

## 993SM - Laboratory of Computational Physics II week October 2, 2023

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# The Numerov's method for the 1D Schroedinger equation

#### codes & notes from: **prof. Paolo Giannozzi (UniUD)**

#### "Numerical methods in Quantum Mechanics"

<u>https://www.fisica.uniud.it/~giannozz/Corsi/MQ/LectureNotes/mq.pdf</u> <u>https://www.fisica.uniud.it/~giannozz/Corsi/MQ/Software/F90/harmonic0.f90</u> <u>https://www.fisica.uniud.it/~giannozz/Corsi/MQ/Software/F90/harmonic1.f90</u>

#### Note:

#### we choose a problem that can be solved exactly (analytically) to check the reliability of the code and the possible problems

### the harmonic oscillator

#### harmonic oscillator: classical

Force F = -Kx

$$m\frac{d^2x}{dt^2} = -Kx$$

Potential 
$$V(x) = V(-x) = \frac{1}{2}Kx^2$$
  
A solution  $x(t) = x_0 \sin(\omega t)$  with  $\omega = \sqrt{\frac{K}{m}}$ 

probability  $\rho(x)dx$  to find the mass between x and x + dx:

$$\rho(x)dx \propto \frac{dx}{v(x)}$$

Since  $v(t) = x_0 \omega \cos(\omega t) = \omega \sqrt{x_0^2 - x_0^2 \sin^2(\omega t)}$ , we have

$$\rho(x) \propto \frac{1}{\sqrt{x_0^2 - x^2}}$$
for  $|x| < x_0$ ; 0 elsewhere

### harmonic oscillator: 1D Schroedinger eq.

In standard notation:

$$\frac{d^2\psi}{dx^2} = -\frac{2m}{\hbar^2} \left( E - \frac{1}{2}Kx^2 \right) \psi(x) \qquad \text{with} \quad \omega = \sqrt{\frac{K}{m}}$$

Note the symmetry of the potential: V(-x) = V(x)

Defining:

$$\xi = \left(\frac{m\omega}{\hbar}\right)^{1/2} x = \left(\frac{mK}{\hbar^2}\right)^{1/4} x = \frac{x}{\lambda} \text{ and } \varepsilon = \frac{E}{\hbar\omega}$$

we rewrite the eq. in adimensional units:

$$\frac{d^2\psi}{d\xi^2} = -2\left(\varepsilon - \frac{\xi^2}{2}\right)\psi(\xi)$$

### harmonic oscillator: 1D Schroedinger eq.

Exact solution (analytical):  $\psi_n(\xi) = H_n(\xi) e^{-\xi^2/2}$ 

odd or even functions n nodes and the same parity as nHermite polynomials.  $H_n(\xi)$ 

The lowest-order Hermite polynomials are

$$H_0(\xi) = 1$$
,  $H_1(\xi) = 2\xi$ ,  $H_2(\xi) = 4\xi^2 - 2$ ,  $H_3(\xi) = 8\xi^3 - 12\xi$ .

corresponding to discretized energies:

$$\varepsilon = n + \frac{1}{2} \quad \square \quad E_n = \left(n + \frac{1}{2}\right) \hbar \omega \quad n = 0, 1, 2, \dots$$

n is a non-negative integer

### harmonic oscillator: 1D Schroedinger eq.

#### Exact solution (plots):



# The Numerov's method

To solve: (g(x), s(x) given)

$$\frac{d^2y}{dx^2} = -g(x)y(x) + s(x)$$

idea: Taylor expansion of y(x), g(x), s(x), followed by a few manipulations

(details in the notes by prof. Giannozzi)

$$\begin{array}{ll} y_{n+1} \left[ 1 + g_{n+1} \frac{(\Delta x)^2}{12} \right] &= 2y_n \left[ 1 - 5g_n \frac{(\Delta x)^2}{12} \right] - y_{n-1} \left[ 1 + g_{n-1} \frac{(\Delta x)^2}{12} \right] \\ &+ (s_{n+1} + 10s_n + s_{n-1}) \frac{(\Delta x)^2}{12} + O[(\Delta x)^6] \end{array}$$

allows to obtain  $y_{n+1}$  starting from  $y_n$  and  $y_{n-1}$ , and recursively go on...

**ID Schroedinger equation:** a form suitable for Numerov's method

The Scroedinger eq.:

$$-\frac{\hbar^2}{2m}\frac{d^2\psi}{dx^2} + V(x)\psi(x) = E\psi(x),$$

has the form:

$$\frac{d^2y}{dx^2} = -g(x)y(x) + s(x)$$

with:

$$g(x) = (2m/\hbar^2)[E - V(x)]$$
 and  $s(x) = 0$ 

## **ID Schroedinger equation:** harmonic oscillator – 1

In adimensional units:

$$\frac{d^2\psi}{d\xi^2} = -2\left(\varepsilon - \frac{\xi^2}{2}\right)\psi(\xi)$$

has the form:

$$\frac{d^2y}{dx^2} = -g(x)y(x) + s(x)$$

with:

$$g(x) = 2\left(\epsilon - \frac{x^2}{2}\right)$$
 and  $s(x) = 0$ 

## **ID Schroedinger equation:** harmonic oscillator - 2

Since s(x) = 0, the Numerov's formula reduces to :

$$y_{n+1} \left[ 1 + g_{n+1} \frac{(\Delta x)^2}{12} \right] = 2y_n \left[ 1 - 5g_n \frac{(\Delta x)^2}{12} \right] - y_{n-1} \left[ 1 + g_{n-1} \frac{(\Delta x)^2}{12} \right] + O[(\Delta x)^6]$$
  
Defining:  $f_n \equiv 1 + g_n \frac{(\Delta x)^2}{12}$ 

we rewrite Numerov's formula as

$$y_{n+1} = \frac{(12 - 10f_n)y_n - f_{n-1}y_{n-1}}{f_{n+1}}$$

The value of the energy is now hidden into  $g_n$  and  $f_n$ .

**1D Schroedinger equation:** harmonic oscillator - 3  $y_{n+1} = \frac{(12 - 10f_n)y_n - f_{n-1}y_{n-1}}{f_{n+1}}$ (\*)

The symmetry of the potential and the parity of the (still unknown) solutions allows to simplify the choice of the starting points

- *n* even choose  $y_0 = 0$  and whatever  $y_1$  you want
- n odd choose whatever  $y_0$  (finite) you want;

 $y_1$  is determined by Numerov's formula :

since  $f_{-1} = f_1$  by symmetry, and  $y_{-1} = y_1$ ; put into (\*) and obtain :

$$y_1 = \frac{(12 - 10f_0)y_0}{2f_1}$$

# harmonic0

The code prompts for some input data:

- the limit  $x_{\text{max}}$  for integration (typical values:  $5 \div 10$ );
- the number N of grid points (typical values range from hundreds to a few thousand); note that the grid point index actually runs from 0 to N, so that  $\Delta x = x_{\text{max}}/N$ ;
- the name of the file where oroutput file written;
- the required number n of nodes (the code will stop if n is negative).

Finally the code prompts for a trial energy. You should answer  $0^{(*)}$  in order to search for an eigenvalue with n nodes. The code will start iterating on the energy.

(\*) It is however possible to specify an energy to force the code to perform an integration at fixed energy useful for testing purposes

new value of the number of nodes answer -1 to stop

### harmonic0

Max value for x (typical value: 10) > 1 Number of grid points (typically a few output file name > 10_500.dat nodes (type -1 to stop) > 0	l0 hundreds) > 500	
1  25.000000000000000000000000000000000000	13	0
2 12.50000000000000	7	0
3 6.250000000000000	3	0
4 3.1250000000000000	2	0

#### Output file contains: $x, \psi(x), |\psi(x)|^2, \rho_{\rm cl}(x), V(x)$

#	х	y(x)
-10	.000	-0.41559380E+11
-9	.980	-0.34101628E+11
-9	.960	-0.27993466E+11
-9	.940	-0.22988666E+11
-9	.920	-0.18886279E+11

y(x)^2	classical p(x)	V
0.17271820E+22	0.0000000E+00	50.000000
0.11629210E+22	0.0000000E+00	49.800200
0.78363416E+21	0.0000000E+00	49.600800
0.52847878E+21	0.0000000E+00	49.401800
0.35669155E+21	0.0000000E+00	49.203200



#### harmonic0





divergences! ... but can be hidden!

#### harmonic0



harmonic0

<u>???</u>

divergences! ... but can be hidden!

what's next?
integration
forward (from x=0 to x\_max) and
backward (from x\_max to 0)
=> harmonic1



harmonic1

harmonic0

# within the Numerov's method: the "shooting" step

