

607SM - Tecniche avanzate di indagine microscopica

Advanced microscopy techniques – 6CFU, 2023/24, 1st semester

Part1:

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Monday 11-13, Aula A -Ed M + Wednesday 9-11, Aula A – Ed M

Part 1 (Dan Cojoc) - 26 h (20 lectures + 6 lab)

1. Optical Microscopy – Physical Principles – 9h

- 1.1. Basics (Image formation, magnification, resolution, image quality)
- 1.2. Digital camera (image acquisition, formats, properties)
- 1.3. Phase imaging (qualitative and quantitative)
- 1.4. Dark field and Polarization microscopy

- 1.5. Non Linear Optical Microscopy
- 1.6. Photoacoustic Microscopy
- 1.7. Super resolution microscopy – STED, PALM/STORM, MINFLUX

Optical microscopy:

- uses light in visible range (wavelength 400 - 700 nm) to image samples with details as small as 200 nm (~ half wavelength).
- recently, new techniques have been developed to improve the resolution and resolve details as small as 1-10 nm in living cells
- moreover, light has been employed to manipulate cells and their components (optical tweezers and scalpels) and measure piconewton forces (force spectroscopy)

What can be seen with a light microscope?

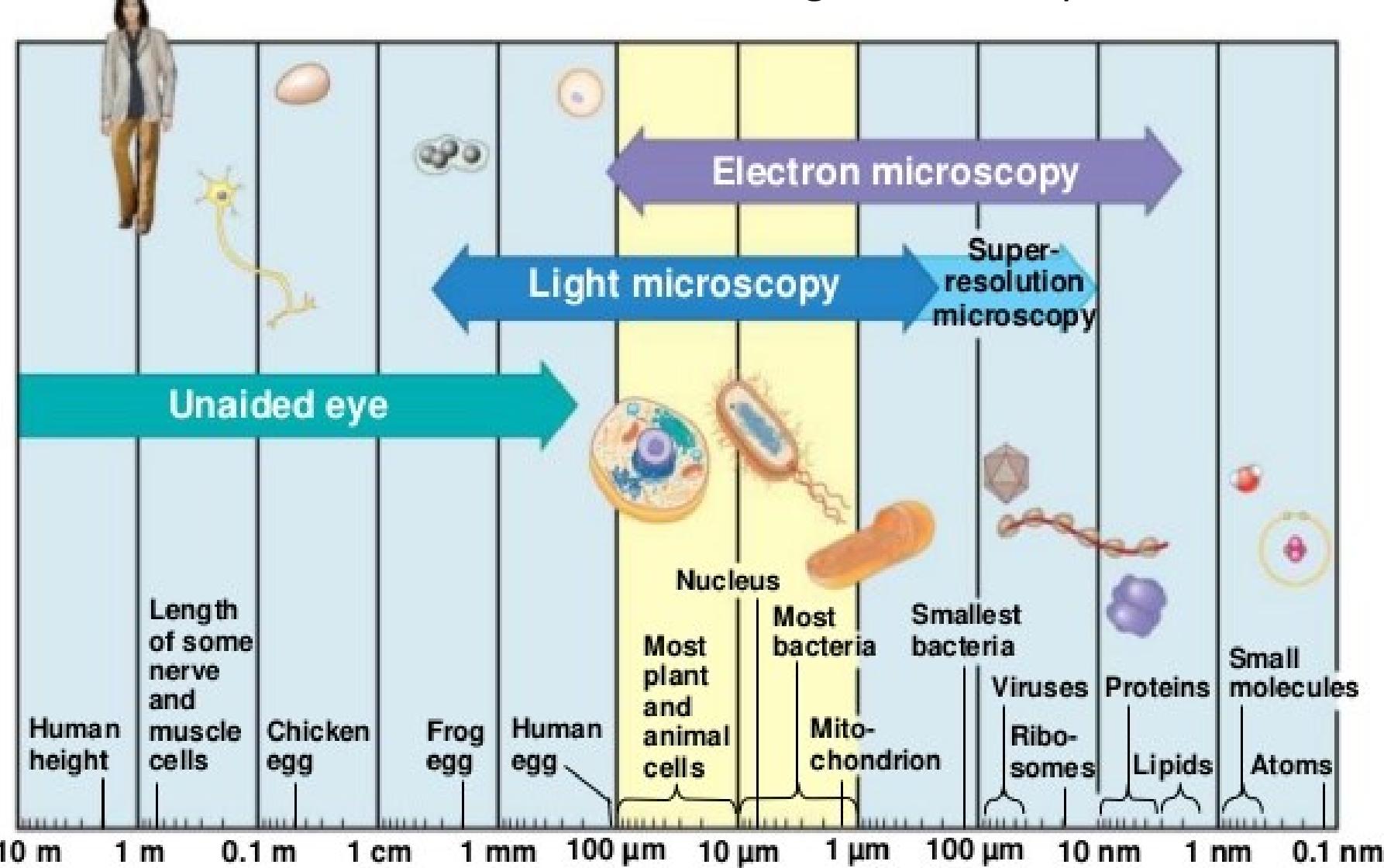
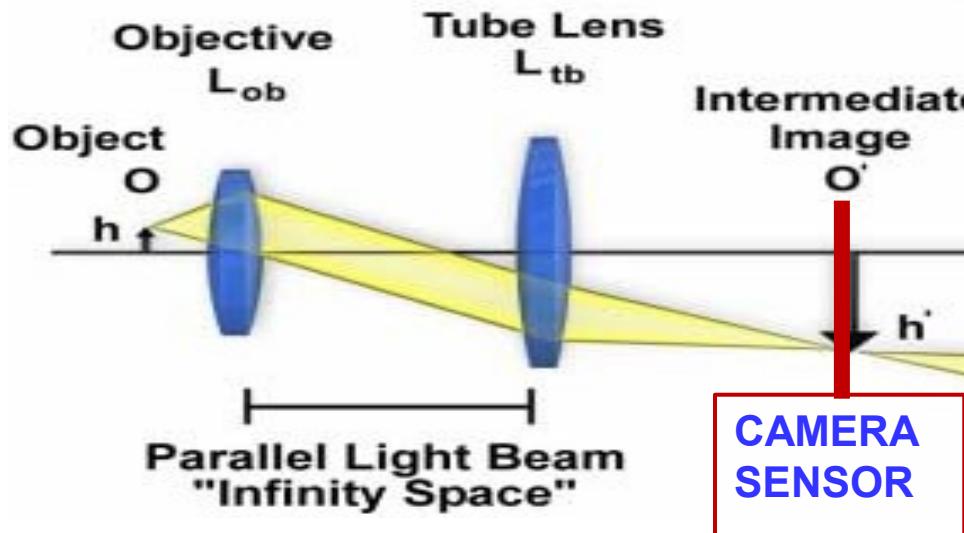


Image formation

Magnification vs resolution in brightfield microscopy

Image formation in the optical microscope



- The object is positioned in the focal plane of the **OBJECTIVE**, hence its image (through the objective) is projected to infinity.
- The **TUBE LENS** 'brings' this image from infinity to its focal plane, forming a magnified image, called intermediate image. in its focal plane.
- The intermediate image can be observed through the EYEPIECE or it is directly captured by a **CAMERA SENSOR** and displayed on a monitor.

Conjugate Planes in the Optical Microscope

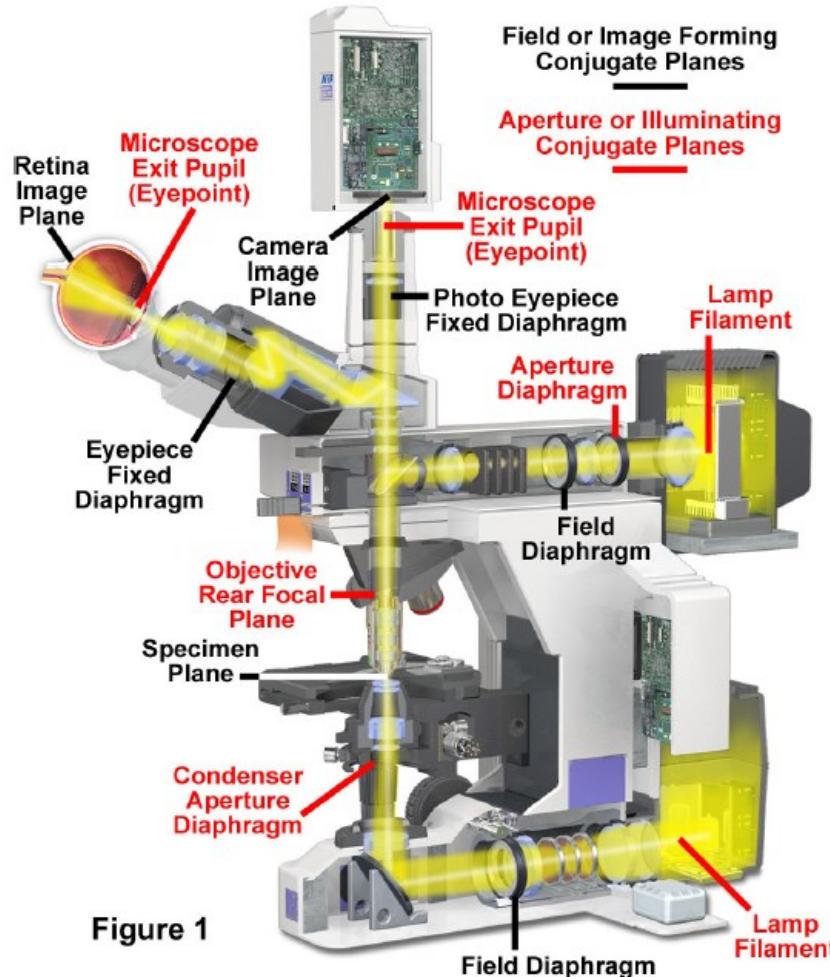


Figure 1

Conjugate Planes in the Optical Microscope

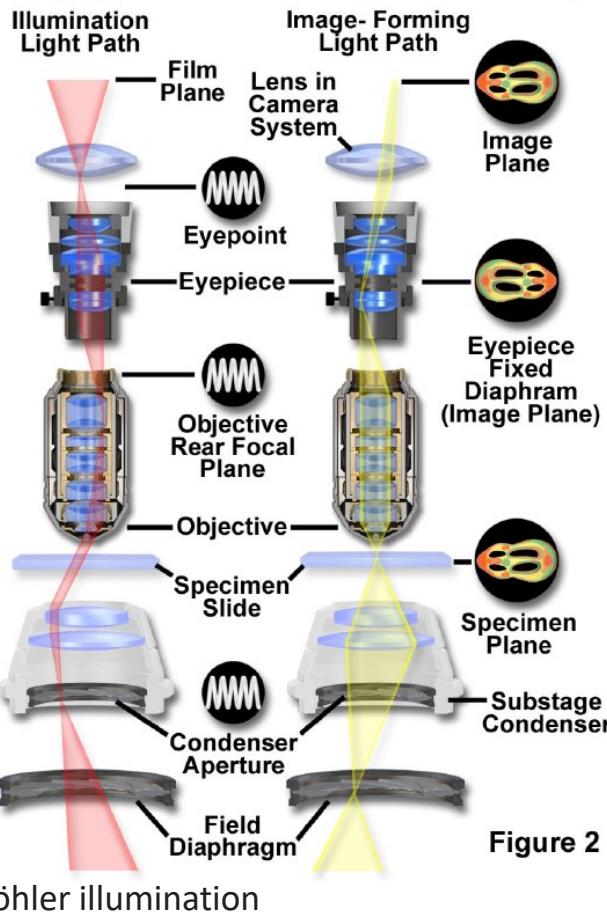


Figure 2

Conjugate planes: one plane is the image of the other plane through a lens (group of lenses)

Good material to read before practical session with optical microscope (moodle or link)

Magnification

Objective magnification

$$M_{OB} = \frac{f_{TL}}{f_{OB}}$$

4X – 120 X

Microscope magnification

$$M_M = M_{OB} M_{EY}$$

50 X – 2000 X

EY – Eyepiece; DC – Digital Camera

$$M_M = M_{OB} M_{DC}$$

Note:

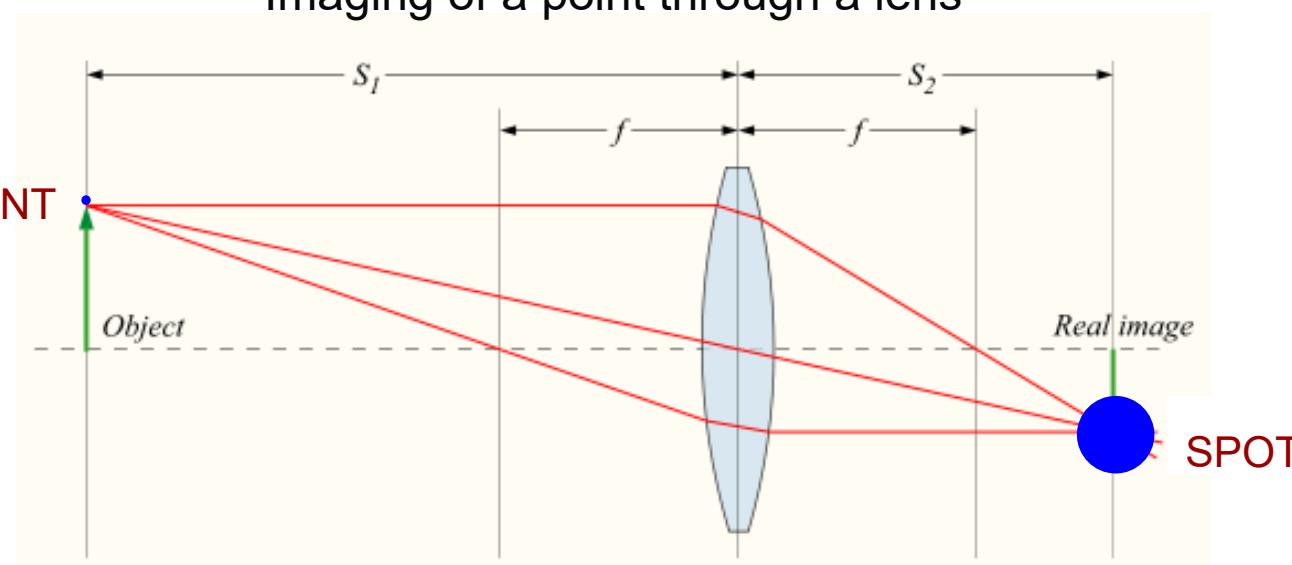
Magnification is different from Resolution !!!

i.e. High magnification does not always mean a better resolved image

Resolution

Resolution describes the ability of an optical system to **resolve details** of the object that is being imaged.

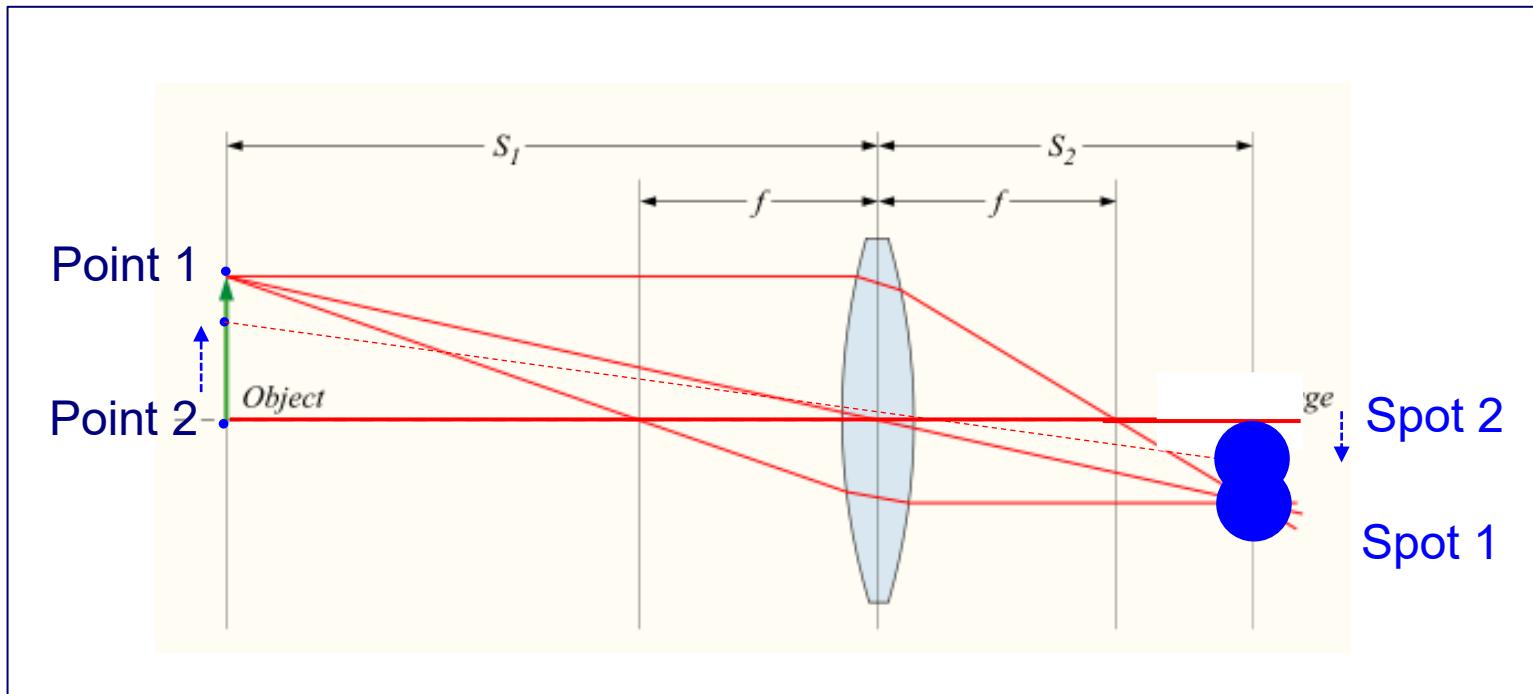
Due to the diffraction of light through an optical system with finite size, a **POINT** object is imaged into a **SPOT** rather than a point.



LARGER the LENS, SMALLER the SPOT !

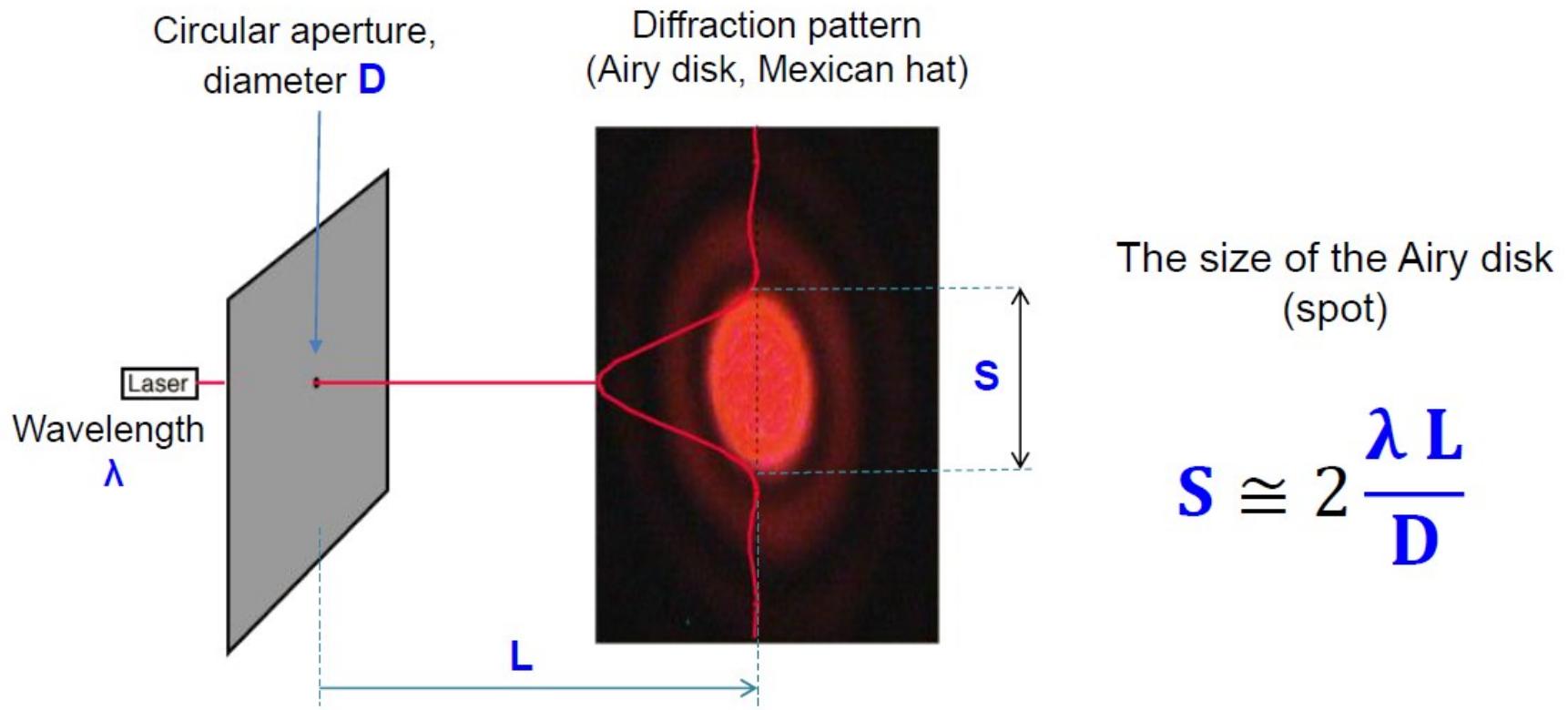
SMALLER the SPOT, BETTER the RESOLUTION !

Imaging TWO points through a lens



When the two points are close each other, their images (two spots) overlap and hence they can not be separated (resolved)!

Diffraction of light through a circular aperture



The size of the Airy disk (spot)

$$S \approx 2 \frac{\lambda L}{D}$$

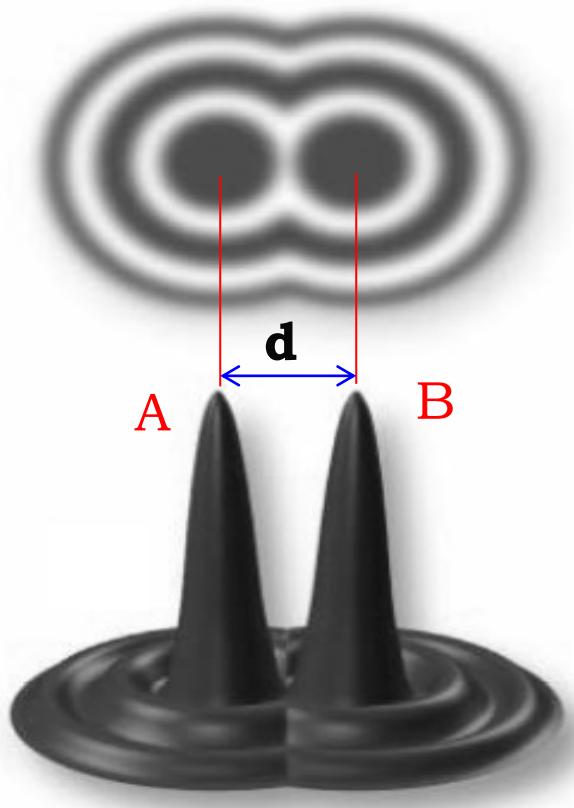
S is proportional to λ and L but it is inverse proportional to the size of the aperture D

LARGER the APERTURE, SMALLER the SPOT !

When the SPOT would be a POINT ?

Resolution criterion

The resolution, r , is defined as the shortest distance between two points on a specimen that can still be distinguished by the observer or camera sensor as separate entities.



Rayleigh criterion

$$r = 0.61 \frac{\lambda}{NA}$$

Abbe criterion

$$r = 0.5 \frac{\lambda}{NA}$$

NA- Numerical Aperture

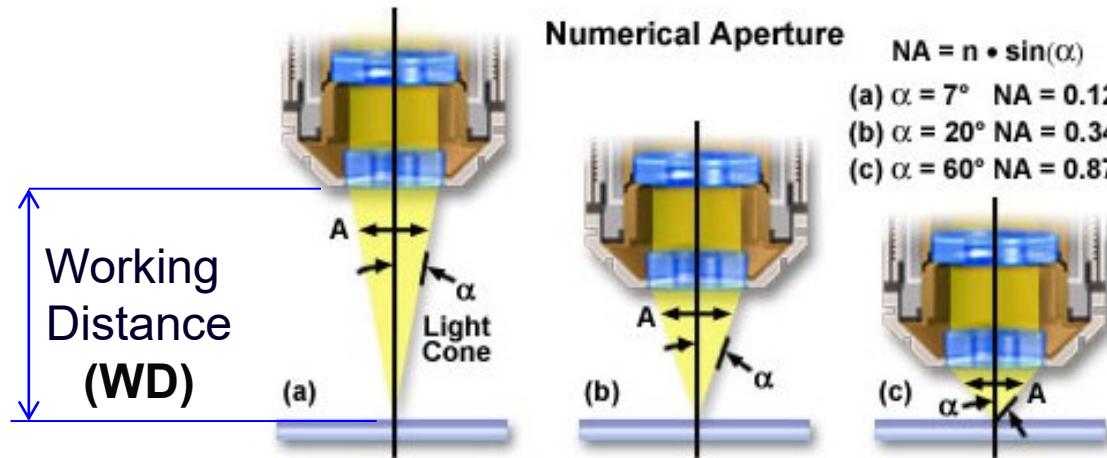
Estimating the lateral resolution
of a microscope objective (lens):

$$NA = 1.5, \lambda = 400 \text{ nm}$$

$$\rightarrow r \sim 200 \text{ nm}$$

A and B are separated if: $d > r$

Resolution – Numerical Aperture – Working Distance



$$NA = n \cdot \sin \alpha$$

Immersed objectives $NA > 1$

Oil ($n=1.515$), Glycerin ($n=1.47$) or Water ($n=1.33$)

Higher NA → better lateral Resolution

Note: WD decreases when NA increases !!!

Microscope objective specs

60x Plan Apochromat Objective



- The magnification of the optical microscope is max 2500X
- The lateral resolution of the optical microscope is diffraction limited to 200 nm (i.e. half of the wavelength: $\lambda / 2$)

Magnification is different from Resolution

i.e. a higher magnification does not guarantee a better resolution

Lateral resolution

$$\Delta x \cong 0.5 \frac{\lambda}{NA}$$

~ 200 nm

Axial resolution

$$\Delta z \cong 0.5 \frac{\lambda n}{NA^2}$$

~ 300 nm

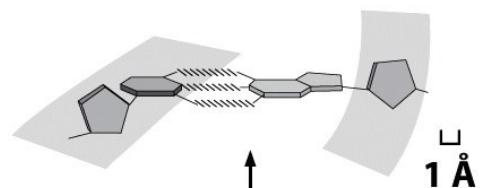
$$\Delta z \approx \frac{n}{NA} \Delta x \longrightarrow \Delta z > \Delta x$$



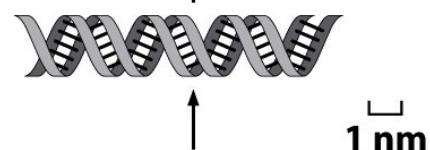
Lateral resolution is better than the axial resolution

Biological Scale and Size

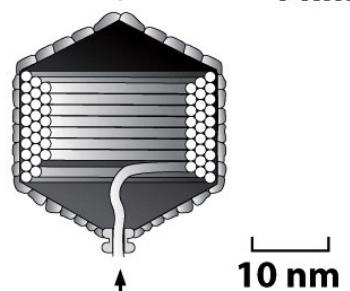
bases



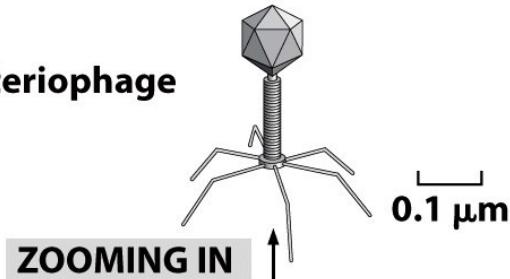
DNA



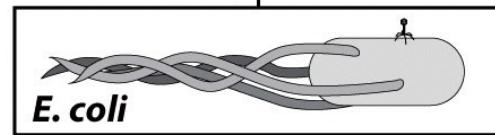
viral capsid



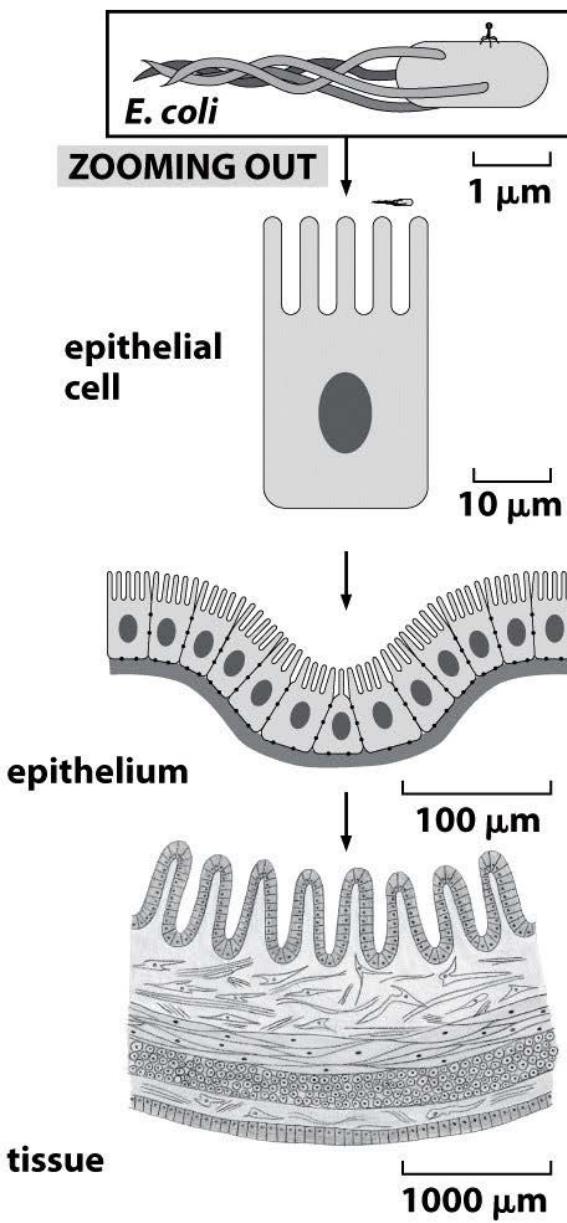
bacteriophage



ZOOMING IN



Optical Microscopy



Optical aberrations and image quality

Optical aberrations

Optical aberration:

deviation of the image from the ideal image given by paraxial approximation

Paraxial optics: propagation of the rays of light close to the optical axis.

Example:

a point object is imaged into a point in paraxial optics but;
in practice the image is a spot (diffraction) deformed (abberations)

There are:

Geometric aberrations:

due to the shape of the optical elements

Chromatic aberrations:

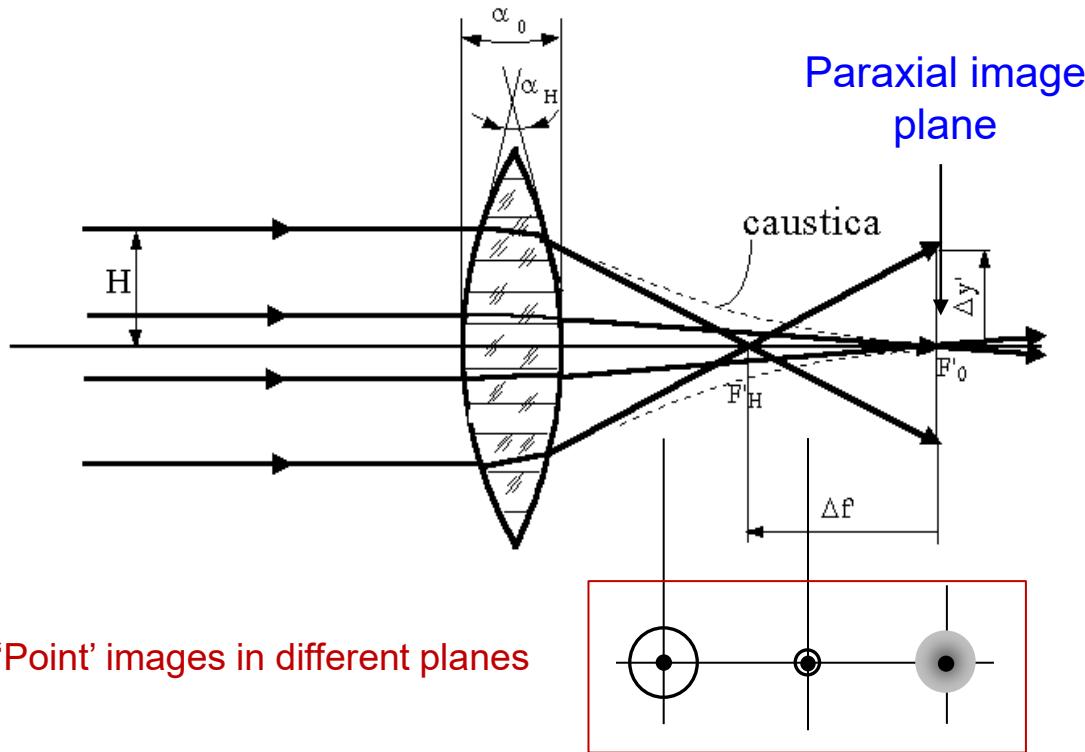
due to the refractive index of the optics, which is function of wavelength

Nice applets:

<http://www.microscopyu.com/tutorials/java/aberrations/astigmatism/index.html>

An example of geometric aberration: spherical aberration

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Rays passing through the outer part of the lens are focused closer than the paraxial rays.

Axial spherical aberration is defined by the difference between the extreme positions of the focal points F'_0 and F'_H

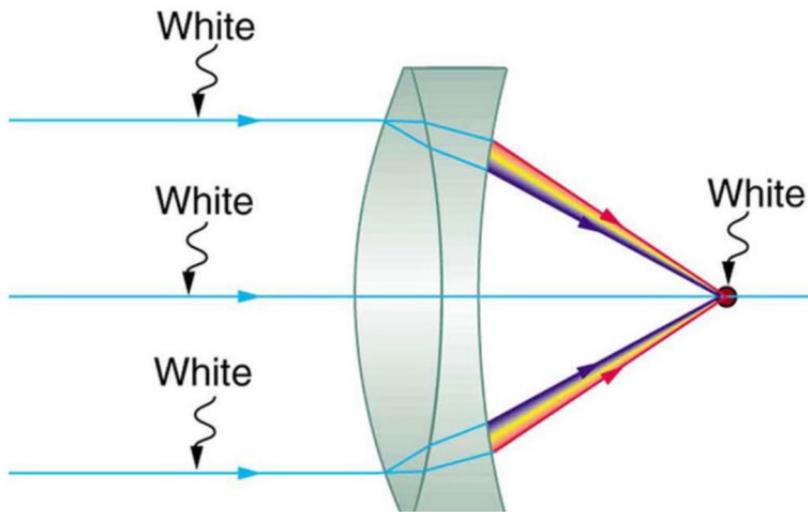
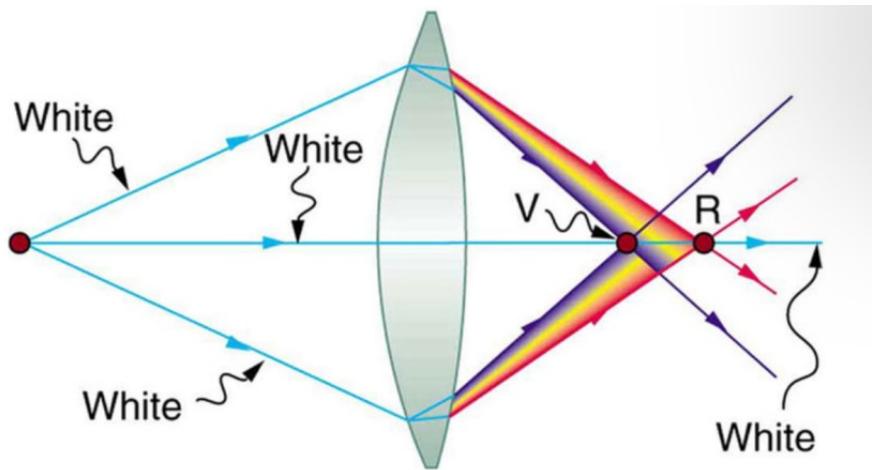
This difference can be explained considering the 'prism' effect:

the emergent ray is deviated by an angle $\delta_i = (n - 1)\alpha_i$ with respect to the incident ray

Chromatic abberation and achromatic doublet

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The image of a point depends of the wavelength (color of light)



Digital camera image acquisition (formats, properties, SNR)

Common requirements for digital cameras in microscopy:

- accurate morphological representation of the specimens
- appropriate spatial resolution
- appropriate signal to noise levels
- accurate capture of dynamic events
- high sensitivity for minimum perturbation of the specimens
- appropriate dynamic range

These requirements depend upon the Sample, Optics, **camera Sensor**

Digital cameras - sensors:

CCD - Charge Coupled Device

EM-CCD – Electronic Multiplying CCD

CMOS - Complementary Metal Oxide Semiconductor

Camera Sensor

- Chip layout (pixel size, frame rate).
- Signal sampling
- Noise and Signal to Noise Ratio (SNR)

Chip Layout

Pixel size: a good sensitivity requires a large pixel size; on the other hand, a good spatial resolution requires a small pixel size → a trade off must be found.

Usual ranges: CCD: 2 to 7.5 μm , EMCCD: 8 to 24 μm , CMOS: 2 to 7.5 μm

Chip size: determines the camera's field of view.

Usual ranges:

CCD: 4.74 x 3.55 mm to 15.16 x15.16 mm (small to medium)

EMCCD: 3.072 x 2.072 mm to 13.3 x13.3 mm (small to medium)

CMOS: 6.97 x 2.23 mm to 16.6 x14 mm (medium to large)

Frame rate: the maximum number of frames that can be captured per second fps

Usual ranges :

CCD: 3-100 fps @ 1362 x1024 pixels

EMCCD: 8-500 fps @ 800x600 pixels

CMOS: 500 fps @ 1700x1600 pixels (up to 5 Mfps for ROI of 128 x 128 pixels)

Note: If a high frame rate is necessary, beside the sensor choice, the PC and data transfer are also very important (e.g. 2 sec recording of images 512x512 pixels with 8 bit depth / pixel at 1000 fps produces a file > 1 GB)

Signal / Image Sampling

It is given by the pixel size and defines the sensor resolution. This should match the optical resolution: $Ro = 0.5 \lambda / NA$ which at the sensor plane becomes:

$$Rs = M \cdot Ro$$

with M being the magnification.

The sampling (Nyquist) theorem says:

sampling frequency should be at least twice higher than the max frequency of the signal.

We get :

$$1/p = f \geq 2 \cdot f_{max} = 2/Rs$$

sampling factor SF :

$$SF = Rs / p \geq 2$$

$$SF = M \times Ro / p \geq 2$$

where: p is the pixel size, f the sampling frequency, f_{max} the max frequency in the image and SF the sampling factor

Example

Table 1: Sampling Factors For Different Cameras (< 2.0 is undersampled)							
MAG	Coupler	NA	Resolution	Resolution	3.5 micron	6.5 micron	14 micron
			Limit (@ sample in microns)	Limit (@ detector in microns)			
10	1	0.45	0.75	7.46	2.13	1.15	0.53
20	1	0.75	0.45	8.95	2.56	1.38	0.64
40	1	0.95	0.35	14.13	4.04	2.17	1.01
40	1	1	0.34	13.42	3.83	2.06	0.96
40	1	1.3	0.26	10.32	2.95	1.59	0.74
60	1	1.2	0.28	16.78	4.79	2.58	1.20
60	1	1.3	0.26	15.48	4.42	2.38	1.11
100	1	1.4	0.24	23.96	6.85	3.69	1.71

Sufficient Sampling
Insufficient Sampling

Sampling factor $SF \geq 2.0 \rightarrow$ good sampling

Sampling factor $SF < 2.0 \rightarrow$ undersampled

Noise (σ) and Signal to Noise ratio (SNR)

Table 1: Types of Noise

Type of Noise	Description
σ_d	Dark Noise Dark current is a time-dependent signal generated on the sensor when no light is present due to heat which causes random generation of holes and electrons in the depletion region of the sensor. Dark noise is the fluctuation in this signal.
σ_r	Read Noise Noise contributed by the amplifier during the conversion of the analog signal to digital signal.
σ_n	Shot Noise Fluctuation in signal due to the quantum properties of photons. The number of photons measured at any given point in time can fluctuate by plus or minus the square-root of the measured signal.

They can be considered as independent,

$$\sigma = \sqrt{\sigma_d^2 + \sigma_r^2 + \sigma_n^2}$$

so the total Noise is given by:

Signal to Noise Ratio

$$SNR = \frac{\Phi}{\sigma} = \frac{S_d}{N}$$

Detected (signal) photons
Noise

CCD, CMOS:

$$\sigma = \sqrt{\sigma_d^2 + \sigma_r^2 + \sigma_n^2}$$

EMCCD:

$$\sigma = \sqrt{\sigma_d^2 + \psi^2 \sigma_r^2 + \psi^2 \sigma_n^2}$$

Excess - noise factor ψ due to the impact ionization through which electronic amplification is achieved

Example: pixel performance of the perfect camera

$$SNR = \frac{QE * S}{\sqrt{F_n^2 * QE * (S + I_b) + (N_r/M)^2}}$$

For a perfect detector (QE=1, Nr=0; Fn=1):

$$SNR = \frac{1 * S}{\sqrt{1^2 * 1 * S + (0/1)^2}} = \sqrt{S}$$

where:

QE: quantum efficiency

S : input signal (photon/pixel)

F_n : noise factor

N_r : readout noise

M : EM gain (=1 for CCD/CMOS)

I_b : background

N_d : dark noise (not included, assumed to be negligible)

Even with a perfect detector we have noise !

This is the photon shot noise and is a function of photon statistics.

In other words, until we have light (photons) we have noise.

To improve SNR we have to collect as many photons as possible.

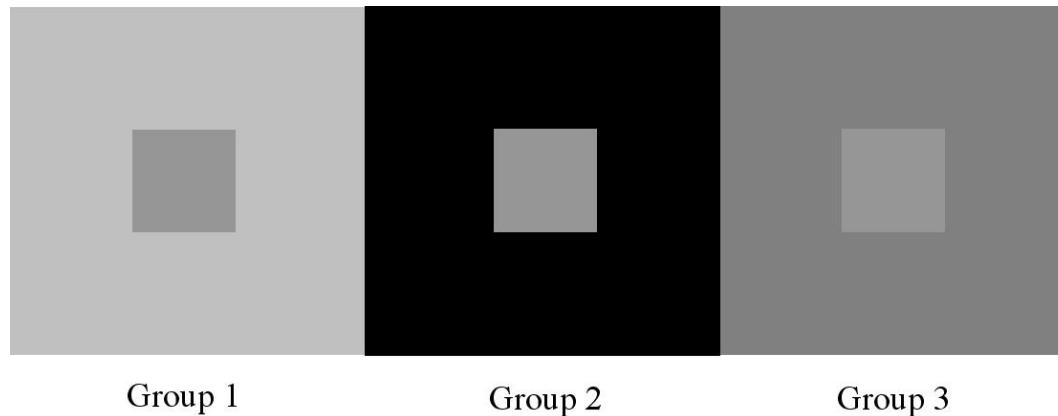
Phase contrast in light microscopy

Image contrast in brightfield microscopy

The contrast is formed by different absorption of light by medium and sample, creating an image with different intensity levels.

An useful image (with features that can be well distinguished) needs to have enough brightness and good contrast:

$$\text{Contrast} = \frac{I_{\max} - I_{\min}}{I_{\max} + I_{\min}}$$



Which center grey box is lighter ?

Contrast mechanisms for image formation

Starting from the solution of the wave equation describing light propagating in materials

$$E(z,t) = E_0 \exp \left[i \left(\omega t - \left(n_{\text{real}} + i n_{\text{imag}} \right) \omega z / c \right) \right] = \underbrace{E_0 \exp \left[i \left(\omega t - n_{\text{real}} \omega z / c \right) \right]}_{\text{Amplitude}} \underbrace{\exp \left(-n_{\text{imag}} \omega z / c \right)}_{\text{Phase change / Absorption / attenuation}}$$

$$n = n_{\text{real}} + j n_{\text{imag}}$$

Refractive index (complex value) is a material characteristic

Electric field

$$\vec{E}(z,t) = \vec{A}(z) \cdot e^{i\omega t} \cdot e^{i\varphi(z)}$$

The contrast mechanisms use

AMPLITUDE, PHASE, POLARIZATION

Cells are almost transparent to visible light, hence it is difficult to obtain a good contrast in brightfield (based on absorption / transmission of light INTENSITY).

However, light is characterized also by PHASE and POLARIZATION which are influenced by the sample, and hence carries information about it.

$$\mathbf{E} = \mathbf{A} \cdot e^{i\varphi}$$

\mathbf{E} – intensity of the electrical (optical) field

Note that \mathbf{E} is a vector; light polarization is related to the vector orientation

\mathbf{A} – amplitude (vector); φ – **phase**

The information carried by **polarization** and **phase** is usually lost because
all existent sensors of light (e.g. retina photoreceptors, photodiodes)
detect only the intensity of light.

$$I - \text{intensity of light} \quad / = |\mathbf{E}|^2 = |\mathbf{A}|^2 \quad \text{because } |e^{i\varphi}| = 1$$

To exploit the **phase** or the **polarization** information
we need to convert it into intensity.

Exploiting the phase of light to enhance image contrast

The phase contrast method exploits **phase retardation of light** by different regions of the sample and the conversion of the phase difference into intensity.

$$E = A \cdot e^{i\varphi}$$

The phase of light:

$$\varphi = 2\pi \frac{OP}{\lambda} \quad \lambda - \text{light wavelength}$$

OP – Optical Path

$$OP = t \cdot n$$

$$\lambda = c T = \frac{2\pi}{\omega} c$$

n – refractive index of the material (sample or medium)

t – geometrical path (e.g. cell height)

Note that the phase φ can be rewritten to obtain the term in the expression of the electric field

$$E(z,t) = E_0 \exp \left[i \left(\omega t - \left(n_{\text{real}} + i n_{\text{imag}} \right) \omega z / c \right) \right] = E_0 \exp \left[i \left(\omega t - n_{\text{real}} \omega z / c \right) \right] \exp \left[\underline{-n_{\text{imag}} \omega z / c} \right]$$

Phase Contrast- principle

A pure phase sample, as a cell, does not absorb light but influence its phase.

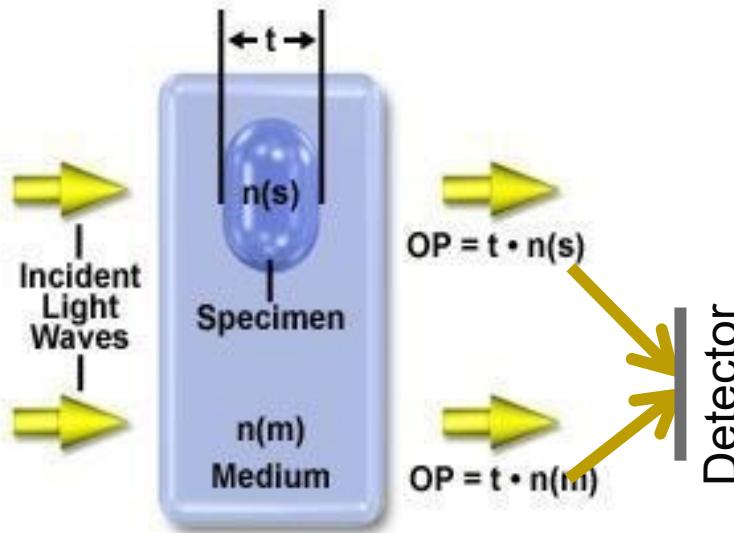
The sample has a refractive index $n(s)$ different from the refractive index $n(m)$ of the medium in which the sample is immersed.

Hence the optical paths (OP) of light through sample and medium are different:

$$OP(s) = t \cdot n(s) \quad ; \quad OP(m) = t \cdot n(m)$$

and so are the phase shifts / retardations: $\Delta\varphi(s) = 2\pi OP(s) / \lambda$; $\Delta\varphi(m) = 2\pi OP(m) / \lambda$

Optical Path Difference in Phase Objects



Unpolarized light and $A(s)=A(m)=1$

$$E(s) = e^{i \Delta\varphi(s)}$$

$$E(m) = e^{i \Delta\varphi(m)}$$

The intensity on the detector:

$$I = |E(s) + E(m)|^2 =$$

$$= |E(s)|^2 + |E(m)|^2 + E(s) * E(m) + E(s) * E(m) * =$$

$$= 2 + 2 \cos(2\pi OPD / \lambda)$$

NOT just 2

as in superposition

OPD - Optical Path Difference:

$$OPD = OP(s) - OP(m) = t(n(s)) - n(m))$$

Phase contrast implementation

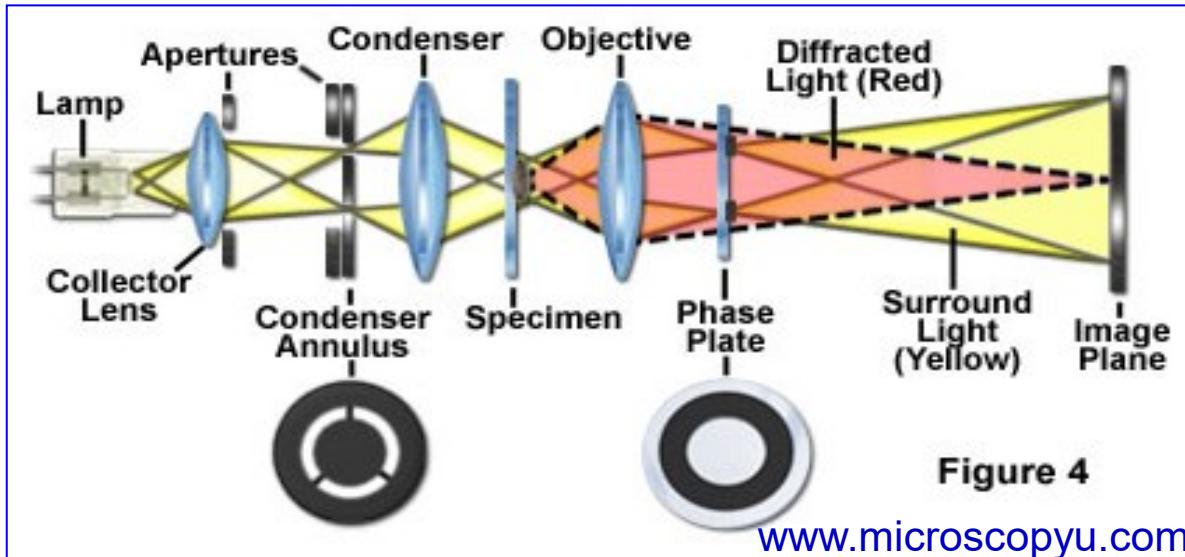


Figure 4

www.microscopyu.com

Using the condenser annulus, the specimen is illuminated obliquely.

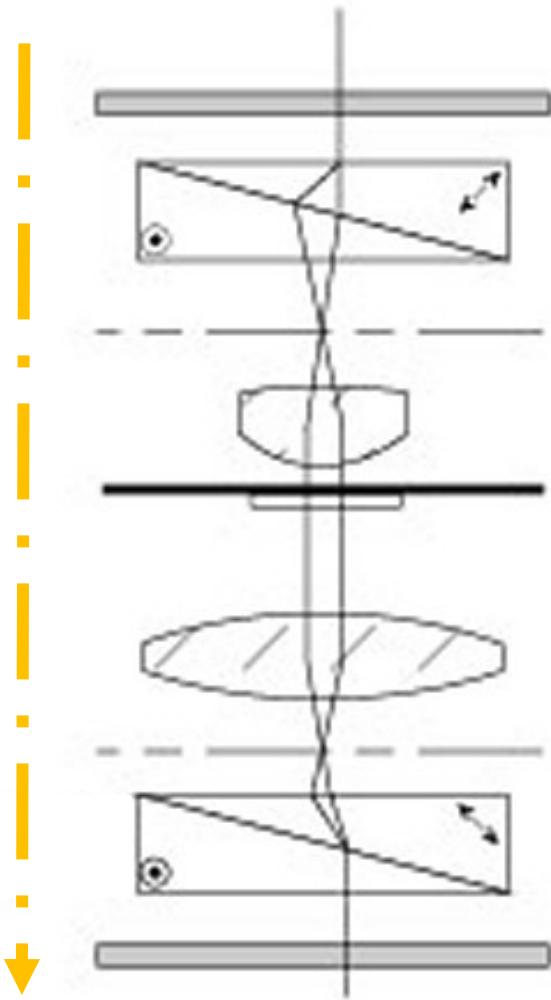
The light which is not intercepted by the specimen is focused onto the ring of the phase plate, which introduces a phase shift. The light reaching the specimen is focused on the image plane. Most of this light does not pass through the phase ring.

Phase contrast is obtained in the image plane from the interference between the light intercepted by the specimen and light not intercepted.

The Nobel Prize in Physics 1953 was awarded to Frits Zernike for "his demonstration of the phase contrast method, especially for his invention of the phase contrast microscope".

Another method: Differential Interference Contrast (DIC) method

DIC is an imaging method generating intensity contrast in the specimen's image by exploiting phase differences between the light passing through pairs of points of the specimen very near to each other.



Pol

W1

fc

Cond

Spec

Obj

f0

W2

Ana

Nomarsky DIC implementation:

Uses polarized light and anisotropic prisms

Pol: polarizer

W1, W2: Nomarski prisms

fc: condenser's focal plane

Cond: condenser lens

Spec: specimen

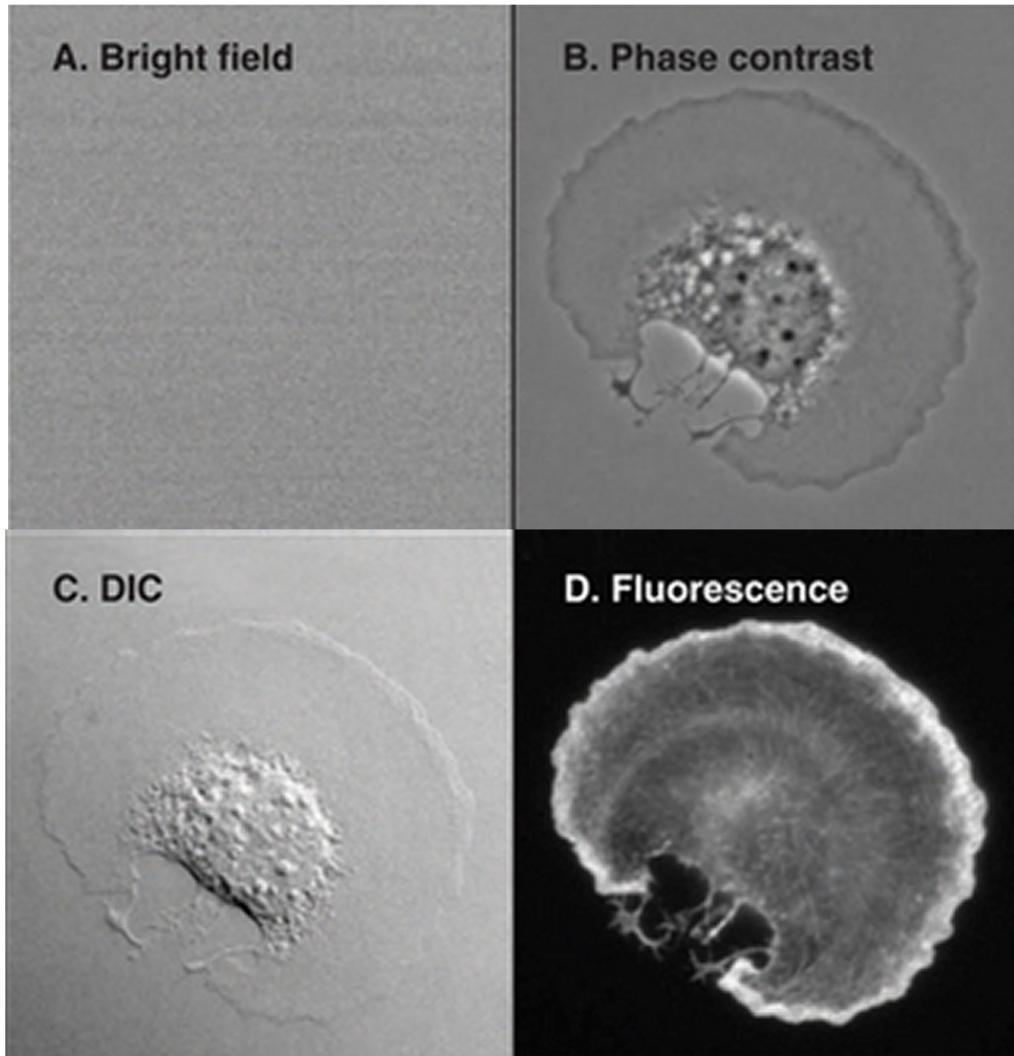
Obj: objective lens

f0: back focal plane of the objective lens

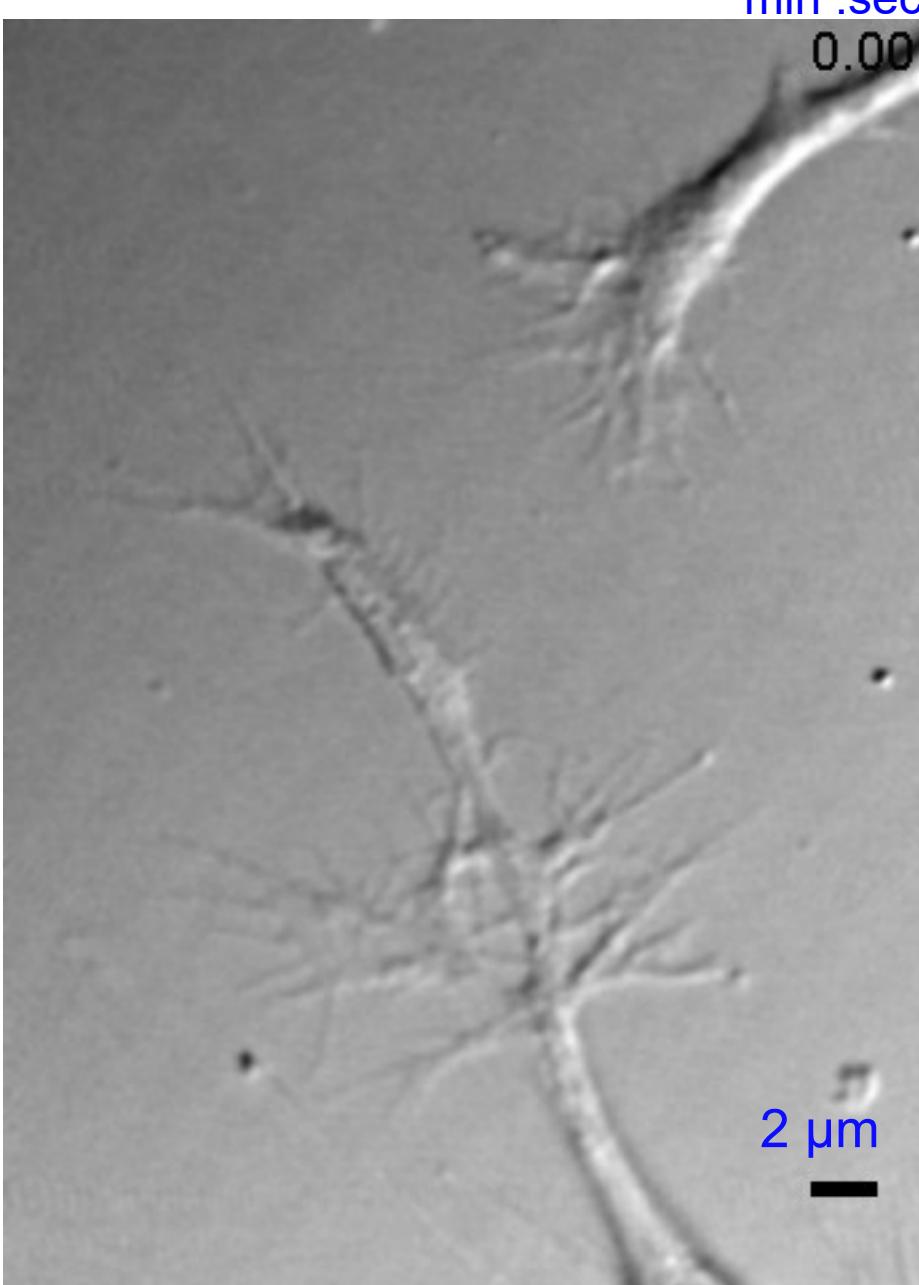
W2: Nomarski-modified Wollaston prism

Ana: polarizer (Analyzer)

Comparison - contrast enhancement by phase imaging



Images of a spread mouse 3T3 cell grown in tissue culture on a microscope slide, then fixed and stained with rhodamine-phalloidin, a fluorescent peptide that binds actin filaments.



Imaging neuronal growth cones
with DIC

movie

F. Difato *et al* (2006) OM-Lab & SISSA

Quantitative phase imaging: digital holographic microscopy

Application Ex: measuring cell height and volume

Phase contrast microscopy is a qualitative technique, using the phase shifts to enhance the contrast and improve the image quality.

Since white light illumination is used ($\lambda : 400 - 800 \text{ nm}$) phase contrast microscopy does not allow to measure the optical paths, which contain information about the sample height (t) and material (n(s)).

$$\Delta\varphi = 2\pi \text{OPD} / \lambda$$

$$\text{OPD} = t(n(s)) - n(m)$$

Optical Path Difference

To get quantitative information on the phase, monochromatic light should be used → coherent sources – lasers

→ quantitative phase microscopy or digital holographic microscopy

Digital holographic microscopy (DHM)

DHM is a modern technology allowing quantitative-phase imaging of phase objects, DHM is non-destructive, marker-free technique, derived from optical holography (OH). DHM includes two steps: RECORDING and RECONSTRUCTION .

RECORDING

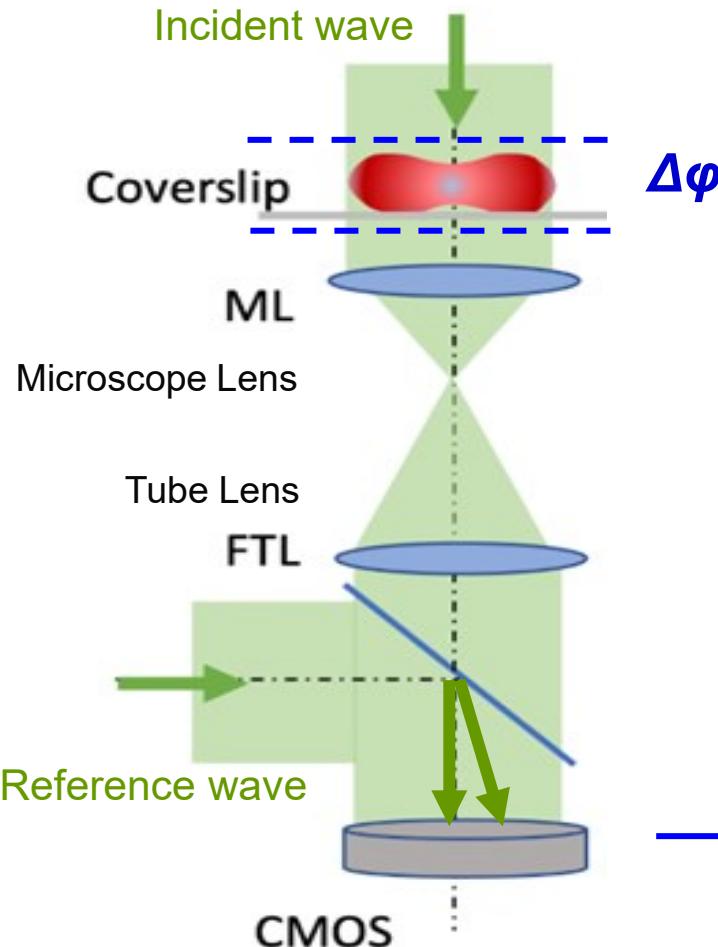
An interference pattern is recorded on the digital camera. A laser is split in two beams: one passes through the specimen and is called object beam and the other, called reference beam, is sent directly to the camera sensor. The interference between the object beam and reference beam forms an interference pattern called digital hologram. This contains both amplitude and phase information about the object.

RECONSTRUCTION

The digital hologram is digitally processed to reconstruct the amplitude and phase information about the object.

Note: the difference between OH and DHM consists in the RECONSTRUCTION step which is DIGITAL in the case of DHM, while it is ANALOGIC (optical) in OH.

DHM – Recording and Reconstruction principles



RECORDING

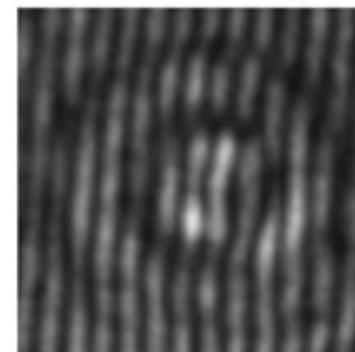
The digital hologram is obtained by interference between the ‘incident’ and ‘reference’ waves and recorded on the CMOS as **intensity pattern** (digital image) $I(x,y)$

RECONSTRUCTION

The intensity pattern I is processed numerically to calculate the phase shift $\Delta\varphi$ introduced by the specimen.

This processing includes Fourier transforms and image spatial filtering, to model light propagation.

$$I = I(\Delta\varphi)$$



Digital Hologram

The result of the numerical reconstruction is the phase shift $\Delta\varphi$:

$$\Delta\varphi(x,y) = 2\pi \text{OPD}(x,y) / \lambda \quad \text{with} \quad \text{OPD}(x,y) = h(x,y) (n_c - n_m)$$

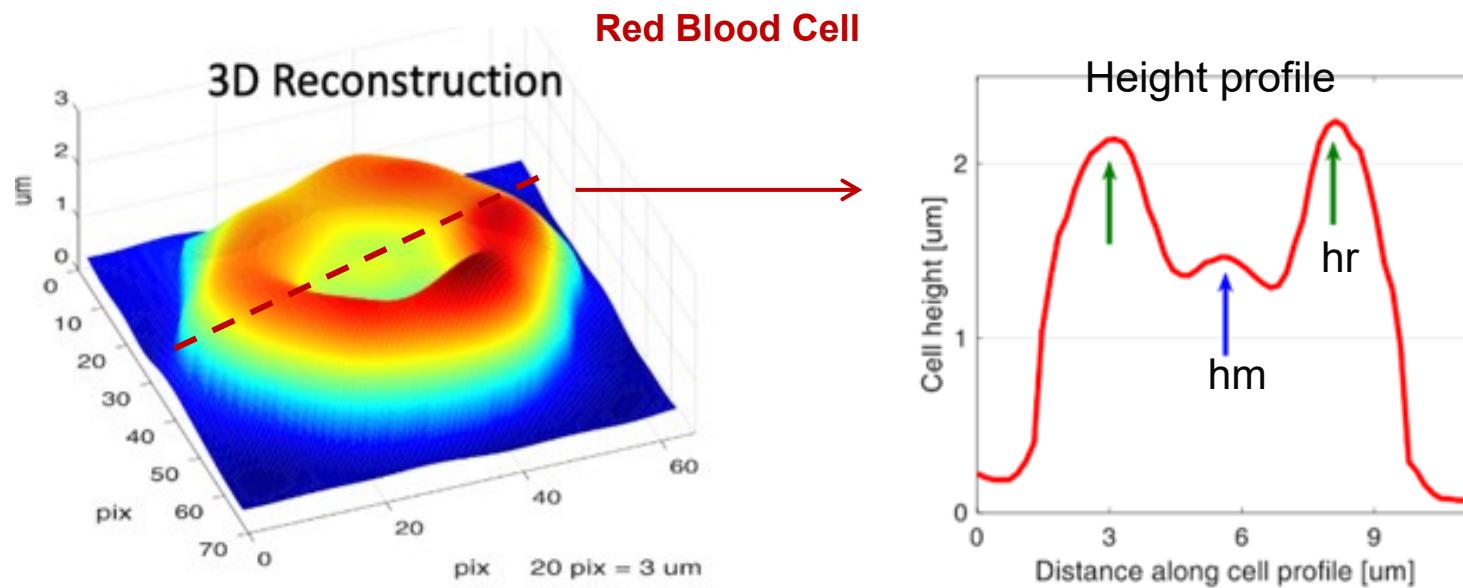
where OPD is the Optical Path Difference. OPD is a function of the geometrical path h , and the refractive indexes of the cell n_c and of the medium n_m

If we know the refractive index n_c and n_m , we can calculate the **height of the cell**, h :

$$h(x,y) = \lambda \Delta\varphi(x,y) / 2\pi (n_c - n_m)$$

the **volume of the cell**: $V = p \sum h(x,y)$, with p being the cell area / pixel

and other morphological parameters as **cell sphericity** $CS = hm / hr$



Example: characterization of *ex vivo* generated red blood cells (RBCs)

Cultured red blood cells (cRBCs) obtained under different conditions are compared with native RBC, from morphology point of view:

Morphology		CA	CV	CS	MCH	hm
		mean ± std	mean ± std	mean ± std	mean ± std	mean
Cells	n	um ²	um ³ (fL)	-	pg	um
nRBC	25	55,42 ± 9,2	95,2 ± 16,6	0,57 ± 0,1	25,24 ± 5	1,72 ± 0,4
cRBC ^{Plasma}	24	41,05 ± 14,4	125,5 ± 43,3	1,04 ± 0,1	31,17 ± 11,7	3,06 ± 0,6
cRBC ^{HPL}	29	70 ± 21,7	107,1 ± 37,8	0,671 ± 0,4	28,1 ± 10,9	1,53 ± 0,3

Cell Area	Cell Volume	Cell Sphericity	Mean Corpuscular Hemoglobin	Cell mean height
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Bernecker et al, *Cells* 2021, 10(3), 552; <https://doi.org/10.3390/cells10030552>

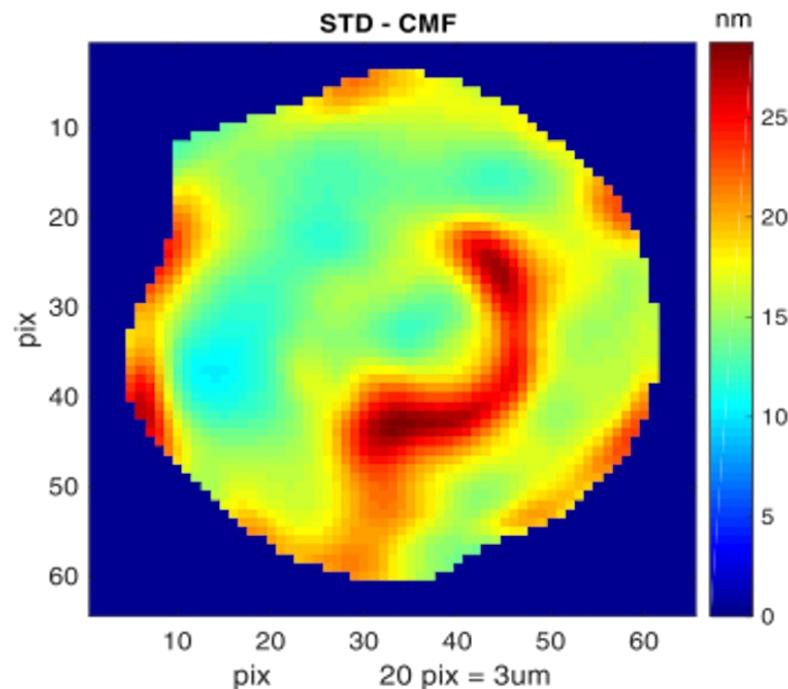
Bernecker et al, *Front Physiol* 2022, <https://doi.org/10.3389/fphys.2022.979298>

Moreover, DHM makes it possible to measure **cell membrane fluctuation (CMF)** which is related to the viscoelastic properties of the membrane.

To determine CMF one calculates, for each pixel within the cell, the fluctuation of the cell height in time at high acquisition rate, then the corresponding standard deviation for each pixel of the cell, STD_{pix_i} .

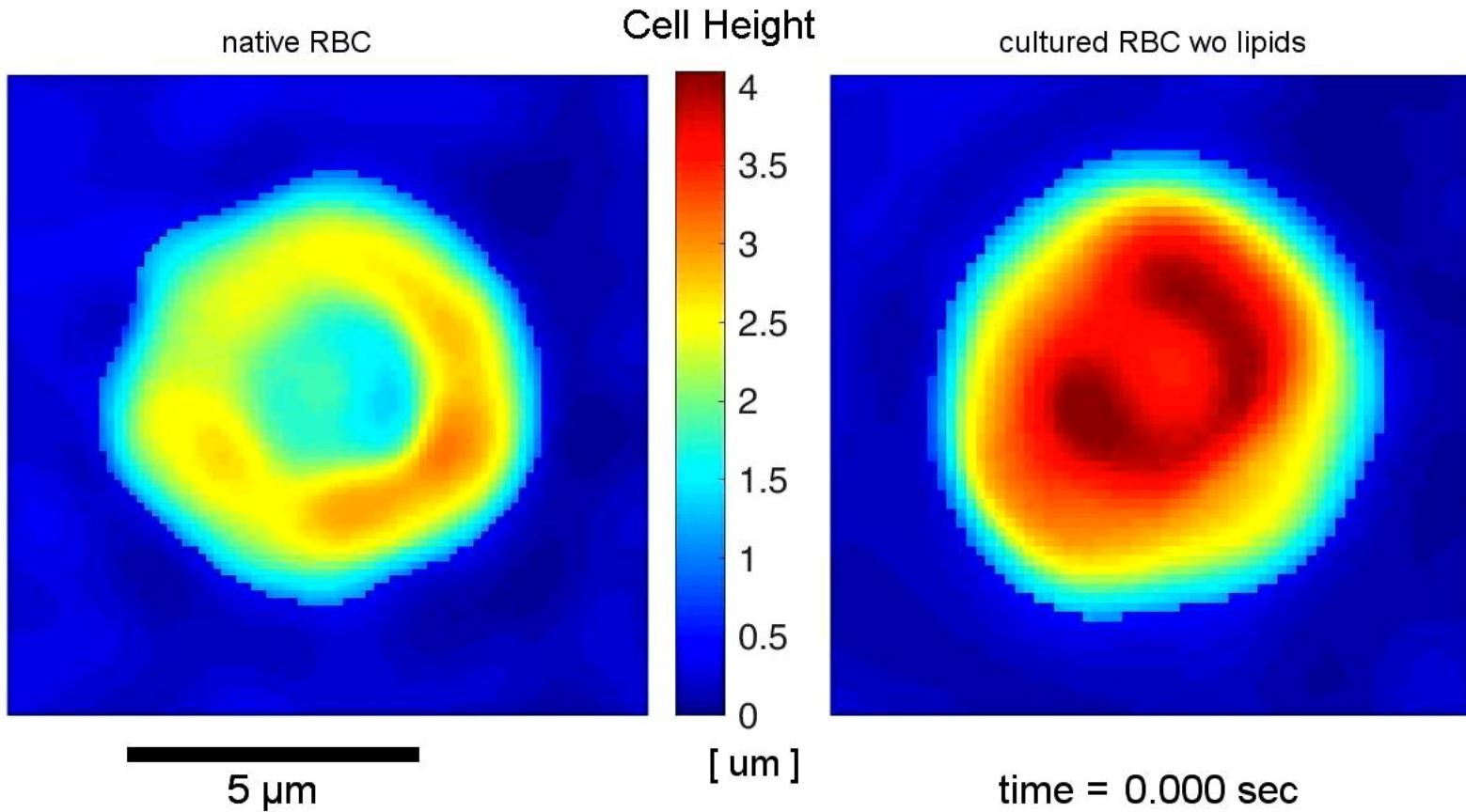
The CMF value is calculated as the mean of STD_{pix} :

$$CMF = 1/N_p \sum_i STD_{pix_i}$$



STD_{pix} - the standard deviation distribution over the cell area.
Note the nanometer sensitivity!

Cell Membrane Fluctuations

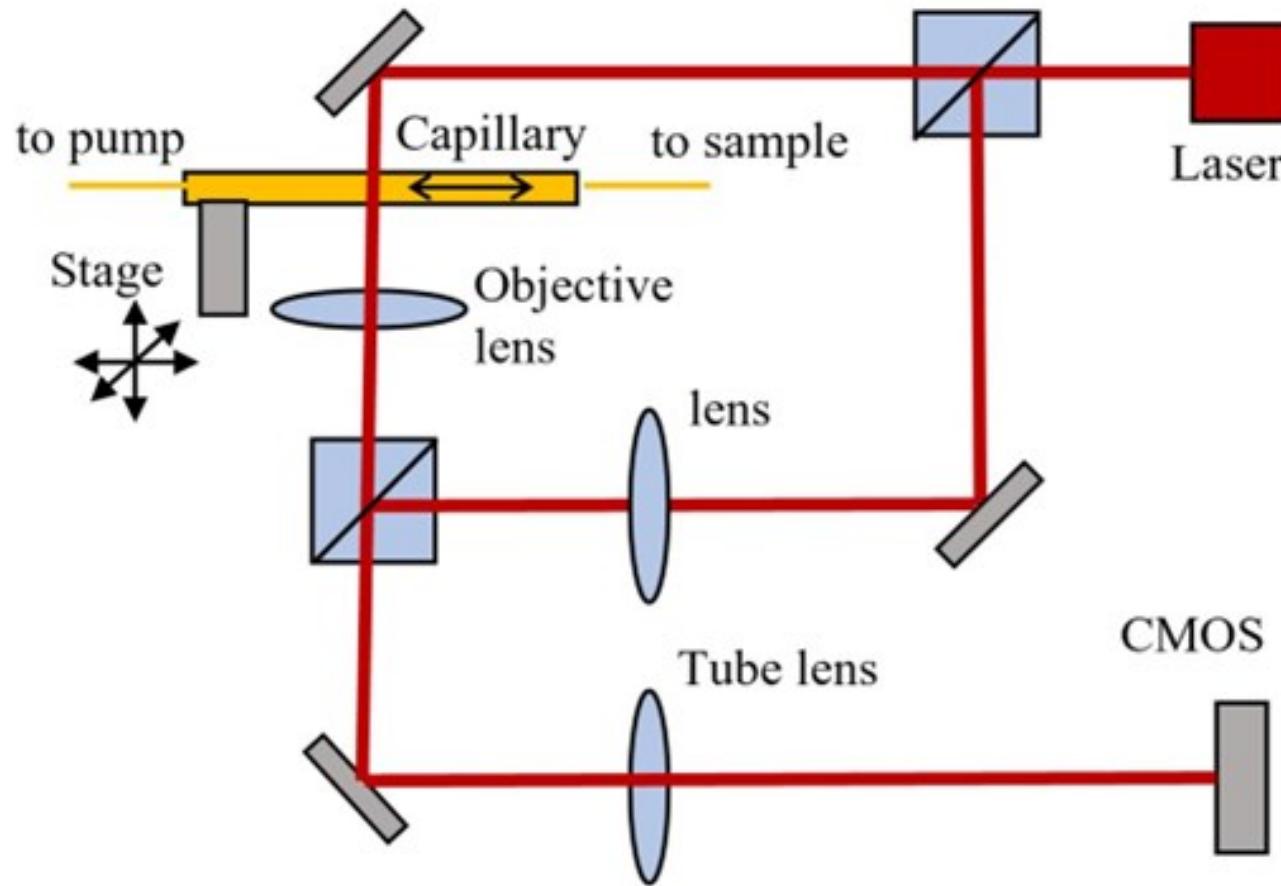


Video shows about 1 second of cell membrane fluctuations (height fluctuations) as measured with DHM

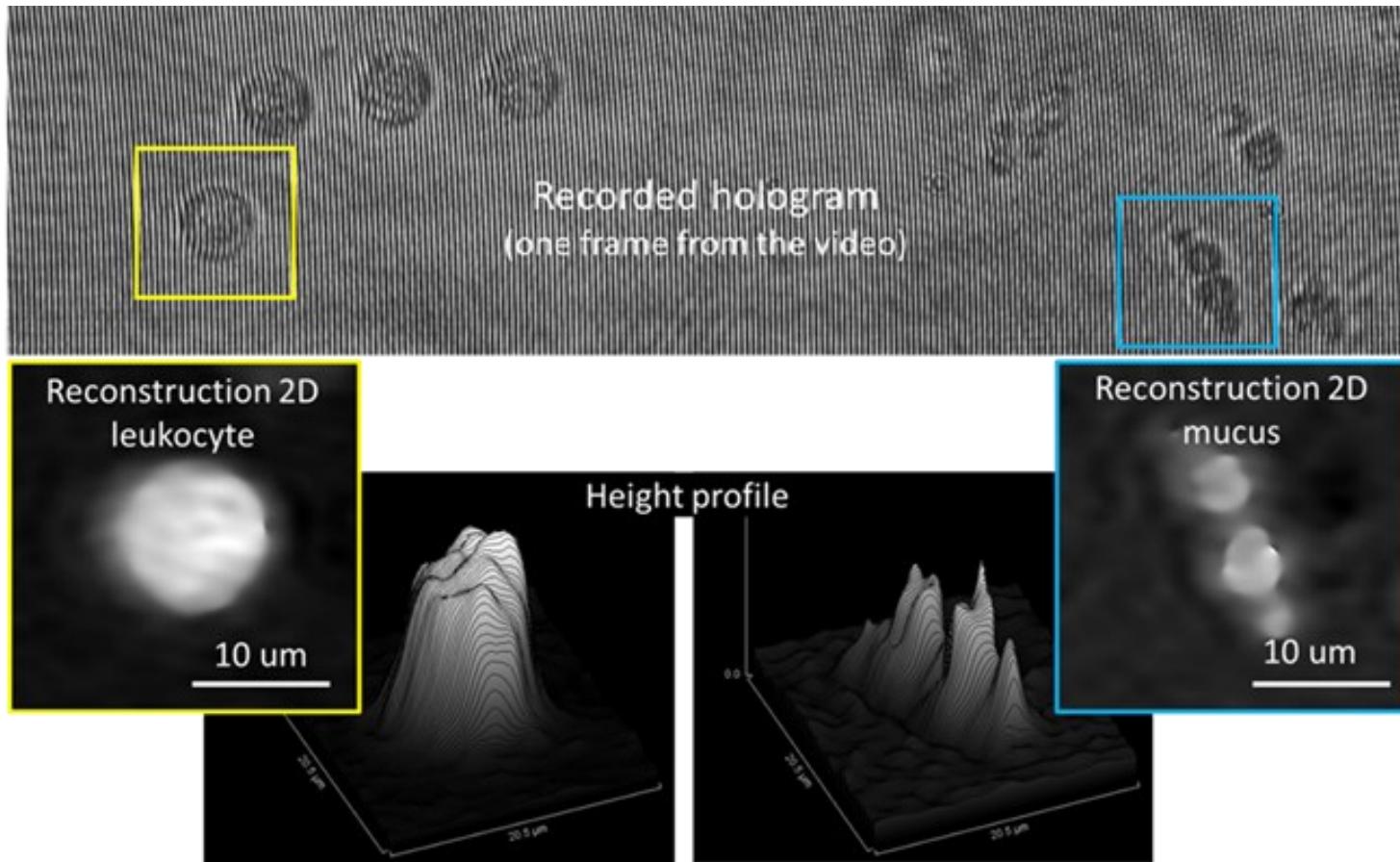
Example 2: Label-Free Analysis of Urine Samples by In-Flow Digital Holographic Microscopy

CNR-IOM & Alifax srl

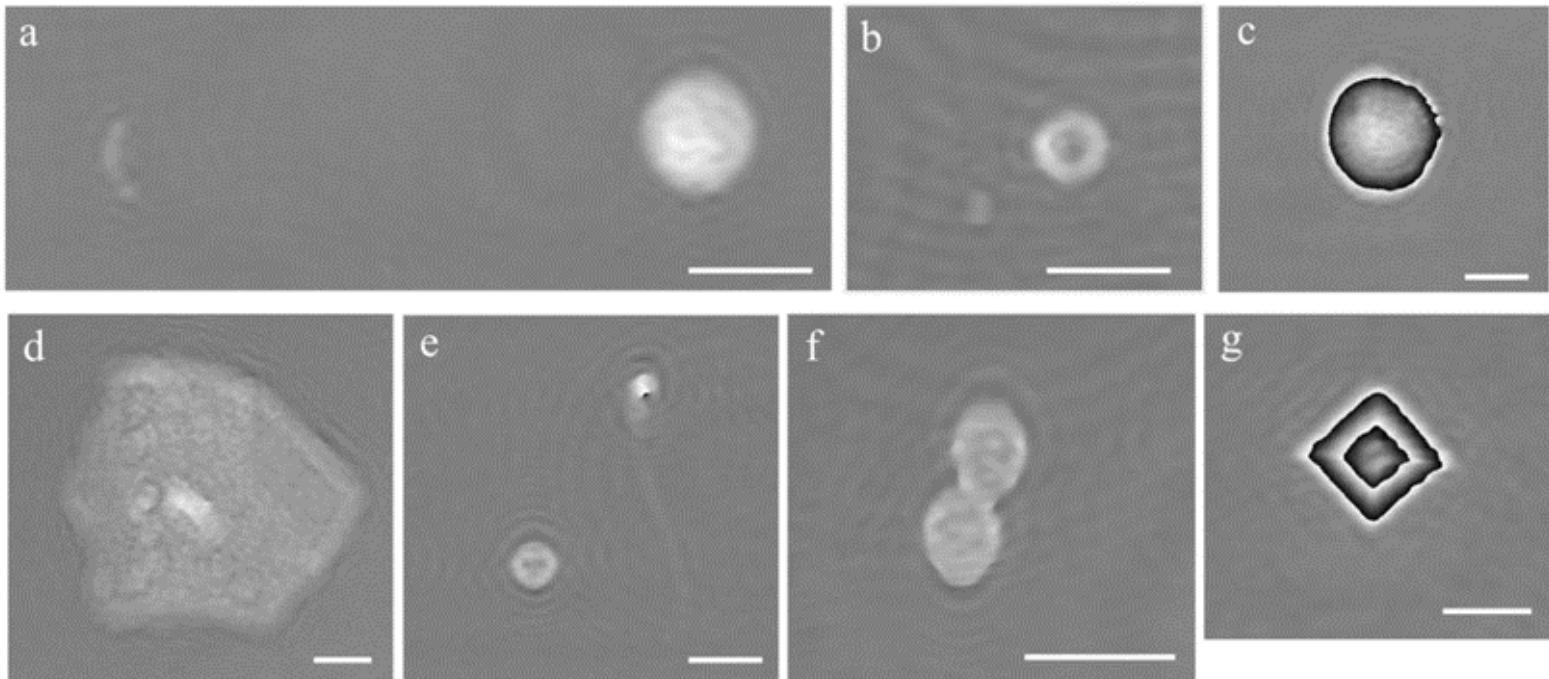
In-flow DHM



DHM Setup: laser beam (red) is split in two and recombined by two cube beam splitters, being directed to the CMOS; the urine is flowing in the capillary and imaged by the objective lens and tube lens on CMOS.



Example of recorded hologram (top image) and reconstructed phase images of a leukocyte (bottom-left yellow inset) and mucus (bottom-right blue inset) with their respective height profiles.



Examples of phase images for different components of the urine samples:

- a) Streptococcus spp chain (left) and leukocyte (right),
- b) *Escherichia coli* (left) and red blood cell (right),
- c) macrophage cell, d) epithelial (squamous) cell,
- e) red blood cell (down-left) and spermatozoa cell (up-right);
- f) fungi and g) crystal. Scale bar 10 μ m

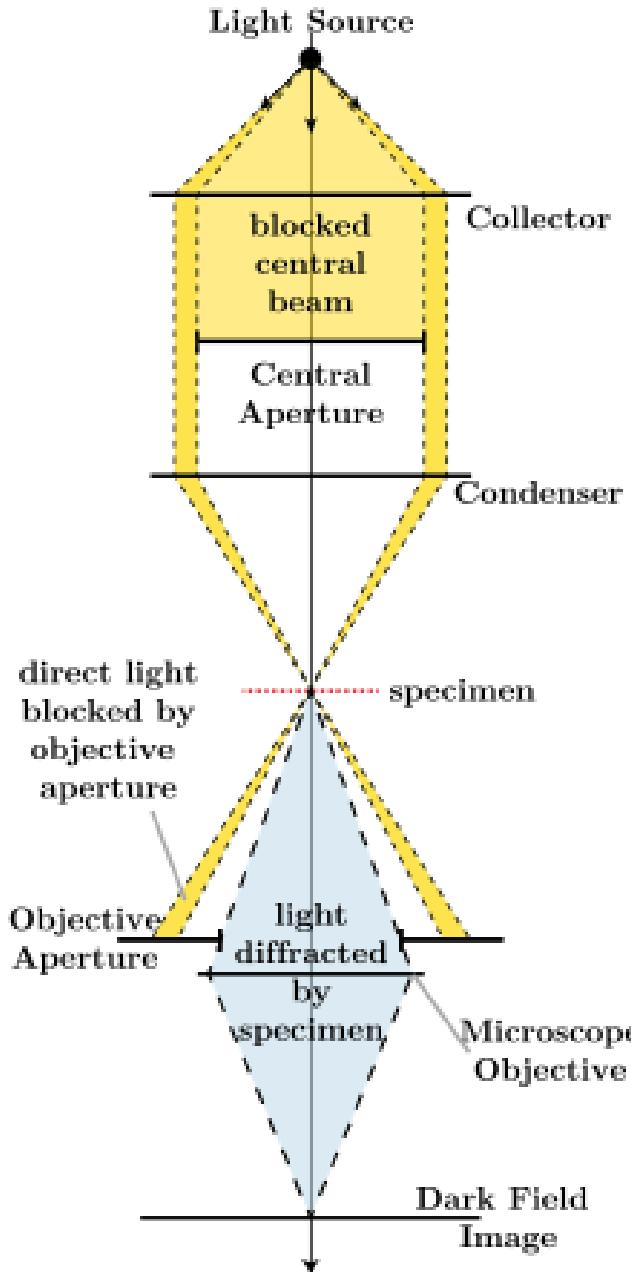
DHM advantages

- It is label free without the need for preliminary preparation steps such as fixation, membrane permeabilization or fluorophore incubation.
- The absence of fluorophores removes any concerns regarding cytotoxicity, phototoxicity through bleaching and molecular oxidation, or phototoxicity from high-power laser intensities.
- Allows measuring the height profile of the cell / refractive index / dry mass
- High axial sensitivity: $10 \text{ nm} \ll 400 \text{ nm}$ (the axial resolution in brightfield microscopy)

DHM drawbacks

- Low lateral resolution: $500 \text{ nm} \sim \lambda / \text{NA}$ instead of $0.5 \lambda / \text{NA}$
- Difficult to separate refractive index contribution from height contribution, if refractive index for sample is not known

Imaging nanoparticles by darkfield microscopy



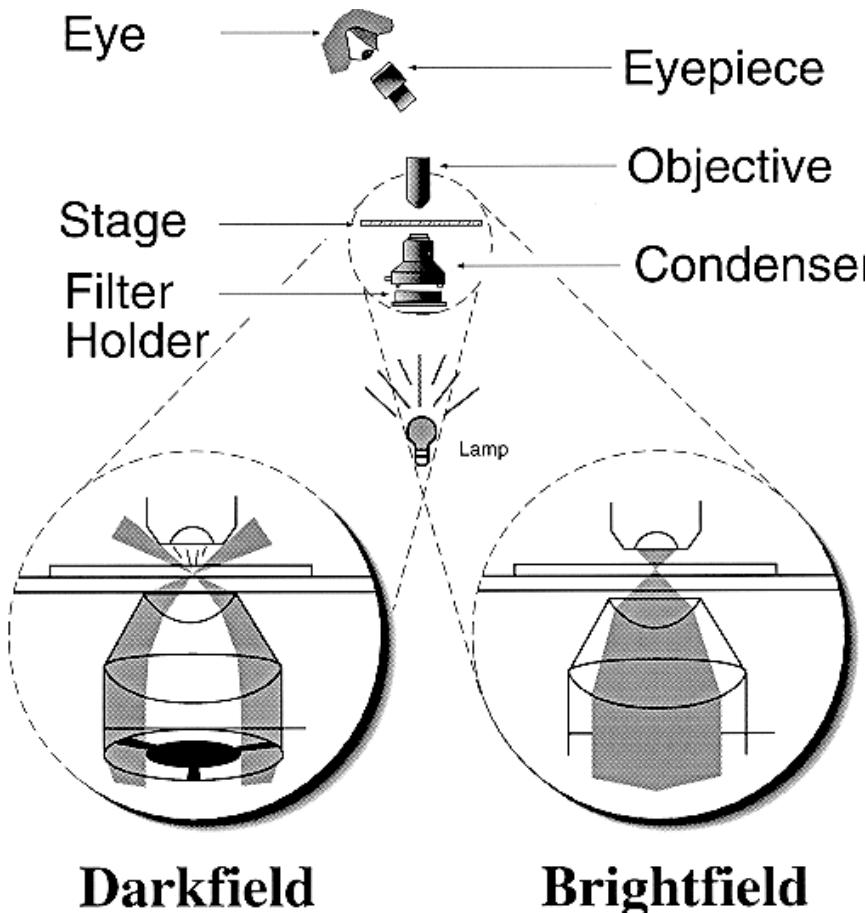
Darkfield Microscopy

The light at the apex of the cone is focused at the plane of the specimen; as this light moves past the specimen plane it spreads again into a hollow cone. The objective lens sits in the dark hollow of this cone; although the light travels around and past the objective lens, no rays enter it.

The entire field appears dark when there is no sample on the microscope stage; when a sample is on the stage, the light at the apex of the cone strikes it; the image is made only by those rays scattered by the sample and captured in the objective and it appears bright against the dark background.

Single nanoparticles can be imaged !

Darkfield vs Brightfield Microscopy (DM vs BM)



DM and BM use
different illumination schemes:

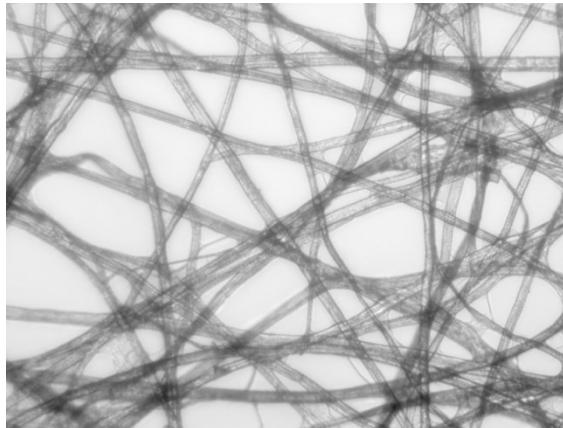
DM uses a darkfield stop to block the center of the beam of light, producing a **hollow cone of light** which does not directly enter the obj lens.

In contrast, a **solid cone of light** illuminates and enters the obj lens in **BM**.

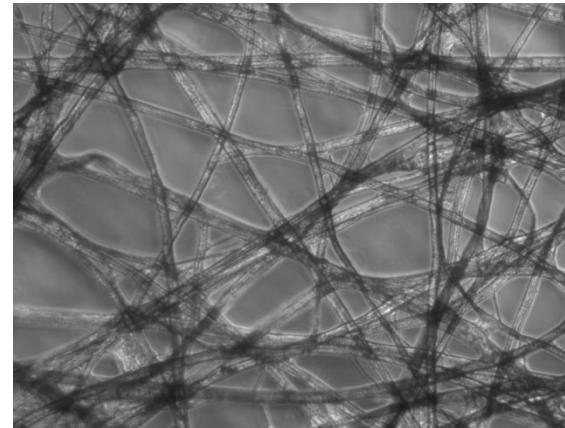
Comparison btw different microscopy techniques

Sample: tissue paper micrograph

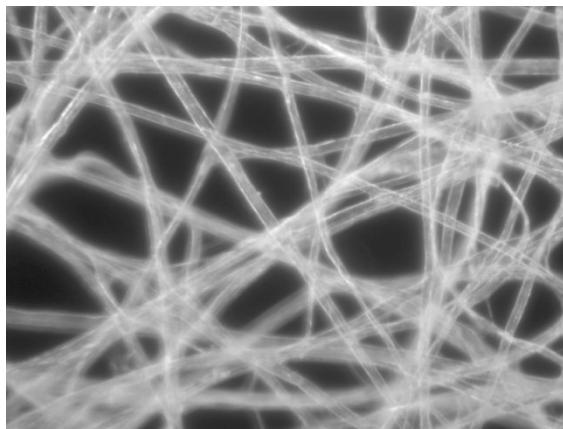
Brightfield



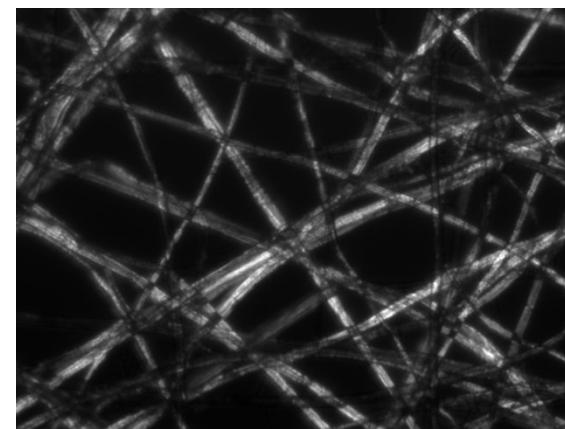
Phase Contrast



Darkfield



Polarization



?

Polarized light microscopy

Polarized light Microscopy (PM)

de Séarnmont Compensator Configuration

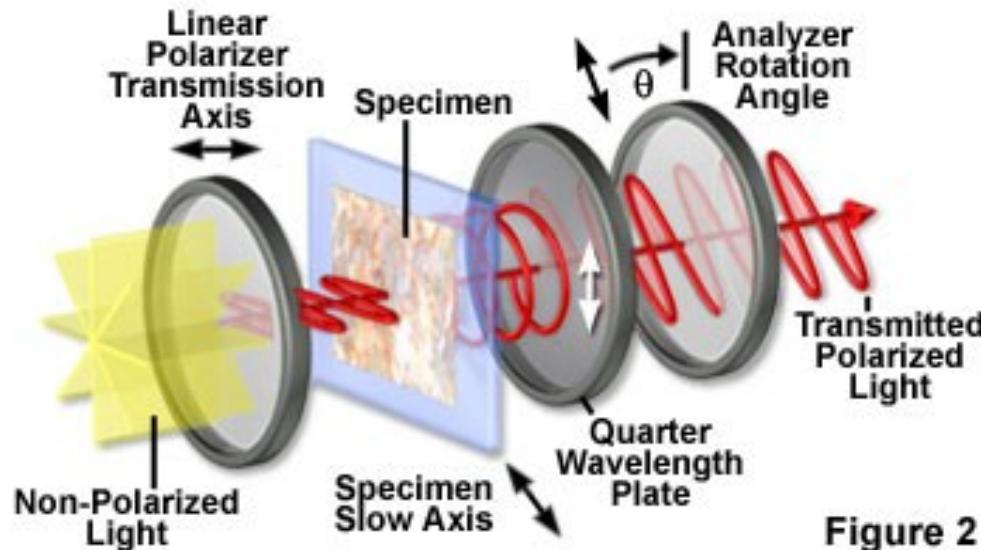


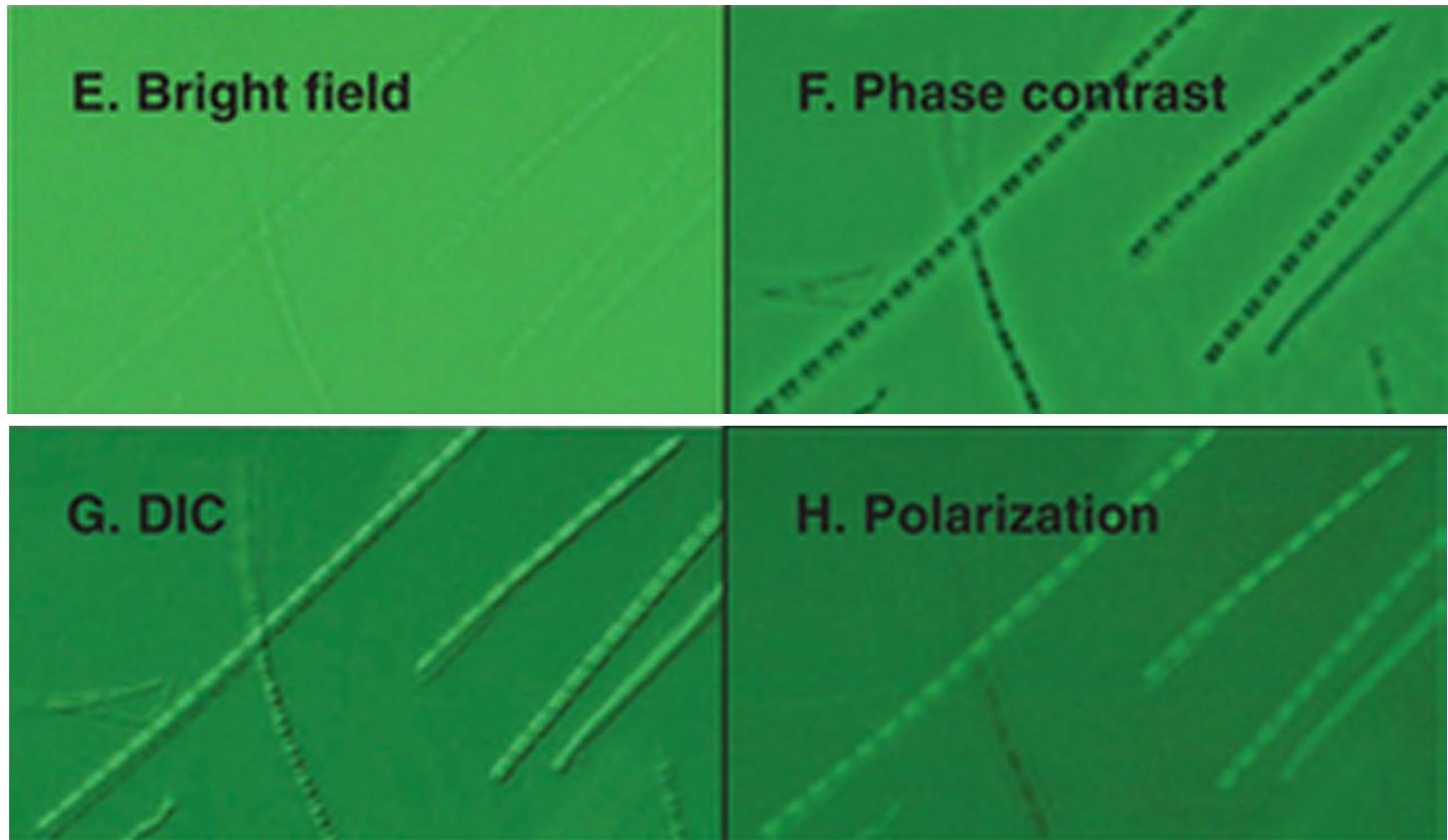
Figure 2

PM requires birefringent samples and polarized light.

A linear polarizer is used to illuminate the sample with linear polarized light.

Due to its **birefringence** the sample changes the **light polarization**. This change is detected with an analyzer + compensator.

Applications: mineralogy, plant biology, blood cells

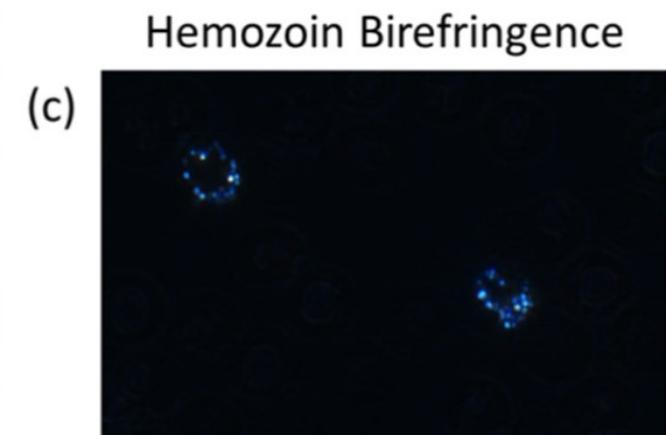
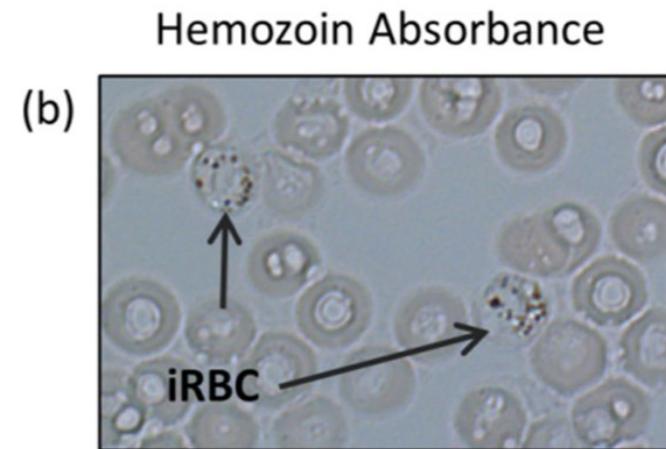
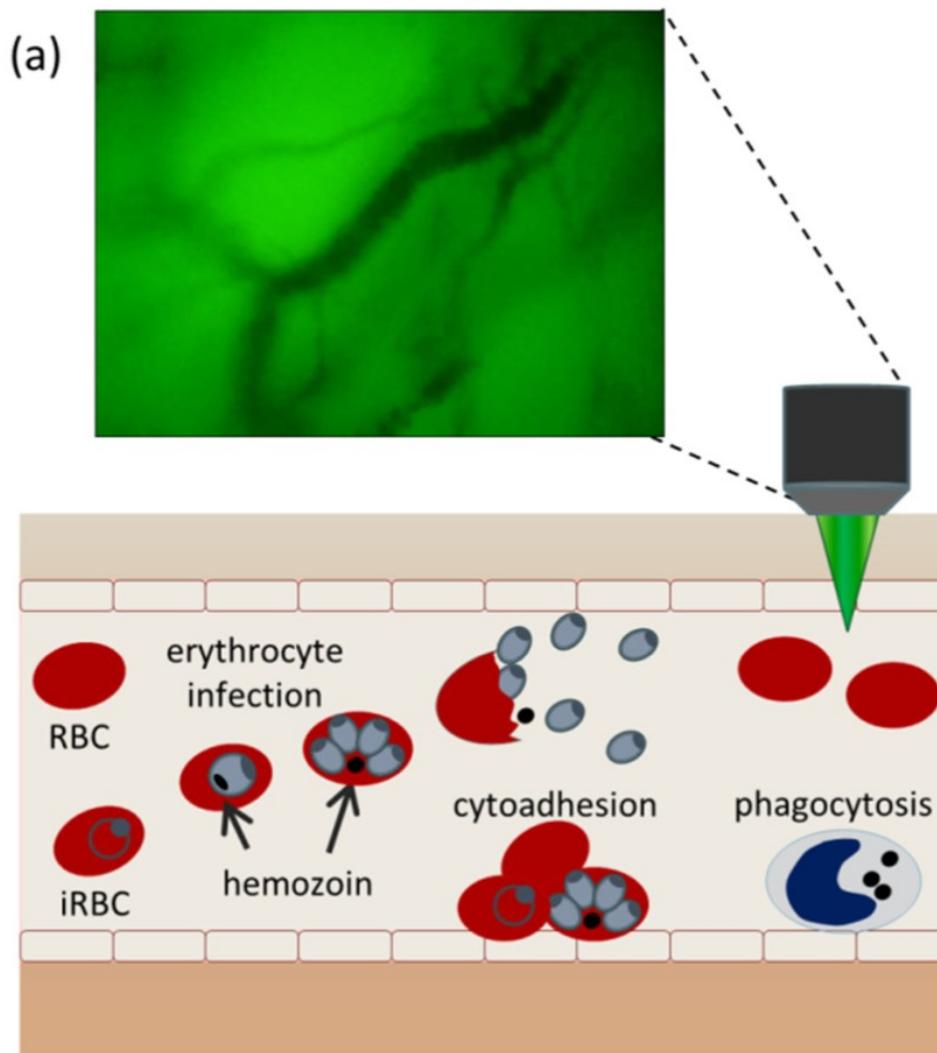


E–H, Micrographs of myofibrils isolated from skeletal muscle.

Contrast methods include bright field (**E**), phase contrast (**F**), differential interference contrast (**G**), and polarization (**H**).

The A-bands, consisting of parallel thick filaments of myosin (see Fig. 39-3), appear as dark bands with phase contrast and are birefringent (either bright or dark, depending on the orientation) with polarization.

Example: Hemozoin imaging for malaria detection



J. Burnett, J. Carns, and R. Richards-Kortum,
"In vivo microscopy of hemozoin: towards a needle free diagnostic for malaria,"
Biomed. Opt. Express **6**, 3462-3474 (2015).

1. Optical Microscopy – Physical Principles – 9h

- 1.1. Basics (Image formation, magnification, resolution, image quality)
- 1.2. Digital camera (image acquisition, formats, properties)
- 1.3. Phase imaging (qualitative and quantitative)
- 1.4. Dark field and Polarization microscopy

Next lecture

- 1.5. Non Linear Optical Microscopy
- 1.6. Photoacoustic Microscopy
- 1.7. Super resolution microscopy – STED, PALM/STORM, MINFLUX