

20.09.2023

$$\begin{aligned}
 & \| f(u_n - u) \|_{L^1([0,T], L^2(\mathbb{R}))} \leq \\
 & \leq \| u \|_{L^2([0,T], L^4(\mathbb{R}^3))} \| u_n - u \|_{L^2([0,T], L^4(\mathbb{R}))} \\
 & \quad \text{with } \| u \|_{L^2([0,T], L^4(\mathbb{R}^3))} \leq 1
 \end{aligned}$$

$$\| \mathcal{S}_{\varepsilon_n} * (u_n - u) \|_{L^2([0,T], L^4(\mathbb{R}))}$$

$$= \left| \int_{\mathbb{R}^3} \mathcal{S}_{\varepsilon_n}(y) (u_n(x-y) - u(x-y)) dy \right|_{L^2([0,T], L^4(\mathbb{R}))}$$

$$\leq \int_{\mathbb{R}^3} \mathcal{S}_{\varepsilon_n}(y) \| u_n(x-y) - u(x-y) \|_{L^2([0,T], L^4(\mathbb{R}))} dy$$

$$\text{supp } \mathcal{S} \subseteq D_{\mathbb{R}^3}(0, 1) \quad \varepsilon_n < 1$$

$$\text{supp } \mathcal{S}_{\varepsilon_n} \subseteq D_{\mathbb{R}^3}(0, 1)$$

$$\leq \int_{\mathbb{R}^3} \mathcal{S}_{\varepsilon_n}(y) \| u_n - u \|_{L^2([0,T], L^4(\mathbb{R}^3))} dy$$

$$\begin{aligned}
 & x-y \in K-y \\
 & x-y \in \bigcup_{|y| \leq 1} K-y \subset \bigcap_{|y| \leq 1} K
 \end{aligned}$$

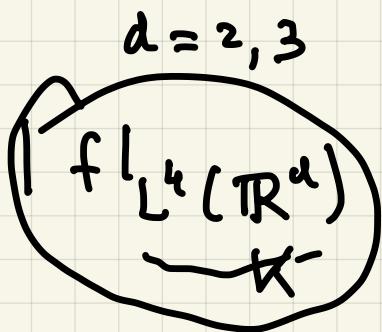
$$= \|u_n - u\|_{L^2([0,T], L^4(K))} \xrightarrow{n \rightarrow \infty} 0$$

$H^1(\mathbb{R}^d)$

$u_n \rightarrow u$  in  $L^2([0,T], L^4(K))$

$u_n \leftarrow u \rightarrow 0$   $\forall K \subset \subset \mathbb{R}^3$  and

$T > 0$ .



$$\|f\|_{L^4(\mathbb{R}^d)} \leq C \|f\|_{L^2(\mathbb{R}^d)}^{1-\frac{4}{d}} \cdot \frac{\|\nabla f\|_{L^2(\mathbb{R}^d)}^{\frac{4}{d}}}{K}$$

$$\chi \in C_c^\infty(\mathbb{R}^d, [0,1])$$

$$\chi \equiv 1 \text{ in } K, \quad \Omega \stackrel{\text{def}}{=} \text{supp } \chi$$

$$\|\nabla \chi\|_{L^\infty(\mathbb{R}^d)} \leq \frac{1}{\Omega} \mathbb{R}^d$$

$$\nabla u \in L^2([0,T] \times \mathbb{R}^d)$$

$$f \in H^1(\mathbb{R}^d)$$

$$|\chi f|_{L^4(\mathbb{R}^d)} \leq C |\chi f|_{L^2(\mathbb{R}^d)}^{1-\frac{4}{d}} \cdot \|\nabla(\chi f)\|_{L^2(\mathbb{R}^d)}^{\frac{4}{d}}$$

$$|f|_{L^4(K)} \leq C |f|_{L^2(\Omega)}^{1-\frac{4}{d}} \cdot$$

$$\cdot \left( |\nabla \chi f|_{L^2(\mathbb{R}^d)} + \|\chi \nabla f\|_{L^2(\mathbb{R}^d)} \right)$$

$$|f|_{L^4(\kappa)} \leq C |f|_{L^2(\Omega)}^{1-\frac{d}{4}} |f|^{\frac{d}{4}}_{H^1(\mathbb{R}^d)}$$

$$\begin{aligned} & |u - u_m|_{L^2([0, \tau], L^4(\kappa))} = \\ &= |u - u_m|_{L^4(\kappa)} |_{L^2([0, \tau])}^{\frac{4-d}{4}} \\ &\lesssim |(u - u_m)|_{H^1(\mathbb{R}^d)}^{\frac{d}{4}} |u - u_m|_{L^2(\Omega)}^{\frac{4-d}{4}} |_{L^2([0, \tau])} \end{aligned}$$

$$\frac{1}{2} = \frac{4-d}{8} + \frac{d}{8} \quad d=2, 3$$

$$\begin{aligned} &\lesssim |u - u_m|_{H^1(\mathbb{R}^d)}^{\frac{d}{4}} |_{L^{\frac{8}{4-d}}([0, \tau])} \\ &\quad |u - u_m|_{L^2(\Omega)}^{\frac{4-d}{4}} |_{L^{\frac{8}{4-d}}([0, \tau])} \\ &= |(u - u_m)|_{H^1(\mathbb{R}^d)}^{\frac{d}{4}} |_{L^2([0, \tau])} \\ &\quad |u - u_m|_{L^2(\Omega)}^{\frac{4-d}{4}} |_{L^2([0, \tau])} \\ &= |u - u_m|_{L^2([0, \tau], H^1(\mathbb{R}^d))}^{\frac{d}{4}} \cdot \\ &\quad |u - u_m|_{L^2([0, \tau] \times \Omega)}^{\frac{4-d}{4}} \xrightarrow{n \rightarrow +\infty} 0 \end{aligned}$$

$$\|u\|_{L^2([0,T], H^1(\mathbb{R}^d))} \leq$$

$$\underbrace{\|u\|_{L^2([0,T], L^2(\mathbb{R}^d))}} + \underbrace{\|\nabla u\|_{L^2([0,T] \times \mathbb{R}^d)}} \leq \|u_0\|_{L^2} \\ \leq \sqrt{T} \underbrace{\|u\|_{L^\infty([0,T], L^2(\mathbb{R}^d))}} \leq \|u_0\|_{L^2}$$

$$\int_0^t \langle \operatorname{div} (\mathcal{G}_{\varepsilon_n} * u_n \otimes u_n), \psi \rangle_{L_x^2} dt \\ \xrightarrow{n \rightarrow +\infty} \int_0^t \langle \operatorname{div}(u \otimes u), \psi \rangle_{L_x^2} ds$$

Remark The above  $u$  satisfies the energy inequality not only with respect to 0, that is

$$\|u(t)\|_{L^2(\mathbb{R}^3)}^2 + 2 \int_1^t \|\nabla u(t')\|_{L^2(\mathbb{R}^3)}^2 dt' \leq \|u(1)\|_{L^2(\mathbb{R}^3)}^2$$

$$\forall t > 1 \geq 0$$

Weak solutions  $u$  of this type are called Leray-Hopf solutions.

# Kato + theory

$$\begin{cases} \partial_t u - \Delta u = -\bar{P} \operatorname{div}(u \otimes u) \\ u|_{t=0} = u_0 \end{cases}$$

$\mathcal{Q}(u, u)$

$\mathcal{Q}(u, v)$  is a symmetrizer of  $-\bar{P} \operatorname{div}(u \otimes v)$

$$\mathcal{Q}(u, v) = -\frac{1}{2} \bar{P} (u \otimes v + v \otimes u)$$

Suppose  $u$  and  $v$  are defined in spacetime such that  $\mathcal{Q}(u, v)$  satisfies the hypothesis of the external force  $f$  in the linear heat equation  $\partial_t w - \Delta w = f$

$$\begin{cases} \partial_t B(u, v) - \Delta B(u, v) = \mathcal{Q}(u, v) \\ B(u, v)|_{t=0} = 0 \end{cases}$$

$$B(u, v)|_{t=0} = 0$$

$$B(u, v) = \int_0^t e^{(t-t')\Delta} \mathcal{Q}(u, v)(t') dt'$$

Then NS

$$\begin{cases} \partial_t u - \Delta u = Q(u, u) \\ u|_{t=0} = u_0 \in H^{\frac{d}{2}-1}(\mathbb{R}^d) \end{cases}$$

Duhamel formula

$$u = e^{t\Delta} u_0 + B(u, u)$$

$$X_T = L^4([0, T], \dot{H}^{\frac{d-1}{2}}(\mathbb{R}^d))$$

Recall

Lemma  $B : X \times X \rightarrow X$  b, linear

bounded,  $X$  Banach space.

$\alpha < \frac{1}{4\|B\|}$ . Then  $\forall x_0 \in \overline{D}(0, \alpha)$

$\exists ! x \in \overline{D}(0, 2\alpha)$  s.t

$$x = x_0 + B(x, x).$$

Remark NS is "invariant" by the scaling  $\lambda > 0$

$$u \rightarrow u_\lambda(t, x) = \lambda u(\lambda^2 t, \lambda x)$$

$$\dot{H}^{\frac{d}{2}-1}(\mathbb{R}^d)$$

$$\begin{aligned} u(0, x) &= u_0(x) \\ u_\lambda(0, x) &= \lambda u_0(\lambda x) \end{aligned}$$

$$\|\lambda u_0(\lambda \cdot)\|_{\dot{H}^{\frac{d}{2}-1}(\mathbb{R}^d)} = \|u_0\|_{\dot{H}^{\frac{d}{2}-1}(\mathbb{R}^d)}$$

$$E_x \quad d=2$$

$$\|\lambda u_0(\lambda \cdot)\|_{L^2(\mathbb{R}^2)} = \lambda \|u_0(\lambda \cdot)\|_{L^2(\mathbb{R}^2)}$$

$$= \lambda \lambda^{-\frac{d}{2}} \|u_0\|_{L^2(\mathbb{R}^2)} = \|u_0\|_{L^2(\mathbb{R}^2)}$$

$$d = 3$$

$$\|\lambda u_0(\lambda \cdot)\|_{\dot{H}^{\frac{1}{2}}(\mathbb{R}^3)} = \|u_0\|_{\dot{H}^{\frac{1}{2}}(\mathbb{R}^3)}$$

$$\dot{H}^{\frac{1}{2}}(\mathbb{R}^3) \hookrightarrow L^3(\mathbb{R}^3)$$

$$\frac{1}{2} - \frac{\frac{1}{2}}{3} = \frac{1}{2} - \frac{1}{6} = \frac{1}{3}$$

$$\left( \|\lambda u_0(\lambda \cdot)\|_{L^3(\mathbb{R}^3)} = \|u_0\|_{L^3(\mathbb{R}^3)} \right)$$

$$\left( = \lambda \|u_0(\lambda \cdot)\|_{L^3(\mathbb{R}^3)} = \cancel{\lambda} \cancel{\frac{\|u_0\|_{L^3(\mathbb{R}^3)}}{3}} \|u_0(\lambda \cdot)\|_{L^3(\mathbb{R}^3)} \right)$$

$u_0$

$T(u_0)$

$T$

$[0, \frac{T}{\lambda^2})$

$u [0, T)$

$u_0$

$u_\lambda(t, x) = u(\lambda^2 t, \lambda x) \quad \lambda$

$u_{0\lambda}$

$$T(\|u_0\|_v) = 1$$

$$\|u(t)\|_{\dot{H}^{\frac{1}{2}}(\mathbb{R}^3)} \xrightarrow{t \rightarrow T} +\infty$$

$\dot{H}^{\frac{1}{2}}(\mathbb{R}^3)$

$L^3(\mathbb{R}^3) \subset \overset{\circ}{\cup} \subseteq \dot{B}_{\infty \infty}^1$

$$\dot{H}^{\frac{1}{2}}(\mathbb{R}^3)$$

$$\dot{H}^{\frac{d}{2}-1}(\mathbb{R}^3)$$

$d=2, 3$

$$X = L^q([0, T], \dot{H}^{\frac{d-1}{2}}(\mathbb{R}^3))$$

$$\begin{cases} \partial_t u - \Delta u = f \\ \nabla \cdot u = 0 \\ u_0 \end{cases}$$

$$\left| \begin{array}{l} \|u\|_{L^p([0, T], \dot{H}^{s+\frac{2}{p}}(\mathbb{R}^d))} \\ \leq \|u_0\|_{\dot{H}^s} + \|f\|_{L^2([0, T], \dot{H}^{\frac{d-1}{2}}(\mathbb{R}^d))} \end{array} \right.$$

$$s = \frac{d}{2} - 1$$

$$p = 4$$

$$s + \frac{2}{p} = \frac{d}{2} - 1 + \frac{2}{4} = \frac{d-1}{2}$$

$$X_T = L^q([0, T], \dot{H}^{\frac{d-1}{2}}(\mathbb{R}^d))$$

$$\partial_\nu B(u, v) - \Delta B(u, v) = Q(u, v)$$

$$B(u, v)|_{\partial} = 0$$

$$\text{We want } (u, v) \mapsto B(u, v) \quad X_T \times X_T \rightarrow X_T$$

step 1 is to show

$$Q : X_T \times X_T \rightarrow L^2([0, T], \dot{H}^{\frac{d}{2}-2}(\mathbb{R}^d))$$