

A.9 Solution to Exercise 4.1

To express the Hadamard gate H as a rotation, we proceed as follows. We consider the general rotation

$$\hat{R}^{\mathbf{n}}(\theta) = \cos(\theta/2)\hat{1} - i \sin(\theta/2)\mathbf{n} \cdot \hat{\sigma}, \quad (\text{A.31})$$

of an angle θ around \mathbf{n} , where the Pauli matrices are

$$X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad Y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}. \quad (\text{A.32})$$

From such a rotation, we want to obtain

$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}. \quad (\text{A.33})$$

First thing, we highlight that the sum $X + Z$ gives

$$X + Z = \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}, \quad (\text{A.34})$$

then it follows that

$$H = \frac{X + Z}{\sqrt{2}} = \left\{ \frac{1}{\sqrt{2}}, 0, \frac{1}{\sqrt{2}} \right\} \cdot \hat{\sigma}. \quad (\text{A.35})$$

Such an expression recalls the last term in Eq. (A.31). Finally, we need to set the angle θ so that the first term in Eq. (A.31) vanishes. This is $\theta = \pi$. Then

$$H = i\hat{R}^{\mathbf{n}}(\pi), \quad \text{where } \mathbf{n} = \left\{ \frac{1}{\sqrt{2}}, 0, \frac{1}{\sqrt{2}} \right\}, \quad (\text{A.36})$$

gives the solution.

A.10 Solution to Exercise 4.2

To prove that, given two fixed non-parallel normalised vectors \mathbf{n} and \mathbf{m} , any unitary \hat{U} can be expressed as

$$\hat{U} = e^{i\alpha} \hat{R}^{\mathbf{n}}(\beta) \hat{R}^{\mathbf{m}}(\gamma) \hat{R}^{\mathbf{n}}(\delta), \quad (\text{A.37})$$

with $\alpha, \beta, \gamma, \delta \in \mathbb{R}$, then one needs to recast \hat{U} in the form

$$\hat{U} = e^{i\alpha} \hat{R}^{\mathbf{t}}(\omega), \quad (\text{A.38})$$

with $\alpha \in \mathbb{R}$ and $\mathbf{t} \in \mathbb{R}^3$ suitably chosen.

The first step of the proof is to write $(\mathbf{m} \cdot \hat{\sigma})(\mathbf{n} \cdot \hat{\sigma})$ in terms of a single Pauli matrix vector $\hat{\sigma}$. We have

$$\begin{aligned} (\mathbf{m} \cdot \hat{\sigma})(\mathbf{n} \cdot \hat{\sigma}) &= (m_1\hat{\sigma}_x + m_2\hat{\sigma}_y + m_3\hat{\sigma}_z)(n_1\hat{\sigma}_x + n_2\hat{\sigma}_y + n_3\hat{\sigma}_z), \\ &= \mathbf{m} \cdot \mathbf{n} + m_1n_2\hat{\sigma}_x\hat{\sigma}_y + m_1n_3\hat{\sigma}_x\hat{\sigma}_z + m_2n_1\hat{\sigma}_y\hat{\sigma}_x + m_2n_3\hat{\sigma}_y\hat{\sigma}_z + m_3n_1\hat{\sigma}_z\hat{\sigma}_x + m_3n_2\hat{\sigma}_z\hat{\sigma}_y. \end{aligned} \quad (\text{A.39})$$

By applying Eq. (A.15), Eq. (A.39) becomes

$$(\mathbf{m} \cdot \hat{\sigma})(\mathbf{n} \cdot \hat{\sigma}) = (\mathbf{m} \cdot \mathbf{n})\hat{1} + i(\mathbf{m} \times \mathbf{n}) \cdot \hat{\sigma}. \quad (\text{A.40})$$

The second step is to consider the composition of two rotations:

$$\begin{aligned}
\hat{R}^{\mathbf{m}}(\gamma)\hat{R}^{\mathbf{n}}(\delta) &= (\cos(\gamma/2)\hat{\mathbf{1}} - i\sin(\gamma/2)\mathbf{m}\cdot\hat{\boldsymbol{\sigma}})(\cos(\delta/2)\hat{\mathbf{1}} - i\sin(\delta/2)\mathbf{n}\cdot\hat{\boldsymbol{\sigma}}), \\
&= \cos(\gamma/2)\cos(\delta/2)\hat{\mathbf{1}} - i\cos(\gamma/2)\sin(\delta/2)\mathbf{n}\cdot\hat{\boldsymbol{\sigma}} - i\cos(\delta/2)\sin(\gamma/2)\mathbf{m}\cdot\hat{\boldsymbol{\sigma}} \\
&\quad - \sin(\gamma/2)\sin(\delta/2)(\mathbf{m}\cdot\hat{\boldsymbol{\sigma}})(\mathbf{n}\cdot\hat{\boldsymbol{\sigma}}).
\end{aligned} \tag{A.41}$$

Substituting Eq. (A.40) in the last expression, we find that

$$\hat{R}^{\mathbf{m}}(\gamma)\hat{R}^{\mathbf{n}}(\delta) = \hat{R}^{\mathbf{h}}(\epsilon) = \cos(\epsilon/2)\hat{\mathbf{1}} - i\sin(\epsilon/2)\mathbf{h}\cdot\hat{\boldsymbol{\sigma}}, \tag{A.42}$$

where ϵ and \mathbf{h} are taken such that

$$\begin{aligned}
\cos(\epsilon/2) &= \cos(\gamma/2)\cos(\delta/2) - \sin(\gamma/2)\sin(\delta/2)\mathbf{m}\cdot\mathbf{n}, \\
\sin(\epsilon/2)\mathbf{h} &= \cos(\gamma/2)\sin(\delta/2)\mathbf{n} + \cos(\delta/2)\sin(\gamma/2)\mathbf{m} + \sin(\gamma/2)\sin(\delta/2)(\mathbf{m}\times\mathbf{n}).
\end{aligned} \tag{A.43}$$

Then, we have

$$\hat{R}^{\mathbf{n}}(\beta)\hat{R}^{\mathbf{m}}(\gamma)\hat{R}^{\mathbf{n}}(\delta) = \hat{R}^{\mathbf{n}}(\beta)\hat{R}^{\mathbf{h}}(\epsilon) = \hat{R}^{\mathbf{t}}(\omega), \tag{A.44}$$

where we applied again the composition of two rotations, which ends the proof.

A.11 Solution to Exercise 4.3

Consider two qubits, where the first is prepared in the superposition

$$|\psi_1\rangle = \frac{|0\rangle + |1\rangle}{\sqrt{2}}, \tag{A.45}$$

while the second is initialised in the ground state $|\psi_2\rangle = |0\rangle$. The total state is

$$|\psi_{12}\rangle = \frac{|0\rangle + |1\rangle}{\sqrt{2}}|0\rangle = \frac{|00\rangle + |10\rangle}{\sqrt{2}}. \tag{A.46}$$

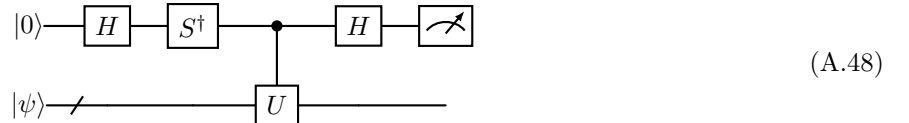
From the first expression, one clearly sees that the state is separable. By applying the CNOT gate, we find that the state becomes

$$|\psi_{12}\rangle = \frac{|00\rangle + |11\rangle}{\sqrt{2}}, \tag{A.47}$$

which is a fully entangled state.

A.12 Solution to Exercise 4.4

Consider the circuit



Its action is the following

$$\begin{aligned}
|0\rangle |\psi\rangle &\xrightarrow{\hat{H} \otimes \hat{\mathbb{1}}} \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle) |\psi\rangle \\
&\xrightarrow{\hat{S}^\dagger \otimes \hat{\mathbb{1}}} \frac{1}{\sqrt{2}}(|0\rangle - i|1\rangle) |\psi\rangle \\
&\xrightarrow{C(U)} \frac{1}{\sqrt{2}}(|0\rangle |\psi\rangle - i|1\rangle \hat{U} |\psi\rangle) \\
&\xrightarrow{\hat{H} \otimes \hat{\mathbb{1}}} \frac{1}{2} \left[(|0\rangle + |1\rangle) |\psi\rangle - i(|0\rangle - |1\rangle) \hat{U} |\psi\rangle \right] \\
&= \frac{1}{2} \left[|0\rangle (\hat{\mathbb{1}} - i\hat{U}) |\psi\rangle + |1\rangle (\hat{\mathbb{1}} + i\hat{U}) |\psi\rangle \right].
\end{aligned} \tag{A.49}$$

Finally, one measures qubit 0, and the probability of finding the qubit in $|0\rangle$ is

$$P(|0\rangle) = \frac{1}{4} \langle \psi | (\hat{\mathbb{1}} + i\hat{U}^\dagger) (\hat{\mathbb{1}} - i\hat{U}) |\psi\rangle = \frac{1}{2} \left(1 + \Im \langle \psi | \hat{U} | \psi \rangle \right), \tag{A.50}$$

which ends the exercise.