

## General form of LWE

$$\frac{\partial^2 \psi(\mathbf{x}, \mathbf{t})}{\partial \mathbf{t}^2} = \mathbf{v}^2 \frac{\partial^2 \psi(\mathbf{x}, \mathbf{t})}{\partial \mathbf{x}^2}$$

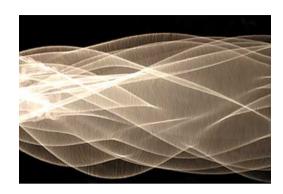
WAVE: organized propagating imbalance, satisfying differential equations of motion



12. . .

## Separation of variables: string





$$\frac{\partial^2 \gamma(x,t)}{\partial x^2} - \frac{1}{c^2} \frac{\partial^2 \gamma(x,t)}{\partial t^2} = 0$$

and if it has separable solutions:

$$y(x,t) = X(x)T(t)$$

$$\frac{d^{2}X(x)}{dx^{2}} + k^{2}X(x) = 0 \qquad T''(t) + c^{2}k^{2}T(t) = 0$$
  

$$X(x) = A\cos(kx) + B\sin(kx) \qquad T(t) = C\cos(\omega t) + D\sin(\omega t)$$
  

$$\omega = ck$$

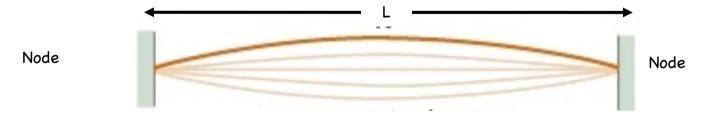
To be determined by **initial** and **boundary** conditions



Consider a string of length L and fixed at both ends

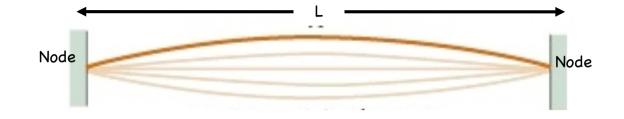
## The string has a number of natural patterns of vibration called **NORMAL MODES**

Each normal mode has a characteristic frequency which we can easily calculate



When the string is displaced at its mid point the centre of the string becomes an antinode.





String is fixed at both ends  $\therefore$  y(x,t) = 0 at x = 0 and L

y(0,t)=0 when x = 0 as sin(kx) = 0 at x = 0

$$y(x,t) = 2A_0 \sin(kx)\cos(\omega t)$$

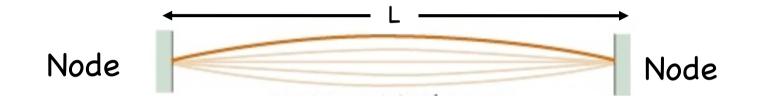
y(L,t) = 0 when sin(kL) = 0 ie  $k_n L = n \pi$  n=1,2,3....

but  $k_n = 2\pi / \lambda$   $\therefore (2\pi / \lambda_n) L = n\pi$  or

$$\lambda_n = 2L/n$$



For first normal mode  $L = \lambda_1 / 2$ 



The next normal mode occurs when the length of the string L = one wavelength, i.e. L =  $\lambda_2$ 

The third normal mode occurs when  $L = 3\lambda_3/2$ 

Generally normal modes occur when  $L = n\lambda_n/2$ 

ie 
$$\lambda_n = \frac{2L}{n}$$
 where  $n = 1, 2, 3$ .....

The natural frequencies associated with these modes can be derived from  $f = v/\lambda$ 

$$f = \frac{v}{\lambda} = \frac{n}{2L}v$$
 with  $n = 1,2,3...$ 

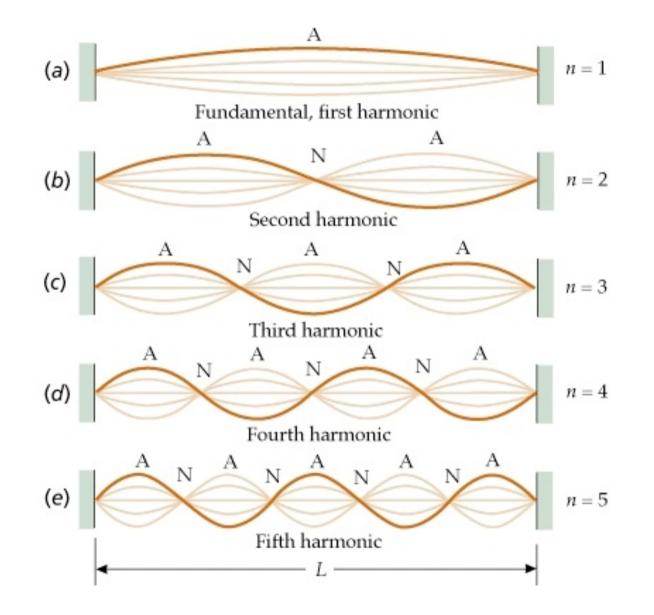
Standing waves in a string fixed at both ends

$$f = \frac{n}{2L} \sqrt{\frac{F}{\mu}}$$
 with  $n = 1, 2, 3....$ 

The lowest frequency (fundamental) corresponds to n = 1 ie f =  $\frac{1}{2L}v$  or f =  $\frac{1}{2L}\sqrt{\frac{F}{\mu}}$ 



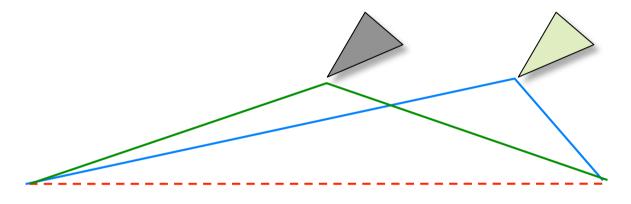








Can one predict the amplitude of each mode (overtone/harmonic?) following plucking?

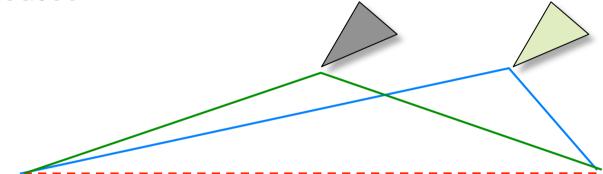


Using the procedure to measure the Fourier coefficients it is possible to predict the amplitude of each harmonic tone.





- •You know the shape just before it is plucked.
- •You know that each mode moves at its own
- frequency
- •The shape when released
- •We rewrite this as



shape = f(x, t = 0)

$$f(x,t=0) = \sum_{n} A_{n} sin(k_{n}x)$$





Each harmonic has its own frequency of oscillation, the m-th harmonic moves at a frequency  $f_m=mf_0$  or m times that of the fundamental mode.

$$f(x, t = 0) = \sum_{n} A_{n} \sin(k_{n}x)$$

$$f(x, t) = \sum_{n} A_{n} \frac{\sin(k_{n}x) \cos(\omega_{n}t)}{n}$$
http://www.falstad.com/loadedstring/





Recall modes on a string:

$$u(x, t) = \sum_{n=0}^{\infty} A_n U_n(x, \omega_n) \cos(\omega_n t)$$

This is the sum of standing waves or *eigenfunctions*,  $U_n(x, \omega_n)$ , each of which is weighted by the amplitude  $A_n$  and vibrates at its *eigenfrequency*  $\omega_n$ .

The eigenfunctions and eigenfrequencies are constants due to the physical properties of the string.

The amplitudes depend on the position and nature of the source that excited the motion.

The eigenfunctions were constrained by the boundary conditions, so that

$$U_n(x, \omega_n) = \sin(n\pi x/L) = \sin(\omega_n x/v)$$
  $\omega_n = n\pi v/L = 2\pi v/\lambda$ 



## Source excitation



$$u(x, t) = \sum_{n=0}^{\infty} \sin(n\pi x_s/L) F(\omega_n) \sin(n\pi x/L) \cos(\omega_n t)$$

The source, at  $x_s = 8$ , is described by

 $F(\omega_n) = \exp[-(\omega_n \tau)^2/4]$ 

with  $\tau = 0.2$ .

