

General form of LWE

$$
\frac{\partial^2 \psi(x,t)}{\partial t^2} = v^2 \frac{\partial^2 \psi(x,t)}{\partial x^2}
$$

WAVE: organized **propagating imbalance**, satisfying differential equations of motion

$$
\frac{\partial^2 y(x,t)}{\partial x^2} - \frac{1}{c^2} \frac{\partial^2 y(x,t)}{\partial t^2} = 0
$$

and if it has separable solutions:

$$
y(x,t) = X(x)T(t)
$$

$$
\frac{d^{2}X(x)}{dx^{2}} + k^{2}X(x) = 0
$$

\n
$$
X(x) = A\cos(kx) + B\sin(kx)
$$

\n
$$
T'(t) = C\cos(\omega t) + D\sin(\omega t)
$$

\n
$$
T(t) = C\cos(\omega t) + D\sin(\omega t)
$$

To be determined by **initial** and **boundary** conditions

Consider a string of length L and fixed at both ends

The string has a number of natural patterns of vibration called **NORMAL MODES**

Each normal mode has a characteristic frequency which we can easily calculate

When the string is displaced at its mid point the centre of the string becomes an antinode.

String is fixed at both ends \therefore y(x,t) = 0 at x = 0 and L

 $y(0,t)=0$ when $x = 0$ as $sin(kx) = 0$ at $x = 0$

$$
y(x,t) = 2A_0 \sin(kx) \cos(\omega t)
$$

 $y(L,t) = 0$ when $sin(kL) = 0$ ie $k_n L = n \pi$ n=1,2,3...

but $k_n = 2\pi / \lambda$: $(2\pi / \lambda_n)L = n\pi$ or

$$
\lambda_n = 2L/n
$$

For first normal mode
$$
L = \lambda_1 / 2
$$

The next normal mode occurs when the length of the string L = one wavelength, i.e. $L = \lambda_2$

The third normal mode occurs when $L = 3\lambda_3/2$

Generally normal modes occur when $L = n\lambda_n/2$

$$
ie \quad \lambda_n = \frac{2L}{n} \text{ where } n = 1, 2, 3, \dots
$$

The natural frequencies associated with these modes can be derived from $f = v/\lambda$

$$
f = \frac{v}{\lambda} = \frac{n}{2L}v
$$
 with $n = 1,2,3,...$

Standing waves in a string fixed at both ends

For a string of mass/unit length µ, under tension F we can replace v by $(F/\mu)^{1/2}$

$$
f = {n \over 2L} \sqrt{\frac{F}{\mu}}
$$
 with $n = 1,2,3,...$

The lowest frequency (**fundamental**) corresponds to n = 1ie $f = \frac{1}{2!}v$ or $f = \frac{1}{2!}\sqrt{\frac{F}{u}}$

Can one predict the amplitude of each mode (overtone/harmonic?) following plucking?

Using the procedure to measure the Fourier coefficients it is possible to predict the amplitude of each harmonic tone.

- •You know the shape just before it is plucked. •You know that each mode moves at its own frequency
- •The shape when released
- •We rewrite this as

shape =
$$
f(x, t = 0)
$$

$$
f(x,t=0)=\sum_n A_n \sin(k_n x)
$$

Each harmonic has its own frequency of oscillation, the m-th harmonic moves at a frequency $f_m=mf_0$ or m times that of the fundamental mode.

$$
f(x, t = 0) = \sum_{n} A_{n} \sin(k_{n}x)
$$

$$
f(x, t) = \sum_{n} A_{n} \sin(k_{n}x) \cos(\omega_{n}t)
$$

Recall modes on a string:

$$
u(x, t) = \sum_{n=0}^{\infty} A_n U_n(x, \omega_n) \cos(\omega_n t)
$$

This is the sum of standing waves or *eigenfunctions*, $U_n(x, \omega_n)$, each of which is weighted by the amplitude A_n and vibrates at its *eigenfrequency* ω_n .

The eigenfunctions and eigenfrequencies are constants due to the physical properties of the string.

The amplitudes depend on the position and nature of the source that excited the motion.

The eigenfunctions were constrained by the boundary conditions, so that

$$
U_n(x, \omega_n) = \sin(n\pi x/L) = \sin(\omega_n x/v) \qquad \omega_n = n\pi v/L = 2\pi v/\lambda
$$

Source excitation

$$
u(x, t) = \sum_{n=0}^{\infty} \sin(n\pi x_s/L) F(\omega_n) \sin(n\pi x/L) \cos(\omega_n t)
$$

The source, at $x_s = 8$, is described by

 $F(\omega_n) = \exp[-(\omega_n \tau)^2/4]$

with $\tau = 0.2$.

