are called **ancillas** or **ancillary qubits**.
- 2. Apply the Walsh-Hadamard transforamtion to the register. Then we have the state

$$|\psi_1\rangle = U_{\mathcal{H}}^{\otimes n+1}|\psi_0\rangle = \frac{1}{\sqrt{2^n}}(|0\rangle + |1\rangle)^{\otimes n} \otimes \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$$
$$= \frac{1}{\sqrt{2^n}} \sum_{x=0}^{2^n-1} |x\rangle \otimes \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle). \tag{5.4}$$

3. Apply  $U_f|x\rangle|c\rangle = |x\rangle|c\oplus f(x)\rangle$ 

The state changes into

$$|\psi_{2}\rangle = U_{f}|\psi_{1}\rangle$$

$$= \frac{1}{\sqrt{2^{n}}} \sum_{x=0}^{2^{n}-1} |x\rangle \frac{1}{\sqrt{2}} (|f(x)\rangle - |\neg f(x)\rangle)$$

$$= \frac{1}{\sqrt{2^{n}}} \sum_{x} (-1)^{f(x)} |x\rangle \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle). \tag{5.5}$$

Although the gate  $U_f$  is applied once for all, it is applied to all the n-qubit states  $|x\rangle$  simultaneously.

4. The Walsh-Hadamard transformation (4.11) is applied on the first n qubits next. We obtain

$$|\psi_3\rangle = (W_n \otimes I)|\psi_2\rangle = \frac{1}{\sqrt{2^n}} \sum_{x=0}^{2^n - 1} (-1)^{f(x)} U_{\mathcal{H}}^{\otimes n} |x\rangle \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle).$$
 (5.6)

#### On the Hadamard gate

It is instructive to write the action of the one-qubit Hadamard gate in the following form,

$$U_{\rm H}|x\rangle = \frac{1}{\sqrt{2}}(|0\rangle + (-1)^x|1\rangle) = \frac{1}{\sqrt{2}} \sum_{y \in \{0,1\}} (-1)^{xy}|y\rangle,$$

where  $x \in \{0, 1\}$ , to find the resulting state. The action of the Walsh-Hadamard transformation on  $|x\rangle = |x_{n-1} \dots x_1 x_0\rangle$  yields

$$W_{n}|x\rangle = (U_{H}|x_{n-1}\rangle)(U_{H}|x_{n-2}\rangle)\dots(U_{H}|x_{0}\rangle)$$

$$= \frac{1}{\sqrt{2^{n}}} \sum_{y_{n-1},y_{n-2},\dots,y_{0}\in\{0,1\}} (-1)^{x_{n-1}y_{n-1}+x_{n-2}y_{n-2}+\dots+x_{0}y_{0}}$$

$$\times |y_{n-1}y_{n-2}\dots y_{0}\rangle$$

$$= \frac{1}{\sqrt{2^{n}}} \sum_{y=0}^{2^{n}-1} (-1)^{x\cdot y}|y\rangle, \qquad (5.7)$$

where  $x \cdot y = x_{n-1}y_{n-1} \oplus x_{n-2}y_{n-2} \oplus \ldots \oplus x_0y_0$ .

#### Coming back to step 4:

4. The Walsh-Hadamard transformation (4.11) is applied on the first n qubits next. We obtain

$$|\psi_3\rangle = (W_n \otimes I)|\psi_2\rangle = \frac{1}{\sqrt{2^n}} \sum_{x=0}^{2^n-1} (-1)^{f(x)} U_{\mathcal{H}}^{\otimes n} |x\rangle \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle). \quad (5.6)$$

$$= \frac{1}{2^n} \left( \sum_{x,y=0}^{2^n - 1} (-1)^{f(x)} (-1)^{x \cdot y} |y\rangle \right) \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle).$$

As we will see, this operation will make the different terms interfere in order to read the desired result

5. The first n qubits are measured. Suppose f(x) is constant. Then  $|\psi_3\rangle$  is put in the form

$$|\psi_3\rangle = \frac{1}{2^n} \sum_{x,y} (-1)^{x \cdot y} |y\rangle \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle)$$

up to an overall phase. Now let us consider the summation

$$\frac{1}{2^n} \sum_{x=0}^{2^n - 1} (-1)^{x \cdot y}$$

with a fixed  $y \in S_n$ . Clearly it vanishes since  $x \cdot y$  is 0 for half of x and 1 for the other half of x unless y = 0. Therefore the summation yields  $\delta_{y0}$ . Now the state reduces to

$$|\psi_3\rangle = |0\rangle^{\otimes n} \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle),$$

and the measurement outcome of the first n qubits is always 00...0.

Example with 3 qubits.

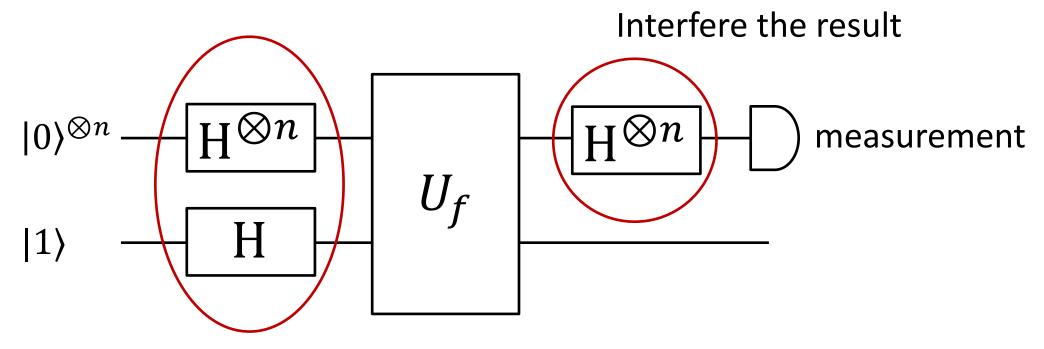
Take y = 110. Then  $x \cdot y = x_2 \oplus x_1$ 

X	$X_2 \oplus X_1$
000	0
001	0
010	1
011	1
100	1
101	1
110	0
111	0

Suppose f(x) is balanced next. The probability amplitude of  $|y=0\rangle$  in  $|\psi_3\rangle$  is proportional to

$$\sum_{x=0}^{2^{n}-1} (-1)^{f(x)} (-1)^{x \cdot 0} = \sum_{x=0}^{2^{n}-1} (-1)^{f(x)} = 0.$$

Therefore the probability of obtaining measurement outcome 00...0 for the first n qubits vanishes. In conclusion, the function f is constant if we obtain 00...0 upon the measurement of the first n qubits in the state  $|\psi_3\rangle$ , and it is balanced otherwise.



Calculate the function on both input values simultaneously

#### Bernstein-Vazirani Algorithm

The **Bernstein-Vazirani algorithm** is a special case of the Deutsch-Jozsa algorithm, in which f(x) is given by  $f(x) = c \cdot x$ , where  $c = c_{n-1}c_{n-2} \dots c_0$  is an n-bit binary number [4]. Our aim is to find c with the smallest number of evaluations of f. If we apply the Deutsch-Jozsa algorithm with this f, we obtain

$$|\psi_3\rangle = \frac{1}{2^n} \left[ \sum_{x,y=0}^{2^n-1} (-1)^{c \cdot x} (-1)^{x \cdot y} |y\rangle \right] \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle).$$

Let us fix y first. If we take y = c, we obtain

$$\sum_{x} (-1)^{c \cdot x} (-1)^{x \cdot c} = \sum_{x} (-1)^{2c \cdot x} = 2^{n}.$$

#### Bernstein-Vazirani Algorithm

If  $y \neq c$ , on the other hand, there will be the same number of x such that  $c \cdot x = 0$  and x such that  $c \cdot x = 1$  in the summation over x and, as a result, the probability amplitude of  $|y \neq c\rangle$  vanishes. By using these results, we end up with

$$|\psi_3\rangle = |c\rangle \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle).$$
 (5.9)

We are able to tell what c is by measuring the first n qubits.

#### Exercise

**EXERCISE 5.1** Let us take n=2 for definiteness. Consider the following cases and find the final wave function  $|\psi_3\rangle$  and evaluate the measurement outcomes and their probabilities for each case.

- $(1) f(x) = 1 \forall x \in S_2.$
- (2) f(00) = f(01) = 1, f(10) = f(11) = 0.
- (3) f(00) = 0, f(01) = f(10) = f(11) = 1. (This function is neither constant nor balanced.)

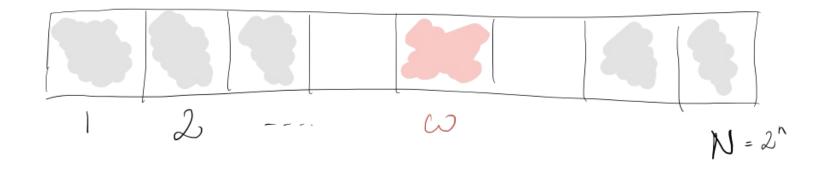
**EXERCISE 5.2** Consider the Bernstein-Vazirani algorithm with n=3 and c=101. Work out the quantum circuit depicted in Fig. 5.2 to show that the measurement outcome of the first three qubits is c=101.

Suppose there is a stack of  $N = 2^n$  files, randomly placed, that are numbered by  $x \in S_n \equiv \{0, 1, ..., N-1\}$ . Our task is to find an algorithm which picks out a particular file which satisfies a certain condition.

In mathematical language, this is expressed as follows. Let  $f: S_n \to \{0, 1\}$  be a function defined by

$$f(x) = \begin{cases} 1 & (x = z) \\ 0 & (x \neq z), \end{cases}$$
 (7.1)

where z is the address of the file we are looking for. It is assumed that f(x) is instantaneously calculable, such that this process does not require any computational steps. A function of this sort is often called an oracle as noted in Chapter 5. Thus, the problem is to find z such that f(z) = 1, given a function  $f: S_n \to \{0,1\}$  which assumes the value 1 only at a single point.



Clearly we have to check one file after another in a classical algorithm, and it will take O(N) steps on average. It is shown below that it takes only  $O(\sqrt{N})$  steps with Grover's algorithm. This is accomplished by amplifying the amplitude of the vector  $|z\rangle$  while cancelling that of the vectors  $|x\rangle$   $(x \neq z)$ .

We first needs to implement the function f(x) quantum mechanically. We define  $U_f$  as follows (**oracle**)

$$U_s |\vec{x}, y\rangle = |\vec{x}\rangle |y \oplus f_s(\vec{x})\rangle, \quad \text{con} \quad f_s(\vec{x}) = \begin{cases} 0 & \text{se } \vec{x} \neq \vec{s} \\ 1 & \text{se } \vec{x} = \vec{s} \end{cases}$$
 (1.35)

si ha:

$$U_s |\vec{x}\rangle \frac{1}{\sqrt{2}} \left[ |0\rangle - |1\rangle \right] = (-1)^{f_s(\vec{x})} |\vec{x}\rangle \frac{1}{\sqrt{2}} \left[ |0\rangle - |1\rangle \right]. \tag{1.36}$$

In questo modo si vede che ignorando l'ultimo qubit, l'operatore  $U_s$  può essere visto come un gate  $V_s$  che agisce sui primi n qubit. Otteniamo dunque la seguente relazione:

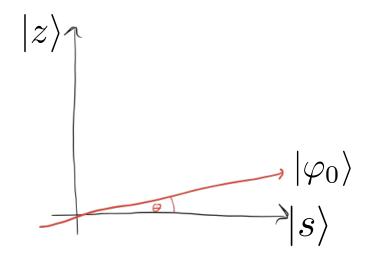
$$U_f|x\rangle = (-1)^{f(x)}|x\rangle$$

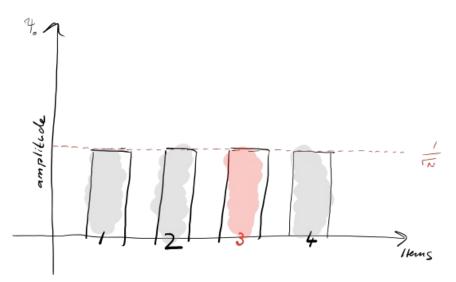
On the computational basis. We see that if x is an unmarked item, the oracle does nothing to the state. It flips the phase for the marked item. It is easy to see that

$$U_f = I - 2|z\rangle\langle z|$$

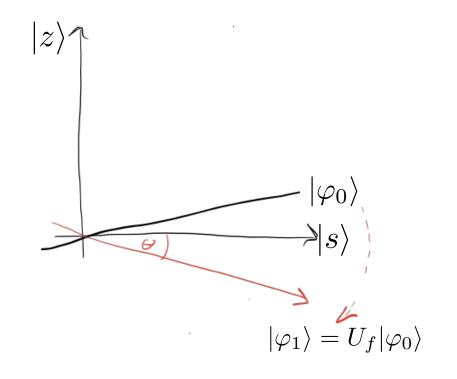
**Step 1:** Create an initially **equal weighted superposition** of all states (this is done with N Hadamard gates):

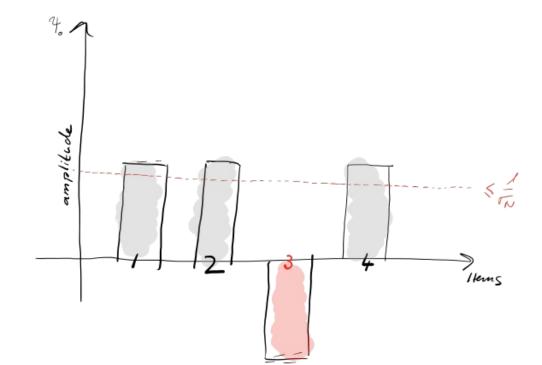
$$|\varphi_0\rangle = \frac{1}{\sqrt{N}} \sum_{x=0}^{N-1} |x\rangle.$$





**Step 2: Apply the oracle U**<sub>f</sub>. Geometrically this corresponds to a reflection of the state  $|z\rangle$  about  $|s\rangle$ . This transformation means that the amplitude in front of the  $|z\rangle$  state becomes negative, which in turn means that the average amplitude has been lowered.





#### **Step 3: Apply the gate**

$$D = -I + 2|\varphi_0\rangle\langle\varphi_0|.$$

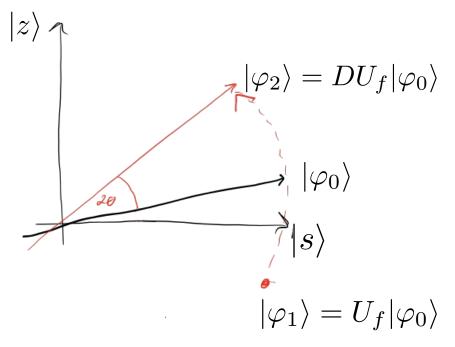
The action of the gate is the following

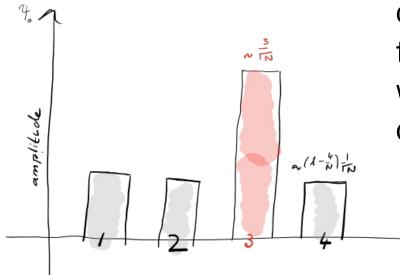
$$\begin{array}{c} \omega_x - \bar{\omega} \\ \overline{\omega} \\ \overline{\omega} - (\omega_x - \bar{\omega}) \\ \overline{\omega} - (\omega_x - \bar{\omega}) \\ = 2\bar{\omega} - \omega_x \\ \end{array}$$
 Reflection around the average

$$|\varphi\rangle = \sum_{x=0}^{N-1} \omega_x |x\rangle \rightarrow D|\varphi\rangle = \left[\frac{2}{N} \sum_{x,y=0}^{N-1} |x\rangle\langle y|\right] \sum_{z=0}^{N-1} \omega_z |z\rangle - \sum_{x=0}^{N-1} \omega_x |x\rangle$$
$$= \frac{2}{N} \left[\sum_{x=0}^{N-1} |x\rangle\right] \left[\sum_{y=0}^{N-1} \omega_y\right] - \sum_{x=0}^{N-1} \omega_x |x\rangle = \sum_{x=0}^{N-1} (2\bar{\omega} - \omega_x)|x\rangle$$

with 
$$ar{\omega} = rac{1}{N} \sum_{x=0}^{N-1} \omega_x$$
 average

In our case we get





Since the average amplitude has been lowered by the first reflection, this transformation boosts the negative amplitude of  $|z\rangle$  to roughly three times its original value, while it decreases the other amplitudes.

**Step 3:** go to step 2 an repeat the application of  $U_f$  and D a sufficient number of times. Let us call  $G_f = D U_f$ .

#### PROPOSITION 7.2 Let us write

$$G_f^k |\varphi_0\rangle = a_k |z\rangle + b_k \sum_{x \neq z} |x\rangle$$
 (7.17)

with the initial condition

$$a_0 = b_0 = \frac{1}{\sqrt{N}}.$$

Then the coefficients  $\{a_k, b_k\}$  for  $k \geq 1$  satisfy the recursion relations

$$a_k = \frac{N-2}{N} a_{k-1} + \frac{2(N-1)}{N} b_{k-1}, \tag{7.18}$$

$$b_k = -\frac{2}{N}a_{k-1} + \frac{N-2}{N}b_{k-1} \tag{7.19}$$

for k = 1, 2, ...

*Proof.* It is easy to see the recursion relations are satisfied for k=1Let  $G_f^{k-1}|\varphi_0\rangle = a_{k-1}|z\rangle + b_{k-1}\sum_{x\neq z}|x\rangle$ . Then  $G_f^k|\varphi_0\rangle = G_f\left(a_{k-1}|z\rangle + b_{k-1}\sum_{x\neq z}|x\rangle\right)$  $= (-I + 2|\varphi_0\rangle\langle\varphi_0|) \left(-a_{k-1}|z\rangle + b_{k-1}\sum_{x\neq z}|x\rangle\right)$  $=-b_{k-1}\sum |x\rangle + a_{k-1}|z\rangle + \frac{2}{\sqrt{N}}(N-1)b_{k-1}|\varphi_0\rangle - \frac{2a_{k-1}}{\sqrt{N}}|\varphi_0\rangle$  $= -b_{k-1} \sum_{n} |x\rangle + a_{k-1}|z\rangle + \frac{2}{N}(N-1)b_{k-1} \sum_{n} |x\rangle - \frac{2a_{k-1}}{N} \sum_{n} |x\rangle$  $= \left[ \frac{N-2}{N} a_{k-1} + \frac{2(N-1)}{N} b_{k-1} \right] |z\rangle + \left[ -\frac{2}{N} a_{k-1} + \frac{N-2}{N} b_{k-1} \right] \sum |x\rangle,$ 

and proposition is proved.

**PROPOSITION 7.3** The solutions of the recursion relations in Proposition 7.2 are explicitly given by

$$a_k = \sin[(2k+1)\theta], \quad b_k = \frac{1}{\sqrt{N-1}}\cos[(2k+1)\theta],$$
 (7.20)

for k = 0, 1, 2, ..., where

$$\sin \theta = \sqrt{\frac{1}{N}}, \quad \cos \theta = \sqrt{1 - \frac{1}{N}}.$$
 (7.21)

*Proof.* Let  $c_k = \sqrt{N-1}b_k$ . The recursion relations (7.18) and (7.19) are written in a matrix form,

$$\begin{pmatrix} a_k \\ c_k \end{pmatrix} = M \begin{pmatrix} a_{k-1} \\ c_{k-1} \end{pmatrix}, M = \begin{pmatrix} (N-2)/N & 2\sqrt{N-1}/N \\ -2\sqrt{N-1}/N & (N-2)/N \end{pmatrix} = \begin{pmatrix} \cos 2\theta & \sin 2\theta \\ -\sin 2\theta & \cos 2\theta \end{pmatrix}.$$

Note that M is a rotation matrix in  $\mathbb{R}^2$ , and its kth power is another rotation matrix corresponding to a rotation angle  $2k\theta$ . Thus the above recursion relation is easily solved to yield

$$\begin{pmatrix} a_k \\ c_k \end{pmatrix} = M^k \begin{pmatrix} a_0 \\ c_0 \end{pmatrix} = \begin{pmatrix} \cos 2k\theta & \sin 2k\theta \\ -\sin 2k\theta & \cos 2k\theta \end{pmatrix} \begin{pmatrix} \sin \theta \\ \cos \theta \end{pmatrix} = \begin{pmatrix} \sin[(2k+1)\theta] \\ \cos[(2k+1)\theta] \end{pmatrix}.$$

Replacing  $c_k$  by  $b_k$  proves the proposition.

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We have proved that the application of  $G_f$  k times on  $|\varphi_0\rangle$  results in the state

$$G_f^{k}|\varphi_0\rangle = \sin[(2k+1)\theta]|z\rangle + \frac{1}{\sqrt{N-1}}\cos[(2k+1)\theta]\sum_{x\neq z}|x\rangle. \tag{7.22}$$

Measurement of the state  $U_f^k|\varphi_0\rangle$  yields  $|z\rangle$  with the probability

$$P_{z,k} = \sin^2[(2k+1)\theta]. \tag{7.23}$$

**STEP 4** Our final task is to find the k that maximizes  $P_{z,k}$ . A rough estimate for the maximizing k is obtained by putting

$$(2k+1)\theta = \frac{\pi}{2} \to k = \frac{1}{2} \left( \frac{\pi}{2\theta} - 1 \right).$$
 (7.24)

**PROPOSITION 7.4** Let  $N \gg 1$  and let

$$m = \left\lfloor \frac{\pi}{4\theta} \right\rfloor, \tag{7.25}$$

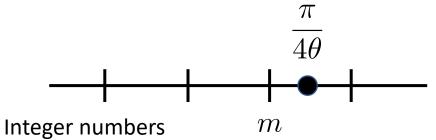
where  $\lfloor x \rfloor$  stands for the floor of x. The file we are searching for will be obtained in  $G_f^m |\varphi_0\rangle$  with the probability

$$P_{z,m} \ge 1 - \frac{1}{N} \tag{7.26}$$

and

$$m = O(\sqrt{N}). \tag{7.27}$$

This is the number of times we repeat the algorithm, which grows with the square root of N



*Proof.* Equation (7.25) leads to the inequality  $\pi/4\theta - 1 < m \le \pi/4\theta$ . Let us define  $\tilde{m}$  by

$$(2\tilde{m}+1)\theta = \frac{\pi}{2} \to \tilde{m} = \frac{\pi}{4\theta} - \frac{1}{2}.$$

Observe that m and  $\tilde{m}$  satisfy

$$|m - \tilde{m}| \le \frac{1}{2},\tag{7.28}$$

from which it follows that

$$|(2m+1)\theta - (2\tilde{m}+1)\theta| = \left| (2m+1)\theta - \frac{\pi}{2} \right| \le \theta.$$
 (7.29)

Considering that  $\theta \sim 1/\sqrt{N}$  is a small number when  $N \gg 1$  and  $\sin x$  is monotonically increasing in the neighborhood of x=0, we obtain

$$0 < \sin|(2m+1)\theta - \pi/2| < \sin\theta$$

or

$$\cos^{2}[(2m+1)\theta] \le \sin^{2}\theta = \frac{1}{N}.$$
 (7.30)

 $\rightarrow = \cos[(2m+1)\theta]$ 

Thus it has been shown that

$$P_{m,z} = \sin^2[(2m+1)\theta] = 1 - \cos^2[(2m+1)\theta] \ge 1 - \frac{1}{N}.$$
 (7.31)

It also follows from  $\theta > \sin \theta = 1/\sqrt{N}$  that

$$m = \left| \frac{\pi}{4\theta} \right| \le \frac{\pi}{4\theta} \le \frac{\pi}{4} \sqrt{N}. \tag{7.32}$$

It is important to note that this quantum algorithm takes only  $O(\sqrt{N})$  steps and this is much faster than the classical counterpart which requires O(N) steps.

Next we will show how to implement the gates