

# Quantum Nonlocality and Special Relativity

## No signalling theorem

Quantum nonlocality cannot be used to send information faster than the speed of light. Actually measurements cannot send information at all



A



B


We have two systems A and B, which in general share an entangled state  $\rho_{AB}$ . They are apart from each other.

Arbitrary measurements can be performed on each of them


# No signalling theorem

Alice performs a measurement of an observable  $\hat{A}$  with eigenprojectors  $P_n^A$ . The state at Bob's side changes to:

$$\rho_{AB} \rightarrow \rho'_{AB} = \sum_n \text{Tr}[(P_n^A \otimes I^B)\rho_{AB}] \frac{(P_n^A \otimes I^B)\rho_{AB}(P_n^A \otimes I^B)}{\text{Tr}[(P_n^A \otimes I^B)\rho_{AB}]} = \sum_n (P_n^A \otimes I^B)\rho_{AB}(P_n^A \otimes I^B)$$



Born rule



Von Neumann collapse

Then the average value of measurements Bob performs are given by:

# No signalling theorem

$$\begin{aligned} \langle O^B \rangle' &= \text{Tr}[(I^A \otimes O^B)\rho'_{AB}] = \sum_n \text{Tr}[(I^A \otimes O^B)(P_n^A \otimes I^B)\rho_{AB}(P_n^A \otimes I^B)] && \text{Cyclicity of trace} \\ &= \sum_n \text{Tr}[(I^A \otimes O^B)(P_n^A \otimes I^B)^2\rho_{AB}] && \text{Idempotent} \\ &= \sum_n \text{Tr}[(I^A \otimes O^B)(P_n^A \otimes I^B)\rho_{AB}] && \text{Linearity of trace + Projectors sum to 1} \\ &= \text{Tr}[(I^A \otimes O^B)\rho_{AB}] = \langle O^B \rangle \end{aligned}$$


The value Bob gets is the same before and after Alice's measurement

# No signalling theorem


Bob does not see any difference in the statistics of the outcomes of his measurements. There is no quantum operation (= unitary evolution or measurement) Alice can do, that allows her to send information to Bob.

If one looks at the reason why it is so, it ultimately rests on the fact that

$$\rho_{AB} \rightarrow \rho'_{AB} = \sum_n \text{Tr}[(P_n^A \otimes I^B)\rho_{AB}] \frac{(P_n^A \otimes I^B)\rho_{AB}(P_n^A \otimes I^B)}{\text{Tr}[(P_n^A \otimes I^B)\rho_{AB}]} = \sum_n (P_n^A \otimes I^B)\rho_{AB}(P_n^A \otimes I^B)$$



Born rule



Von Neumann collapse

In measurements, the Born rule and the von Neumann collapse are just the right recipes that avoid superluminal communication

# Teleportation

The teleportation protocol begins with a quantum state or qubit  $|\psi\rangle$ , in Alice's possession, that she wants to convey to Bob. This qubit can be written generally, in bra-ket notation, as:

$$|\psi\rangle_C = \alpha|0\rangle_C + \beta|1\rangle_C.$$

The subscript  $C$  above is used only to distinguish this state from  $A$  and  $B$ , below.

Next, the protocol requires that Alice and Bob share a maximally entangled state. This state is fixed in advance, by mutual agreement between Alice and Bob, and can be any one of the four Bell states shown. It does not matter which one.

$$\begin{aligned} |\Phi^+\rangle_{AB} &= \frac{1}{\sqrt{2}}(|0\rangle_A \otimes |0\rangle_B + |1\rangle_A \otimes |1\rangle_B), \\ |\Psi^+\rangle_{AB} &= \frac{1}{\sqrt{2}}(|0\rangle_A \otimes |1\rangle_B + |1\rangle_A \otimes |0\rangle_B), \\ |\Psi^-\rangle_{AB} &= \frac{1}{\sqrt{2}}(|0\rangle_A \otimes |1\rangle_B - |1\rangle_A \otimes |0\rangle_B), \\ |\Phi^-\rangle_{AB} &= \frac{1}{\sqrt{2}}(|0\rangle_A \otimes |0\rangle_B - |1\rangle_A \otimes |1\rangle_B). \end{aligned}$$

In the following, assume that Alice and Bob share the state  $|\Phi^+\rangle_{AB}$ . Alice obtains one of the particles in the pair, with the other going to Bob. (This is implemented by preparing the particles together and shooting them to Alice and Bob from a common source.) The subscripts  $A$  and  $B$  in the entangled state refer to Alice's or Bob's particle.

# Teleportation

At this point, Alice has two particles ( $C$ , the one she wants to teleport, and  $A$ , one of the entangled pair), and Bob has one particle,  $B$ . In the total system, the state of these three particles is given by

$$|\psi\rangle_C \otimes |\Phi^+\rangle_{AB} = (\alpha|0\rangle_C + \beta|1\rangle_C) \otimes \frac{1}{\sqrt{2}}(|0\rangle_A \otimes |0\rangle_B + |1\rangle_A \otimes |1\rangle_B).$$

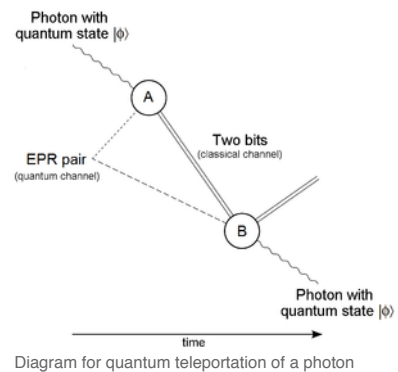
Alice will then make a local measurement in the Bell basis (i.e. the four Bell states) on the two particles in her possession. To make the result of her measurement clear, it is best to write the state of Alice's two qubits as superpositions of the Bell basis. This is done by using the following general identities, which are easily verified:

$$|0\rangle \otimes |0\rangle = \frac{1}{\sqrt{2}}(|\Phi^+\rangle + |\Phi^-\rangle),$$

$$|0\rangle \otimes |1\rangle = \frac{1}{\sqrt{2}}(|\Psi^+\rangle + |\Psi^-\rangle),$$

$$|1\rangle \otimes |0\rangle = \frac{1}{\sqrt{2}}(|\Psi^+\rangle - |\Psi^-\rangle),$$

$$|1\rangle \otimes |1\rangle = \frac{1}{\sqrt{2}}(|\Phi^+\rangle - |\Phi^-\rangle).$$



# Teleportation

One applies these identities with  $A$  and  $C$  subscripts. The total three particle state, of  $A$ ,  $B$  and  $C$  together, thus becomes the following four-term superposition:

$$\begin{aligned} |\psi\rangle_C \otimes |\Phi^+\rangle_{AB} = & \\ \frac{1}{2} \left[ |\Phi^+\rangle_{CA} \otimes (\alpha|0\rangle_B + \beta|1\rangle_B) + |\Phi^-\rangle_{CA} \otimes (\alpha|0\rangle_B - \beta|1\rangle_B) \right. & \text{[38]} \\ \left. + |\Psi^+\rangle_{CA} \otimes (\alpha|1\rangle_B + \beta|0\rangle_B) + |\Psi^-\rangle_{CA} \otimes (\alpha|1\rangle_B - \beta|0\rangle_B) \right]. & \end{aligned}$$

The above is just a change of basis on Alice's part of the system. No operation has been performed and the three particles are still in the same total state. The actual teleportation occurs when Alice measures her two qubits  $A, C$ , in the Bell basis

$$|\Phi^+\rangle_{CA}, |\Phi^-\rangle_{CA}, |\Psi^+\rangle_{CA}, |\Psi^-\rangle_{CA}.$$

Experimentally, this measurement may be achieved via a series of laser pulses directed at the two particles. Given the above expression, evidently the result of Alice's (local) measurement is that the three-particle state would collapse to one of the following four states (with equal probability of obtaining each):

- $|\Phi^+\rangle_{CA} \otimes (\alpha|0\rangle_B + \beta|1\rangle_B)$
- $|\Phi^-\rangle_{CA} \otimes (\alpha|0\rangle_B - \beta|1\rangle_B)$
- $|\Psi^+\rangle_{CA} \otimes (\alpha|1\rangle_B + \beta|0\rangle_B)$
- $|\Psi^-\rangle_{CA} \otimes (\alpha|1\rangle_B - \beta|0\rangle_B)$

# Teleportation

- $|\Phi^+\rangle_{CA} \otimes (\alpha|0\rangle_B + \beta|1\rangle_B)$
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- $|\Psi^+\rangle_{CA} \otimes (\alpha|1\rangle_B + \beta|0\rangle_B)$
- $|\Psi^-\rangle_{CA} \otimes (\alpha|1\rangle_B - \beta|0\rangle_B)$

Alice's two particles are now entangled to each other, in one of the four Bell states, and the entanglement originally shared between Alice's and Bob's particles is now broken. Bob's particle takes on one of the four superposition states shown above. Note how Bob's qubit is now in a state that resembles the state to be teleported. The four possible states for Bob's qubit are unitary images of the state to be teleported.

The result of Alice's Bell measurement tells her which of the above four states the system is in. She can now send her result to Bob through a classical channel. Two classical bits can communicate which of the four results she obtained.

# Teleportation

After Bob receives the message from Alice, he will know which of the four states his particle is in. Using this information, he performs a unitary operation on his particle to transform it to the desired state  $\alpha|0\rangle_B + \beta|1\rangle_B$ :

- If Alice indicates her result is  $|\Phi^+\rangle_{CA}$ , Bob knows his qubit is already in the desired state and does nothing. This amounts to the trivial unitary operation, the identity operator.
- If the message indicates  $|\Phi^-\rangle_{CA}$ , Bob would send his qubit through the unitary quantum gate given by the Pauli matrix

$$\sigma_3 = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$$

to recover the state.

- If Alice's message corresponds to  $|\Psi^+\rangle_{CA}$ , Bob applies the gate

$$\sigma_1 = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$

to his qubit.

- Finally, for the remaining case, the appropriate gate is given by

$$\sigma_3\sigma_1 = -\sigma_1\sigma_3 = i\sigma_2 = \begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix}.$$

Teleportation is thus achieved. The above-mentioned three gates correspond to rotations of  $\pi$  radians ( $180^\circ$ ) about appropriate axes (X, Y and Z) in the Bloch sphere picture of a qubit.

# Teleportation

Some remarks:

- After this operation, Bob's qubit will take on the state  $|\psi\rangle_B = \alpha|0\rangle_B + \beta|1\rangle_B$ , and Alice's qubit becomes an (undefined) part of an entangled state. Teleportation does not result in the copying of qubits, and hence is consistent with the no cloning theorem.
- There is no transfer of matter or energy involved. Alice's particle has not been physically moved to Bob; only its state has been transferred. The term "teleportation", coined by Bennett, Brassard, Crépeau, Jozsa, Peres and Wootters, reflects the indistinguishability of quantum mechanical particles.
- For every qubit teleported, Alice needs to send Bob two classical bits of information. These two classical bits do not carry complete information about the qubit being teleported. If an eavesdropper intercepts the two bits, she may know exactly what Bob needs to do in order to recover the desired state. However, this information is useless if she cannot interact with the entangled particle in Bob's possession.

Role of the collapse of the wave function in the process

Role of classical communication (teleportation protocol is subluminal)

## FLASH—A superluminal communicator based upon a new kind of measurement

As usual, there are Alice and Bob sharing a singlet state and perform distant spin measurements, as in a standard Bell setup.

The basis we will consider are  $|\uparrow\rangle$ ,  $|\downarrow\rangle$  and  $|+\rangle$ ,  $|-\rangle$ .

The FLASH protocol goes as follows.

1. Alice performs measurements in one of the two basis indicated above. Bob will receive the opposite state.

$\uparrow / \downarrow$  measurements. Alice obtains 50%  $|\uparrow\rangle$  and 50%  $|\downarrow\rangle$ . The states Bob receives are 50%  $|\downarrow\rangle$  and 50%  $|\uparrow\rangle$ .

$+/-$  measurements. Alice obtains 50%  $|+\rangle$  and 50%  $|-\rangle$ . The states Bob receives are 50%  $|-\rangle$  and 50%  $|+\rangle$ .

## FLASH—A superluminal communicator based upon a new kind of measurement

2. Bob amplifies the signal:

$$\begin{aligned} |\uparrow\rangle &\rightarrow |\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\rangle \\ |\downarrow\rangle &\rightarrow |\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\rangle \end{aligned}$$

in case Alice makes  $\uparrow / \downarrow$  measurements.

$$\begin{aligned} |+\rangle &\rightarrow |+++++\rangle \\ |-\rangle &\rightarrow |-----\rangle \end{aligned}$$

in case Alice makes +/- measurements.

## FLASH—A superluminal communicator based upon a new kind of measurement

3. Bob divides the states in two subsets. For half of them he performs a  $\uparrow / \downarrow$  measurement; for the other half he performs a +/- measurement.

In case Alice made  $\uparrow / \downarrow$  measurements:

$$\begin{aligned} |\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\rangle &\quad \rightarrow \quad \text{Half to } \uparrow / \downarrow : 100\% |\uparrow\rangle \text{ or } 100\% |\downarrow\rangle \\ |\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\rangle &\quad \rightarrow \quad \text{Half to } +/- : 50\% |+\rangle \text{ and } 50\% |-\rangle \end{aligned}$$

In case Alice made +/- measurements:

$$\begin{aligned} |+++++\rangle &\quad \rightarrow \quad \text{Half to } \uparrow / \downarrow : 50\% |\uparrow\rangle \text{ and } 50\% |\downarrow\rangle \\ |-----\rangle &\quad \rightarrow \quad \text{Half to } +/- : 100\% |+\rangle \text{ or } 100\% |-\rangle \end{aligned}$$

**Bob can understand what Alice measured. Faster than light**

# The No Cloning Theorem

The theorem says that it is not possible to clone an arbitrary quantum state.

Let us consider a unitary operator  $U$  such that:

$$U|\psi\rangle \otimes |s\rangle \rightarrow |\psi\rangle \otimes |\psi\rangle \quad \forall \psi \in \mathcal{H}$$

The state  $\psi$  has been duplicated. In particular we have, for two given states:

$$U|\psi_1\rangle \otimes |s\rangle \rightarrow |\psi_1\rangle \otimes |\psi_1\rangle$$

$$U|\psi_2\rangle \otimes |s\rangle \rightarrow |\psi_2\rangle \otimes |\psi_2\rangle$$

# The No Cloning Theorem

Then:

$$\langle \psi_1 | \psi_2 \rangle = \langle \psi_1 | \otimes \langle s | s \rangle \otimes \langle \psi_2 | = \langle \psi_1 | \otimes \langle s | U^\dagger U | s \rangle \otimes \langle \psi_2 | = \langle \psi_1 | \psi_2 \rangle^2$$

So we have the equation:  $x^2 = x$ , whose solution is  $x = 0, 1$ . This means that the two states  $\psi_1$  and  $\psi_2$  are either the same or orthogonal to each other.

The conclusion is that it is possible to copy orthogonal states, but it is not possible to copy arbitrary non-orthogonal states. This violates the unitarity of quantum evolutions.





# Why FLASH does not work

The suppose Alice prepared in the  $\uparrow / \downarrow$  so that Bob's machine generates

$$|\uparrow\rangle \rightarrow |\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\rangle$$

$$|\downarrow\rangle \rightarrow |\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\rangle$$

Bob divides the set un two subsets. For half of them he performs a  $\uparrow / \downarrow$  measurement; for the other half he performs a +/- measurement.

$$|\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\rangle$$



Half to  $\uparrow / \downarrow$  : 100%  $|\uparrow\rangle$  or 100%  $|\downarrow\rangle$

$$|\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\rangle$$

Half to +/- : 50%  $|+\rangle$  and 50%  $|-\rangle$

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$$|\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\rangle$$



Half to  $\uparrow / \downarrow$  : 100%  $|\uparrow\rangle$  or 100%  $|\downarrow\rangle$

$$|\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\rangle$$

Half to +/- : 50%  $|+\rangle$  and 50%  $|-\rangle$

# Why FLASH does not work

The suppose Alice prepared in the  $+/-$  so that Bob's machine generates

$$|+\rangle = \frac{1}{\sqrt{2}}[|\uparrow\rangle + |\downarrow\rangle] \rightarrow \frac{1}{\sqrt{2}}[|\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\rangle + |\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\rangle]$$

$$|-\rangle = \frac{1}{\sqrt{2}}[|\uparrow\rangle - |\downarrow\rangle] \rightarrow \frac{1}{\sqrt{2}}[|\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\rangle - |\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\rangle]$$

Bob divides the set un two subsets. For half of them he performs a  $\uparrow / \downarrow$  measurement; for the other half he performs a  $+/-$  measurement. It is evident that as soon as he performs a  $\uparrow / \downarrow$  measurement on the first system, the whole state collapses to

$$50\% \quad |\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\uparrow\rangle \qquad 50\% \quad |\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\rangle$$

Therefore the same statistics as in the previous case is recovered

# Why FLASH does not work

Exercise: Repeat the calculation assuming that Bob's machine does the following

$$|+\rangle \rightarrow |++++++\rangle$$

$$|-\rangle \rightarrow |-----\rangle$$

# Cryptography

Classical cryptography can be divided into two major branches; **secret or symmetric key cryptography** and **public key cryptography**, which is also known as **asymmetric cryptography**.

**Secret key cryptography** represents the most traditional form of cryptography in which two parties both encrypt and decrypt their messages using the **same shared secret key**. While some secret key schemes, such as one-time pads, are **perfectly secure** against an attacker with arbitrary computational power, they have the **major practical disadvantage** that before two parties can communicate securely they must somehow establish a secret key.

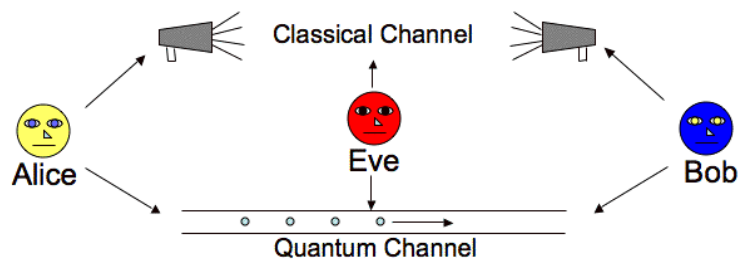
In order to establish a secret key over an insecure channel, **key distribution schemes** based on public key cryptography, such as Diffie-Hellman, are typically employed.

# Cryptography

In contrast to secret key cryptography, a shared secret key does not need to be established prior to communication in **public key cryptography**. Instead **each party has a private key**, which remains secret, **and a public key**, which they may distribute freely. If one party, say Alice, wants to send a message to another party, Bob, **she would encrypt her message with Bob's public key after which only Bob could decrypt the message using his private key**. While there is no need for key exchange, the security of public key cryptography algorithms are currently all based on the **unproven assumption of the difficulty of certain problems** such as integer factorization or the discrete logarithm problem. This means that **public key cryptography algorithms are potentially vulnerable** to improvements in computational power or the discovery of efficient algorithms to solve their underlying problems. Indeed algorithms have already been proposed to perform both integer factorization and solve the discrete logarithm problem in polynomial time on a quantum computer

# QKD

The basic model for Quantum Key Distribution (QKD) protocols involves two parties, referred to as Alice and Bob, wishing to exchange a key both with access to a classical public communication channel and a quantum communication channel. This is shown in the figure. An eavesdropper, called Eve, is assumed to have access to both channels and no assumptions are made about the resources at her disposal. With this basic model established, we describe in layman's terms the necessary quantum principles needed to understand the QKD protocols.

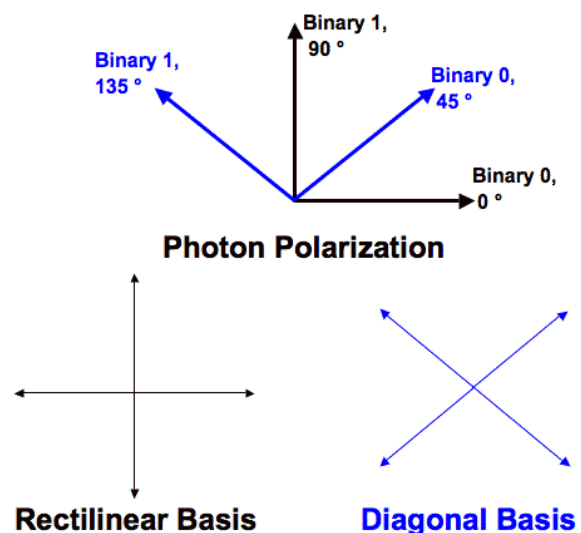


# QKD - BB84

The Figure shows how a bit can be encoded in the polarization state of a photon in BB84.

We define a binary 0 as a polarization of 0 degrees in the rectilinear bases or 45 degrees in the diagonal bases. Similarly a binary 1 can be 90 degrees in the rectilinear bases or 135 in diagonal bases.

Thus a bit can be represented by polarizing the photon in either one of two bases.



# QKD - BB84

1. Alice begins by choosing a random string of bits.
2. For each bit, Alice will randomly choose a basis, rectilinear or diagonal, by which to encode the bit.
3. She will transmit a photon for each bit with the corresponding polarization, as just described, to Bob.
4. For every photon Bob receives, he will measure the photon's polarization by a randomly chosen basis. If, for a particular photon, Bob chose the same basis as Alice, then in principle, Bob should measure the same polarization and thus he can correctly infer the bit that Alice intended to send. If he chose the wrong basis, his result, and thus the bit he reads, will be random.

# QKD - BB84

5. Bob will notify Alice over any insecure channel what basis he used to measure each photon. Alice will report back to Bob whether he chose the correct basis for each photon.
6. Alice and Bob will discard the bits corresponding to the photons which Bob measured with a different basis. On the average, only half of the photons have to be disregarded. Provided no errors occurred or no one manipulated the photons, Bob and Alice should now both have an identical string of bits which is called a sifted key.

# QKD - BB84

Alice's bit	0	1	1	0	1	0	0	1
Alice's basis	+	+	X	+	X	X	X	+
Alice's polarization	↑	→	↖	↑	↖	↗	↗	→
Bob's basis	+	X	X	X	+	X	+	+
Bob's measurement	↑	↗	↖	↗	→	↗	→	→
Public discussion								
Shared Secret key	0		1			0		1

# QKD - BB84 - Eve

Assume that Eve tries to intercept the basis. She will do that by measuring the photon's state. In this way, she will introduce an error with probability 25%

A sends bit  
0 in basis +

The best Eve can do is:  
50% +: outcome 0  
50% x: outcome 0 or 1

Bob measures in basis +  
→ Outcome 0  
→ 50 % 0 and 50% 1

So 25% of the times Bob gets a different result from Alice, in spite they have measured in the same basis.

# QKD - BB84 - Eve

If now Alice and Bob publicly compare  $n$  bits (then disregarding them as key bits, since they are no longer secret) the probability of finding a disagreement is

$$\mathbb{P}_D^{(n)} = 1 - (3/4)^n \quad (\text{where } 3/4 \text{ is the probability that they all match})$$

Then for  $n = 72$ :  $\mathbb{P}_D^{(n)} = 0,999999999$  (nine 9)

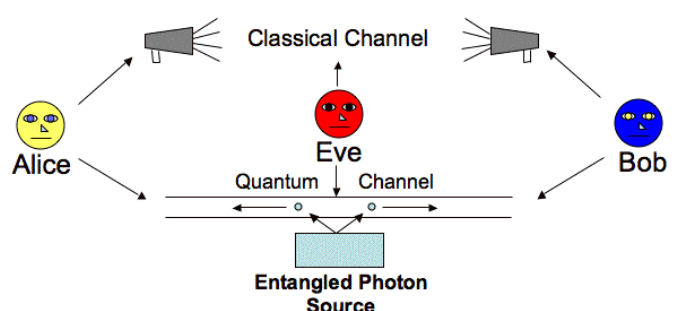
Almost immediately Alice and Bob realize that Eve tried to copy the key and abort the operation of key distribution.

**In general, if there are too many errors when comparing the bits, the quantum channel is considered insecure and the protocol is aborted.**

# QKD - E91

Eckert describes a channel where there is a single source that emits pairs of entangled particles, which could be polarized photons. The particles are separated and Alice and Bob each receive one particle from each pair as shown in figure 5. Alice and Bob would each choose a random bases on which to measure their received particles. As in BB84, they would discuss in the clear which bases they used for their measurements. For each measurement where Alice and Bob used the same bases, they should expect opposite results due to the principle of quantum entanglement as described earlier.

This means that if Alice and Bob both interpret their measurements as bits as before, they each have a bit string which is the binary complement of the other. Either party could invert their key and they would thus share a secret key.





# QKD - E91

The presence of an eavesdropper can be detected by examining the photons for which Alice and Bob chose different bases for measurement. Alice and Bob can measure these photons in a third basis and discuss their results. With this information they can test Bell's Inequality which should not hold for entangled particles. If the inequality does hold, it would indicate that the photons were not truly entangled and thus there may be an eavesdropper present.