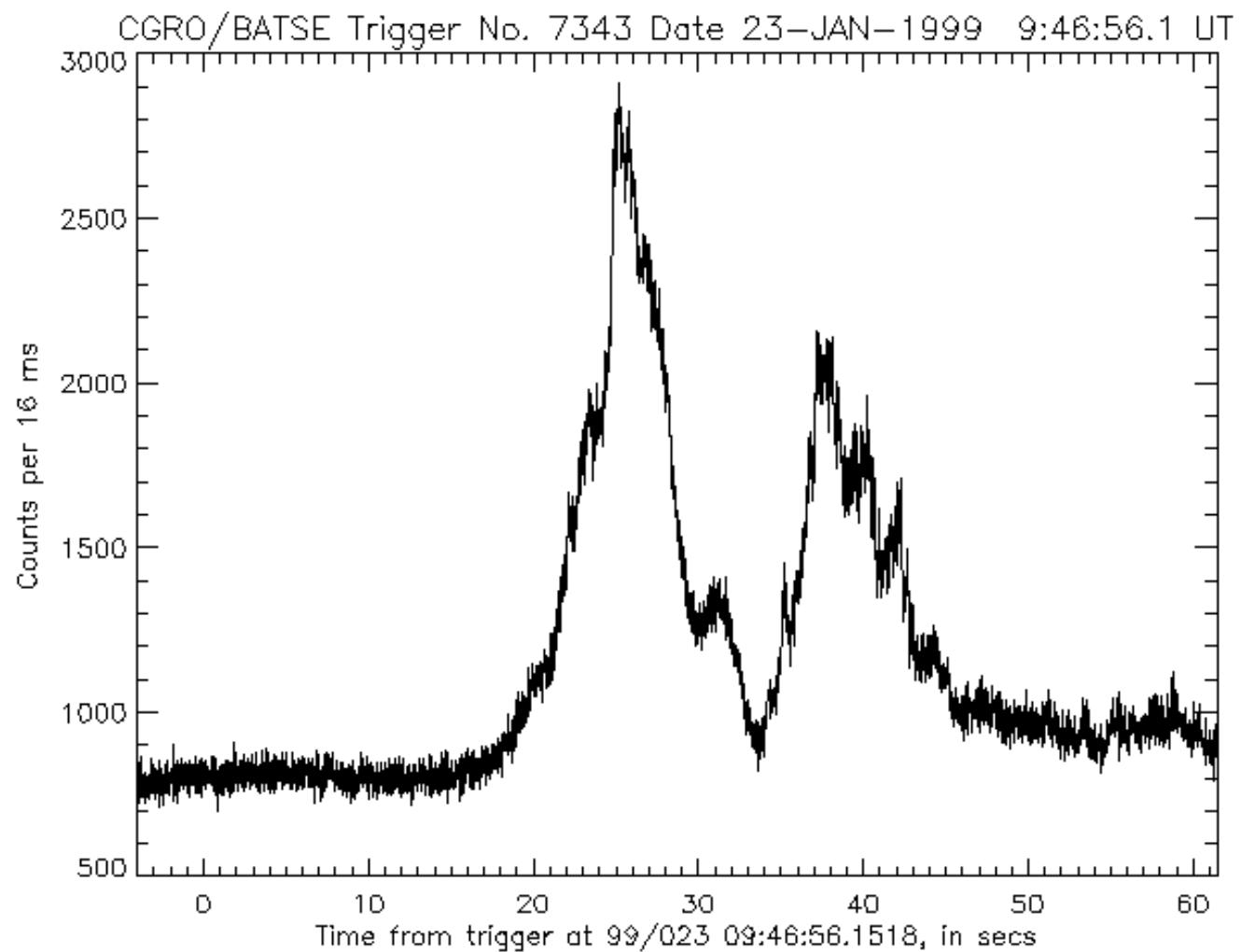


Astrofisica Nucleare e Subnucleare

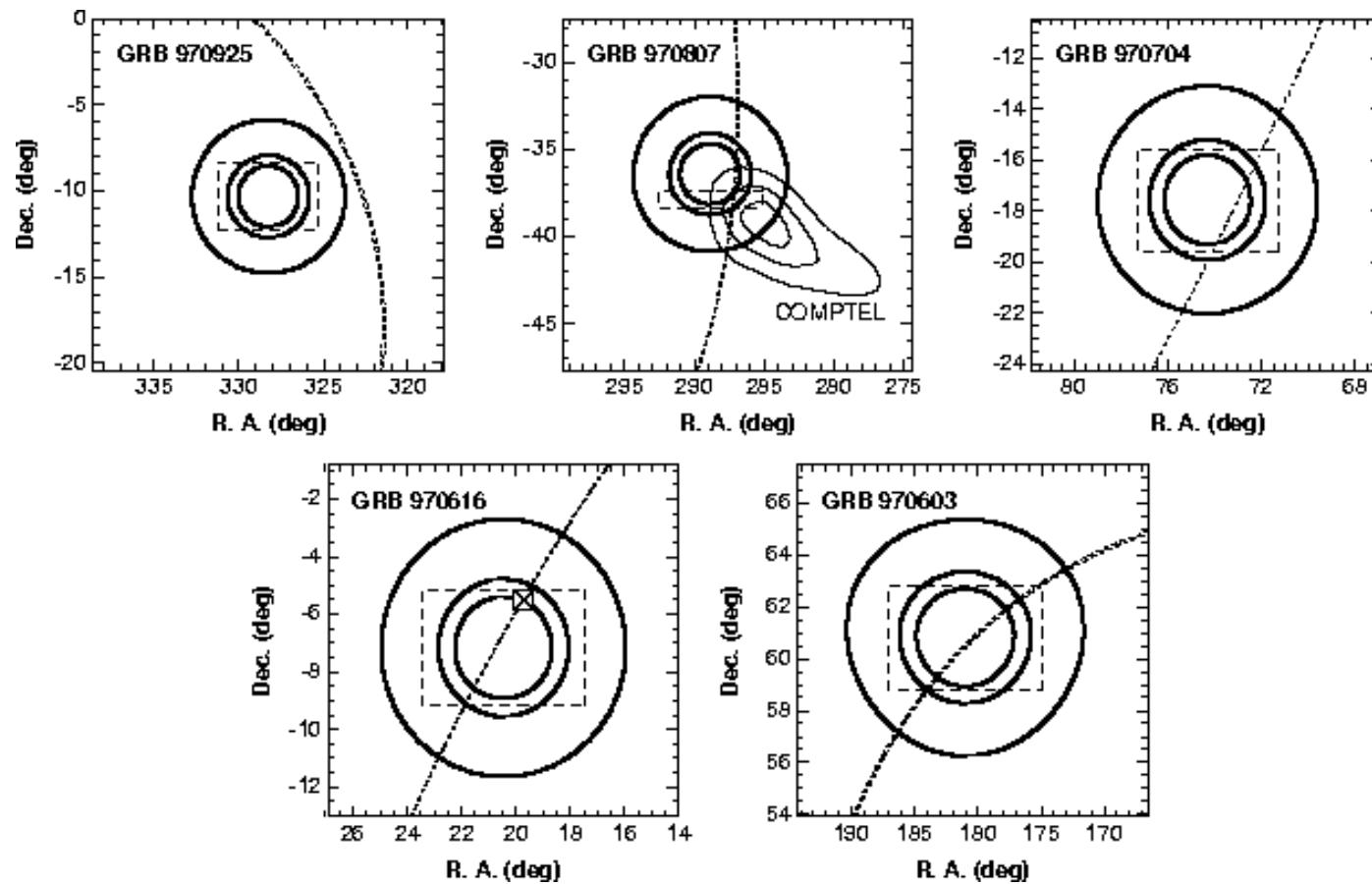
keV-MeV Astrophysics

GRB detectors I

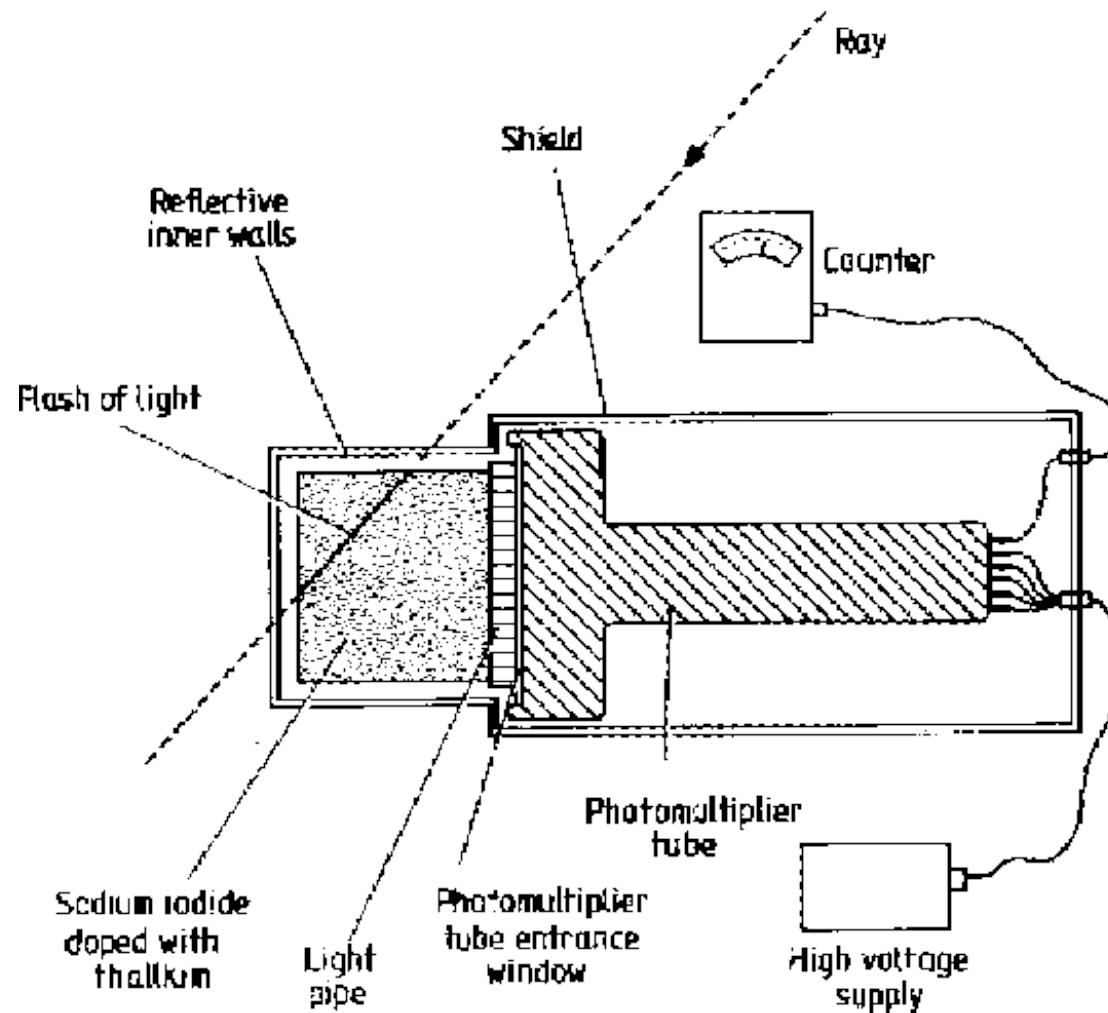
A GRB ...



some GRB locations ...

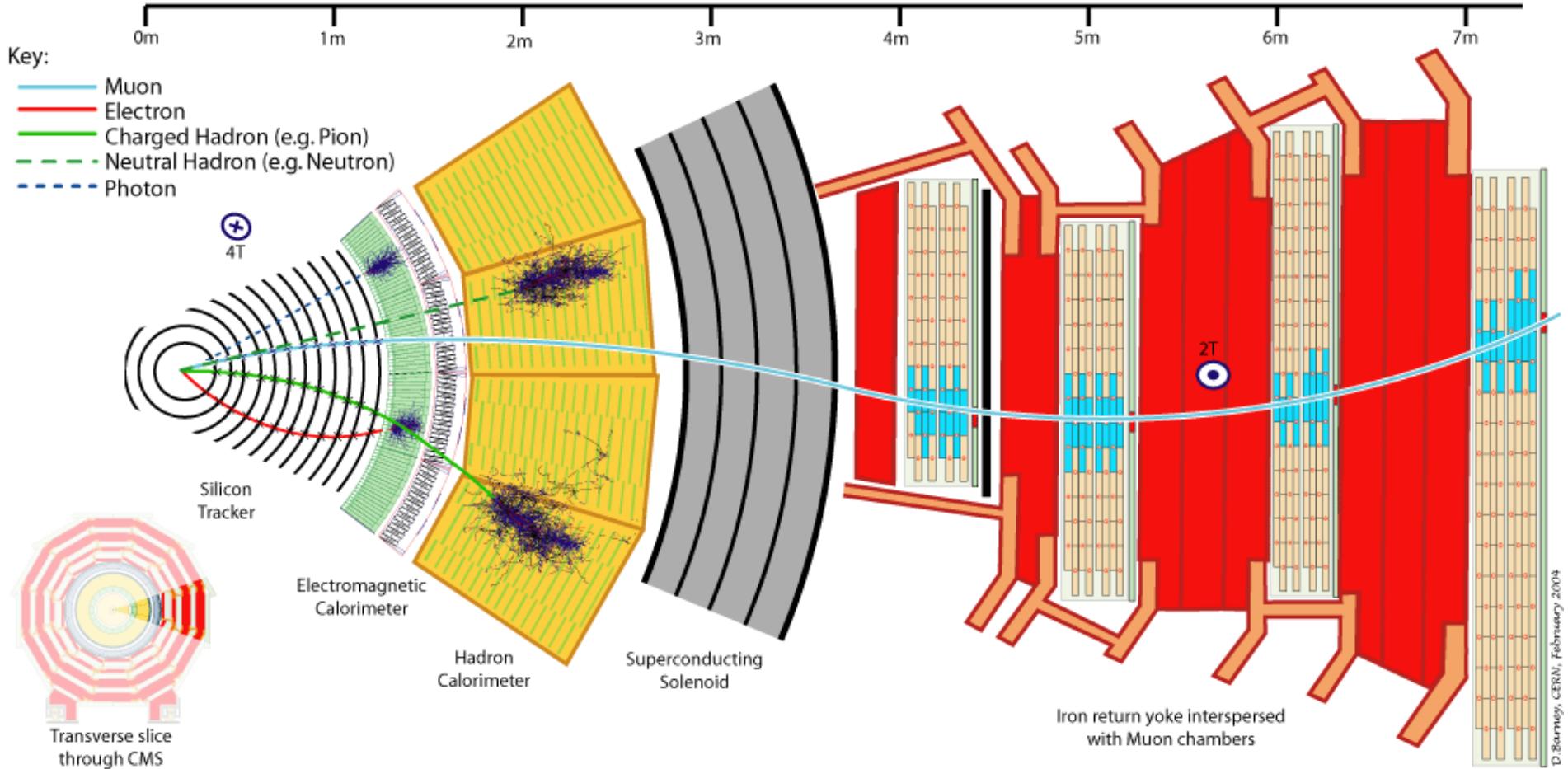


GRB Detectors



Astrofisica Nucleare e Subnucleare

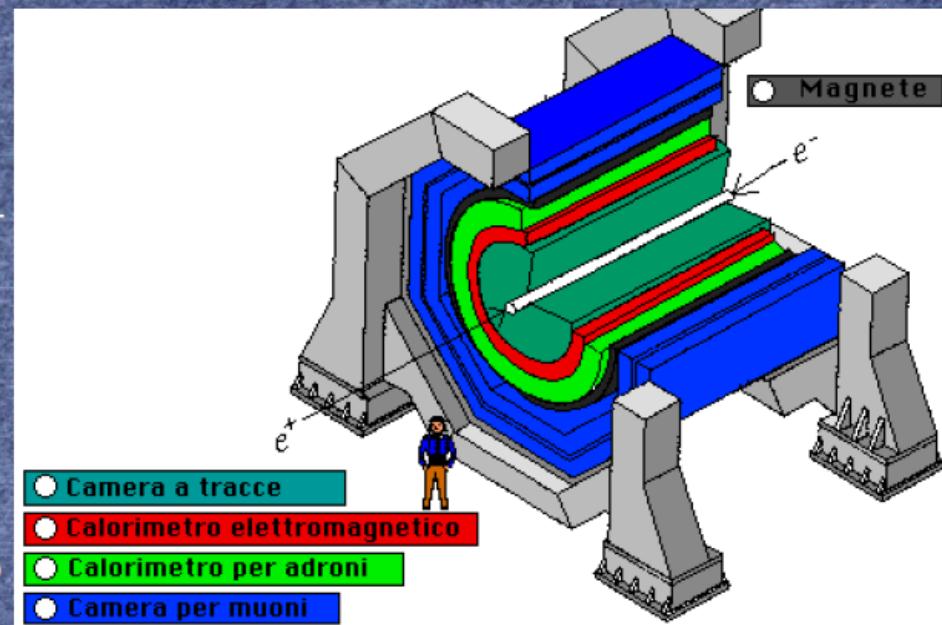
Interazione Radiazione Materia



Rivelazione delle particelle elementari

Struttura tipica dei rivelatori

- ◆ I rivelatori di particelle sono costituiti da un insieme di sotto-rivelatori diversi.
- ◆ Le tecniche di rivelazione descritte precedentemente dettano una struttura a strati cilindrici concentrici dei rivelatori.
- ◆ Procedendo dall'asse del cilindro verso l'esterno tipicamente abbiamo:
 - ◆ un tracciatore immerso in un campo magnetico uniforme;
 - ◆ un "calorimetro elettromagnetico" per misurare l'energia degli elettroni e dei fotoni;
 - ◆ un "calorimetro adronico" per misurare l'energia degli adroni;
 - ◆ "rivelatori di posizione" per rivelare il passaggio dei muoni.

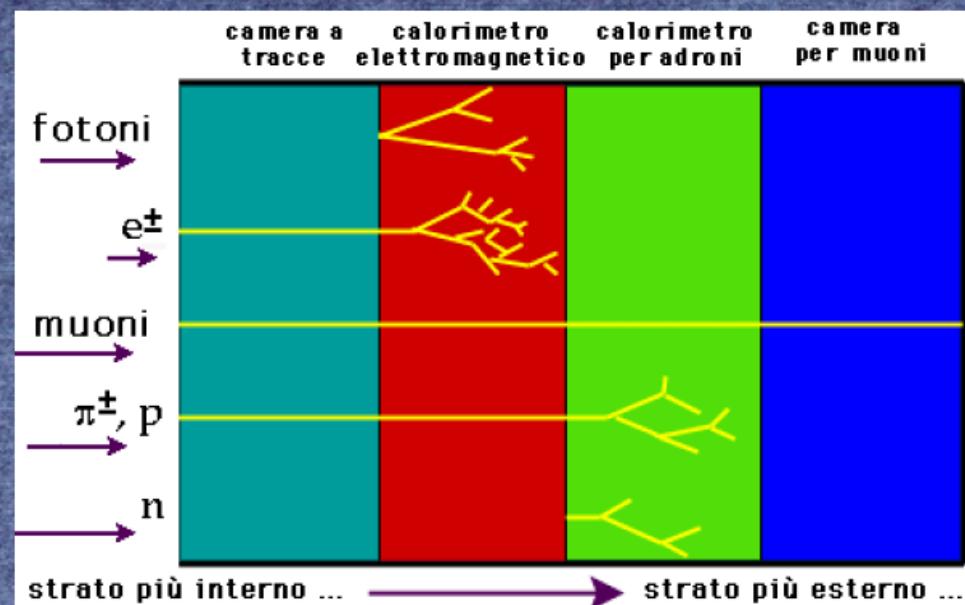


Come?

Identificazione delle particelle

◆ Quindi, disponendo i rivelatori in un certo ordine e combinandone le informazioni, è possibile identificare diversi tipi di particelle:

- ◆ fotone: energia nel calorimetro elettromagnetico;
- ◆ elettrone: traccia + energia nel calorimetro elettromagnetico;
- ◆ muone: traccia + segnale nei rivelatori di muoni;
- ◆ adrone carico: traccia + energia nel calorimetro adronico;
- ◆ adrone neutro: energia nel calorimetro adronico.



Particle Detectors

Summer Student Lectures 2010

Werner Riegler, CERN, werner.riegler@cern.ch

- ◆ **History of Instrumentation ↔ History of Particle Physics**
- ◆ **The ‘Real’ World of Particles**
- ◆ **Interaction of Particles with Matter**
- ◆ **Tracking Detectors, Calorimeters, Particle Identification**
- ◆ **Detector Systems**

Interaction of Particles with Matter

Any device that is to detect a particle must interact with it in some way → almost ...

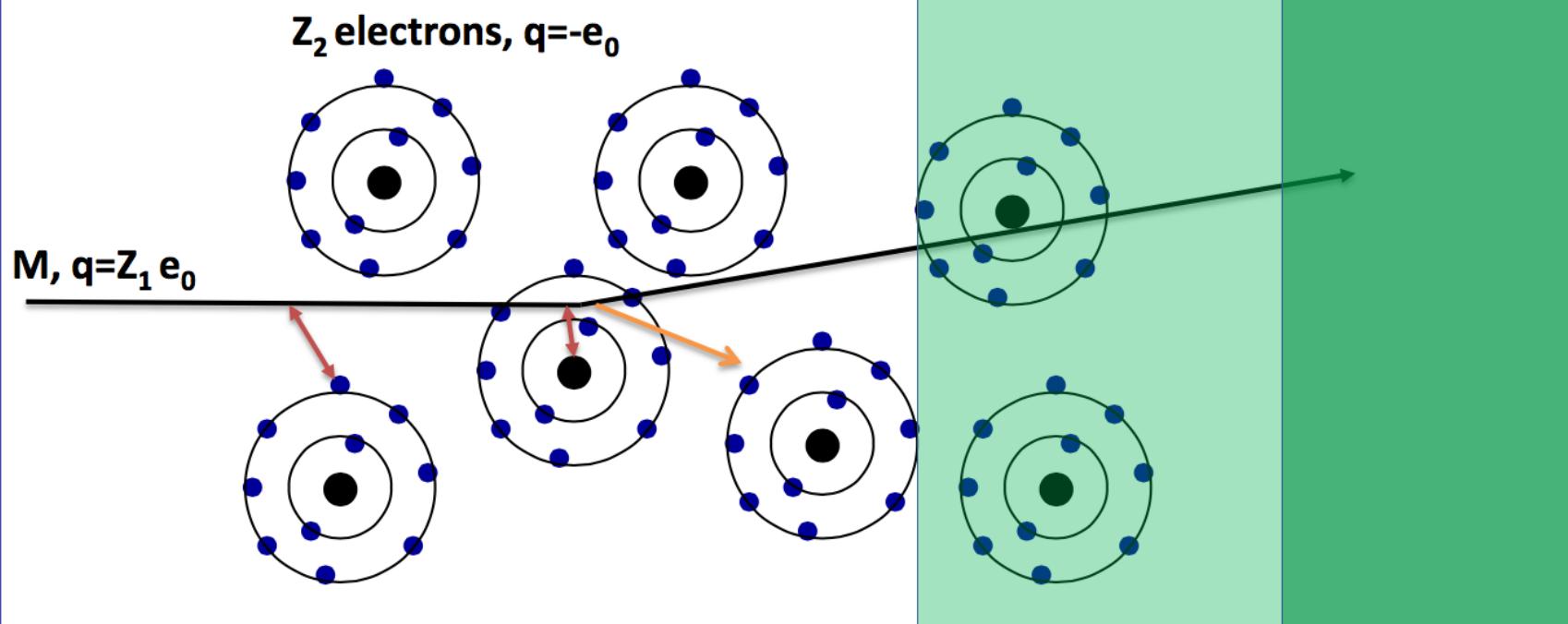
In many experiments neutrinos are measured by missing transverse momentum.

**E.g. e^+e^- collider. $P_{tot}=0$,
If the Σp_i of all collision products is $\neq 0 \rightarrow$ neutrino escaped.**



“Did you see it?”
“No nothing.”
“Then it was a neutrino!”

Electromagnetic Interaction of Particles with Matter



Interaction with the atomic electrons. The incoming particle loses energy and the atoms are excited or ionized.

Interaction with the atomic nucleus. The particle is deflected (scattered) causing multiple scattering of the particle in the material. During this scattering a Bremsstrahlung photon can be emitted.

In case the particle's velocity is larger than the velocity of light in the medium, the resulting EM shockwave manifests itself as Cherenkov Radiation. When the particle crosses the boundary between two media, there is a probability of the order of 1% to produce an X ray photon, called Transition radiation.

Astrofisica Nucleare e Subnucleare

Ionizzazione

Bethe Bloch Formula

$$\frac{1}{\rho} \frac{dE}{dx} = -4\pi r_e^2 m_e c^2 \frac{Z_1^2}{\beta^2} N_A \frac{Z}{A} \left[\ln \frac{2m_e c^2 \beta^2 \gamma^2 F}{I} - \beta^2 - \frac{\delta(\beta\gamma)}{2} \right]$$

Für $Z > 1$, $I \approx 16Z^{0.9}$ eV

For Large $\beta\gamma$ the medium is being polarized by the strong transverse fields, which reduces the rise of the energy loss \rightarrow density effect

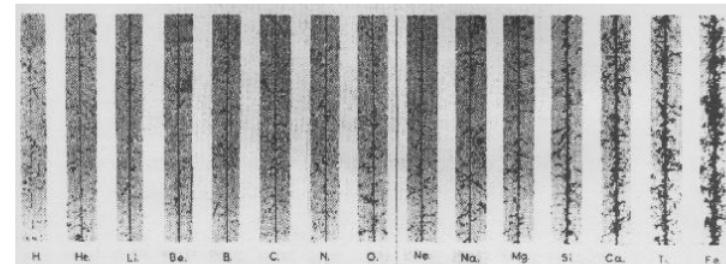
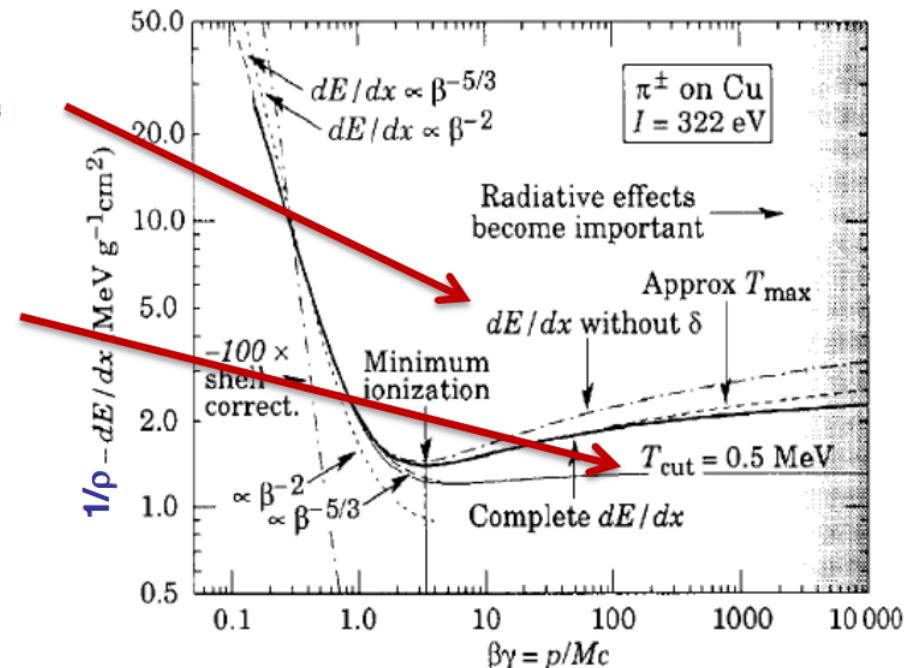
At large Energy Transfers (delta electrons) the liberated electrons can leave the material.

In reality, E_{max} must be replaced by E_{cut} and the energy loss reaches a plateau (Fermi plateau).

Characteristics of the energy loss as a function of the particle velocity ($\beta\gamma$)

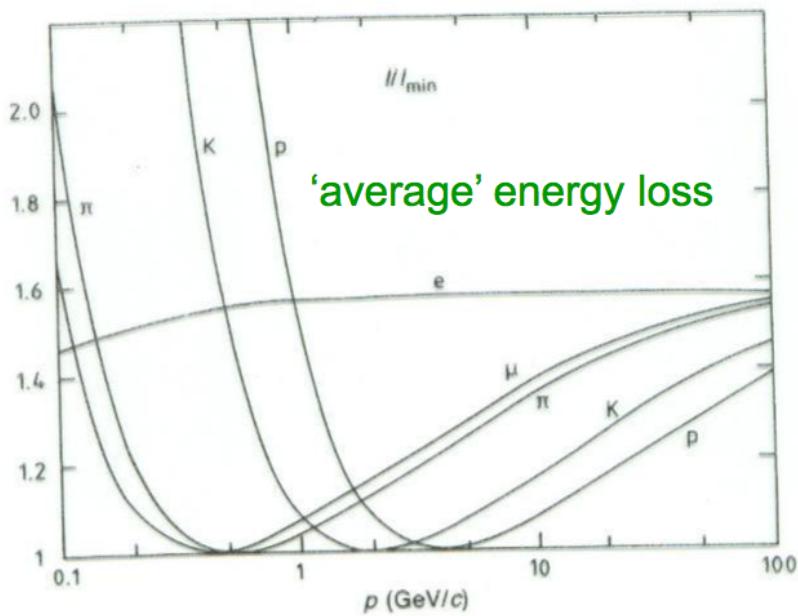
The specific Energy Loss $1/\rho dE/dx$

- first decreases as $1/\beta^2$
- increases with $\ln \gamma$ for $\beta=1$
- is \approx independent of M ($M \gg m_e$)
- is proportional to Z_1^2 of the incoming particle.
- is \approx independent of the material ($Z/A \approx \text{const}$)
- shows a plateau at large $\beta\gamma$ (>100)
- $dE/dx \approx 1-2 \times \rho [\text{g/cm}^3] \text{ MeV/cm}$

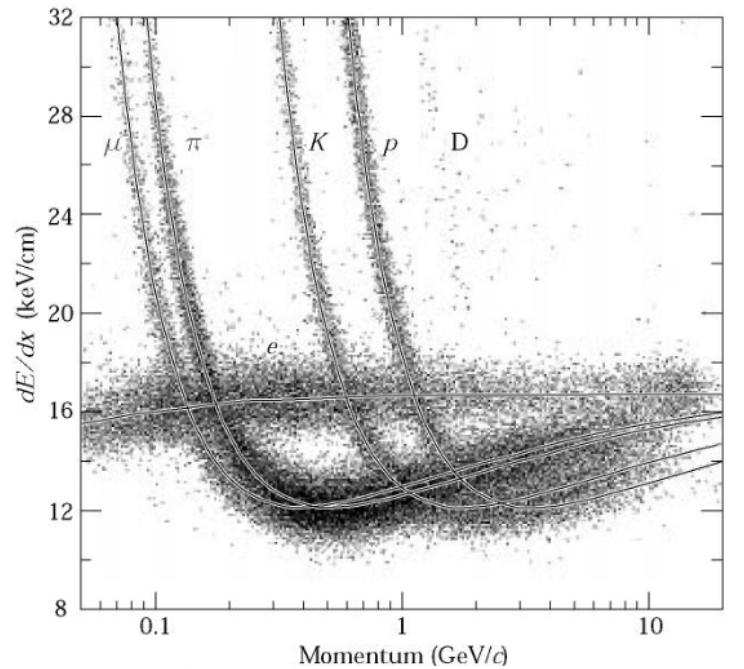


Particle Identification

Measured energy loss

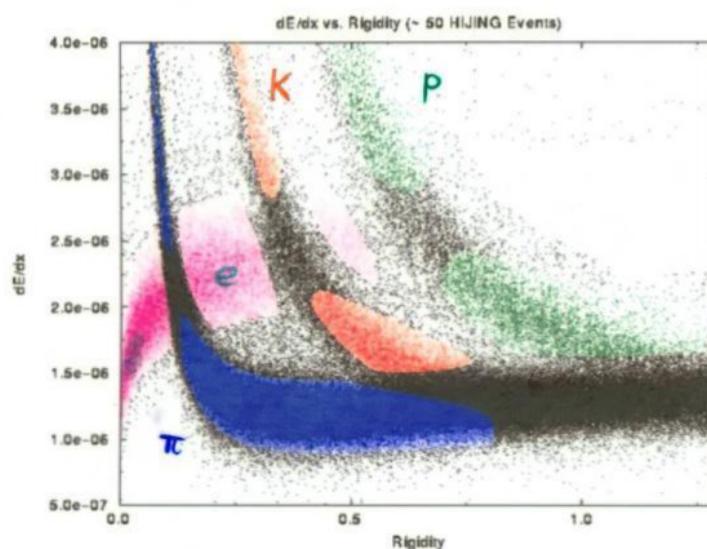


'average' energy loss

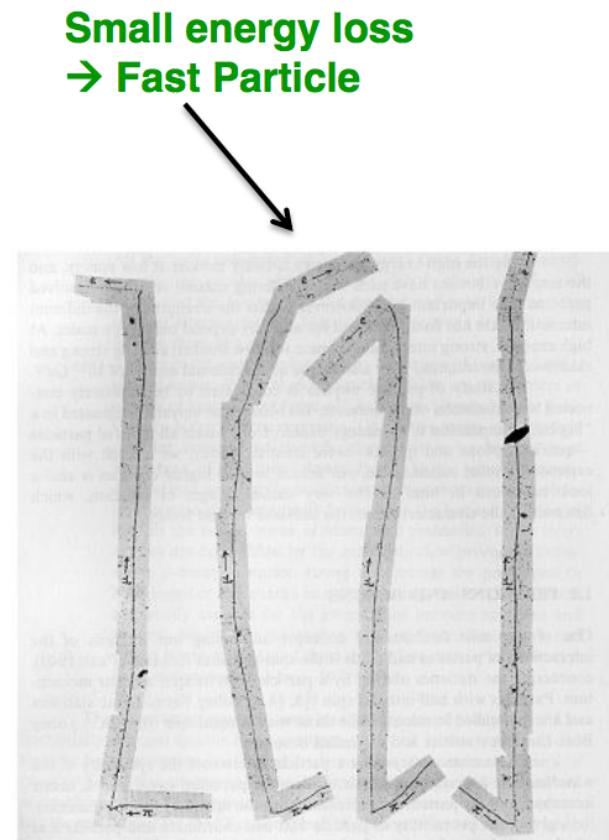


BLUE => PIONS RED => KAONS GREEN => PROTONS MAGENTA => ELECTRONS BLACK => NO ID POSSIBLE

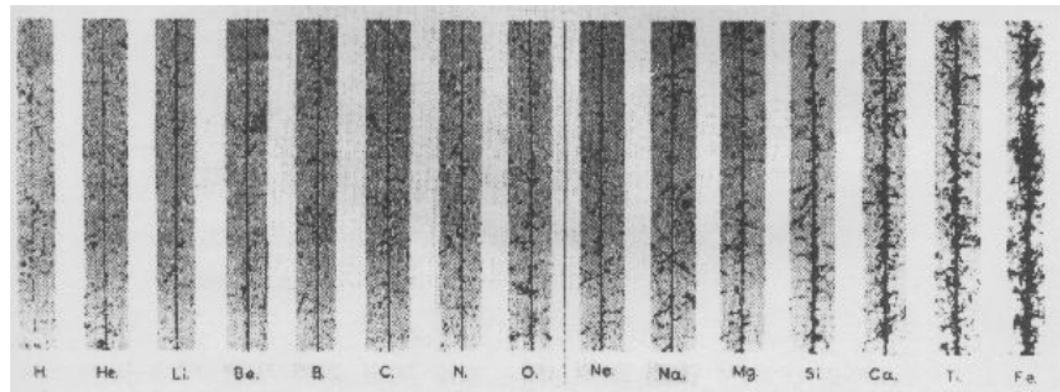
STAR
TPC



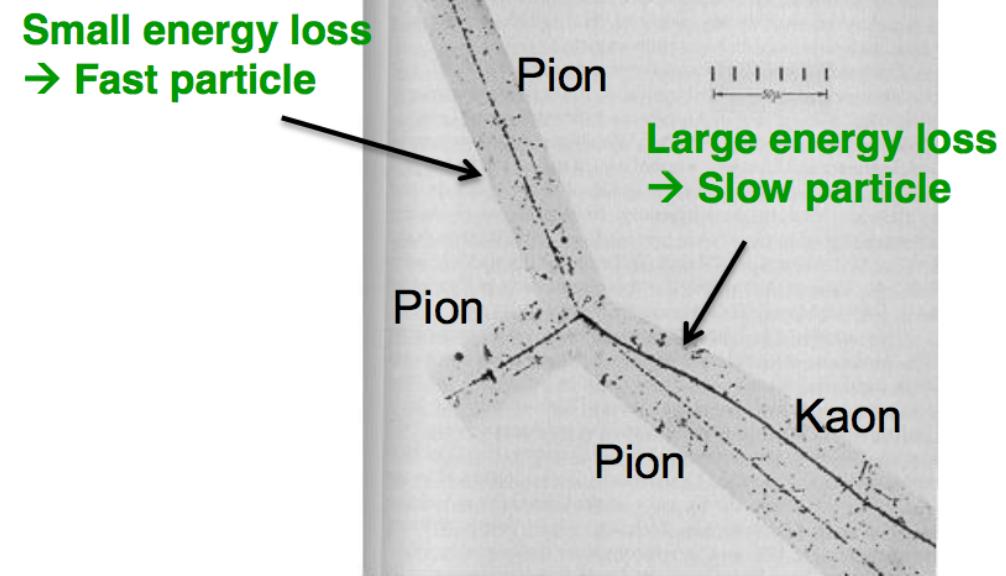
In certain momentum ranges,
particles can be identified by
measuring the energy loss.



Discovery of muon and pion



Cosmics rays: $dE/dx \propto Z^2$

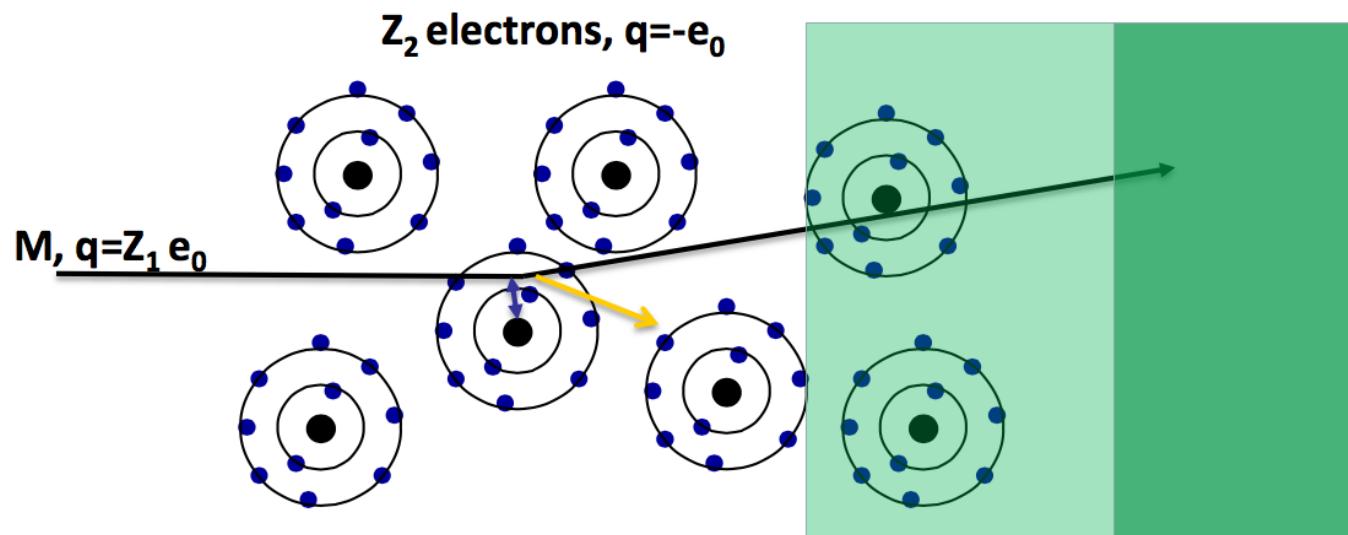


Astrofisica Nucleare e Subnucleare

Bremsstrahlung

Bremsstrahlung

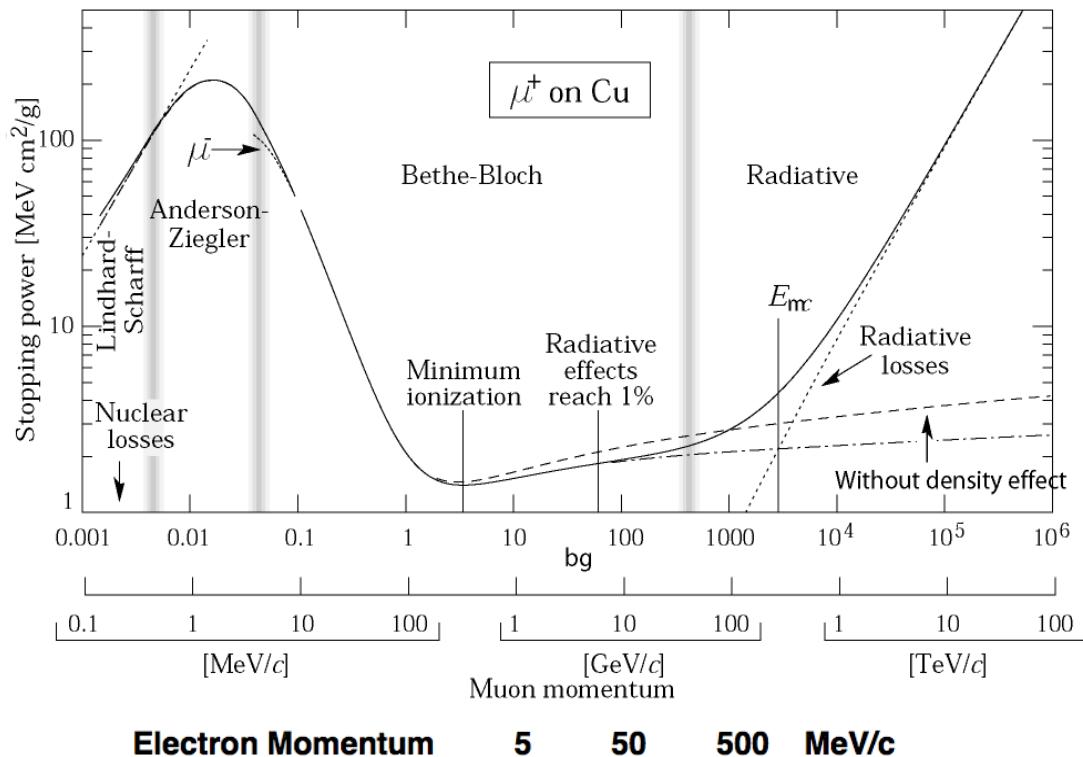
A charged particle of mass M and charge $q=Z_1 e_0$ is deflected by a nucleus of charge Ze which is partially 'shielded' by the electrons. During this deflection the charge is 'accelerated' and it therefore radiated \rightarrow Bremsstrahlung.



7/15/2010

Critical Energy

such as copper to about 1% accuracy for energies between about 6 MeV and 6 GeV



For the muon, the second lightest particle after the electron, the critical energy is at 400GeV.

The EM Bremsstrahlung is therefore only relevant for electrons at energies of past and present detectors.

Critical Energy: If dE/dx (Ionization) = dE/dx (Bremsstrahlung)

Myon in Copper: $p \approx 400\text{GeV}$
Electron in Copper: $p \approx 20\text{MeV}$

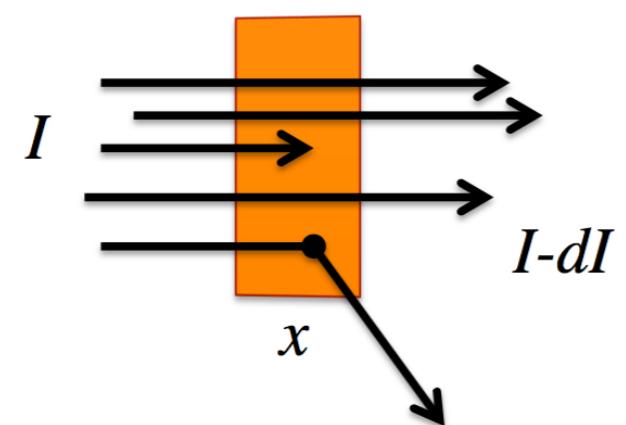
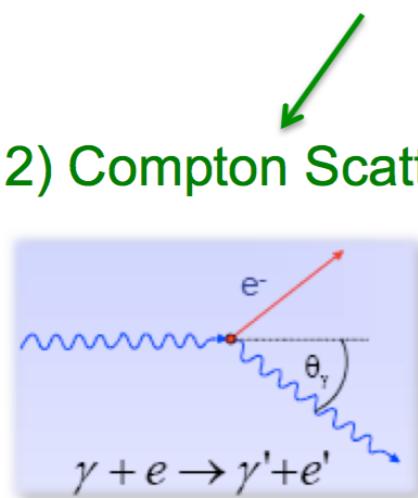
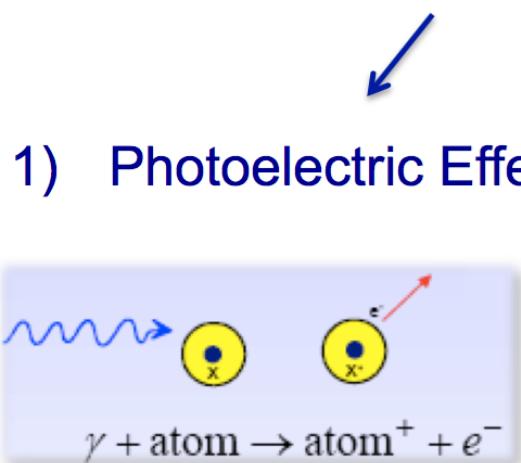
Astrofisica Nucleare e Subnucleare

Interazione di Foton

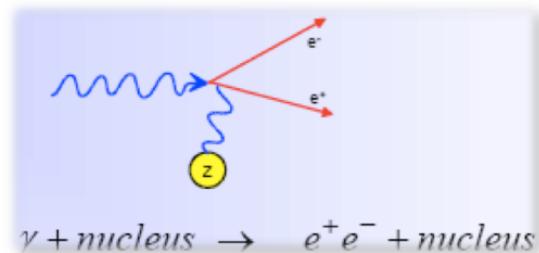
Interactions of photons with matter

Characteristic for interactions of photons with matter:

A photon is removed from the beam after one single interaction either because of **total absorption** or **scattering**



3) Pair Production

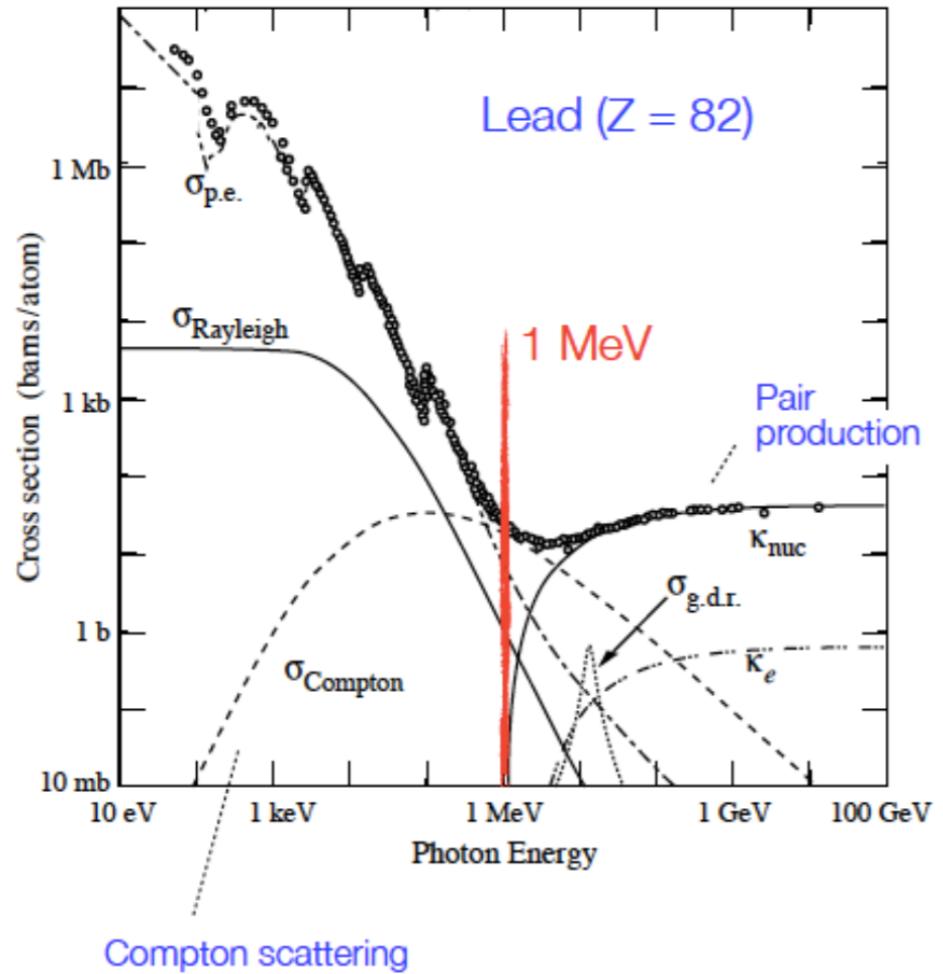
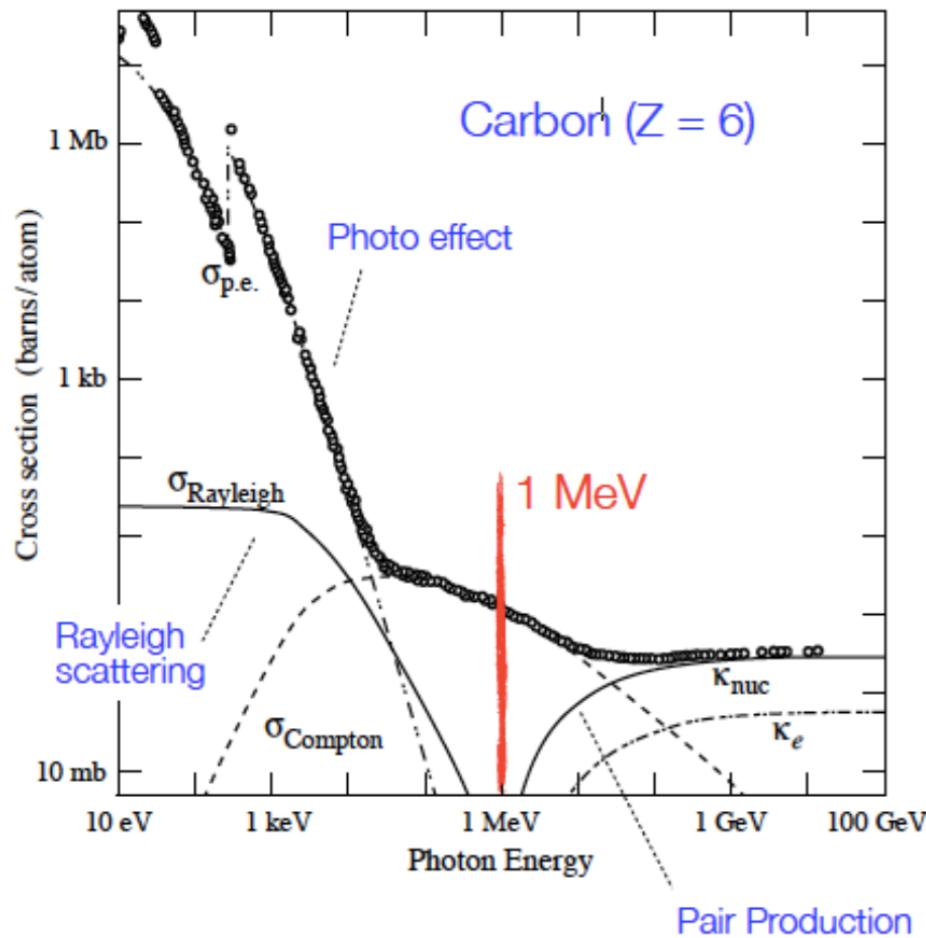


$$I(x) = I_0$$

$$\lambda = 1 / \mu \quad \text{Mean free path}$$

Interactions of photons with matter

Photon Total Cross Sections



Photoelectric effect

From energy conservation:

$$E_e = E_\gamma - E_N = h\nu - I_b$$

I_b = Nucleus binding energy
introduces strong Z dependence

Cross-section largest for $E_\gamma \approx$ K-shell energy

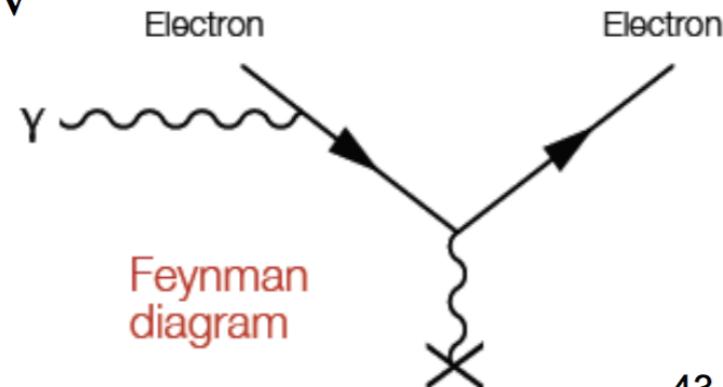
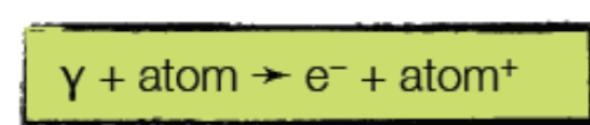
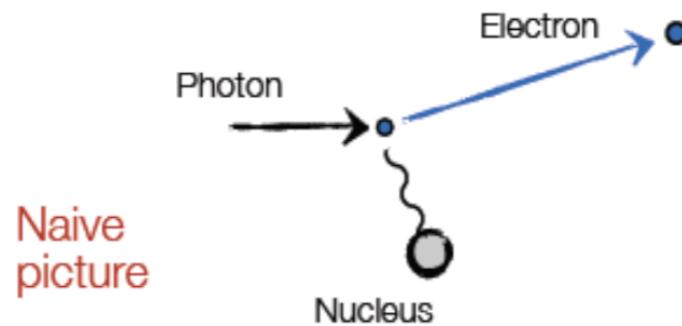
Strongest E dependence for $I_0 < E_\gamma < m_e c^2$

$$\sigma_{ph} = \alpha \pi a_B Z^5 (I_0 / E_\gamma)^{7/2}$$

$a_B = 0.53 \text{ \AA}$
 $I_0 = 13.6 \text{ eV}$

E-dependence softer for $E_\gamma > m_e c^2$

$$\sigma_{ph} = 2\pi r_e^2 \alpha^4 Z^5 (mc)^2 / E_\gamma$$



Compton scattering

Best known electromagnetic process
(Klein–Nishina formula)

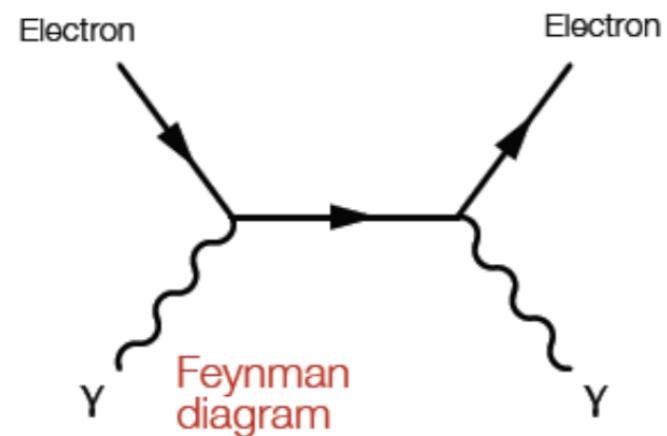
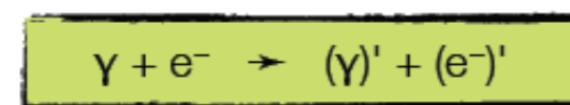
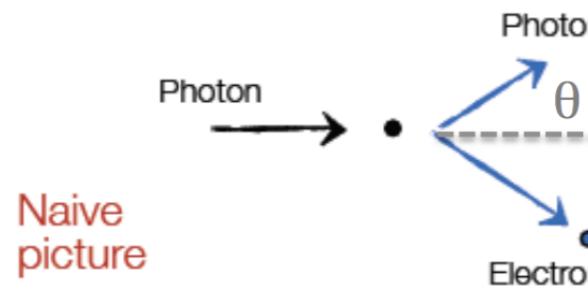
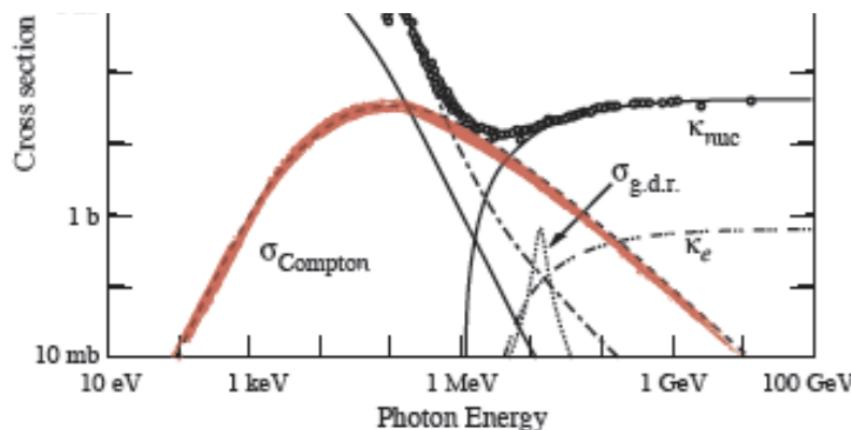
for $E_\lambda \ll m_e c^2$ $\sigma_c \propto \sigma_{Th}(1 - 2\varepsilon)$

Thompson cross-section:

$$\sigma_{Th} = 8\pi/3 r_e^2 = 0.66 \text{ barn}$$

$$\varepsilon = \frac{E_\lambda}{m_e c^2}$$

for $E_\lambda \gg m_e c^2$ $\sigma_c \propto \frac{\ln \varepsilon}{\varepsilon} Z$



Compton scattering

From E and p conservation get the energy of the scattered photon

$$E_\gamma' = \frac{E_\gamma}{1 + \varepsilon(1 - \cos\theta)} \quad \varepsilon = \frac{E_\lambda}{m_e c^2}$$

Kinetic energy of the outgoing electron:

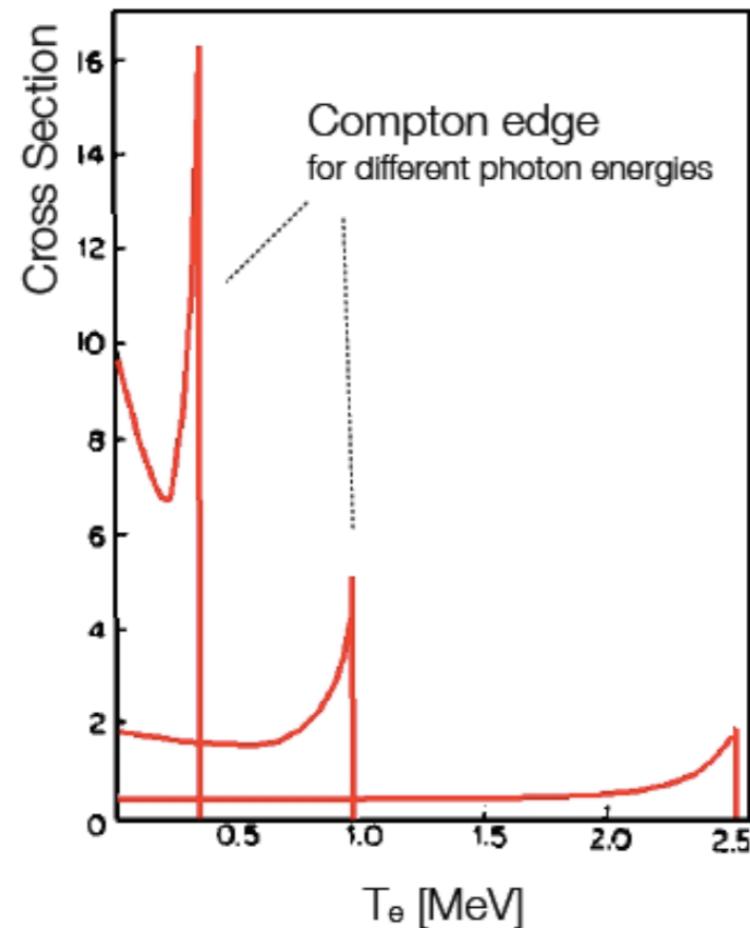
$$T_e = E_\gamma - E_\gamma' = E_\gamma \frac{\varepsilon(1 - \cos\theta)}{1 + \varepsilon(1 - \cos\theta)}$$

Max. electron recoil energy for $\theta = \pi$:

$$T_{\max} = E_\gamma \frac{2\varepsilon}{1 + 2\varepsilon}$$

Transfer of complete γ -energy via Compton scattering not possible:

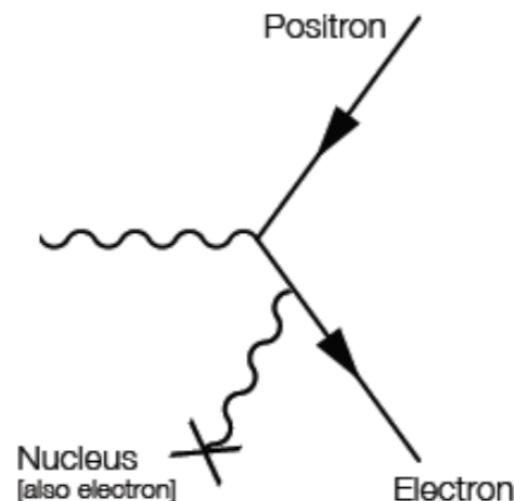
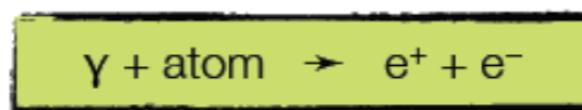
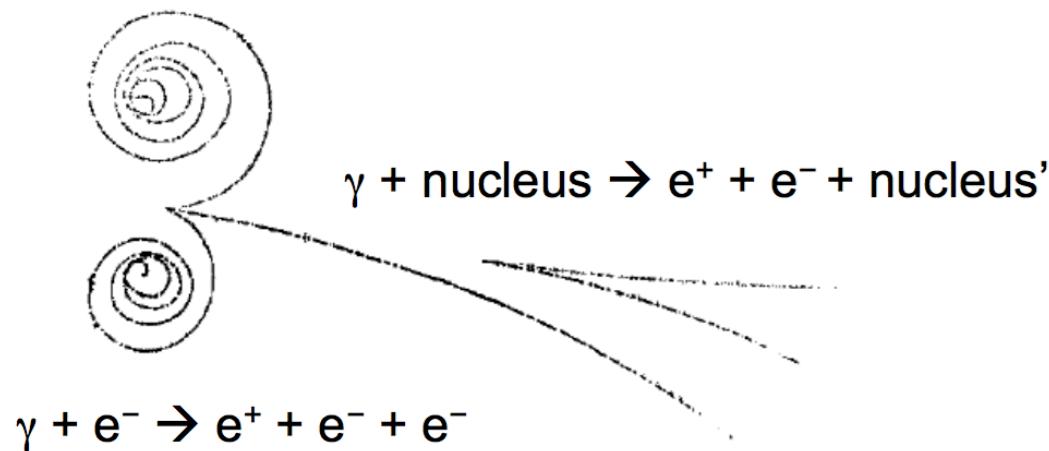
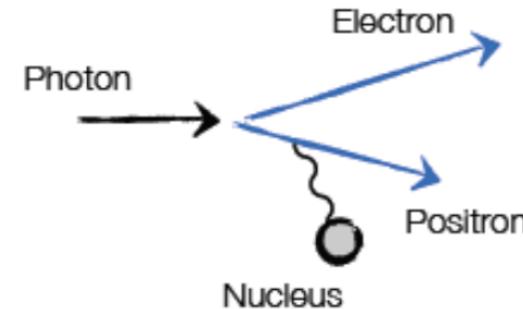
$$\Delta E = E_\gamma - T_{\max} = E_\gamma \frac{1}{1 + 2\varepsilon}$$



Pair production

Minimum energy required for this process
 $2 m_e + \text{Energy transferred to the nucleus}$

$$E_\gamma \geq 2m_e c^2 + \frac{2m_e c^2}{m_{\text{Nucleus}}}$$

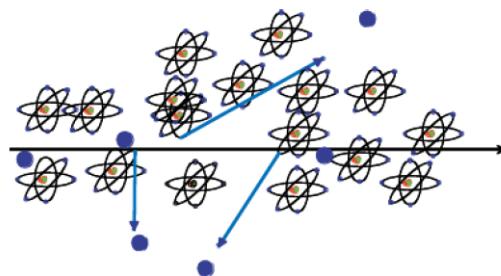


Astrofisica Nucleare e Subnucleare

Scintillation Detectors

Creation of the Signal

Charged particles traversing matter leave excited atoms, electron-ion pairs (gases) or electrons-hole pairs (solids) behind.

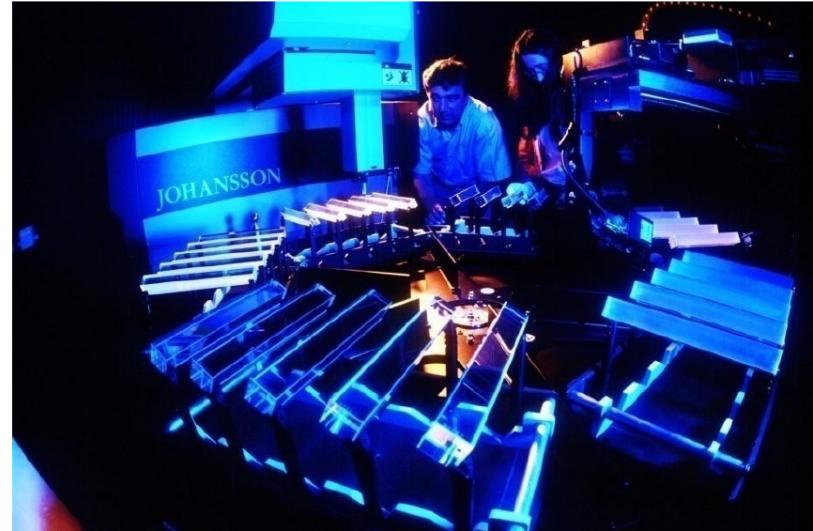


Excitation:

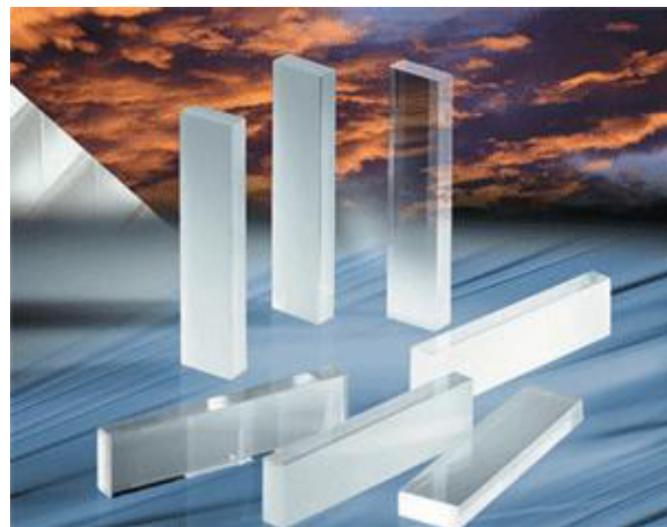
The photons emitted by the excited atoms in transparent materials can be detected with photon detectors like photomultipliers or semiconductor photon detectors.

Ionization:

By applying an electric field in the detector volume, the ionization electrons and ions are moving, which induces signals on metal electrodes. These signals are then read out by appropriate readout electronics.



**Detectors based on registration of
excited Atoms → Scintillators**



Detectors based on Registration of excited Atoms → Scintillators

Emission of photons of by excited Atoms, typically UV to visible light.



a) Observed in Noble Gases (even liquid !)

b) Inorganic Crystals

→ Substances with largest light yield. Used for precision measurement of energetic Photons. Used in Nuclear Medicine.

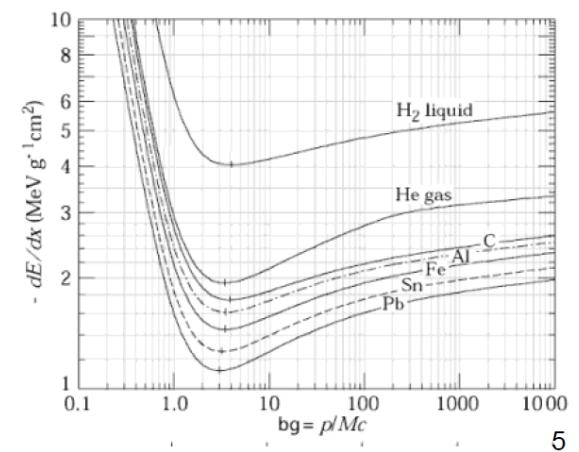
c) Polycyclic Hydrocarbons (Naphtalen, Anthrazen, organic Scintillators)

→ Most important category. Large scale industrial production, mechanically and chemically quite robust. Characteristic are one or two decay times of the light emission.

Typical light yield of scintillators:

Energy (visible photons) ≈ few % of the total energy Loss.

z.B. 1cm plastic scintillator, $\rho \approx 1$, $dE/dx = 1.5$ MeV, ~15 keV in photons; i.e. ~ 15 000 photons produced.



Detectors based on Registration of excited Atoms → Scintillators

Organic ('Plastic') Scintillators

Low Light Yield

Type	Light ^a output	$\lambda_{\text{max}}^{\text{b}}$ (nm)	Attenuation ^c length (cm)	Risetime (ns)	Decay ^d time (ns)	Pulse FWHM (ns)
NE 102A	58–70	423	250	0.9	2.2–2.5	2.7–3.2
NE 104	68	406	120	0.6–0.7	1.7–2.0	2.2–2.5
NE 104B	59	406	120	1	3.0	3
NE 110	60	434	400	1.0	2.9–3.3	4.2
NE 111	40–55	375	8	0.13–0.4	1.3–1.7	1.2–1.6
NE 114	42–50	434	350–400	~1.0	4.0	5.3
Pilot B	60–68	408	125	0.7	1.6–1.9	2.4–2.7
Pilot F	64	423	300	0.9	2.1	3.0–3.3
Pilot U	58–67	391	100–140	0.5	1.4–1.5	1.2–1.9
BC 404	68	408	—	0.7	1.8	2.2
BC 408	64	423	—	0.9	2.1	~2.5
BC 420	64	391	—	0.5	1.5	1.3
ND 100	60	434	400	—	3.3	3.3
ND 120	65	423	250	—	2.4	2.7
ND 160	68	408	125	—	1.8	2.7

Fast: 1-3ns

Inorganic (Crystal) Scintillators

Large Light Yield Slow: few 100ns

	Relative light output	λ_{max} emission (nm)	Decay time (ns)	Density (g/cm ³)
<i>Inorganic crystals</i>				
Nal(Tl)	230	415	230	3.67
CsI(Tl)	250	560	900	4.51
Bi ₄ Ge ₃ O ₁₂ (BGO)	23–86	480	300	7.13
<i>Organic crystals</i>				
Anthracene	100	448	22	1.25
Trans-stilbene	75	384	4.5	1.16
Naphthalene	32	330–348	76–96	1.03
p,p'-Quaterphenyl	94	437	7.5	1.20
<i>Primary activators</i>				
2,5-Diphenyl-oxazole (PPO)	75	360–416	5*	
2-Phenyl-5-(4-biphenylyl)-1,3,4-oxadiazole (PBD)	96	360–5		
4,4"-Bis(2-butyloctyloxy)-p-quaterphenyl (BIBUQ)	60	365,393	1.30 ^b	

LHC bunchcrossing 25ns

LEP bunchcrossing 25μs

Scintillators

Inorganic Scintillators – Properties

Scintillator material	Density [g/cm ³]	Refractive Index	Wavelength [nm] for max. emission	Decay time constant [μs]	Photons/MeV
Nal	3.7	1.78	303	0.06	$8 \cdot 10^4$
Nal(Tl)	3.7	1.85	410	0.25	$4 \cdot 10^4$
CsI(Tl)	4.5	1.80	565	1.0	$1.1 \cdot 10^4$
Bi ₄ Ge ₃ O ₁₂	7.1	2.15	480	0.30	$2.8 \cdot 10^3$
CsF	4.1	1.48	390	0.003	$2 \cdot 10^3$
LSO	7.4	1.82	420	0.04	$1.4 \cdot 10^4$
PbWO ₄	8.3	1.82	420	0.006	$2 \cdot 10^2$
LHe	0.1	1.02	390	0.01/1.6	$2 \cdot 10^2$
LAr	1.4	1.29*	150	0.005/0.86	$4 \cdot 10^4$
LXe	3.1	1.60*	150	0.003/0.02	$4 \cdot 10^4$

* at 170 nm

Scintillators

Organic Scintillators – Properties

Scintillator material	Density [g/cm ³]	Refractive Index	Wavelength [nm] for max. emission	Decay time constant [ns]	Photons/MeV
Naphtalene	1.15	1.58	348	11	4·10 ³
Antracene	1.25	1.59	448	30	4·10 ⁴
p-Terphenyl	1.23	1.65	391	6-12	1.2·10 ⁴
NE102*	1.03	1.58	425	2.5	2.5·10 ⁴
NE104*	1.03	1.58	405	1.8	2.4·10 ⁴
NE110*	1.03	1.58	437	3.3	2.4·10 ⁴
NE111*	1.03	1.58	370	1.7	2.3·10 ⁴
BC400**	1.03	1.58	423	2.4	2.5·10 ²
BC428**	1.03	1.58	480	12.5	2.2·10 ⁴
BC443**	1.05	1.58	425	2.2	2.4·10 ⁴

* Nuclear Enterprises, U.K.

** Bicron Corporation, USA

Scintillators

Scintillators – Basic Counter Setup

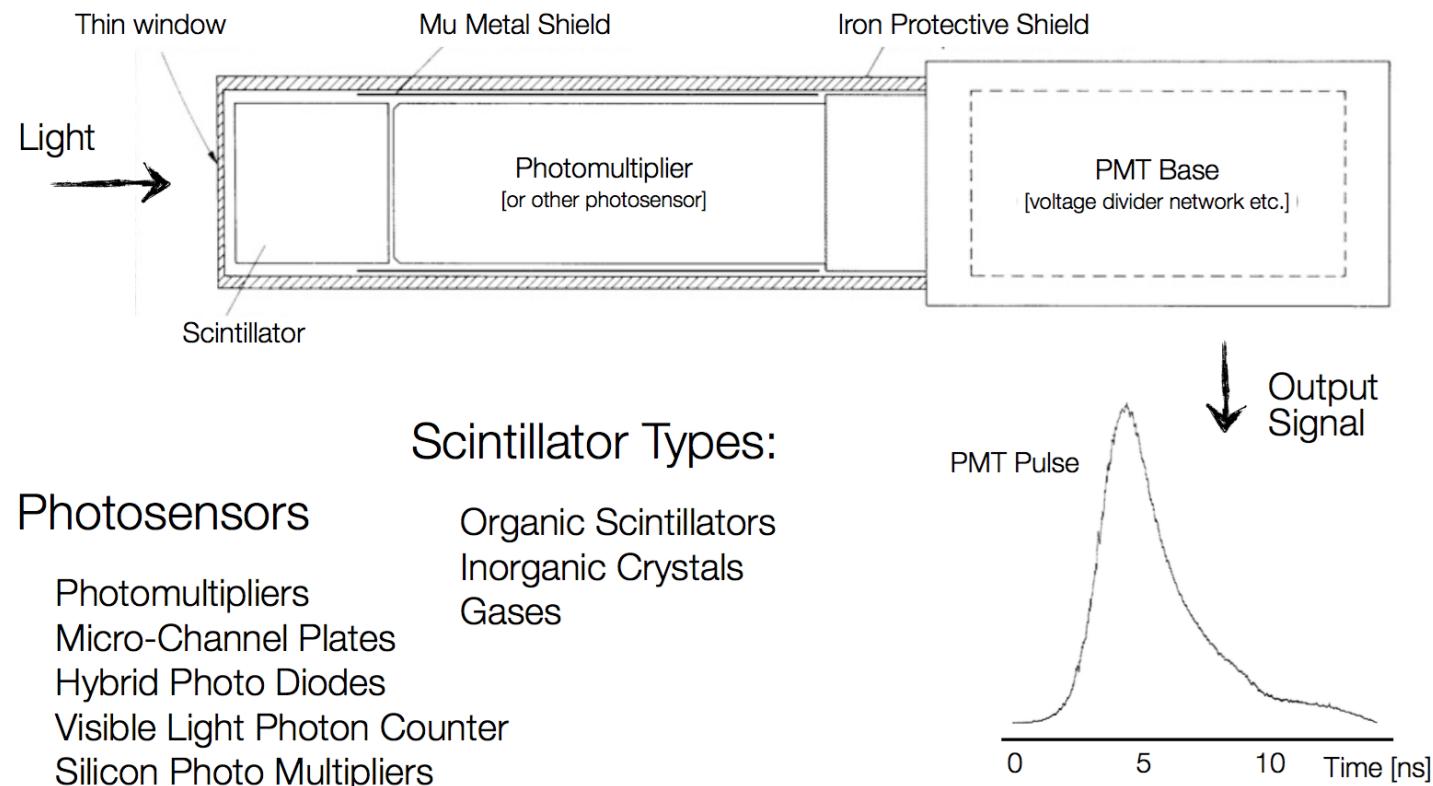
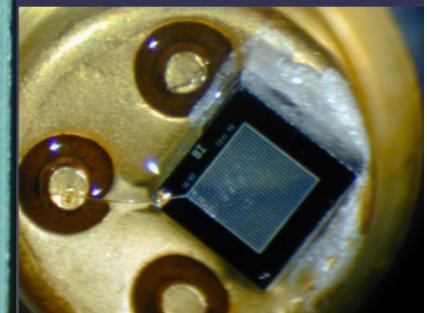
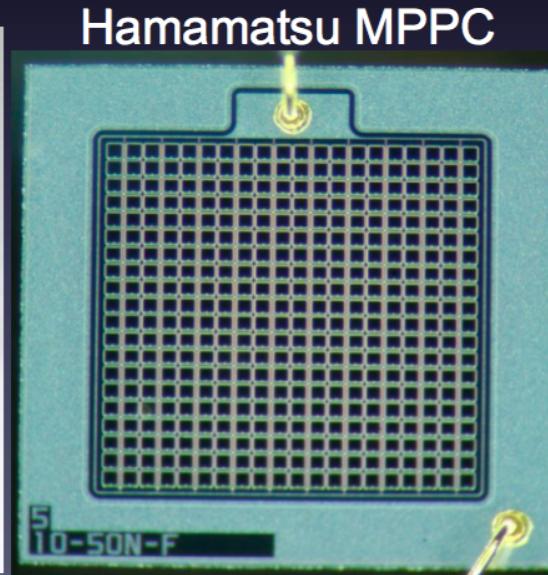


Photo-detectors

- Convert light into an electronic signal by using the photo-electric effect to convert photons into photo-electrons (p.e.)
- Requirement :
 - High Photon Detection Efficiency (PDE) or
 - Quantum Efficiency; Q.E. = $N_{\text{p.e.}}/N_{\text{photons}}$
- Photomultipliers
- SiPM



One of the first
SiPM
Pulsar, Moscow

PMTs

Photomultipliers

Principle:

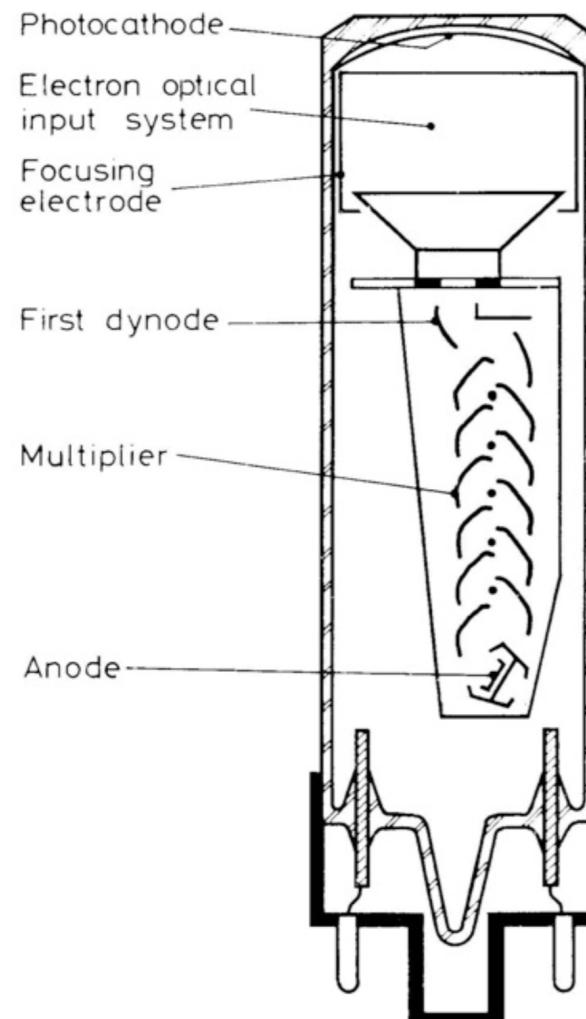
Electron emission
from photo cathode

Secondary emission
from dynodes; dynode gain: 3-50 [$f(E)$]

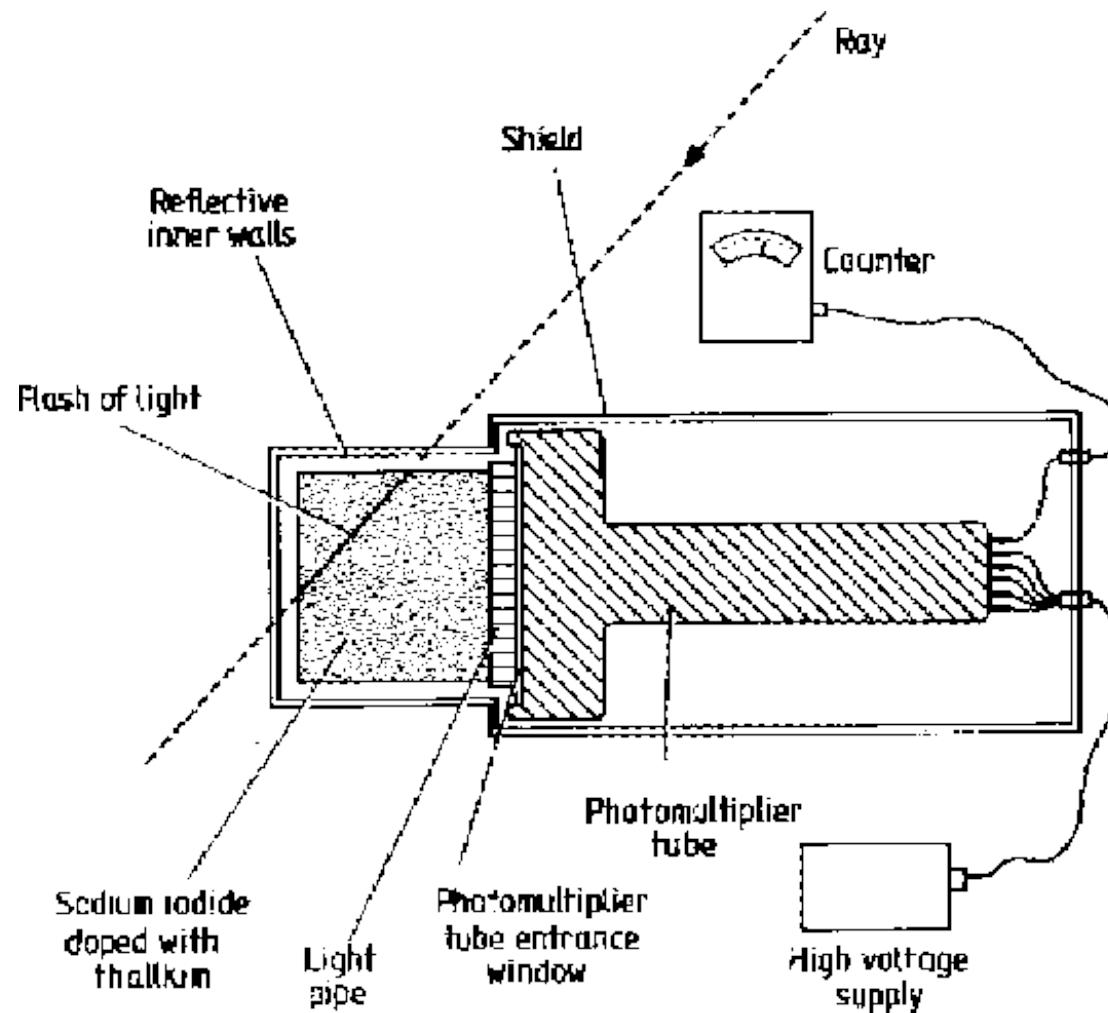
Typical PMT Gain: $> 10^6$
[PMT can see single photons ...]



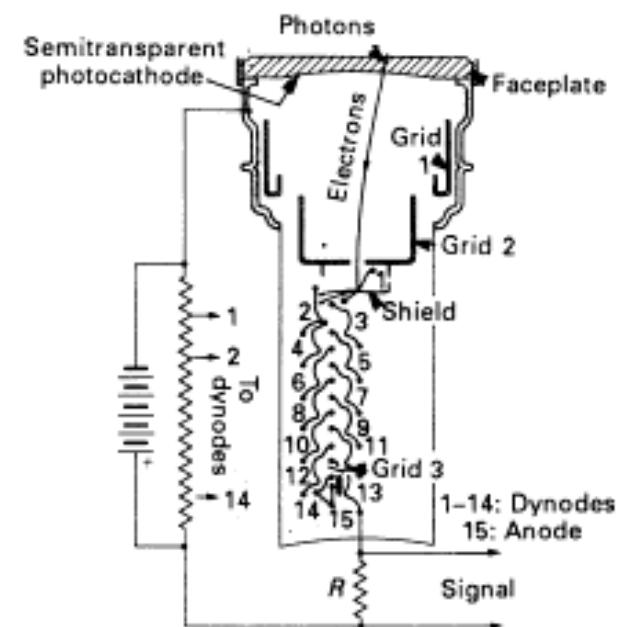
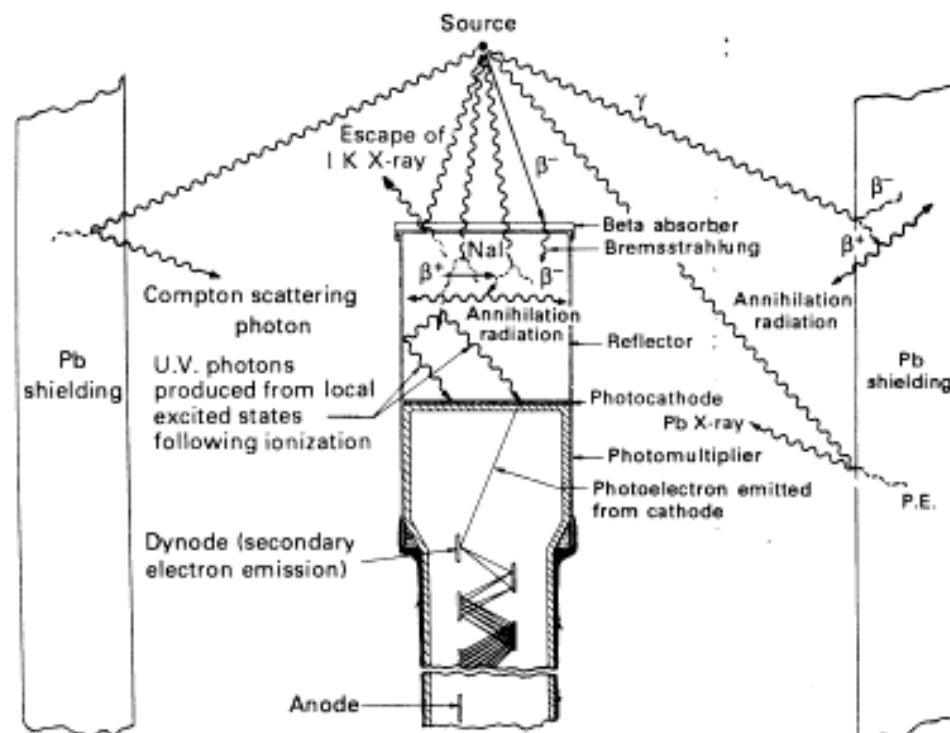
PMT Collection



Scintillator Detectors



Scintillation Detectors



Risposta del rivelatore - 1

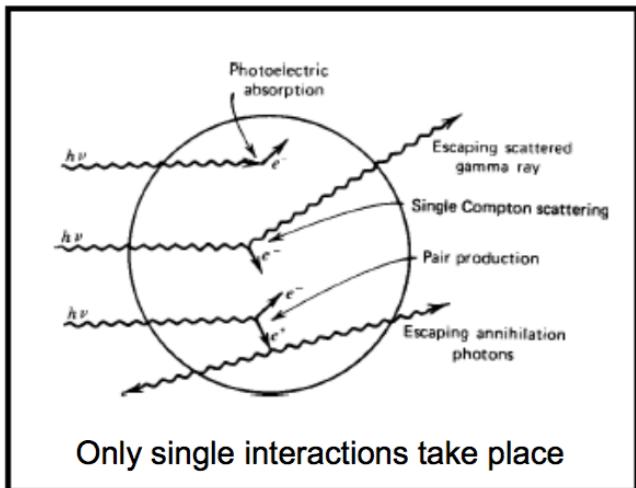


Figure 9: "Small" detector

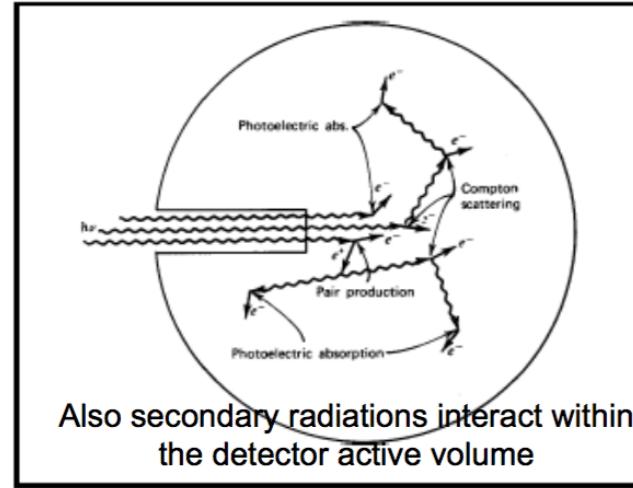


Figure 10: "Large" detector

most of the "secondary products" remain in the detector

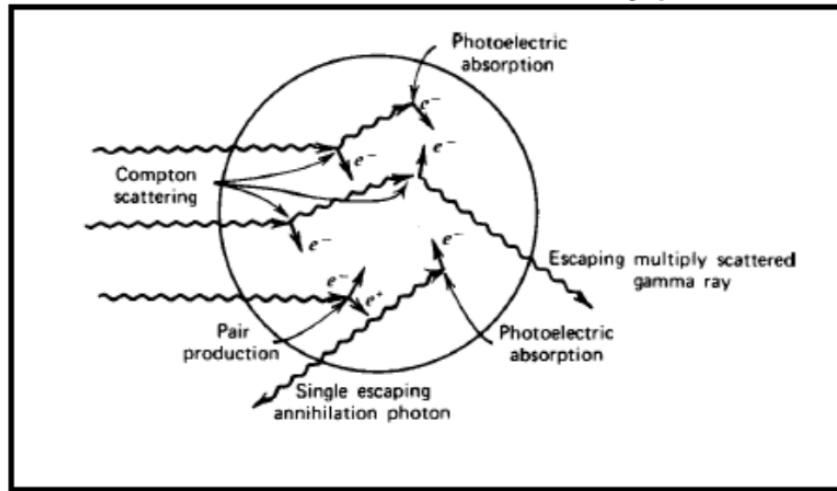
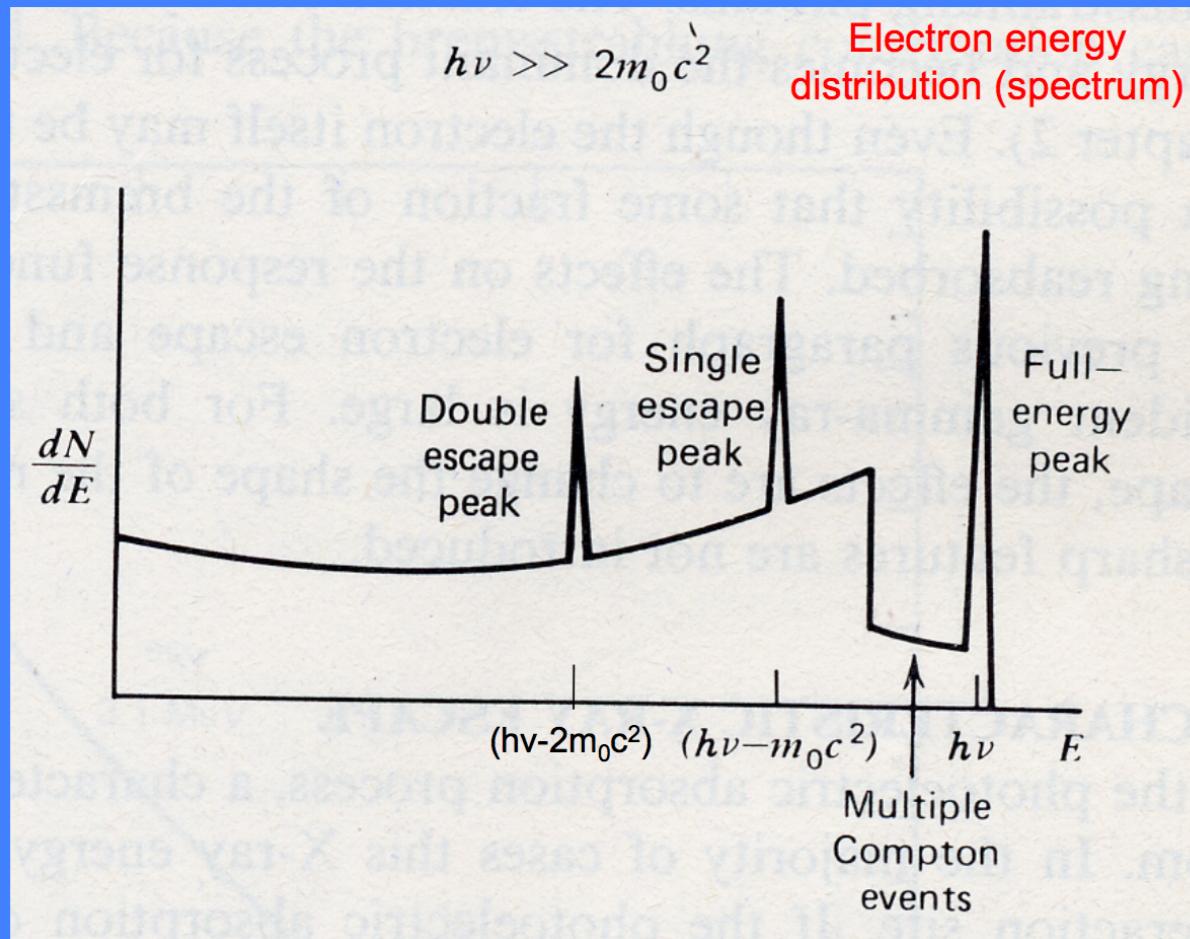


Figure 11: Intermediately sized detector

F. Knoll

Risposta del rivelatore - 2



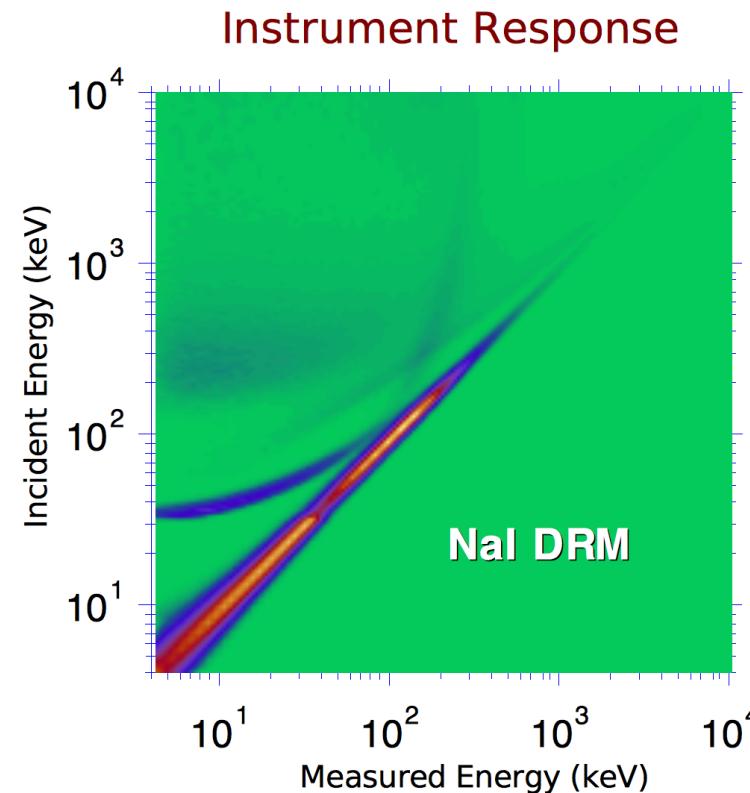
Case of intermediate-size detector (Knoll)

Photo-peak (full-energy peak): all photoelectric events remain in the detector and produce an energy deposit at the energy of the incoming photon

Single-escape peak: one annihilation photon leaves the detector without further interaction

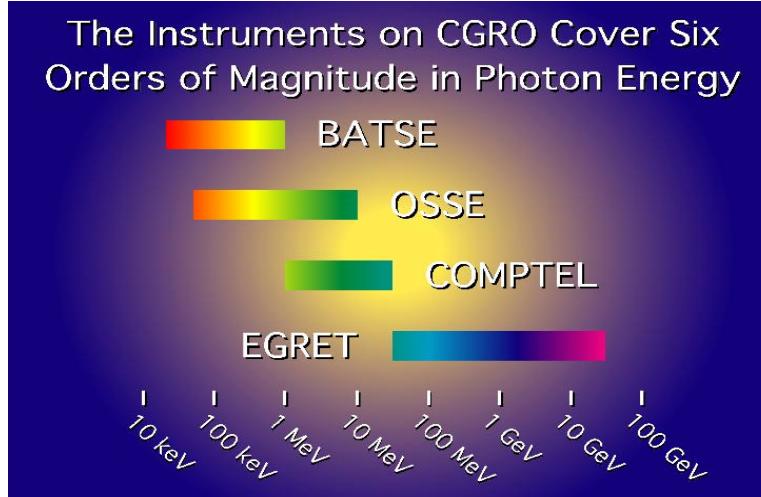
Double-escape peak: both annihilation photons leave the detector (escape)

Detector Response Matrix

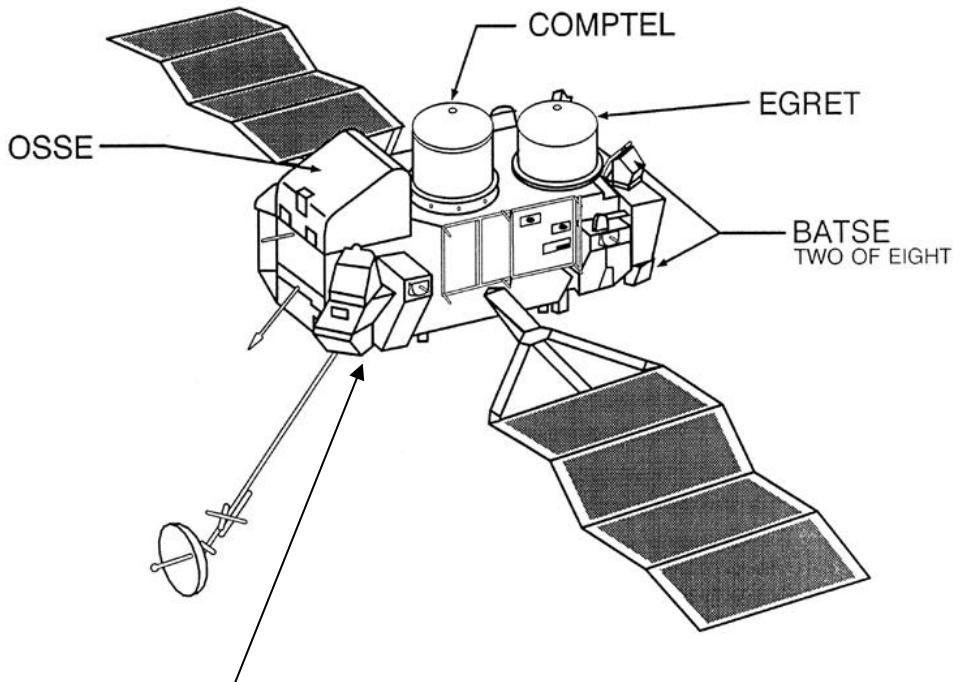


The response of a detector, whose signal depends on the energy of an incoming photon, distributes the photon of a certain energy over many pulse height channels according to the gain and energy resolution of the detector. Usually this resolution function is relatively complicated and depends on the photon energy. Since the energy acceptance and resolution of a given detector is determined by its design it is convenient to tabulate this function while the photon energy serves as a parameter. This procedure leads directly to a form of a matrix and gives the whole data set the name *detector response matrix*.

CGRO-BATSE (1991-2000)



COMPTON OBSERVATORY INSTRUMENTS



CGRO/BATSE (20 keV÷10 MeV)

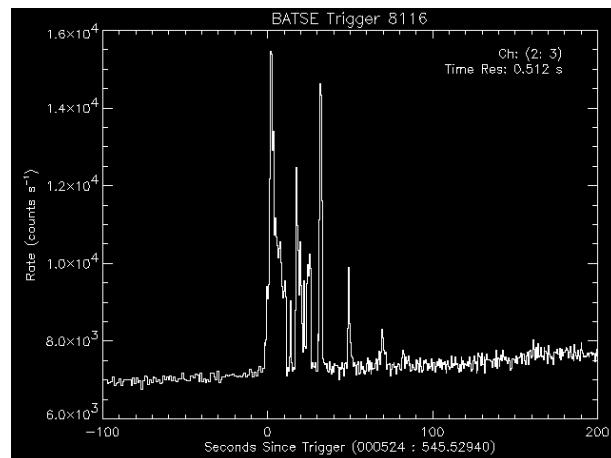
The BATSE instrument

- NaI scintillators
- 20 keV – 2 MeV
- FoV 4π (LAD)
- SD spectroscopy

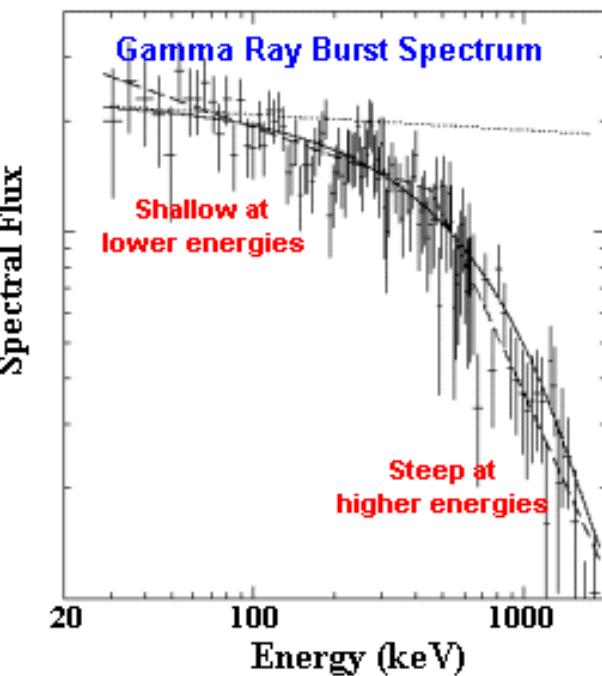


Gamma-Ray Bursts

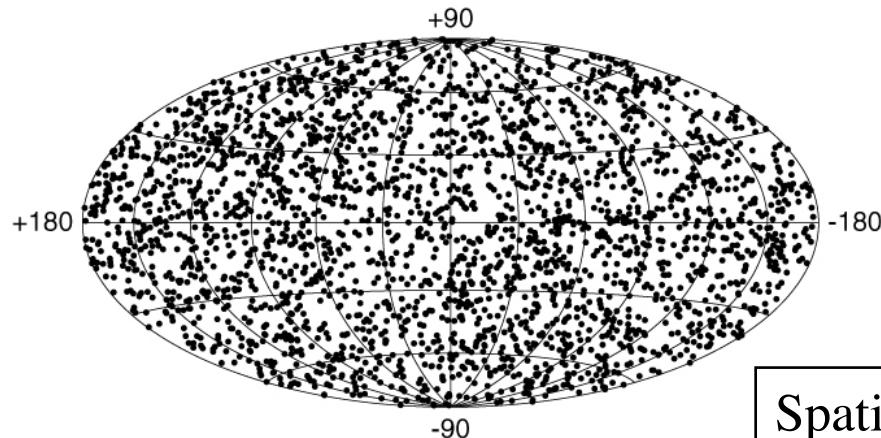
Temporal behaviour



Spectral shape



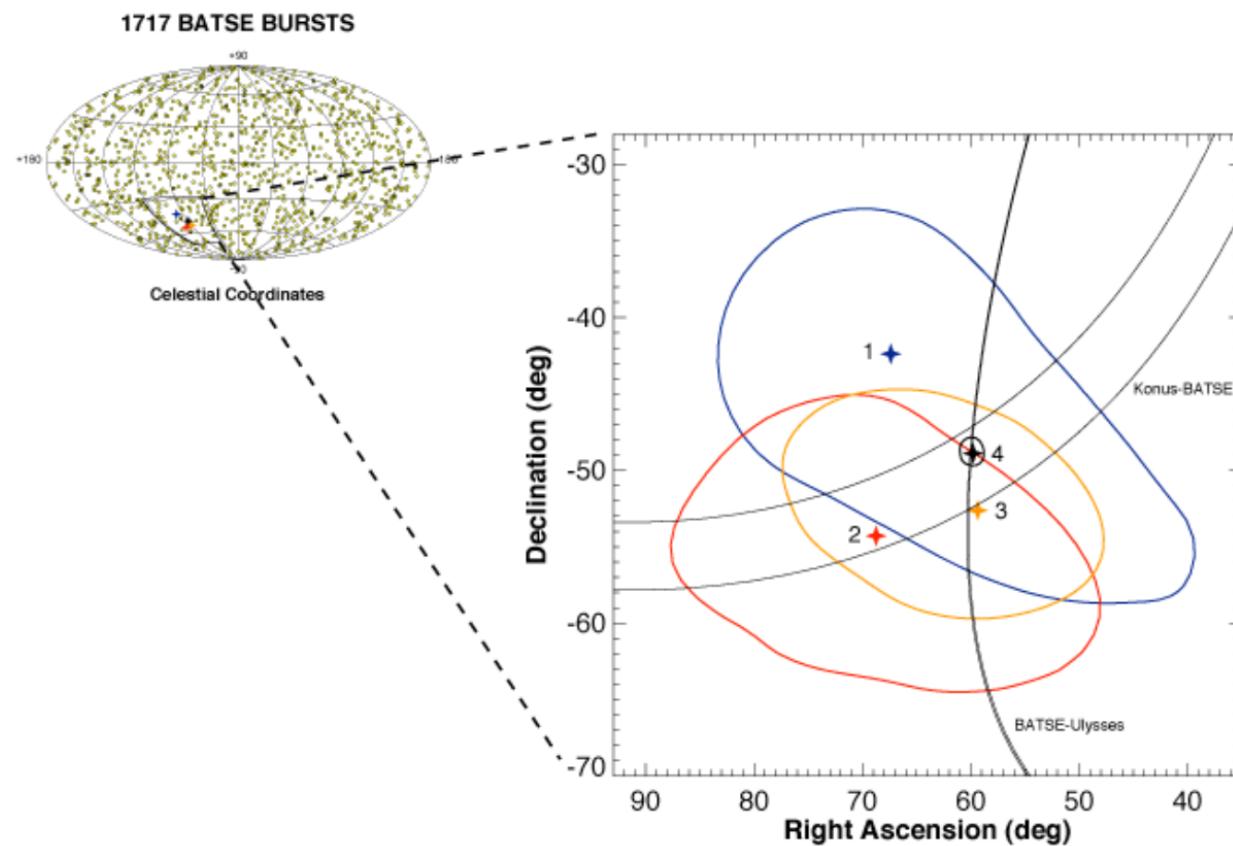
2704 BATSE Gamma-Ray Bursts



Spatial distribution

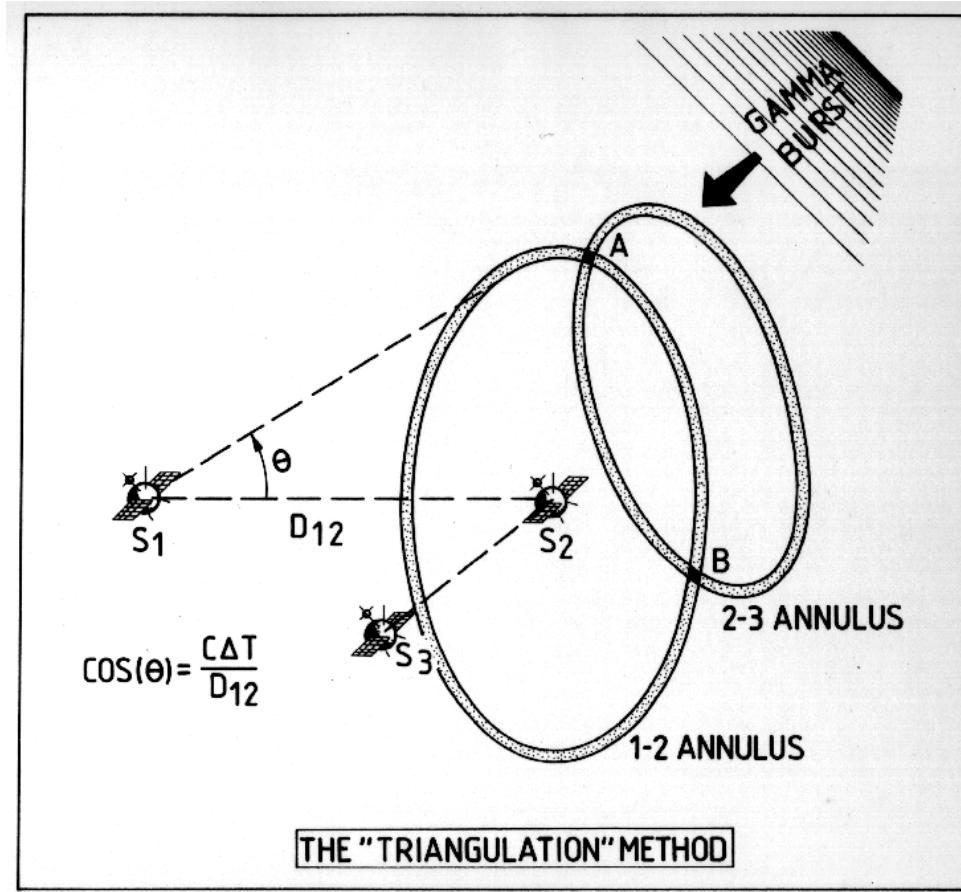
GRB localisation

BATSE GAMMA-RAY BURST CLUSTER



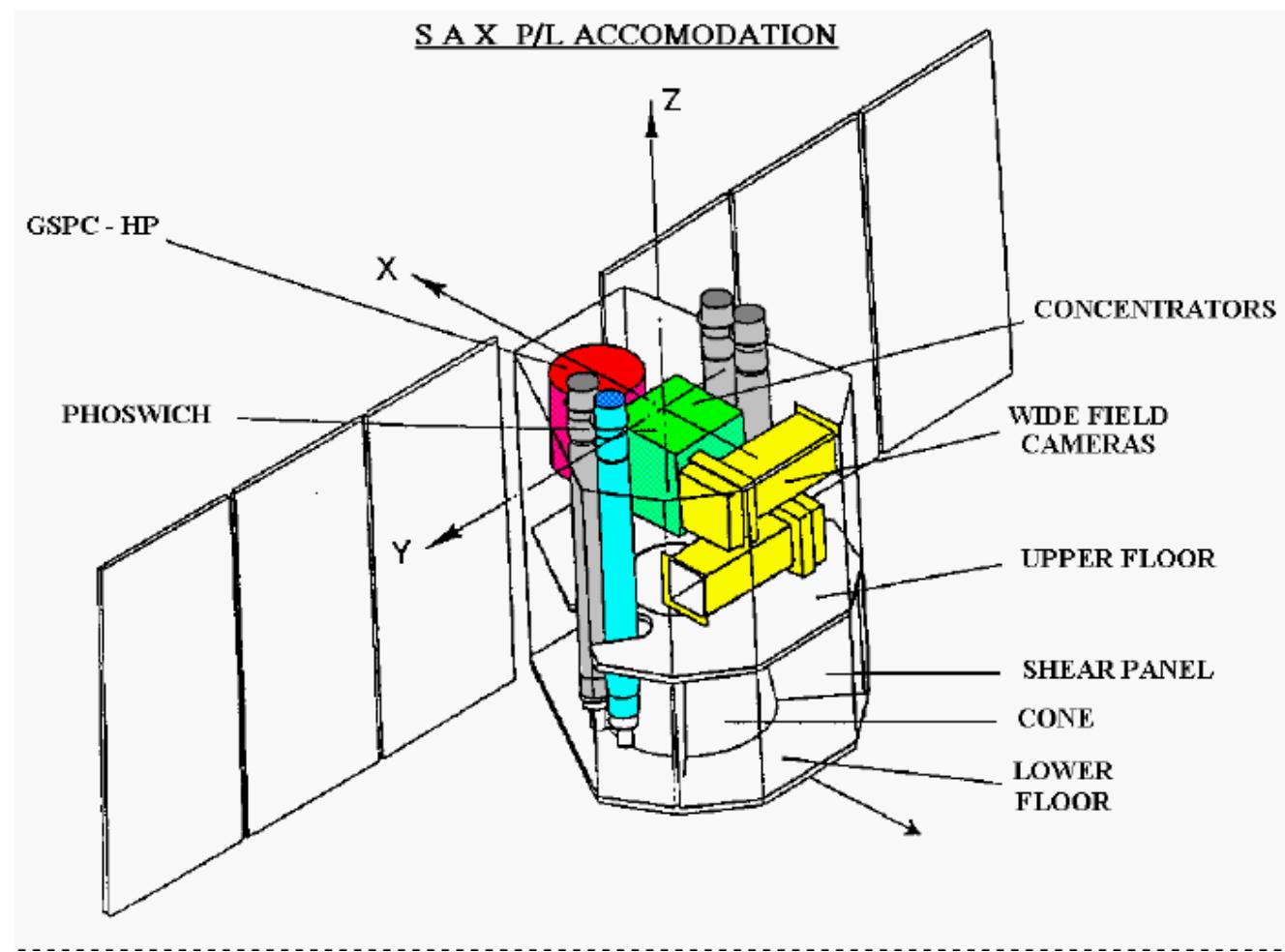
GRB History

- Interplanetary Network (IPN)

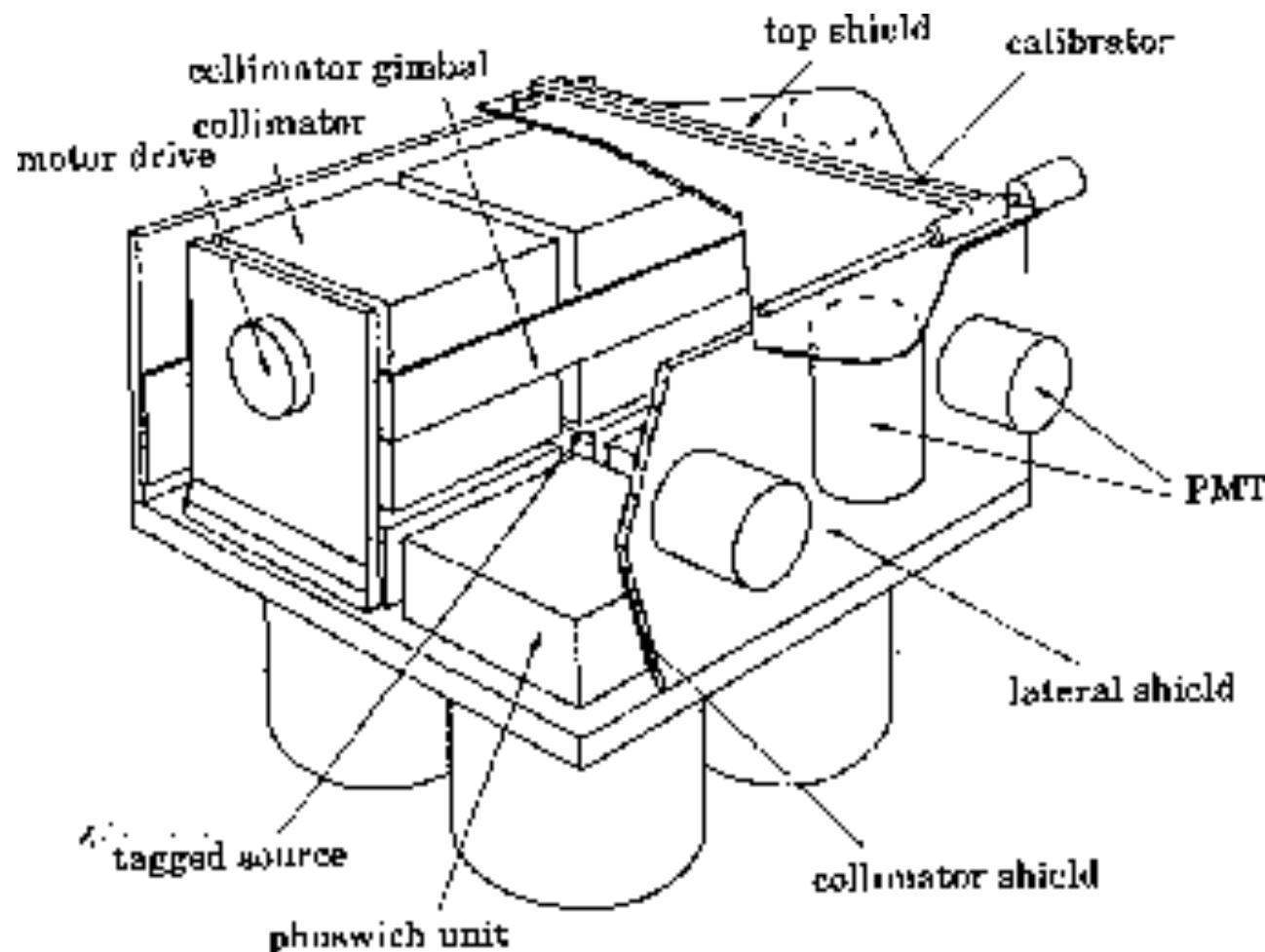


<http://www.ssl.berkeley.edu/ipn3/>

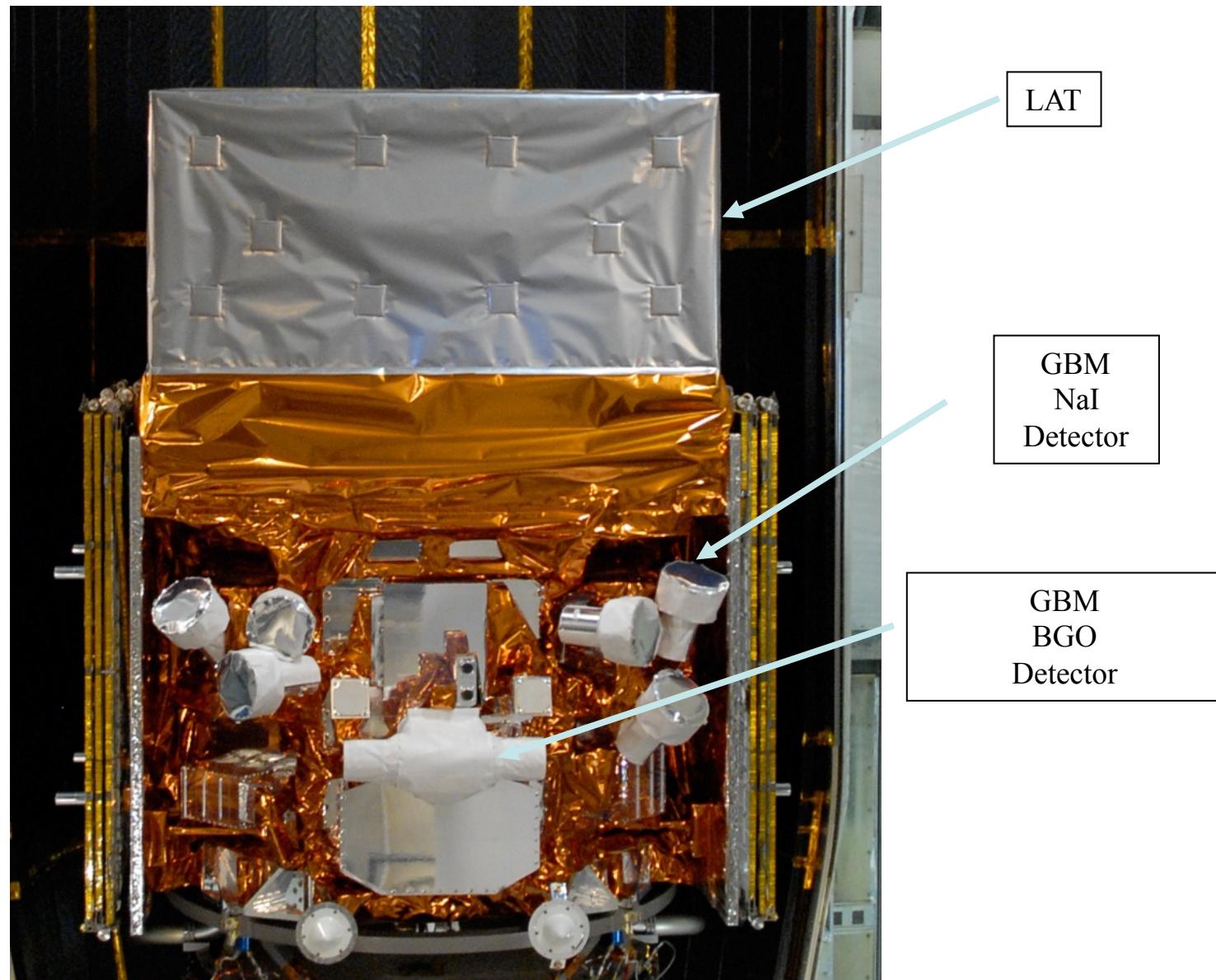
BeppoSAX (1995 - 2002)



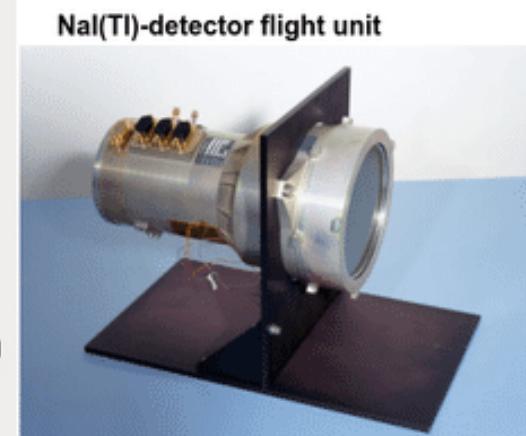
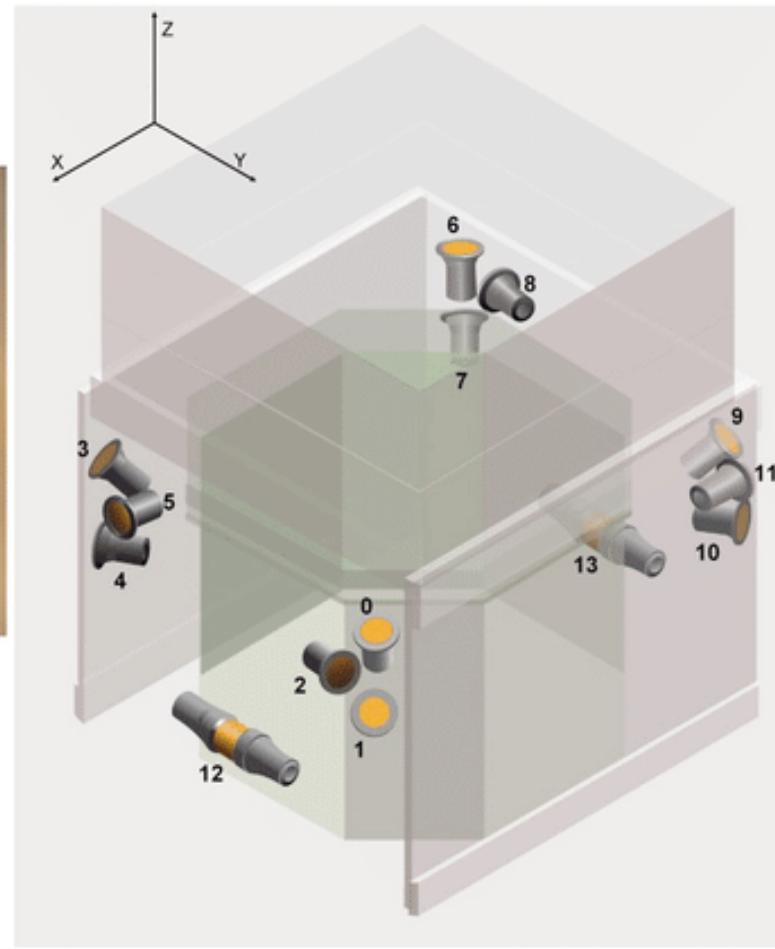
BeppoSAX (1995 - 2002)



Fermi/GBM detector (2008 -- ..)



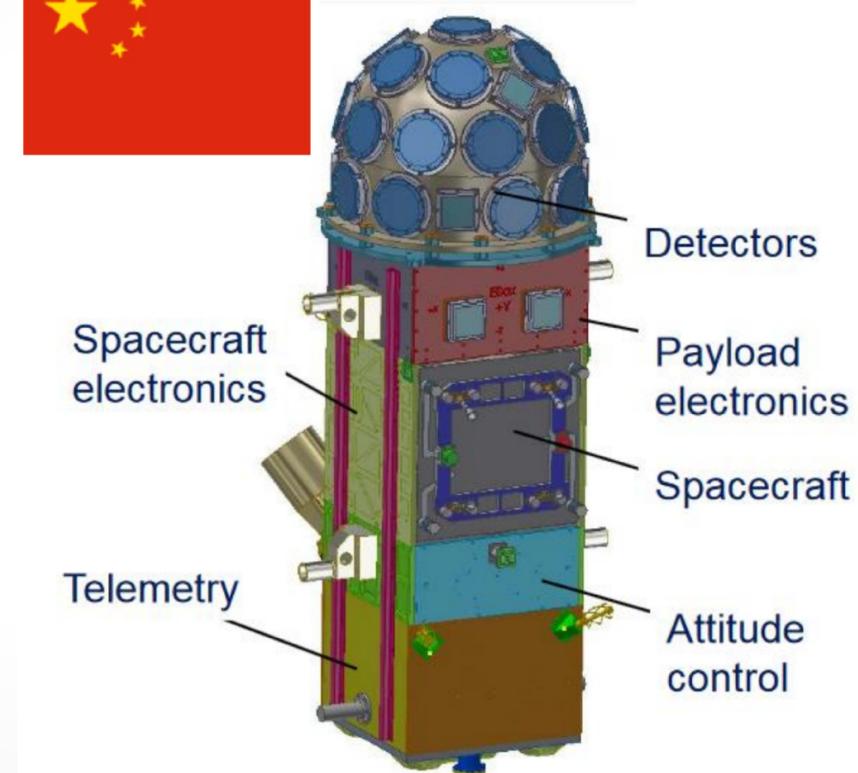
Fermi/GBM detector (2008 -- ..)



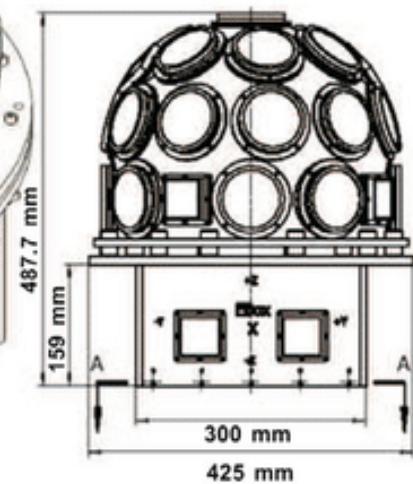
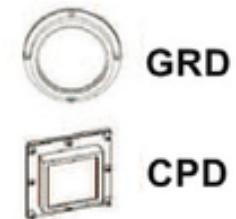
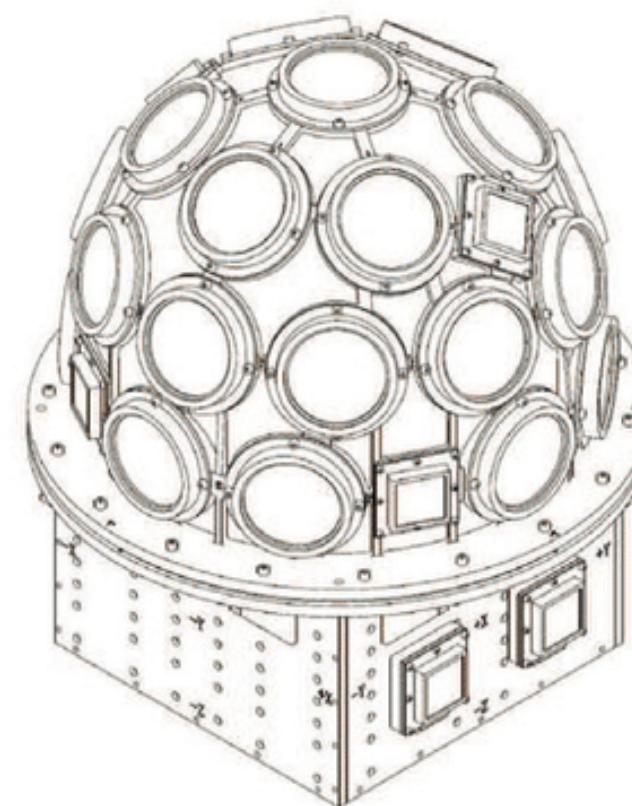
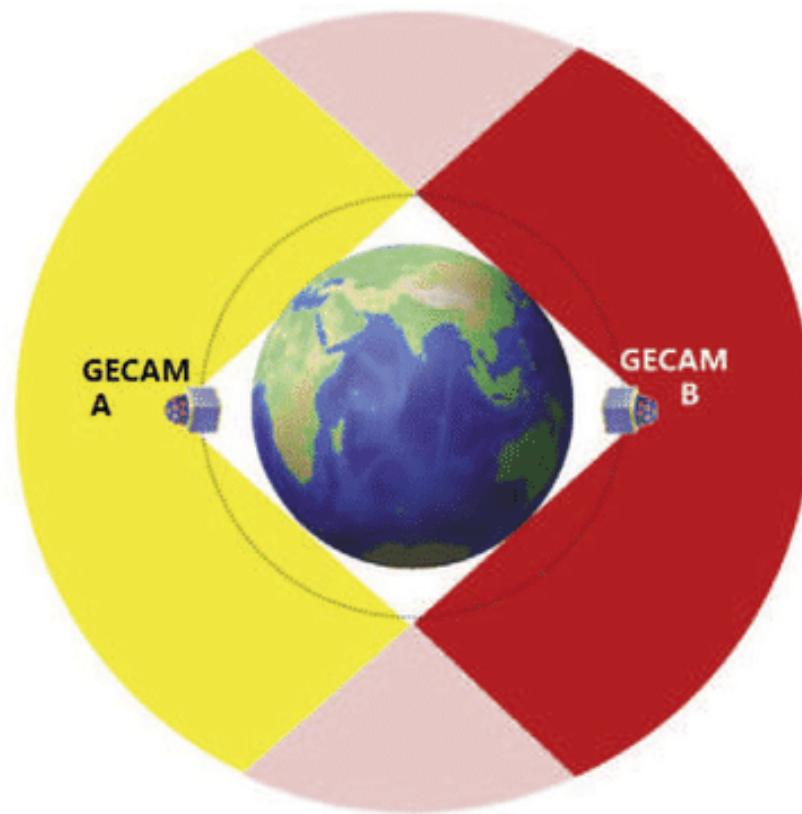
GECAM

GCAM

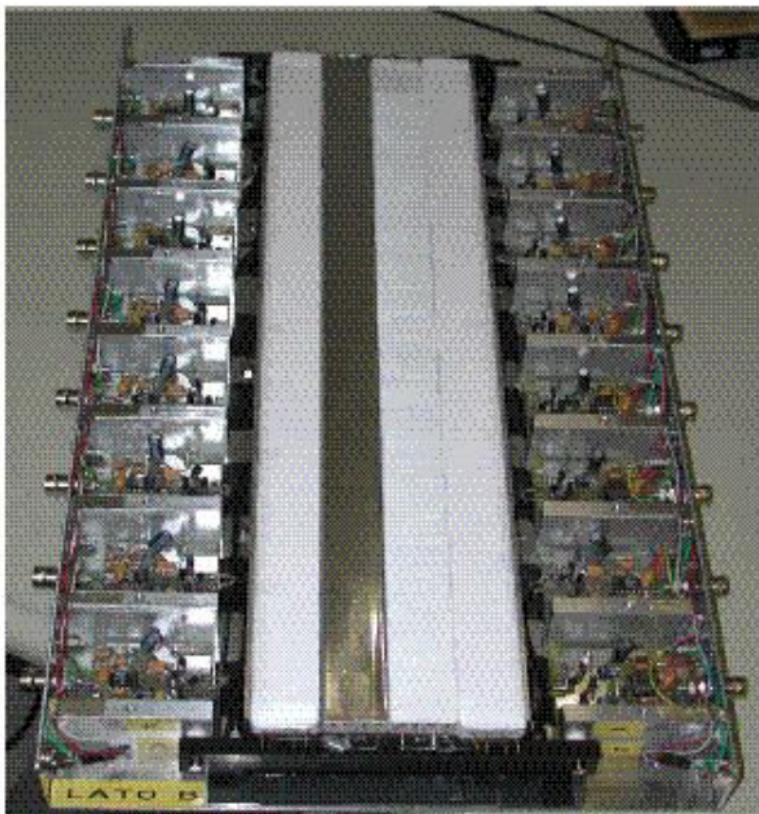
Chinese mission
Launch 2021
2 satellites
100% sky coverage.
Very similar to GBM.
Positioning by triangulation
(need 3 participants)
Some directionality.
Cannot do the work to alert CTA all
by itself.



GECAM



AGILE MCAL



MINI-CALORIMETER

DETECTOR

- 30 CsI bars wrapped with tight diffusion material organized in 2 orthogonal trays
 - bar dimension: 40x2.3x1.5 cm³
 - total radiation length: 1.5X₀ (in axis)

FRONTEND ELECTRONICS

- 1 photodiode on each side of the bar
 - optically coupled

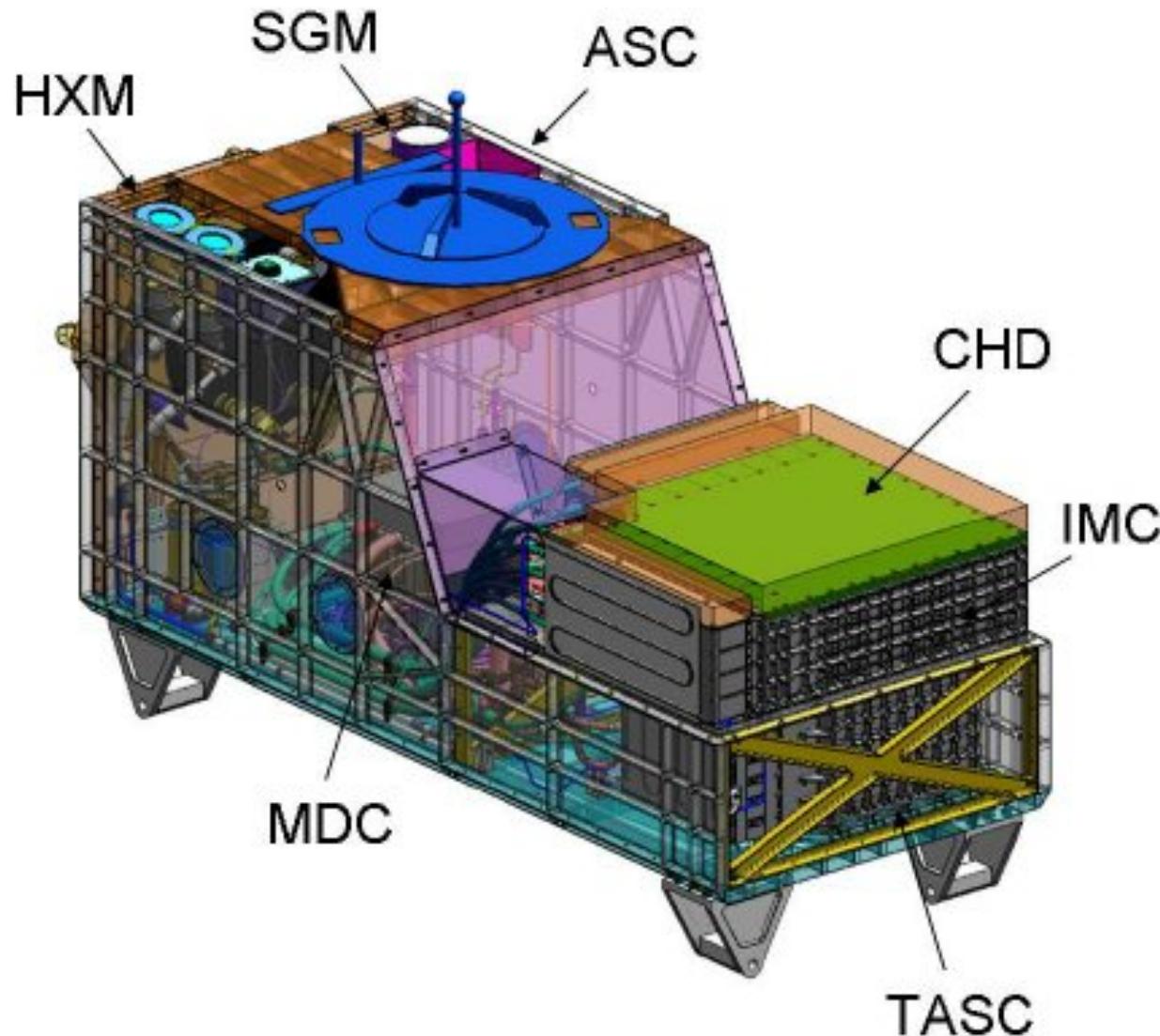
GOAL

- measure energy deposit of the photon conversion pair (GRID mode)
- detect GRBs and transients in the range 0.25-250MeV (BURST mode)

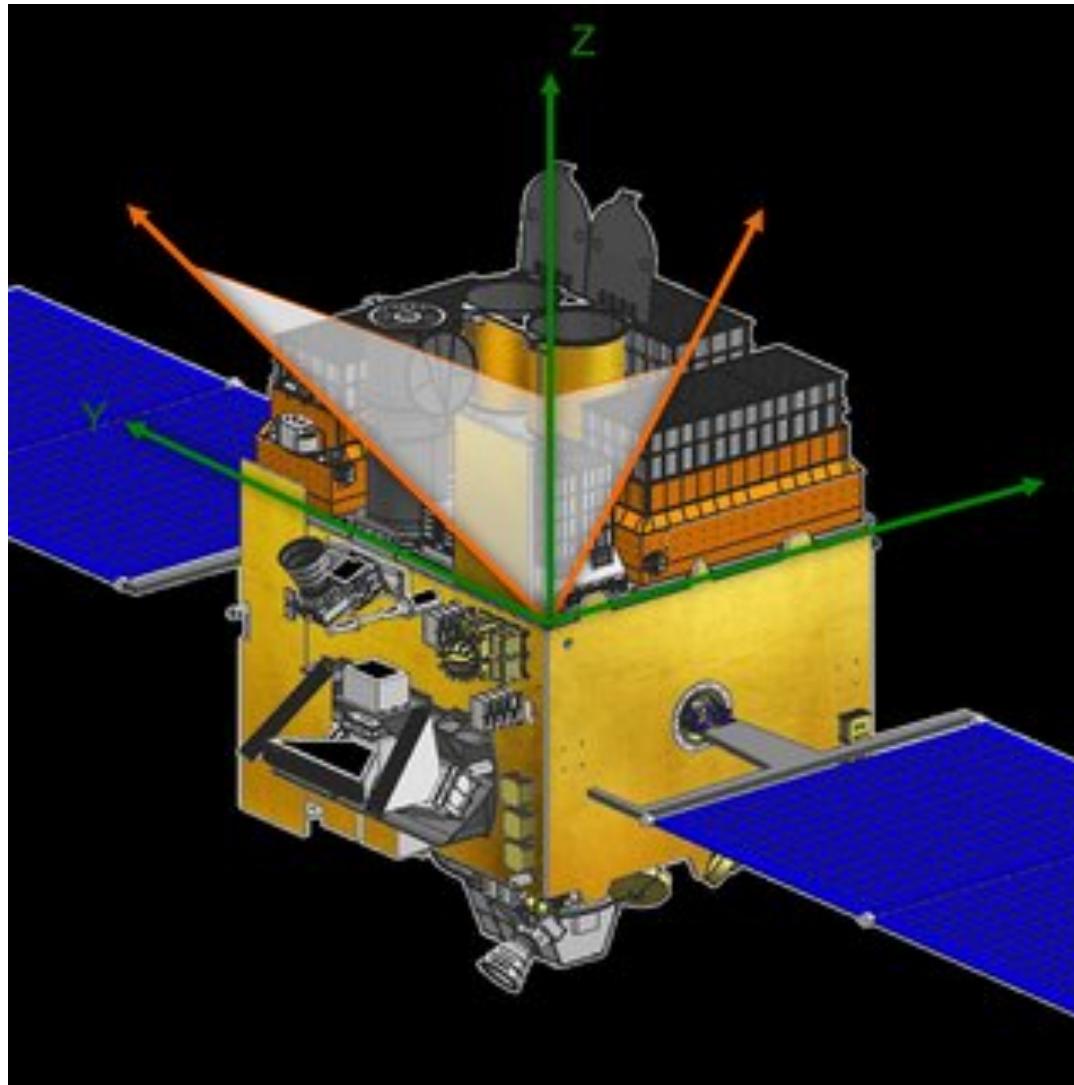
SCIENTIFIC FEATURES

- energy resolution: 22-24%(FWHM) @ 1MeV
0.7% @ 100MeV
- spatial resolution: 15mm @ 1MeV
2mm @ 100MeV
- timing resolution: 2_μs (BURST mode)

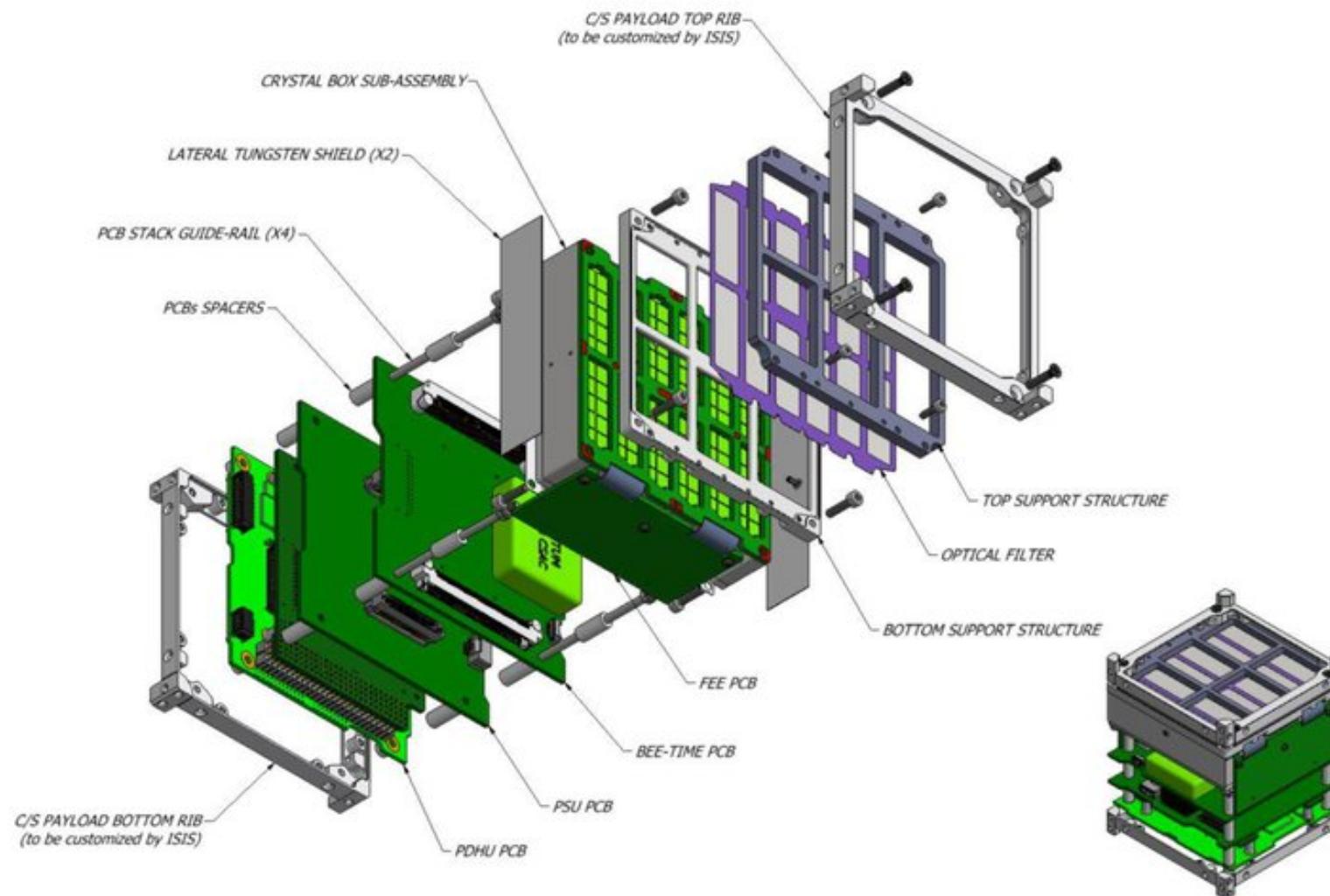
CALET



AstroSAT



HERMES



Exercise #1

- Find the web sites of BATSE
- Find the web site (if any) of BeppoSAX
- Find the web site of Fermi/GBM
- Find the web site of AGILE/MCAL GRB catalog
- Find the web site of CALET GRBM
- Find the web site of AstroSAT CZTI GRB