

A.6 Solution to Exercise 3.1

Consider the map

$$\mathcal{T}_s[\hat{\rho}(0)] = \hat{\rho}(s) = e^{-4s}\hat{\rho}(0) + \frac{(1 - e^{-4s})}{2}\hat{\mathbb{1}}. \quad (\text{A.20})$$

If we set $s = 0$, we trivially find

$$\mathcal{T}_0[\hat{\rho}(0)] = \hat{\rho}(0), \quad (\text{A.21})$$

thus $\mathcal{T}_0 = \text{id}$. Then, we apply \mathcal{T}_t to the expression in Eq. (A.20):

$$\begin{aligned} \mathcal{T}_t[\hat{\rho}(s)] &= e^{-4t}\hat{\rho}(s) + \frac{(1 - e^{-4t})}{2}\hat{\mathbb{1}}, \\ &= e^{-4t} \left(e^{-4s}\hat{\rho}(0) + \frac{(1 - e^{-4s})}{2}\hat{\mathbb{1}} \right) + \frac{(1 - e^{-4t})}{2}\hat{\mathbb{1}}, \\ &= e^{-4(t+s)}\hat{\rho}(0) + \frac{e^{-4t}(1 - e^{-4s})}{2}\hat{\mathbb{1}} + \frac{(1 - e^{-4t})}{2}\hat{\mathbb{1}}, \\ &= e^{-4(t+s)}\hat{\rho}(0) + \frac{(1 - e^{-4(t+s)})}{2}\hat{\mathbb{1}}, \end{aligned} \quad (\text{A.22})$$

which covers $\mathcal{T}_t\mathcal{T}_s = \mathcal{T}_{t+s}$. Finally, $\lim_{h \rightarrow 0} \mathcal{T}_{t+h}[\hat{\rho}(0)] = \mathcal{T}_t[\hat{\rho}(0)]$ can be obtained by applying the limit to Eq. (A.22).

A.7 Solution to Exercise 3.2

For $\epsilon = i\epsilon_0$ with $\epsilon_0 \in \mathbb{R}$, we have that Eq. (A.27) becomes

$$d|\psi_t\rangle = \left[-\frac{i}{\hbar}\hat{H}dt + \sum_k \left(-\epsilon_0\hat{L}_k dW_{k,t} - \frac{1}{2}\epsilon_0^2\hat{L}_k^\dagger\hat{L}_k dt \right) \right] |\psi_t\rangle, \quad (\text{A.23})$$

where, notably, in the last term one has ϵ_0^2 in place of $\epsilon^2 = -\epsilon_0^2$. This is due to the constraint on the conservation of the norm for $|\psi_t\rangle$, which reflects in the preservation of the trace of $\hat{\rho}_t$. Similarly, the equation for $\langle\psi_t|$ reads

$$d\langle\psi_t| = \langle\psi_t| \left[\frac{i}{\hbar}\hat{H}dt + \sum_k \left(-\epsilon_0\hat{L}_k^\dagger dW_{k,t} - \frac{1}{2}\epsilon_0^2\hat{L}_k^\dagger\hat{L}_k dt \right) \right]. \quad (\text{A.24})$$

By following the calculations already performed, we find the following master equation

$$d\hat{\rho}_t = -\frac{i}{\hbar} \left[\hat{H}, \hat{\rho}_t \right] dt + \epsilon_0^2 \sum_k \left(\hat{L}_k \hat{\rho}_t \hat{L}_k^\dagger - \frac{1}{2} \left\{ \hat{L}_k^\dagger \hat{L}_k, \hat{\rho}_t \right\} \right) dt, \quad (\text{A.25})$$

which has the same form of Eq. (3.79) with ϵ substituted by ϵ_0 .

A.8 Solution to Exercise 3.3

Let us consider Eq. (3.79) with $\hat{L}_k^\dagger = \hat{L}_k$. In such a case, the master equation can be rewritten as

$$d\hat{\rho}_t = -\frac{i}{\hbar} [\hat{H}, \hat{\rho}_t] dt - \frac{\epsilon^2}{2} \sum_k [\hat{L}_k, [\hat{L}_k, \hat{\rho}_t]] dt, \quad (\text{A.26})$$

which can be obtained from the unravelling

$$d|\psi_t\rangle = \left[-\frac{i}{\hbar} \hat{H} dt + \sum_k \left(i\epsilon \hat{L}_k dW_{k,t} - \frac{1}{2} \epsilon^2 \hat{L}_k^2 dt \right) \right] |\psi_t\rangle. \quad (\text{A.27})$$

The relation between the two holds also for the specific case of $\hat{L}_k = \hat{A}_k - \langle \psi_t | \hat{A}_k | \psi_t \rangle$, for which the unravelling becomes non-linear in $|\psi_t\rangle$:

$$d|\psi_t\rangle = \left\{ -\frac{i}{\hbar} \hat{H} dt + \sum_k \left[i\epsilon \left(\hat{A}_k - \langle \psi_t | \hat{A}_k | \psi_t \rangle \right) dW_{k,t} - \frac{1}{2} \epsilon^2 \left(\hat{A}_k - \langle \psi_t | \hat{A}_k | \psi_t \rangle \right)^2 dt \right] \right\} |\psi_t\rangle. \quad (\text{A.28})$$

However, in such a case, in the commutators appearing in Eq. (A.26) have the following structure:

$$\left[\hat{A}_k - \langle \psi_t | \hat{A}_k | \psi_t \rangle, \hat{X} \right] = \left[\hat{A}_k, \hat{X} \right], \quad (\text{A.29})$$

which is valid for any arbitrary operator \hat{X} . It follows that the master equation becomes

$$d\hat{\rho}_t = -\frac{i}{\hbar} [\hat{H}, \hat{\rho}_t] dt - \frac{\epsilon^2}{2} \sum_k \left[\hat{A}_k, [\hat{A}_k, \hat{\rho}_t] \right] dt, \quad (\text{A.30})$$

which notably is linear in the state $\hat{\rho}_t$.