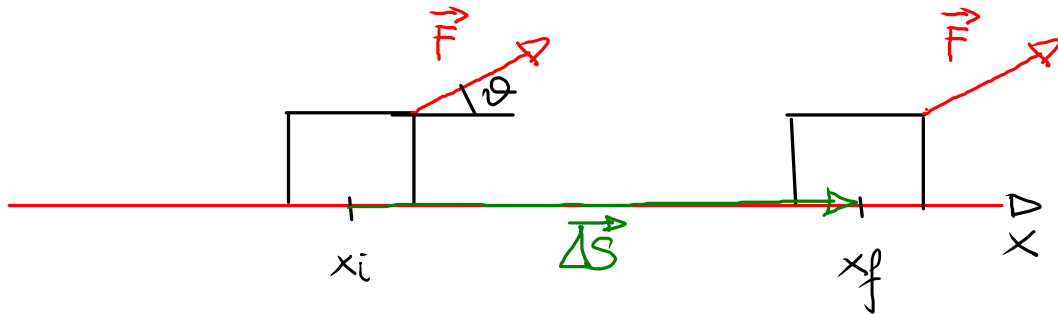


LAVORO di \vec{F} su uno $\Delta\vec{S}$



livello base : $\Delta\vec{S}$ è rettilineo

\vec{F} è costante

θ è l'angolo tra $\vec{\Delta S}$ e \vec{F}

Definisco $\mathcal{L} = \vec{F} \cdot \vec{\Delta S} = |\vec{F}| \cdot |\vec{\Delta S}| \cdot \cos\theta$

{	+	$0 \leq \theta < \frac{\pi}{2}$
	0	per $\theta = \frac{\pi}{2}$
	-	$\frac{\pi}{2} < \theta \leq \pi$

$$[\mathcal{L}] = [\vec{F}] \cdot [\vec{\Delta S}] = \text{N} \cdot \text{m} = \text{J} \quad (\text{Joule})$$

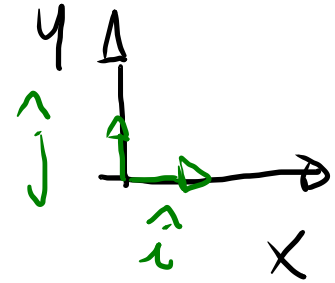
L'UOMO CON LA VALIGIAA

$$\vec{h} = h\hat{j}$$

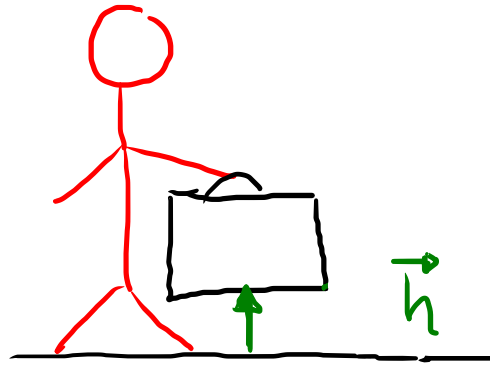
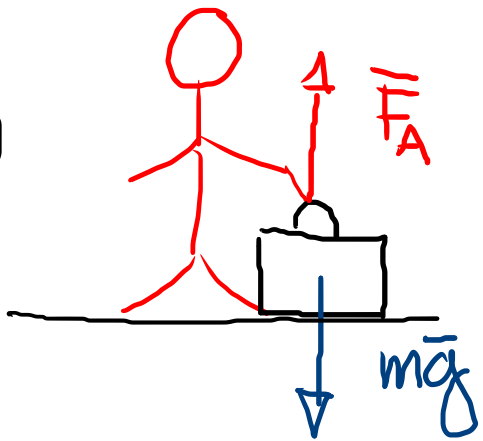
$$\sum \vec{F} = \vec{F}_A + m\vec{g} = 0$$

$$L_u = \vec{F}_A \cdot \vec{h} = m\vec{g} \cdot \vec{h} = mgh$$

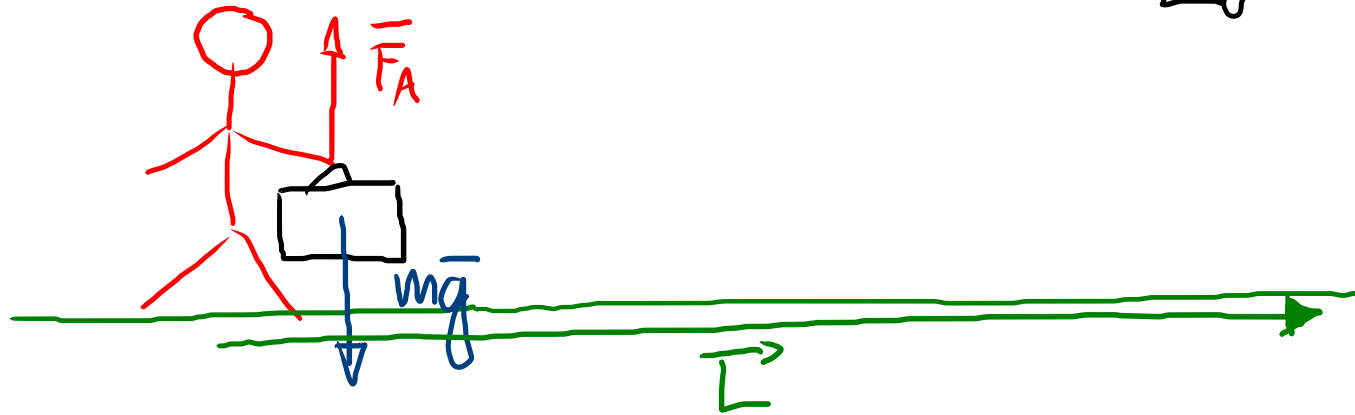
$$L_u = \vec{F}_A \cdot \vec{L} = 0$$



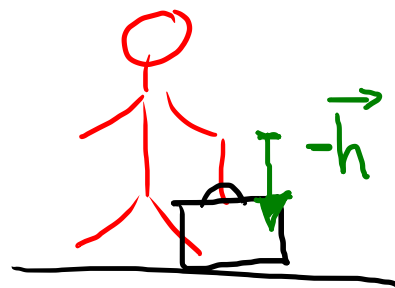
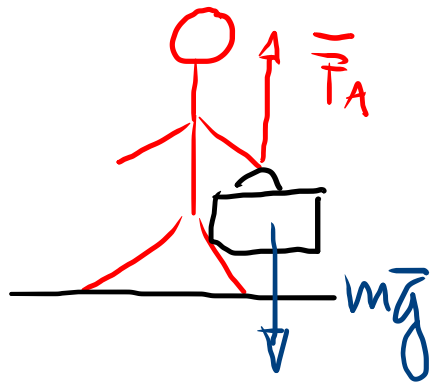
①



②



③



$$L_G = m\vec{g} \cdot (-\vec{h}) = -m\vec{g} \cdot \vec{h}$$

$$= -m|\vec{g}||\vec{h}| \cancel{\sin \pi} \cos \pi = mgh$$

LAVORO DI UNA \bar{F} VARIABILE (caso 1D)

1D lungo asse x : $x_i \rightarrow x_f$

F è variabile: $F(x) \rightarrow F_x$

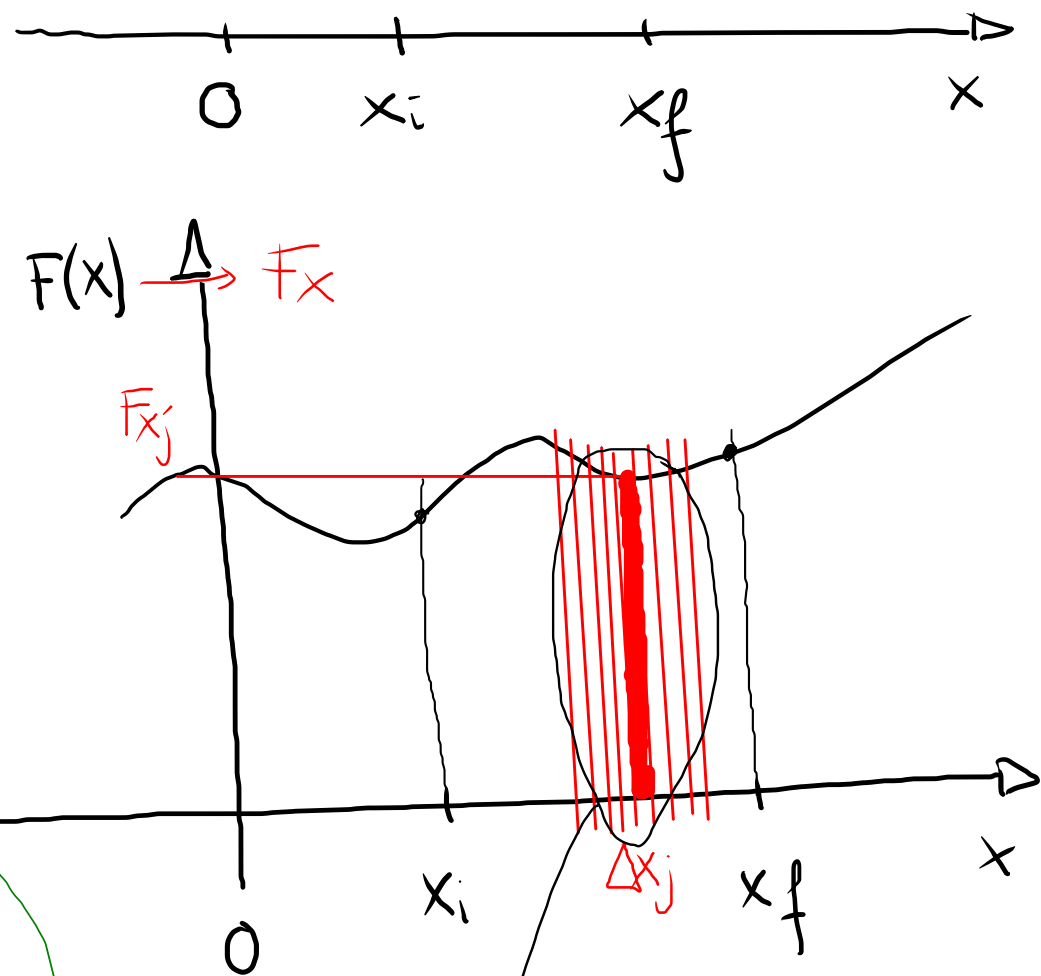
$$\mathcal{L}_j = F_{x_j} \cdot \Delta x_j$$

$$\mathcal{L} = \sum_j \mathcal{L}_j = \sum_j \underbrace{F_{x_j} \cdot \Delta x_j}_{A_j}$$

$= \sum_j A_j = \text{area sotto la curva}$

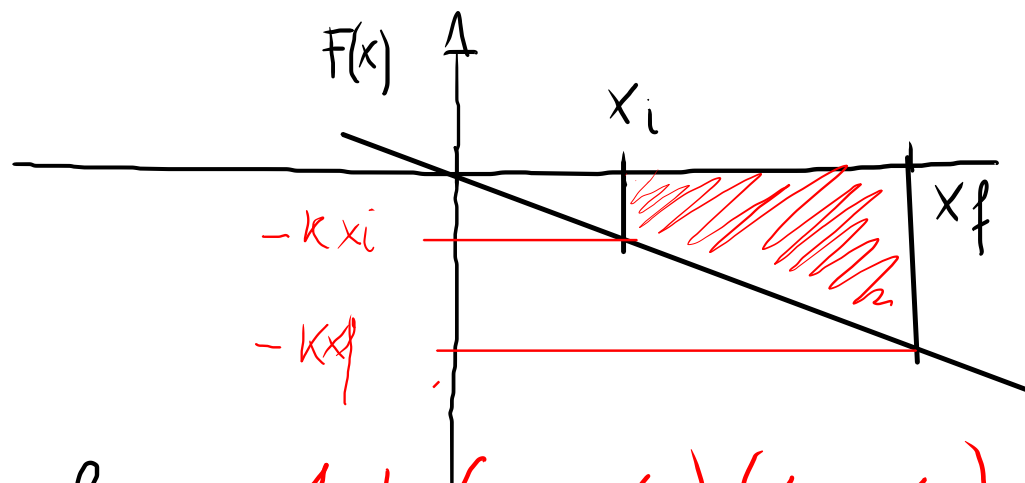
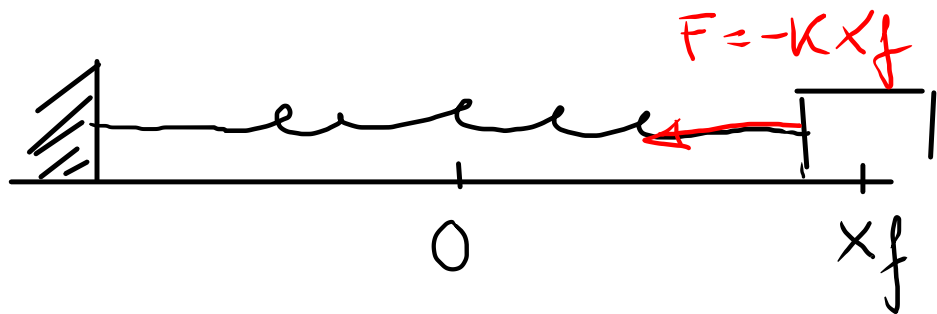
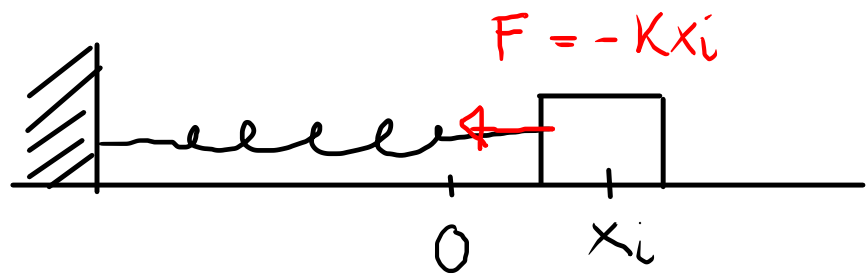
$$\mathcal{L} = \lim_{\Delta x_j \rightarrow 0} \sum_j \mathcal{L}_j = \int_{x_i}^{x_f} F(x) dx$$

$= \text{area sotto la curva}$



$$A_j = F_{x_j} \cdot \Delta x_j = \mathcal{L}_j$$

ESEMPIO: $\vec{F} = -K\vec{x}$ (faza elastica)



$$\begin{aligned} \mathcal{L} &= -\frac{1}{2}k(x_i + x_f)(x_f - x_i) \\ &= \frac{1}{2}k(x_i^2 - x_f^2) \end{aligned}$$

$$\begin{aligned} \mathcal{L} &= \int_{x_i}^{x_f} F(x) dx = \int_{x_i}^{x_f} -Kx dx \\ &= -K \int_{x_i}^{x_f} x dx = -K \left. \frac{1}{2}x^2 \right|_{x_i}^{x_f} \\ &= -\frac{1}{2}k(x_f^2 - x_i^2) = \frac{1}{2}k(x_i^2 - x_f^2) \end{aligned}$$

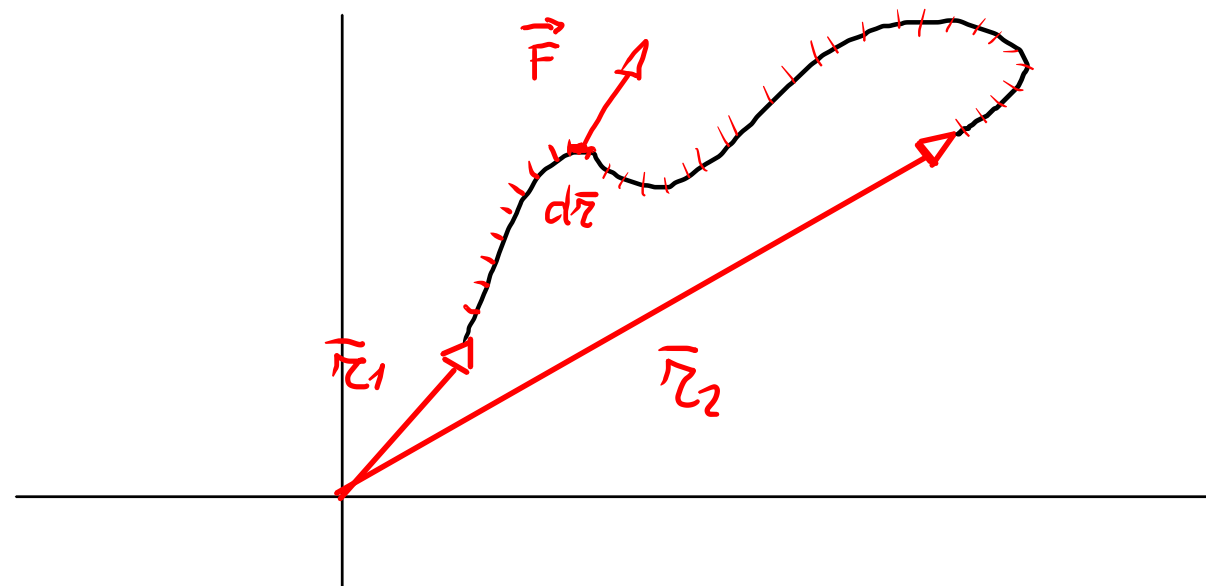
DEFINIZIONE GENERALE di L

$$L = \int_{\vec{r}_1}^{\vec{r}_2} \vec{F}(\vec{r}) \cdot d\vec{r}$$

↑
prodotto scalare!

integrale di linea!

ovvero va calcolato lungo la particolare
traiettoria del moto



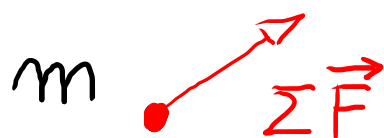
Se ho molte Forze, posso valutare $\Sigma \vec{F}$ e

$$L = \int_{\vec{r}_1}^{\vec{r}_2} \Sigma \vec{F} \cdot d\vec{r}$$

ENERGIA

Energia cinetica: m \vec{v} $K = \frac{1}{2} m v^2$

TEOREMA LAVORO - ENERGIA



$$\mathcal{L} = \Delta K = K_f - K_i$$

$$\mathcal{L} \text{ di } \Sigma \vec{F} : \mathcal{L} = \frac{1}{2} m v_f^2 - \frac{1}{2} m v_i^2 = \frac{1}{2} m (v_f^2 - v_i^2)$$

(\mathcal{L} è anche il lavoro totale, ovvero la somma dei lavori delle singole forze prese singolarmente)

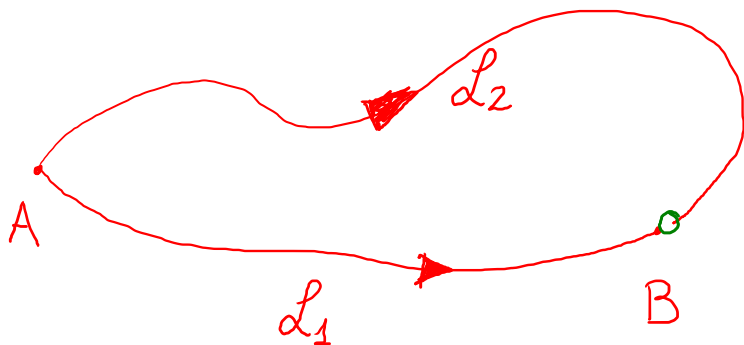
Dimostrazione di $\mathcal{L} = \Delta K$

II principio dinamica

$$\begin{aligned} \mathcal{L} &= \int_{x_i}^{x_f} \underbrace{\sum F(x)}_{ma} dx = \int_{x_i}^{x_f} m a dx = \int_{x_i}^{x_f} m \frac{dv}{dt} dx \\ &= \int_i^f m dv \frac{dx}{dt} = \int_{v_i}^{v_f} m v dv = m \int_{v_i}^{v_f} v dv = \\ &= m \left. \frac{1}{2} v^2 \right|_{v_i}^{v_f} = \frac{1}{2} m (v_f^2 - v_i^2) = K_f - K_i = \Delta K \end{aligned}$$

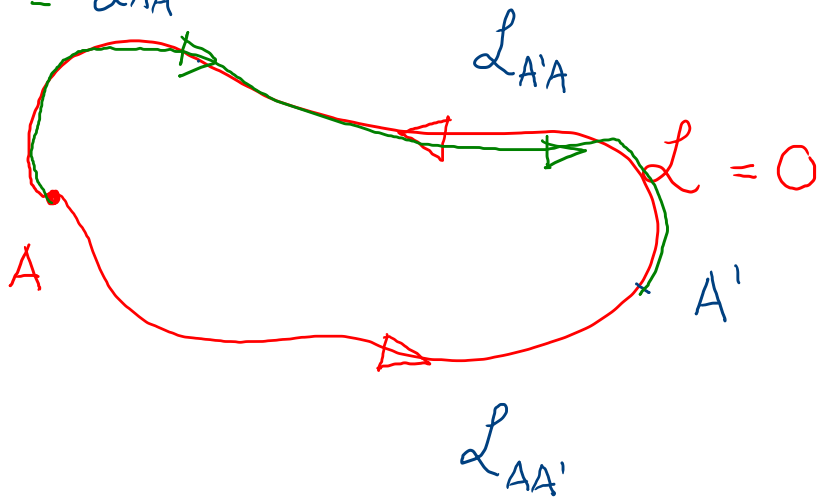
$$\vec{F} = m\vec{a} \iff \mathcal{L} = \Delta K$$

FORZE CONSERVATIVE



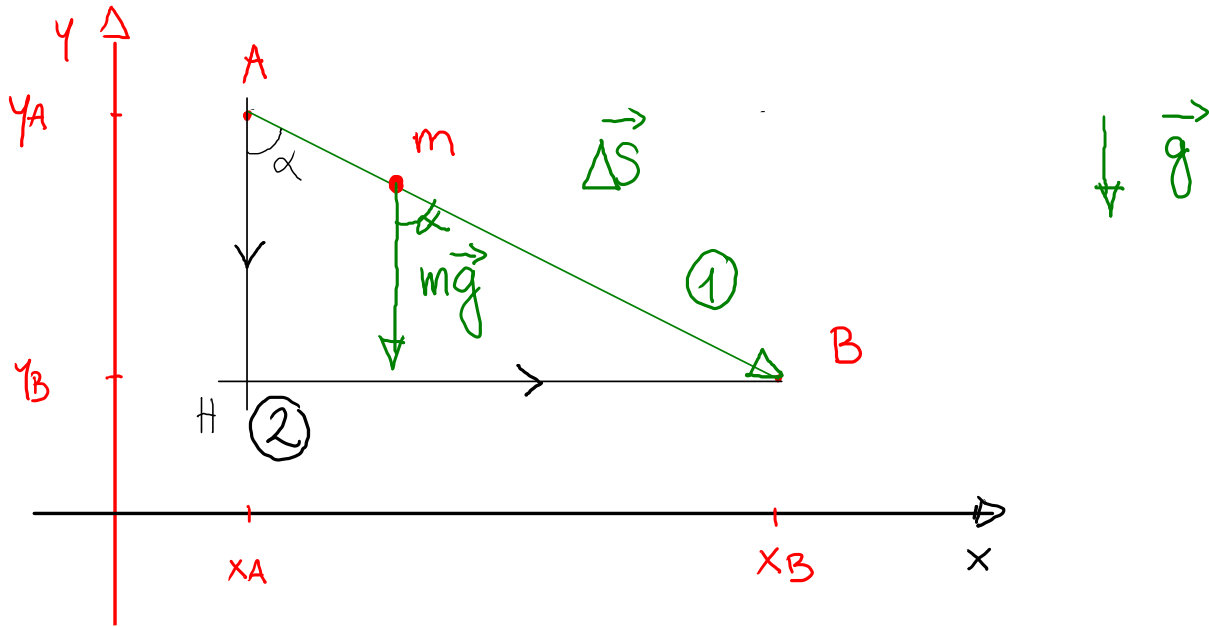
$$L_1 = L_2 = L_{AB}$$

$$L'_{AA'} = -L_{A'A}$$



$$\begin{aligned} L &= L_{AA'} + L_{A'A} \\ &= L_{AA'} - L'_{AA'} \\ &= 0 \end{aligned}$$

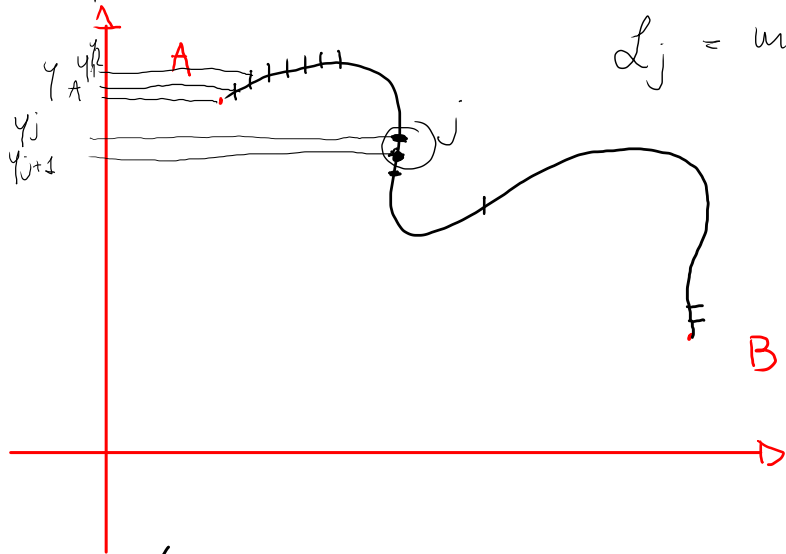
ESEMPIO: FORZA PESO



$$\textcircled{1} \quad \mathcal{L} = \vec{F} \cdot \Delta \vec{S} = |m\vec{g}| \cdot |\Delta \vec{S}| \cos \alpha = mg \Delta S \cos \alpha = mg (y_A - y_B)$$

$$\textcircled{2} \quad \mathcal{L} \stackrel{!}{=} \mathcal{L}_{AH} + \mathcal{L}_{HB} = mg (y_A - y_B) + 0 = mg (y_A - y_B)$$

Percorso qualsiasi... (approfondimento, non in programma)

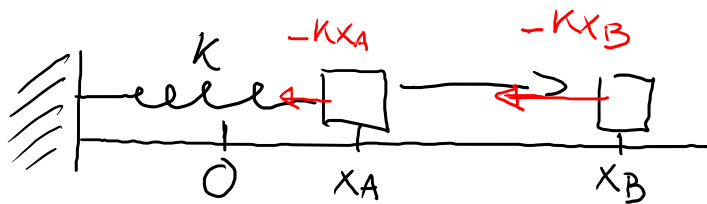


$$\begin{aligned} \mathcal{L} &= \sum_{j=1}^N \mathcal{L}_j = mg(y_A - y_2) + mg(y_1 - y_2) + \dots \\ &+ \dots mg(y_j - y_{j+1}) + mg(y_{j+1} - y_{j+2}) \\ &+ mg(y_n - y_B) = mg(y_A - y_B) \end{aligned}$$

ULTERIORE ESEMPIO: Forza elastica

$$\vec{F} = -k\vec{x}$$

Però abbiamo calcolato il lavoro della forza elastica per andare da x_A a x_B :



$$d = \frac{1}{2} k (x_A^2 - x_B^2)$$

Il lavoro dipende solo da x_A e x_B

\Rightarrow La forza elastica è conservativa!

ENERGIA POTENZIALE

Per ogni forza CONSERVATIVA posso definire

$$U(\vec{r}) \quad \text{tale che} \quad \mathcal{L}_{AB} = U(\vec{r}_A) - U(\vec{r}_B)$$
$$= U_A - U_B = -\Delta U$$

Esempi precedenti

Forza peso : $\mathcal{L}_{AB} = mg(y_A - y_B)$

$$U_g = mgy$$

gravitazionale

$$\mathcal{L}_{AB} = U_A - U_B$$
$$= mgy_A - mgy_B$$

Forza elastica $\mathcal{L}_{AB} = \frac{1}{2}k(x_A^2 - x_B^2)$

$$U_e = \frac{1}{2}kx^2$$

elastica

$$\mathcal{L}_{AB} = U_A - U_B$$
$$= \frac{1}{2}kx_A^2 - \frac{1}{2}kx_B^2$$

ENERGIA MECCANICA

$$E_{mecc} = K + U$$

Sistema conservativo: il \mathcal{L} è fatto da forze conservative

① $\mathcal{L} = \Delta K = K_B - K_A$ (vale sempre)

② $\mathcal{L} = -\Delta U = U_A - U_B$ (solo per sistemi conservativi)

①② $0 = \Delta K + \Delta U = K_B + U_B - (K_A + U_A)$

$$\Delta K + \Delta U = 0$$

$$\Delta(K + U) = 0$$

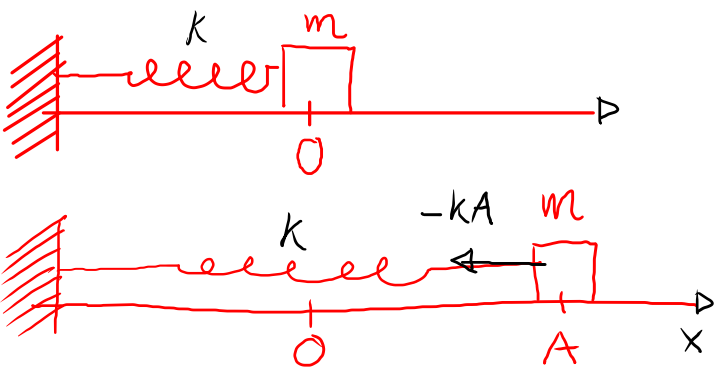
$$\Delta E_{mecc} = 0$$

$$K_B + U_B - (K_A + U_A) = 0$$

$$E_{mecc B} - E_{mecc A} = 0$$

$$\Delta E_{mecc} = 0$$

L' E_{mecc} si conserva.



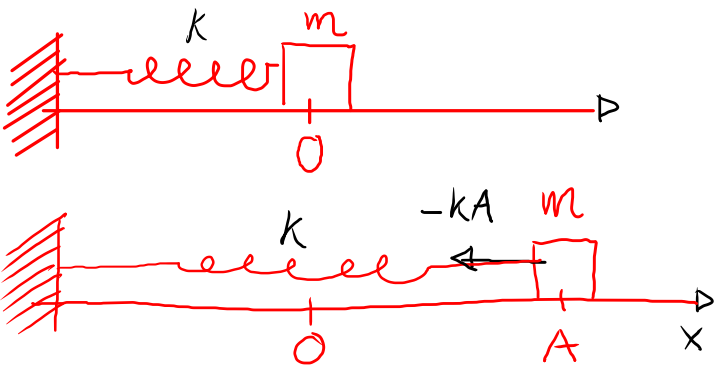
cinematica

$$\begin{cases}
 * & x(t) = A \cos(\omega t) & \omega = \frac{2\pi}{T} \\
 ** & v(t) = -A\omega \sin(\omega t) \\
 & a(t) = -A\omega^2 \cos(\omega t) = -\omega^2 x(t)
 \end{cases}$$

dinamica : $\omega = \sqrt{\frac{k}{m}}$ $\omega^2 = \frac{k}{m}$

energia : $U_e = \frac{1}{2} k x^2$
potenziale elastica

$$\begin{aligned}
 E_{mecc} = K + U_e &= \frac{1}{2} m v^2 + \frac{1}{2} k x^2 \\
 &= \frac{1}{2} m \left[-A\omega \sin(\omega t) \right]^2 + \frac{1}{2} k \left[A \cos(\omega t) \right]^2 \\
 &= \frac{1}{2} m \underbrace{A^2 \omega^2}_{***} \sin^2(\omega t) + \frac{1}{2} k A^2 \cos^2(\omega t) \\
 &= \frac{1}{2} k A^2 \sin^2(\omega t) + \frac{1}{2} k A^2 \cos^2(\omega t) \\
 &= \frac{1}{2} k A^2 \left[\underbrace{\sin^2(\omega t) + \cos^2(\omega t)}_1 \right] = \frac{1}{2} k A^2
 \end{aligned}$$



$$E_{mecc} = K + U_e = \frac{1}{2} k A^2$$

$$\begin{aligned}
 x=0 &\Rightarrow U_e = 0 \\
 K &= \frac{1}{2} k A^2 \\
 \frac{1}{2} m v^2 &= \frac{1}{2} k A^2 \\
 v^2 &= \frac{k}{m} A^2 = \omega^2 A^2 \\
 v &= \pm \omega A
 \end{aligned}$$

$$\begin{aligned}
 x \neq 0 \\
 \frac{1}{2} k A^2 &= \frac{1}{2} m v^2 + \frac{1}{2} k x^2 \\
 \omega^2 A^2 &= v^2 + \omega^2 x^2
 \end{aligned}$$

cinematica

$$\begin{aligned}
 * &\left\{ \begin{aligned} x(t) &= A \cos(\omega t) & \omega &= \frac{2\pi}{T} \\ v(t) &= -A\omega \sin(\omega t) \\ a(t) &= -A\omega^2 \cos(\omega t) = -\omega^2 x(t) \end{aligned} \right. \\
 ** &
 \end{aligned}$$

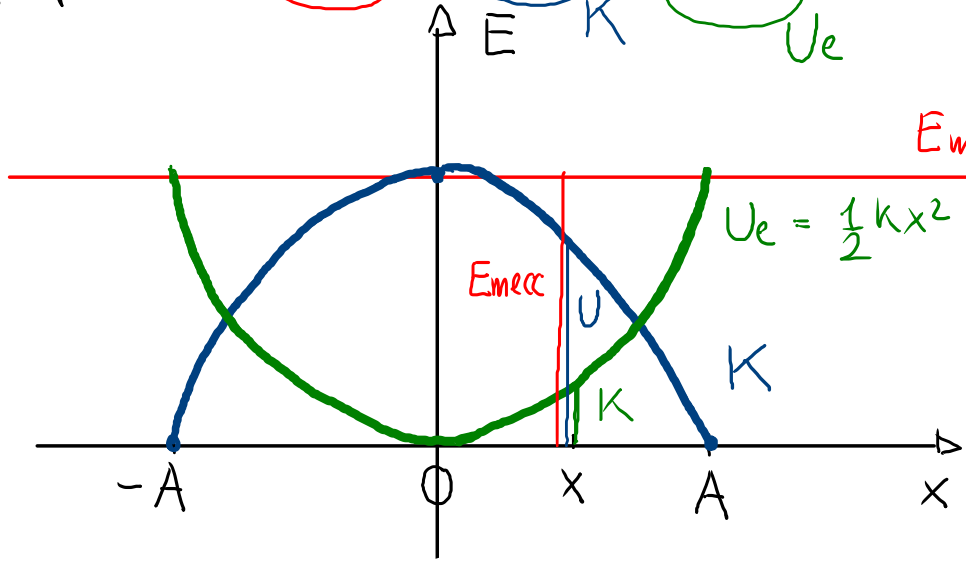
$$\begin{aligned}
 \text{dinamica: } \omega &= \sqrt{\frac{k}{m}} & \omega^2 &= \frac{k}{m} \\
 \text{energia potenziale elastica: } U_e &= \frac{1}{2} k x^2
 \end{aligned}$$

$$\begin{aligned}
 v^2 &= \omega^2 (A^2 - x^2) \\
 v &= \pm \omega \sqrt{A^2 - x^2}
 \end{aligned}$$

In forma grafica:

$$E_{\text{mecc}} = \underbrace{\frac{1}{2}mv^2}_E + \underbrace{\frac{1}{2}kx^2}_{U_e} = \frac{1}{2}kA^2$$

$$K = E_{\text{mecc}} - U_e$$



$$E_{\text{mecc}} = \frac{1}{2}kA^2$$

$$U_e = \frac{1}{2}kx^2$$

• $\mathcal{L} = \Delta K$ Vale SEMPRE

$\mathcal{L} = -\Delta U$ vale solo x sistemi conservativi

In un sistema conservativo valgono entrambe

$$\begin{cases} \mathcal{L} = \Delta K \\ \mathcal{L} = -\Delta U \end{cases} \Rightarrow \Delta K = -\Delta U \quad \Delta K + \Delta U = 0$$

Definiamo ora $K + U = E_{mecc}$

$$\Delta K + \Delta U = 0$$

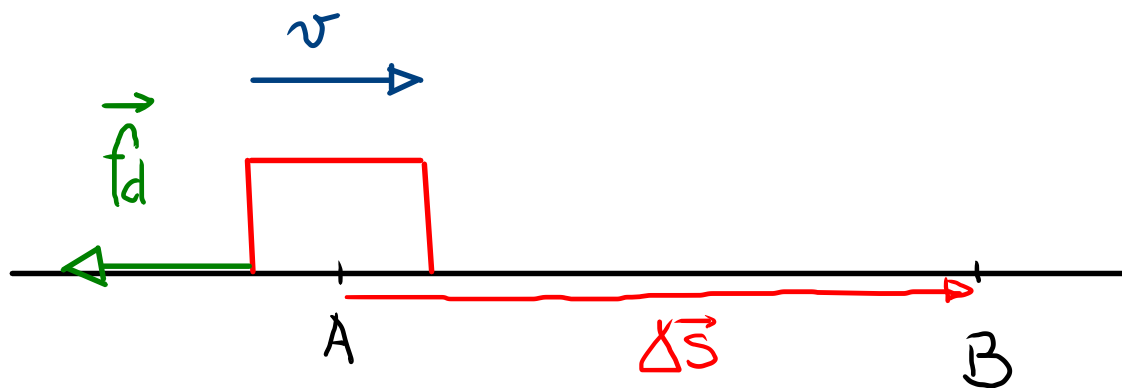
$$\Delta(K + U) = 0$$

$$\Delta E_{mecc} = 0$$

E_{mecc} si conserva! ▽

FORZE DISSIPATIVE (ovvero non conservative)

Es: attrito dinamico

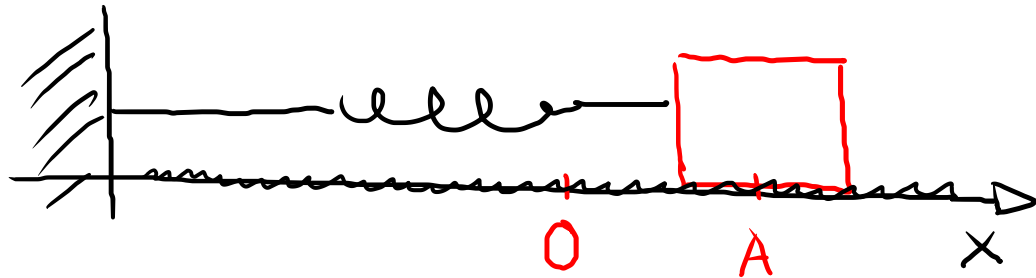


$$\begin{aligned} \mathcal{L}_d &= \vec{f}_d \cdot \Delta \vec{S} = f_d \Delta S \cdot \cos \pi \\ &= -f_d \Delta S \end{aligned}$$

↑ negativo!
↑ dipende dal percorso!

SISTEMA CON FORZE CONSERVATIVE E DISSIPATIVE

Es.



piano liscio

⇒ conservativo

$$U = \frac{1}{2} K x^2$$

piano ruvido ⇒ c'è attrito
non conservativo

$$\mathcal{L} = \Delta K$$

↑ somma dei lavori fatti da ciascuna forza
assumiamo: una forza conservativa \mathcal{L}_C
una forza dissipativa \mathcal{L}_D

$$\mathcal{L} = \mathcal{L}_C + \mathcal{L}_D$$

$$\mathcal{L}_C + \mathcal{L}_D = \Delta K$$

$$\mathcal{L}_C = -\Delta U$$

$$-\Delta U + \mathcal{L}_D = \Delta K$$

$$\mathcal{L}_D = \Delta K + \Delta U = \Delta E_{mecc} *$$

$$\Delta E_{mecc} = \mathcal{L}_D < 0$$

E_{mecc} NON si conserva ma diminuisce ($\Delta E_{mecc} < 0$)

SISTEMI ISOLATI

Non scambia né materia né energia con l'ambiente che lo circonda

Definisco E_{int} , energia interna del sistema.

E_{TOT} , energia totale del sistema. \leftarrow si conserva!

$$E_{TOT} = E_{mecc} + E_{int}$$

$$\Delta E_{TOT} = 0 = \Delta E_{mecc} + \Delta E_{int}$$

$$\Delta E_{int} = - \underbrace{\Delta E_{mecc}}_{< 0} > 0$$

Quindi da *

$$\Delta K + \Delta U = \Delta E_{mecc}$$

$$\Delta K + \Delta U - \Delta E_{mecc} = 0$$

$$\Delta K + \Delta U + \Delta E_{int} = 0$$

$$\underline{\Delta E_{TOT} = 0}$$

\rightarrow si conserva!

POTENZA

$$P_m = \frac{L}{\Delta t} \quad \frac{J}{s} = \text{Watt} \quad (\text{potenza media})$$

$$P = \lim_{\Delta t \rightarrow 0} \frac{L}{\Delta t} \quad (\text{potenza istantanea})$$

$$L = \vec{F} \cdot \Delta \vec{x}$$

$$= \lim_{\Delta t \rightarrow 0} \frac{\vec{F} \cdot \Delta \vec{x}}{\Delta t} = \lim_{\Delta t \rightarrow 0} \vec{F} \cdot \frac{\Delta \vec{x}}{\Delta t} = \vec{F} \cdot \vec{v}$$

$$P = \vec{F} \cdot \vec{v}$$

Nota: $1 \text{ kWh} = \underbrace{1 \text{ kW}}_{\frac{\text{kJ}}{\text{s}}} \cdot 1 \text{ h} = 3600 \text{ s} = 3600 \text{ kJ}$

RENDIMENTO

$$\eta = \frac{L}{E}$$

← lavoro compiuto dalla macchina
← energia necessaria a far funzionare

in percentuale

$$\eta = \frac{L}{E} \cdot 100 \quad (\%)$$

(fine programma obbligatoria meccanica)

seguono approfondimenti