

**Esercitazioni VIII lezione:  
algoritmo di Metropolis - Monte Carlo  
uso nel campionamento d'importanza**

**1. Campionamento di quantità fisiche con una distribuzione gaussiana: campionamento diretto e Metropolis**

- (a) *Campionamento diretto.* Valutare energia cinetica, potenziale, e i primi momenti  $\langle x^i \rangle$  dello stato fondamentale dell'oscillatore armonico, descritto dalla funzione d'onda  $\psi(x) = Ae^{-\beta x^2}$  (sempre con  $\sigma = 1/\sqrt{4\beta}$ ) con un campionamento diretto usando ad esempio la subroutine `gasdev` (algoritmo Box-Muller) vista precedentemente. Un esempio è il codice `diretto.f90`. Studiare l'accuratezza numerica e la convergenza delle quantità suddette in funzione del numero di punti di campionamento.
- (b) E' importante la costante A per i nostri scopi?
- (c) *Campionamento Metropolis.* Ripetere il campionamento usando il metodo di Metropolis. Un esempio è il codice `metropolis.f90`, che riprende l'esercizio (1). Si studino anche in questo caso l'accuratezza numerica e la convergenza delle quantità più rilevanti in funzione del numero di punti di campionamento.

**2. Correlations**

- (a) Calculate the autocorrelation function  $C(j) = \frac{\langle x_i x_{i+j} \rangle - \langle x_i \rangle^2}{\langle x_i^2 \rangle - \langle x_i \rangle^2}$  for a sequence of random numbers with a gaussian distribution using the Metropolis method, with different values of  $\delta/\sigma$ : 1, 5, 10, 25, 50. Comment the results.
- (b) For a fixed value of  $\sigma$  compare the autocorrelation function for two sequences of random numbers with a gaussian distribution (i) using the Metropolis method and (ii) using some ad-hoc routine, like for instance `gasdev` based on the Box-Muller algorithm. Discuss the results.

### 3. Verification of the Boltzmann distribution

We can verify directly that the Metropolis algorithm yields the Boltzmann distribution. We consider a **single classical particle** in one dimension in equilibrium with a heat bath (*canonical ensemble*). We fix therefore the temperature  $T$ , which labels a *macrostate*. The energy  $E$  can vary according to the particular *microstate* (in this particular case, it is enough to label a microstate, a part from the sign of the velocity).

- (a) Write a code (see e.g. `boltzmann_metropolis.f90`) to determine the form of the probability distribution  $P(E)$  that is generated by the Metropolis algorithm. Let for instance  $T=1$ , the initial velocity  $v_{initial}=0$ , the number of Monte Carlo steps  $nmcs=1000$ , and the maximum variation of the velocity  $dvmax=2$ . Calculate the mean energy, the mean velocity, and the probability density  $P(E)$ .
- (b) Consider  $\ln P(E)$  as a function of  $E$ . Can you recognize the expected behavior? (see slides for the analytic derivation of  $P(E)$ ) You should recognise that the asymptotic behavior is a straight line whose slope is  $-1/T$ .
- (c) How many  $nmcs$  do you need to have a reasonable estimate of the mean energy and mean velocity?
- (d) Verify that your results do not depend from the initial conditions by changing  $v_{initial}$ . What does it change? What does it change by changing instead  $dvmax$ ?
- (e) Modify the program to simulate an ideal gas of **N particles** in one dimension. [Hint: modify the subroutine *Metropolis* inserting a loop over the particles] Consider for instance  $N=20$ ,  $T=100$ ,  $nmcs=200$ . Assume all particles to have the same initial velocity  $v_{initial}=10$ . Determine the value of  $dvmax$  so that the acceptance ratio is about 50%? What are the mean energy  $\langle E \rangle$  (i.e., total energy of the system  $\langle E_{tot} \rangle$  divided by the number of particles) and the mean velocity? [the symbol  $\langle \rangle$  indicates temporal(statistical) averages]
- (f) Calculate  $P(E)$  ( $E$  now indicates the mean energy per particle), make a plot and describe its behaviour. Is it similar to the case  $N=1$ ? Comment on that.
- (g) Calculate the total energy  $E_{tot}$  for  $T=10, 20, 30, 90, 100$ , and  $110$ , and estimate the heat capacity as the numerical derivative of the energy with respect to the temperature,  $C = \partial \langle E_{tot} \rangle / \partial T$ . [C is the heat capacity, i.e. referred to the whole system; you may consider, alternatively, the specific heat, referred to a single particle...]
- (h) Calculate the mean square energy fluctuation  $\langle \Delta E_{tot}^2 \rangle = \langle E_{tot}^2 \rangle - \langle E_{tot} \rangle^2$  for  $T=10$  and  $T=40$ . Compare the magnitude of the ratio  $C = \langle \Delta E_{tot}^2 \rangle / T^2$  numerically estimated from the mean square energy fluctuation with that obtained in (f).

#### 4. MC simulation of a simple N-particles model

Consider an ideal gas of  $N$  non interacting, distinguishable particles, **confined** in a box (fixed  $\mathbf{V}$ ) and **isolated** (fixed  $\mathbf{E}$ ), divided into left/right with the possibility for one particle at a time to pass through the separation wall, with equal probability from the left to the right or viceversa.

A **macrostate** is specified for instance by the number of particles on the left side, say  $n$ , that can correspond to different **microstates** depending on the list of the specific particles there. A Monte Carlo approach consists in generating a certain number of movements, randomly, and consider them as representative of all the possible movements. The program `box.f90` is a possible implementation of the algorithm describing the time evolution of the system in terms of macrostates, i.e. –given an initial number of particles on the left,  $n$ – the approach to equilibrium and which is the equilibrium macrostate.

- (a) Choose  $N=4, 10, 20, 40, 80$ , and  $n=N$  initially. Make a plot of  $n$  (or, better, of  $n/N$ ) with respect to time. What is the equilibration time  $\tau_{eq}$  (=how many MC steps)?
- (b) Modify the program so that at each time step  $t$  it calculates the number of particles  $\langle n(t) \rangle$  averaged over different runs (e.g. 5 runs). Make a plot to compare  $n(t)$  over the individual runs and averaged  $\langle n(t) \rangle$ .
- (c) (*Optional; do it at home!*) Compare the numerical value of  $\langle n(t) \rangle$  with the exact analytic results for a simple case, for instance  $N=4$ .
- (d) (*Optional*) Consider only one run. Modify the program to calculate numerically the probability  $P_n$  of having *at equilibrium* a macrostate with  $n$  particles on the left, by simply counting the number of occurring microstates that correspond to the macrostate  $n$  and dividing for the total number of microstates generated in the time evolution. Plot the histogram  $P_n$  for  $N=20, 40, 80$  and a “sufficiently” long run. Comment.
- (e) Modify the program to measure the statistical fluctuations at the equilibrium, by calculating the variance  $\sigma^2 = \langle n^2 \rangle - \langle n \rangle^2$ , where the average is done over a time interval *after* reaching the equilibrium.
- (f) Determine  $\langle n \rangle$  and  $\sigma / \langle n \rangle$  at equilibrium for  $N=20, 40, 80$ . Which is the dependence of these quantities on  $N$ ?



```

        call random_number(rnd)          !
        if (p > rnd) then                !
            x = xp                       !
        !cccccccccccccccccccccccccccccc
            acc=acc+1.0_dp
        endif
    enddo

write(unit=*,fmt=*)"acceptance ratio = ",acc/n
write(unit=*,fmt=*)"Risultati (simulazione vs.risultato esatto):"
write(unit=*,fmt=format1)"etot = ",etot/n,1.0_dp/(8.0_dp*sigma**2)&
    +0.5_dp*sigma**2
write(unit=*,fmt=format1)"ekin = ",ekin/n,1.0_dp/(8.0_dp*sigma**2)
write(unit=*,fmt=format1)"epot = ",epot/n,0.5_dp*sigma**2
write(unit=*,fmt=format1)"<x> = ",x1/n,0.0_dp
write(unit=*,fmt=format1)"<x^2>= ",x2/n,sigma**2
write(unit=*,fmt=format1)"<x^3>= ",x3/n,0.0_dp
write(unit=*,fmt=format1)"<x^4>= ",x4/n,3.0_dp*sigma**4

end program metropolis

```



```

program diretto
  use gaussian
  implicit none
  integer, parameter :: dp=selected_real_kind(13)
  integer :: i,n
  integer, dimension(1) :: seed
  real :: rnd
  real(kind=dp):: sigma,beta,etot,ekin,epot
  real(kind=dp):: x,x1,x2,x3,x4
  character(len=13), save :: format1 = "(a7,2x,2f9.5)"

  x1 = 0.0_dp
  x2 = 0.0_dp
  x3 = 0.0_dp
  x4 = 0.0_dp
  ekin = 0.0_dp
  epot = 0.0_dp
  print*, "seed, n, beta ="
  read*, seed(1),n,beta
  call random_seed(put=seed)
  sigma=1.0_dp/sqrt(4.0_dp*beta)
  do i=1,n
    !cccccccccccccccccccccccccccccccc
    call gasdev(rnd) !
    x=rnd*sigma ! campionamento diretto
    !cccccccccccccccccccccccccccccccc!
    ekin = ekin - 0.5_dp * ((2*beta*x)**2 - 2*beta)
    epot = epot + 0.5_dp * x**2
    etot = ekin + epot
    x1 = x1 + x
    x2 = x2 + x**2
    x3 = x3 + x**3
    x4 = x4 + x**4
  end do

  write(unit=*,fmt=*)"Risultati (simulazione verso risultato esatto):"
  write(unit=*,fmt=format1)"etot = ",etot/n,1.0_dp/(8.0_dp*sigma**2)&
    +0.5_dp*sigma**2
  write(unit=*,fmt=format1)"ekin = ",ekin/n,1.0_dp/(8.0_dp*sigma**2)
  write(unit=*,fmt=format1)"epot = ",epot/n,0.5_dp*sigma**2
  write(unit=*,fmt=format1)"<x> = ",x1/n,0.0_dp
  write(unit=*,fmt=format1)"<x^2>= ",x2/n,sigma**2
  write(unit=*,fmt=format1)"<x^3>= ",x3/n,0.0_dp
  write(unit=*,fmt=format1)"<x^4>= ",x4/n,3.0_dp*sigma**4

end program diretto

```





```

print *,"# deltaE  :",del_E
print *,"# nbin   :",nbin
open(unit=9,file="boltzmann.dat",status="replace",action="write")
write(unit=9,fmt=*)"# T      :",T
write(unit=9,fmt=*)"# <E0>  :",E
write(unit=9,fmt=*)"# <v0>  :",vel
write(unit=9,fmt=*)"# dvmax  :",dvmax
write(unit=9,fmt=*)"# nMCsteps:",nmcs
write(unit=9,fmt=*)"# deltaE  :",del_E
write(unit=9,fmt=*)"# nbin   :",nbin
allocate (P(0:nbin))
ecum = 0.0
e2cum = 0.0
vcum = 0.0
P = 0.0
accept= 0.0
end subroutine initial

subroutine Metropolis()
real :: dv,vtrial,de,rnd
call random_number(rnd)
dv = (2*rnd - 1) * dvmax           ! trial variation for v
vtrial = vel + dv                 ! trial velocity v
de = 0.5 * (vtrial*vtrial - vel*vel) ! corresponding variation of E
call random_number(rnd)
if (de >= 0.0) then
    if ( exp(-beta*de) < rnd ) return ! trial step not accepted
end if
vel = vtrial
accept = accept + 1
E = E + de
end subroutine Metropolis

subroutine data(vcum,ecum,e2cum)
real, intent(inout) :: vcum,ecum,e2cum
Ecum = Ecum + E
E2cum = E2cum + E*E
vcum = vcum + vel
call probability()
end subroutine data

subroutine probability()
integer :: ibin
ibin = int(E/del_E)
if ( ibin <= nbin ) P(ibin) = P(ibin) + 1
end subroutine probability

```

```

subroutine averages(nequil,vcum,Ecum,E2cum)
  integer, intent(in) :: nequil
  real, intent(in) :: vcum,Ecum,E2cum
  real :: znorm, Eave, E2ave, vave, sigma2
  integer :: ibin
  znorm = 1.0/nmcs
  accept = accept / (nmcs+nequil) ! acceptance ratio
  Eave = Ecum * znorm ! average energy
  E2ave = E2cum * znorm !
  vave = vcum * znorm ! average velocity
  sigma2 = E2ave - Eave*Eave
  print *, "# <E2>num.:", E2ave
  print *, "# <E> num.:", Eave
  print *, "# <E> th. :", T/2
  print *, "# <v> :", vave
  print *, "# accept. :", accept
  print *, "# sigma :", sqrt(sigma2)
  write(unit=9,fmt=*) "# <E2>num:", E2ave
  write(unit=9,fmt=*) "# <E> num.:", Eave
  write(unit=9,fmt=*) "# <E> th. :", T/2
  write(unit=9,fmt=*) "# <v> :", vave
  write(unit=9,fmt=*) "# accept. :", accept
  write(unit=9,fmt=*) "# sigmaE :", sqrt(sigma2)
  write(unit=9,fmt=*) "# ibin*del_E, P(E)"
  do ibin = 0, nbin
    write(unit=9,fmt=*) ibin*del_E, P(ibin) * znorm
  end do
  close(unit=9)
end subroutine averages
end module common

program Boltzmann
  use common
  real :: vcum, ecum, e2cum
  integer :: imcs, nequil
  ! parameters and variable initialization
  call initial(nequil,vcum,ecum,e2cum)
  do imcs = 1, nmcs + nequil
    call Metropolis()
    ! data accumulation after each Metropolis step
    if ( imcs > nequil ) call data(vcum,ecum,e2cum)
  end do
  call averages(nequil,vcum,Ecum,E2cum)
  deallocate(P)
end program Boltzmann

```



```
! elsewhere from the right to the left  
call initial()  
call move()  
end program box
```