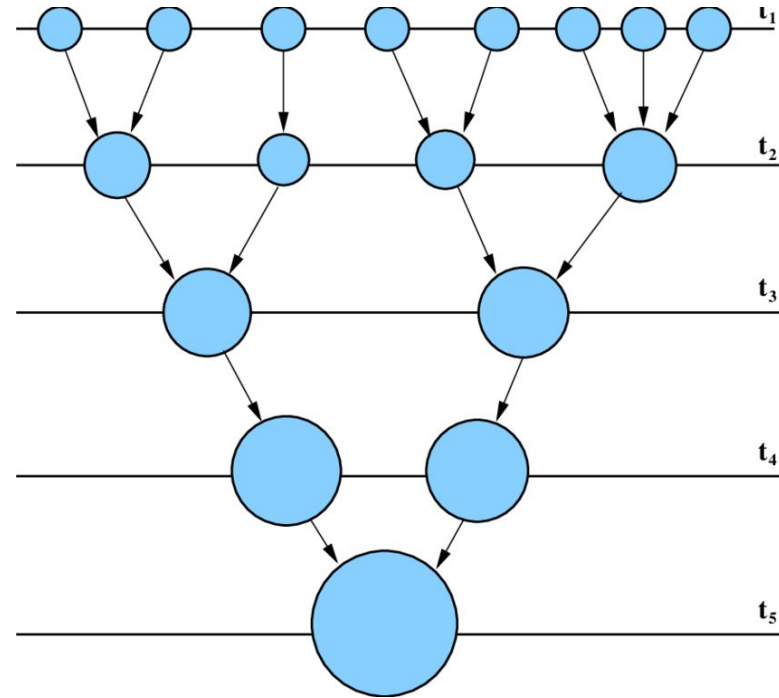
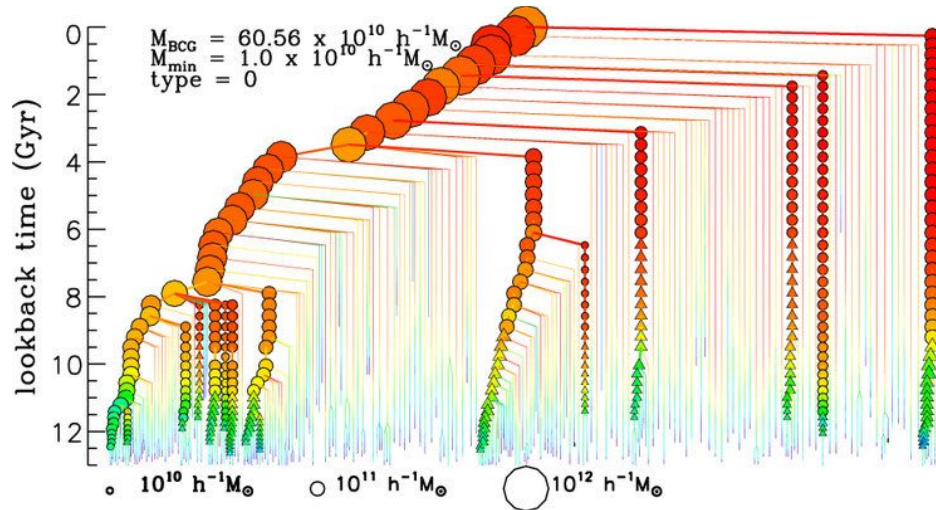


**STATISTICAL PROPERTIES OF THE LARGE SCALE STRUCTURES:
CLUSTER NUMBER COUNTS**

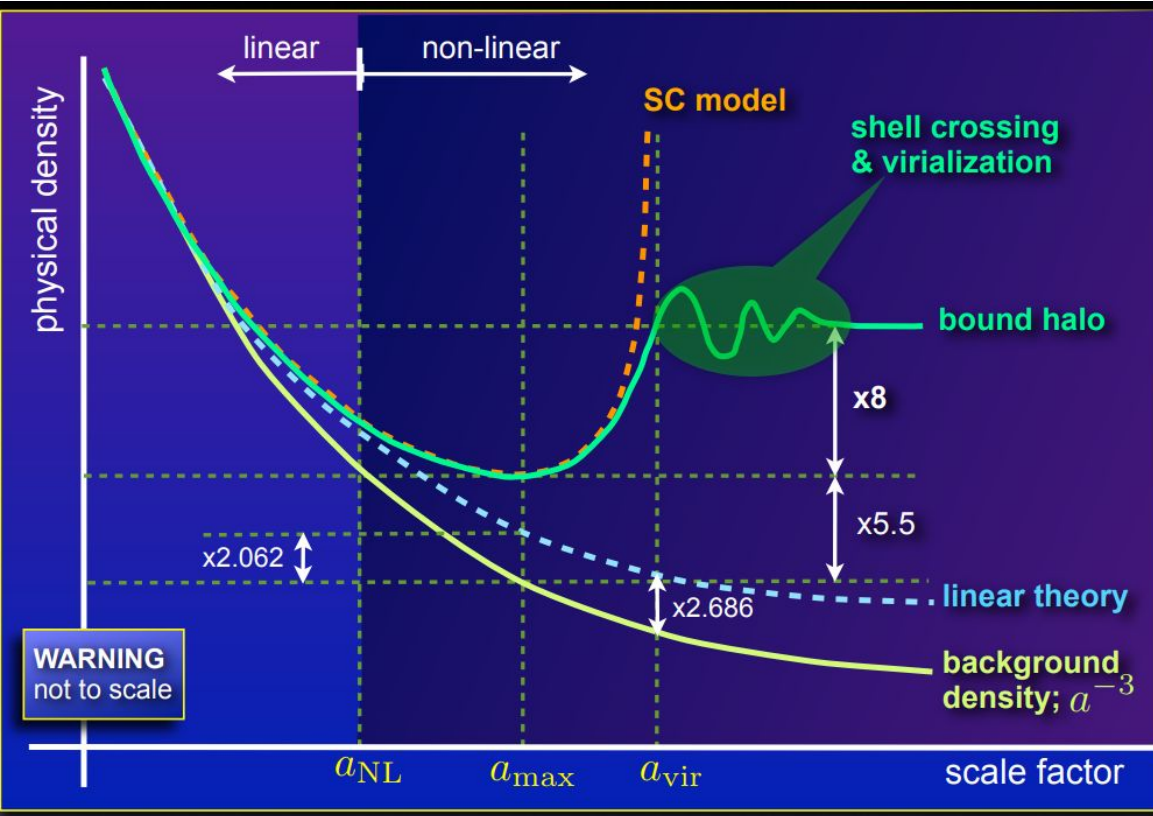
For a review: [Allen+2011](#) or [Kravtsov+2012](#)

STRUCTURE FORMATION: DARK MATTER HALOS

In the LCDM scenario, structures grow *hierarchically*: Small overdensities are able to overcome the cosmological expansion and collapse first, and the resulting dark matter "halos" merge together to form larger halos which serve as sites of galaxy and galaxy cluster formation



STRUCTURE FORMATION: SPHERICAL COLLAPSE MODEL



We can follow the collapse of a spherical overdensity in a homogeneous universe. SC model becomes inaccurate (brakes down) shortly after turn-around it is still a useful model to identify important epochs in the linearly evolved density field.

- The linearly extrapolated density field collapses when $\delta_{lin} = \delta_c = 1.686$
- Virialized dark matter haloes have an average overdensity of $\Delta_{vir} = 178$

STRUCTURE FORMATION: SPHERICAL COLLAPSE MODEL

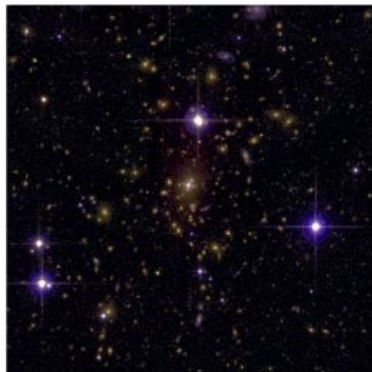


According to the spherical collapse model, regions with $\delta(\mathbf{x}, t) > \delta_c \approx 1.686$ will have collapsed to produce dark matter haloes by time t

GALAXY CLUSTERS

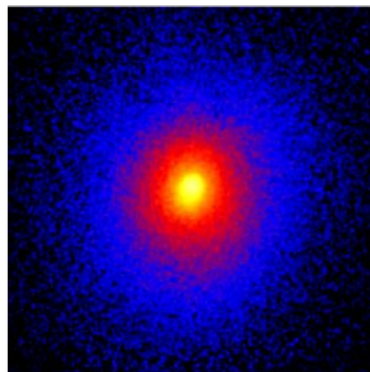
- **Most massive bound objects in the Universe:**
 $M \approx 10^{13} - 10^{15} M_{\odot}$ and $R \approx 1 - 5$ Mpc
- **Multi-component systems:**
Galaxies and stars (~5%), ICM (~15%), DM (~80%)

OPTICAL



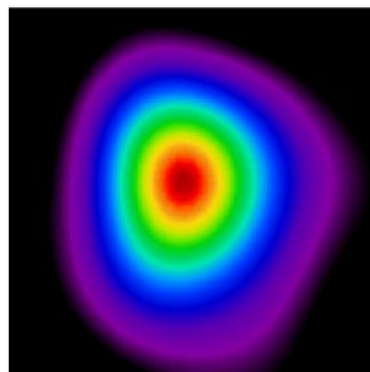
RICHNESS, LENSING EFFECTS

X-RAYS

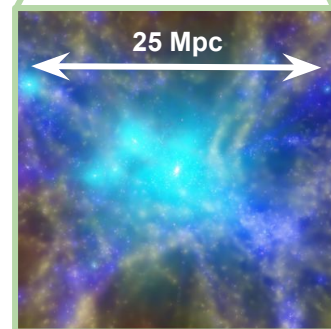
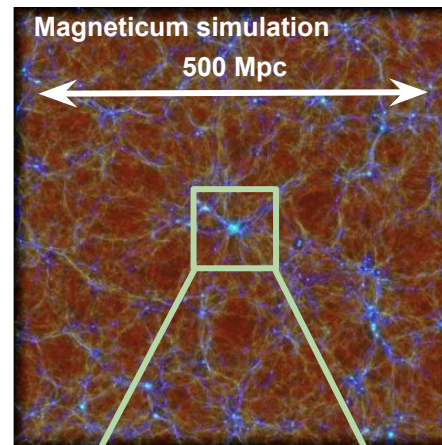


LUMINOUS AND EXTENDED X-RAY SOURCES

MICROWAVES



SUNYAEV-ZEL'DOVICH EFFECT



From Hirschmann+2014

GALAXY CLUSTERS AS COSMOLOGICAL PROBE

The abundance and spatial distribution of galaxy clusters are sensitive to the **growth rate** of cosmic structures and **expansion history** of the Universe

σ_8 : Amplitude of the matter power spectrum

Ω_m : Present-day total matter density

$$S_8 = \sigma_8 (\Omega_m / 0.3)^{0.5}$$

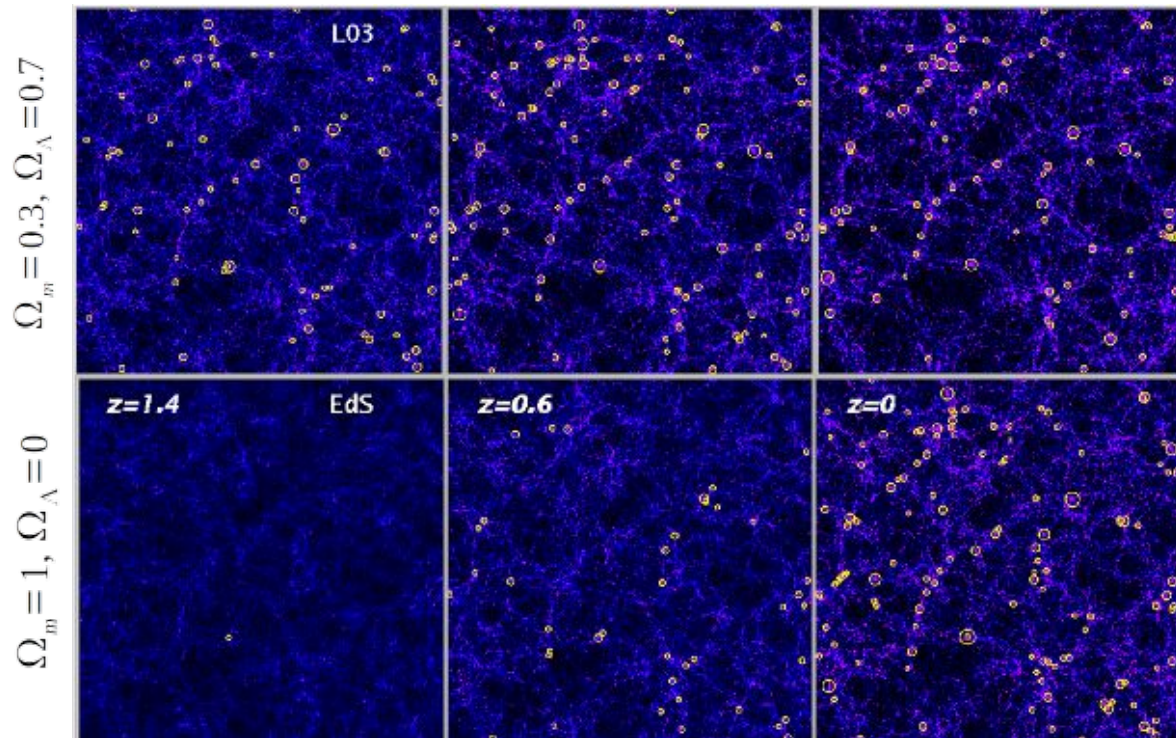
Dark energy equation of state parameter w

Total neutrino mass

Deviation from GR

....

Evolution of the clusters population in 2 N-body simulations



time

From Borgani, Guzzo 2001

THE HALO MASS FUNCTION

Cluster abundance:

$$\frac{dN}{dzd\Omega} = \frac{dV}{dzd\Omega} n(M, z)$$

● geometry
● growth

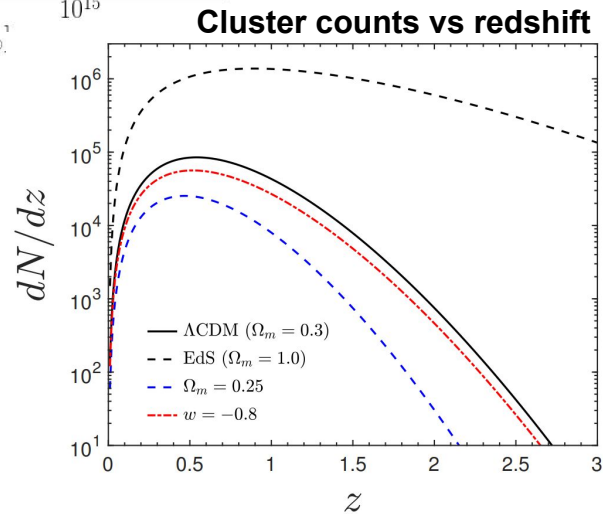
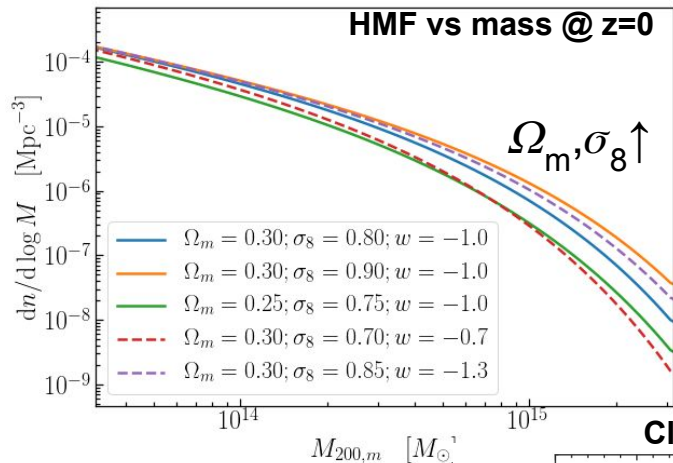
The halo mass function:

$$n(z, M) = \frac{\rho_m}{M} f(\sigma) \frac{d \ln(\sigma^{-1})}{dM}$$

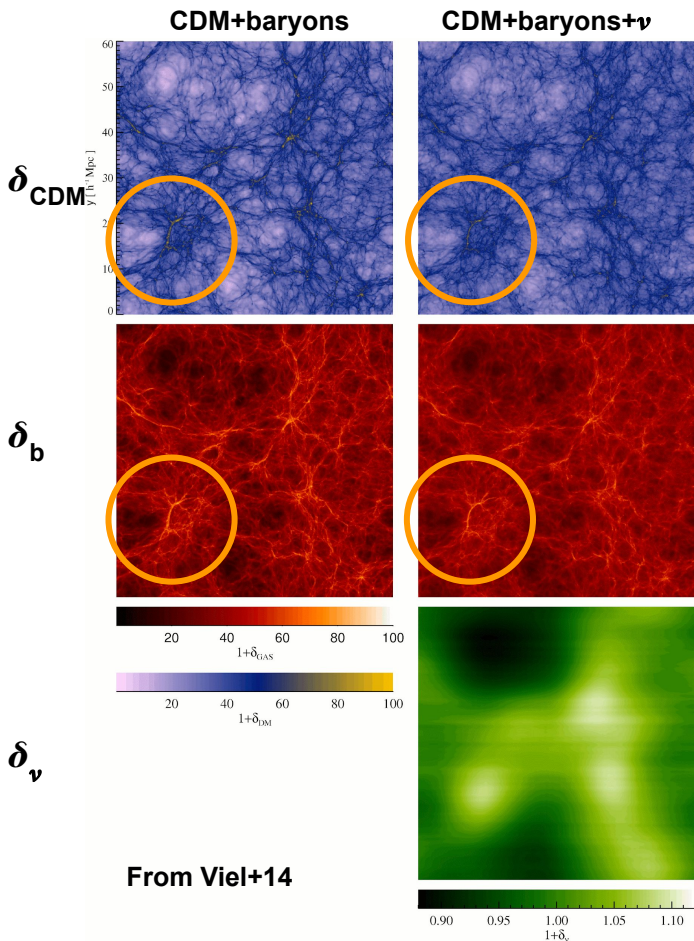
Variance of the density field:

$$\sigma(z, R) = \frac{1}{2\pi^2} \int_0^\infty dk k^2 P_m(z, k) |W(kR)|^2$$

● Matter power spectrum



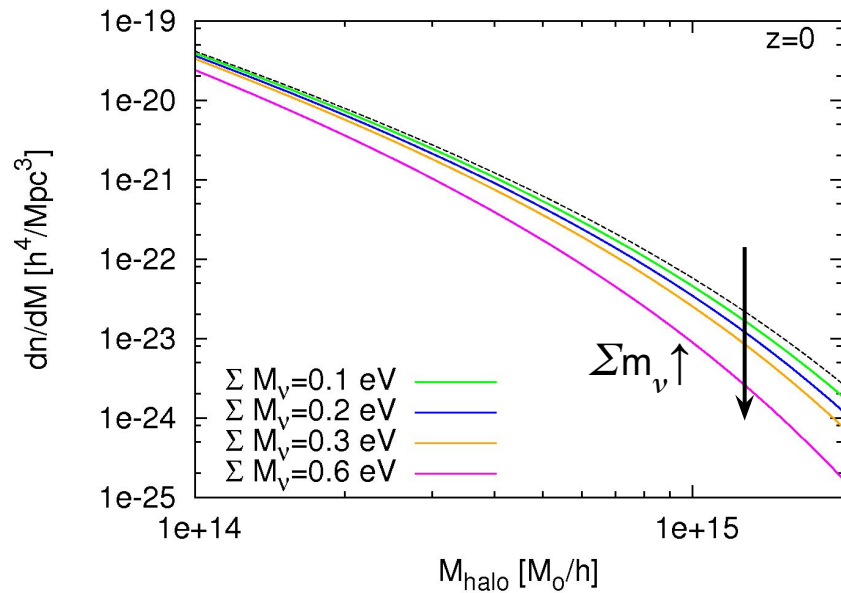
THE HALO MASS FUNCTION: MASSIVE NEUTRINOS



Massive neutrinos:

- Delay the epoch of matter-radiation equality
- Suppress the growth of density fluctuation on scale smaller than the free-streaming length

Effects on the number density of halos as a function of mass



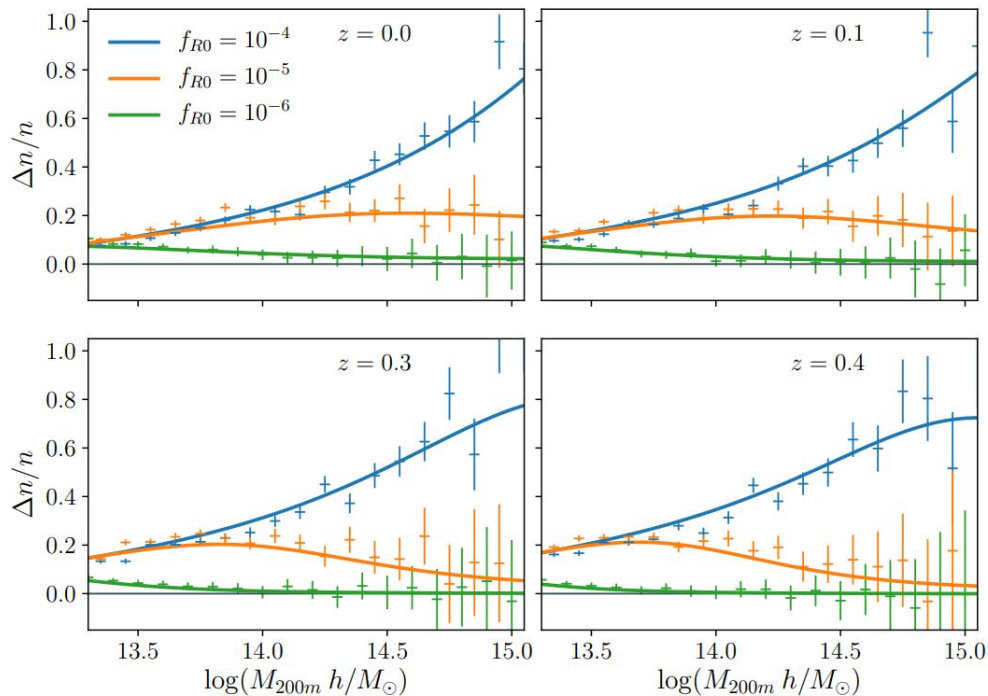
THE HALO MASS FUNCTION: MODIFIED GRAVITY

Modified gravity models, e.g. $f(R)$:

$$S = \frac{1}{16\pi G} \int \sqrt{-g}[R + f(R)]d^4x.$$

- Give rise to accelerated expansion and enhance gravity
- Introduce screening mechanism that restores GR in high density environments

Relative effect on the Halo Mass Function compared to Λ CDM



From Hagstotz+18

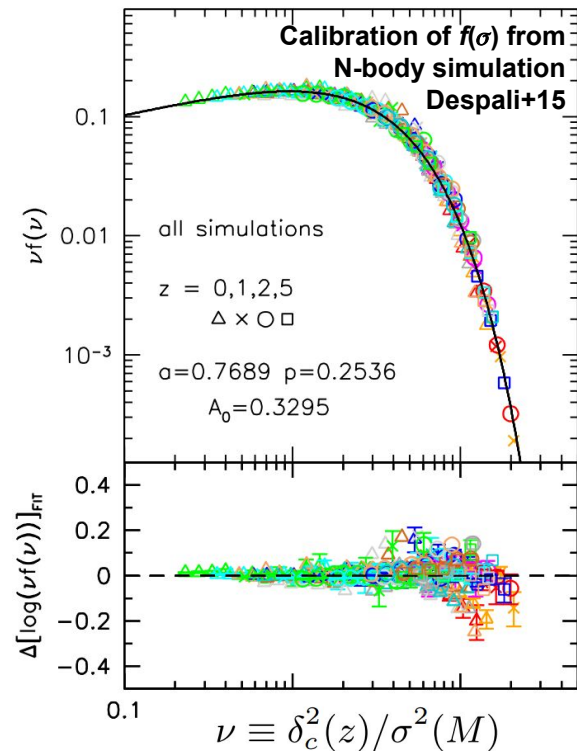
THE MULTIPLICITY FUNCTION: $f(\sigma)$

Halo mass function:
$$n(z, M) = \frac{\rho_m}{M} f(\sigma) \frac{d \ln(\sigma^{-1})}{dM}$$

- $f(\sigma)$ “universal” function:

- Press & Schechter (1974) approximated from spherical collapse of Gaussian density field
- Improved modeling using ellipsoidal collapse, e.g. Sheth & Tormen (1999)
- Nowadays calibrated against N-body simulations

Reference	Functional form
Press & Schechter (1974)	$f_{\text{PS}}(\sigma) = \sqrt{\frac{2}{\pi}} \frac{\delta_c}{\sigma} \exp\left(-\frac{\delta_c^2}{2\sigma^2}\right)$
Sheth & Tormen (1999)	$f_{\text{ST}}(\sigma) = A \sqrt{\frac{2a}{\pi}} \frac{\delta_c}{\sigma} \exp\left(-\frac{a \delta_c^2}{2\sigma^2}\right) \left[1 + \left(\frac{\sigma^2}{a \delta_c^2}\right)^p\right]$
Jenkins et al. (2001)	$f_{\text{J}}(\sigma) = A \exp(- \ln \sigma^{-1} + B ^p)$
Reed et al. (2003)	$f_{\text{R}}(\sigma) = f_{\text{ST}}(\sigma) \exp\left(\frac{-a}{\sigma(\cosh 2\sigma)^b}\right)$
Warren et al. (2006)	$f_{\text{W}}(\sigma) = A \left(\sigma^{-a} + b\right) \exp\left(-\frac{c}{\sigma^2}\right)$
Tinker et al. (2008)	$f_{\text{T}}(\sigma) = A \left[\left(\frac{\sigma}{b}\right)^{-a} + 1\right] \exp\left(-\frac{c}{\sigma^2}\right)$



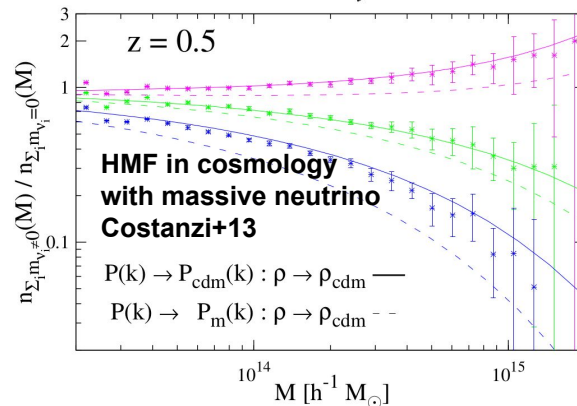
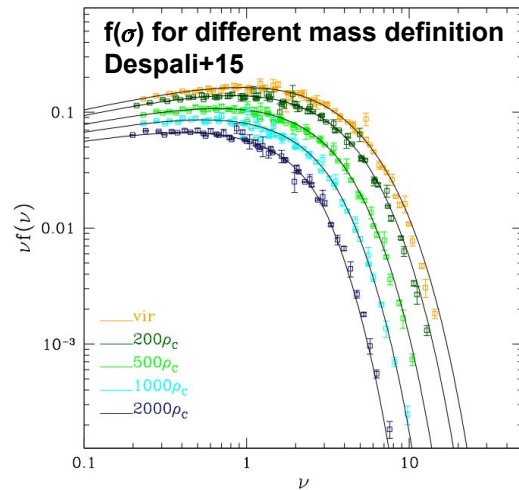
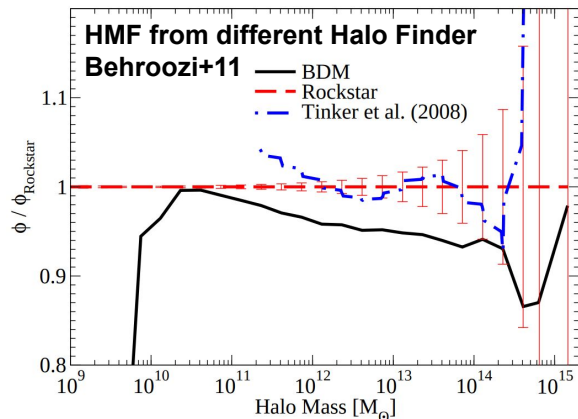
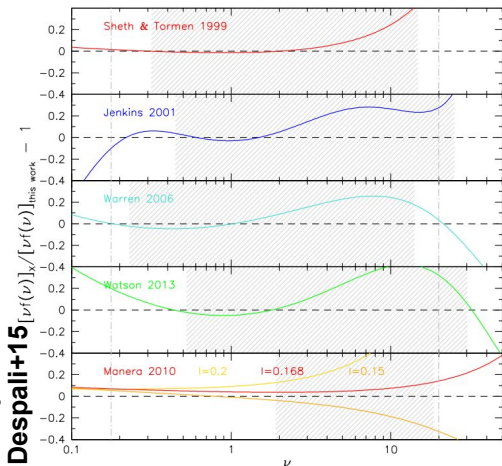
See also: Despali+15 ; Castro+22

HALO MASS FUNCTION: UNIVERSALITY

How accurate is the calibration of $f(\sigma)$? Is it universal?

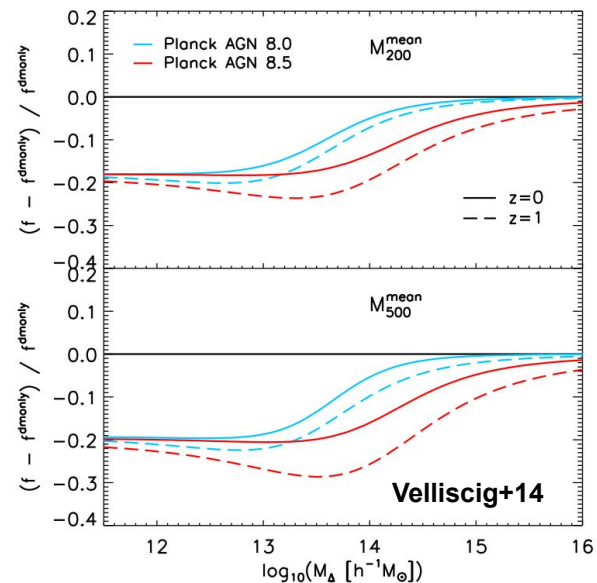
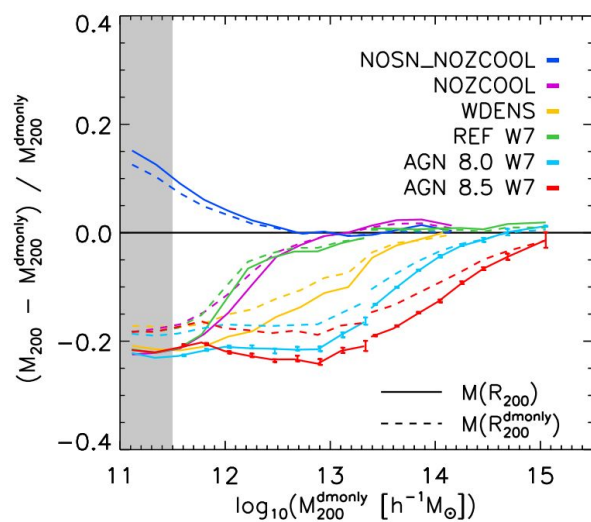
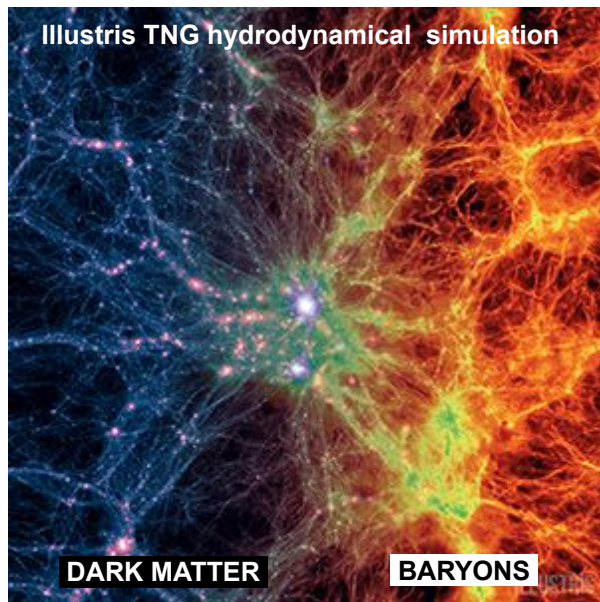
- Specific of the simulation (e.g. box size, number of particles, softening length)
- Halo finder (e.g. linking length, FoF, SO)
- Mass definition (e.g. $M_{200,m}$, $M_{500,c}$)
- Redshift dependence
- Cosmological model (e.g. LCDM, Λ CDM, massive neutrino)

Comparison different HMF calibration



HALO MASS FUNCTION: BARYONIC EFFECTS

Baryonic feedbacks (radiative cooling, star formation, AGN feedback) redistribute and expel mass from galaxy clusters

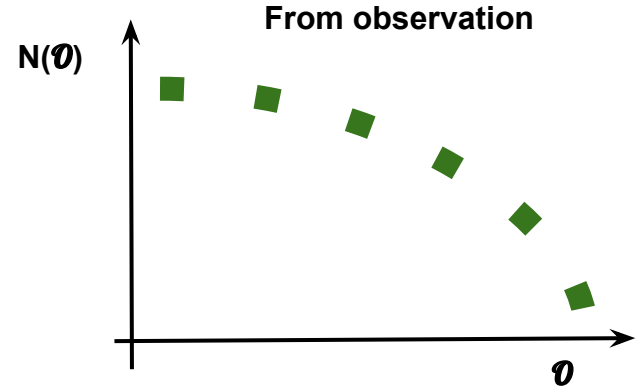
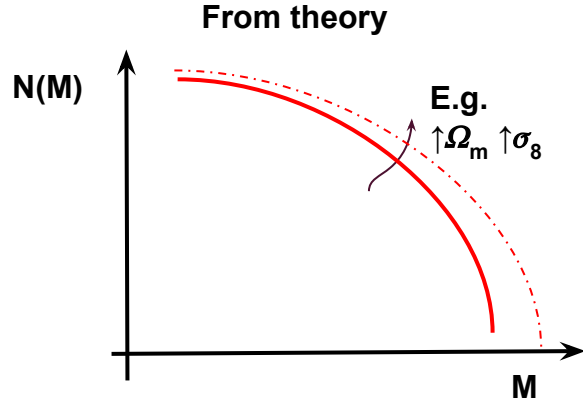


Baryonic feedbacks most effective in the inner the regions of the halo and in low mass systems

See also Castro+21

FROM THEORY TO OBSERVATION

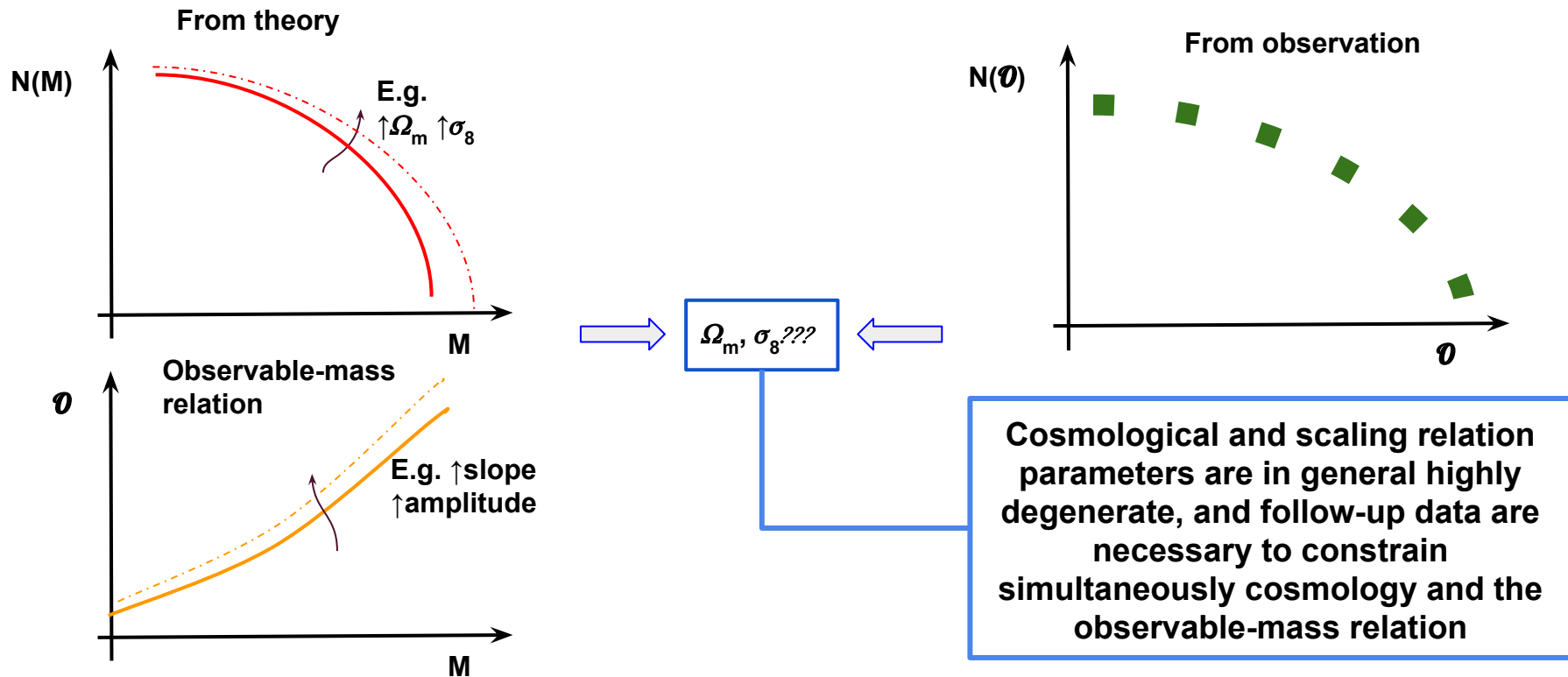
- Masses are not directly observable. Galaxy clusters are selected according to some observable, in general related to the observational technique, which correlate with the mass.



ϕ : Observable used to detect/select clusters (e.g. number of galaxies, X-ray luminosity, SZ signal)

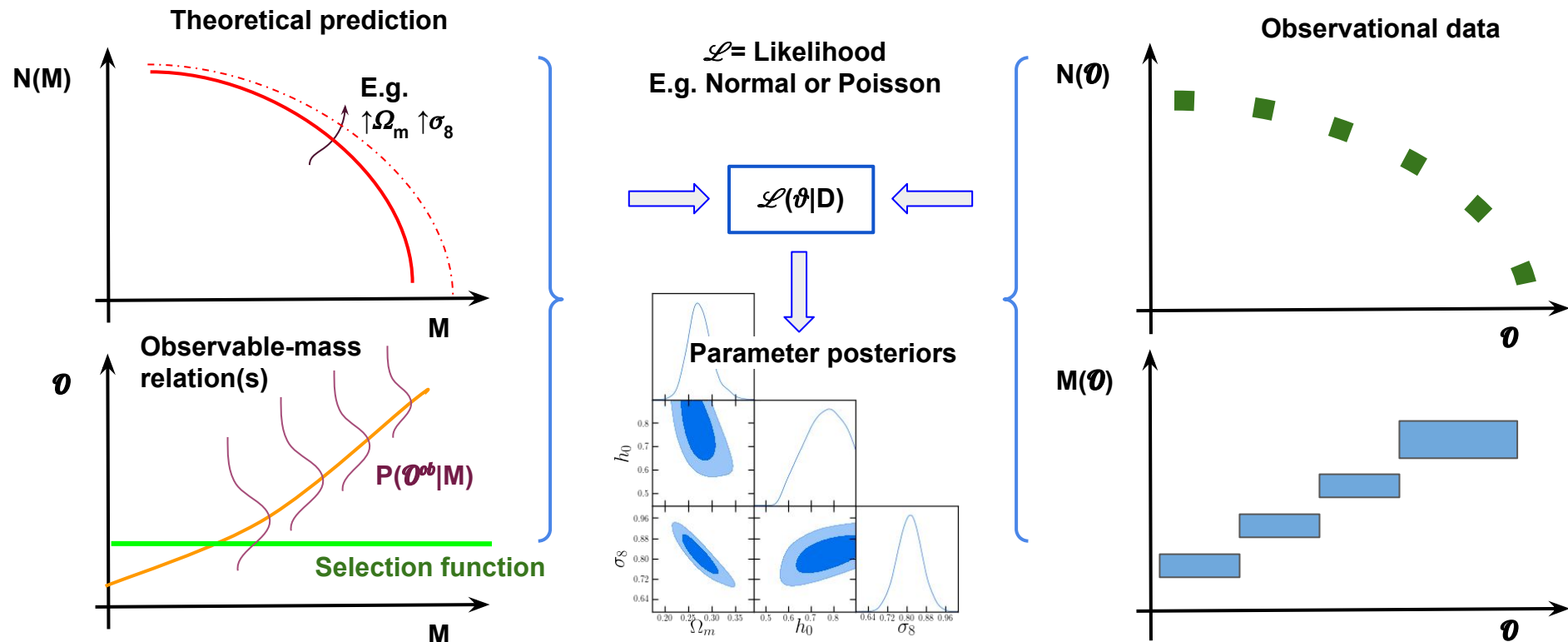
FROM THEORY TO OBSERVATION

- Individual mass measurements are expensive and not feasible for cluster survey. We need to rely on mass proxies which are tightly correlated with the halo mass.



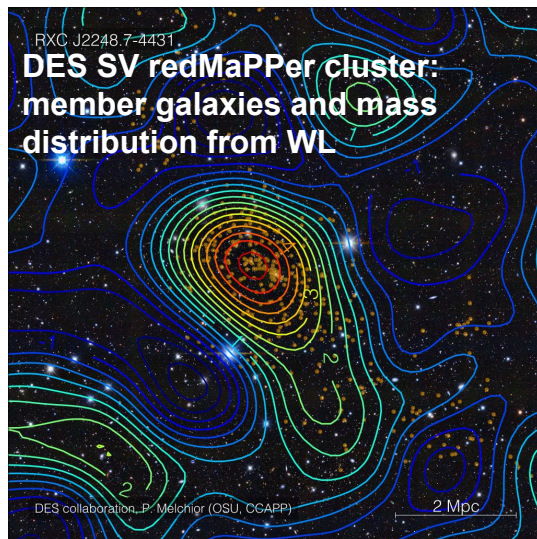
FROM THEORY TO OBSERVATION: CONSTRAINTS

- Combine cluster abundance and cluster mass estimates data to **simultaneously** constrain cosmology and the observable-mass relation(s)

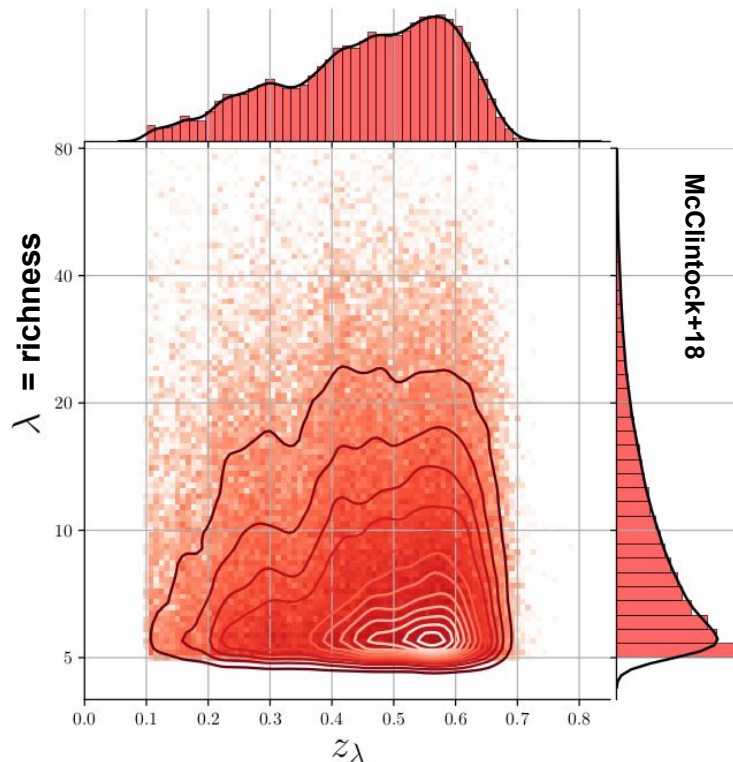


CLUSTER DETECTION: PHOTOMETRIC SURVEY

- **Detection:**
 - Overdensity of (red-sequence) galaxies
 - Lensing effect
- **Observable/Mass proxy:**
 - Richness (# member galaxies)
 - Luminosity
 - Lensing signal
 - Velocity dispersion (with spectra)



**z - λ distribution of redMaPPer clusters
in DES Y1**



CLUSTER DETECTION: X-RAY SURVEY

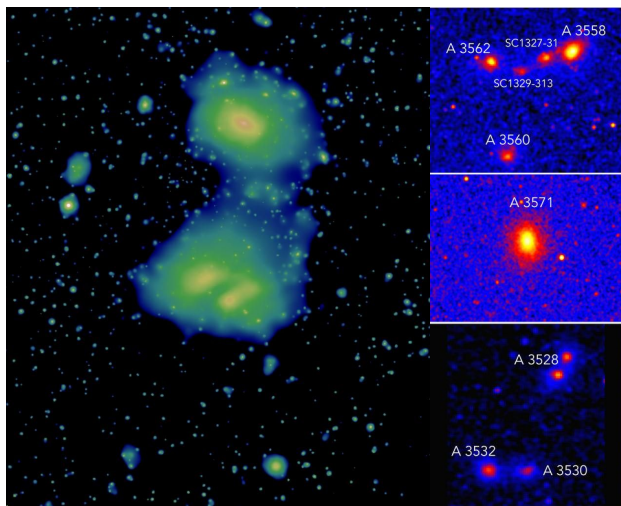
- **Detection:**

- Extended x-ray sources

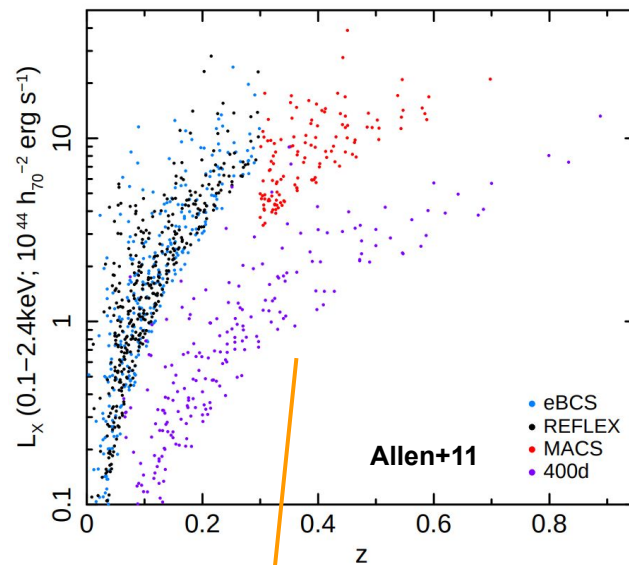
- **Observable/Mass proxy:**

- L_X
- T_X
- Flux
- $Y_x = M_{\text{gas}} T_X$ (gas thermal energy)

X-ray images of clusters from eROSITA



L_X, z distribution of X-ray selected catalogs



- **X-ray emissivity from *bremsstrahlung* radiation of the ICM:**

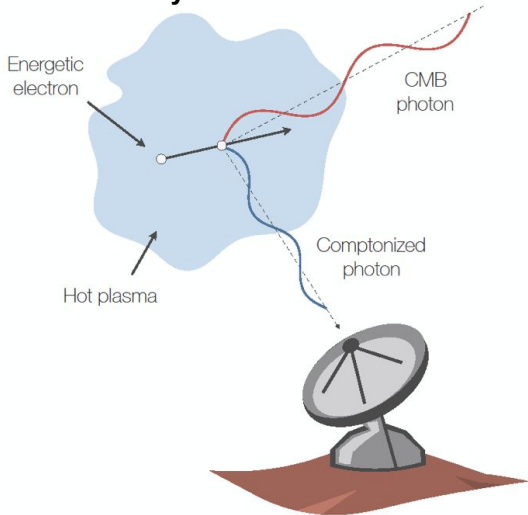
$$\epsilon_\nu \equiv \frac{dL}{dVd\nu} \propto n_e^2 g(\nu, T) T^{-1/2} \exp(-h\nu/k_B T)$$

Not very sensitive to projections

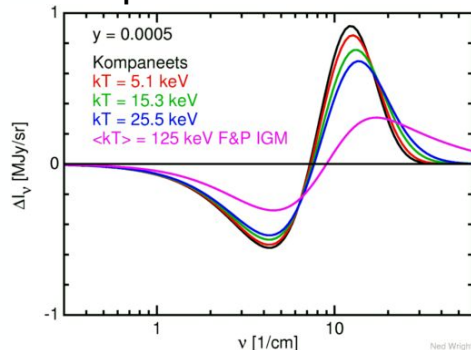
CLUSTER DETECTION: SZ SURVEY

- **Detection:**
Thermal Sunyaev-Zel'dovich effect
(mm-wavelength)
- **Observable/Mass proxy:**
SZ signal

Thermal Sunyaev-Zel'dovich effect



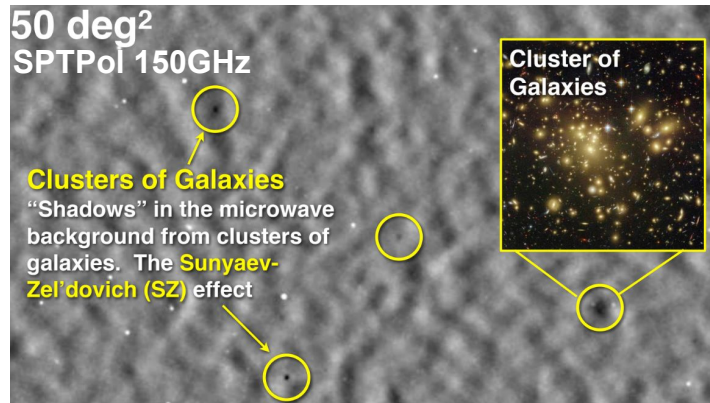
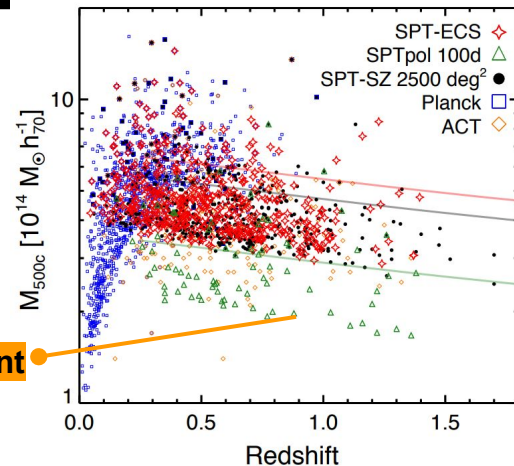
SZ spectral distortion



Compton-y parameter

$$y = \int dl \frac{k_B T_e}{m_e c^2} n_e \sigma_T$$

Mass and redshift distribution of SZ-selected cluster catalogs (Bleem+19)



Credits B. Benson, SPT Collaboration

CLUSTER CATALOGS

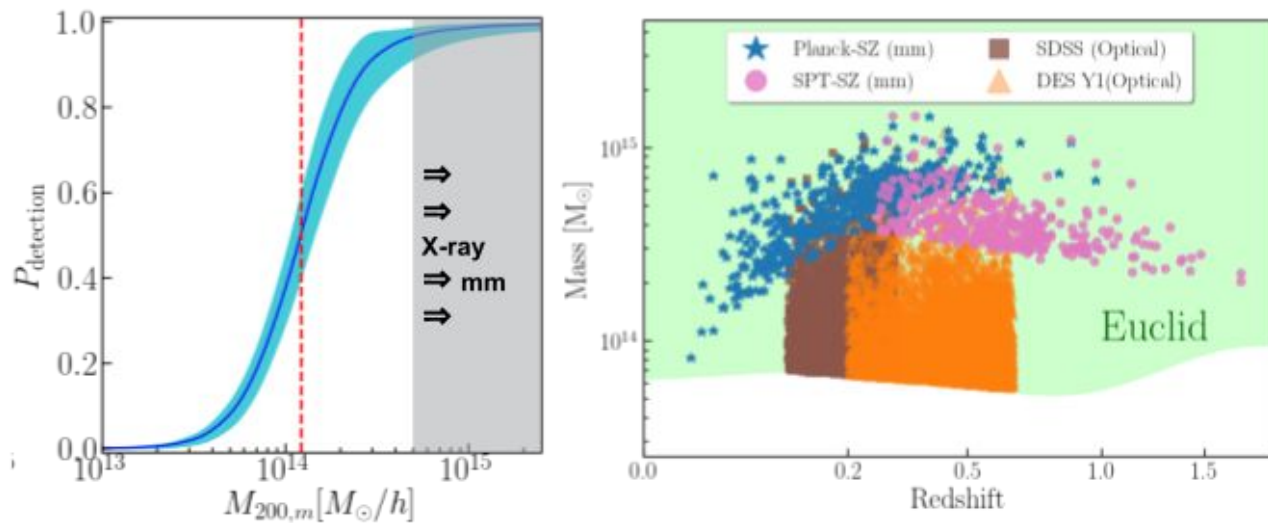


Figure 1: *Left:* redMaPPer DES Y1 cluster catalog detection probability as a function of mass: systems down to $\sim 5 \cdot 10^{13} M_{\odot}$ have a non-negligible chance to be included in the optical catalog ($\lambda > 20$), while clusters selected at different wavelength (X-ray, mm) have masses typically above $5 \cdot 10^{14} M_{\odot}$ (*gray area*; adapted from [Ab20]). *Right:* Mass and redshift ranges probed by current optical (SDSS, DES Y1) and millimeter (Planck-SZ, SPT-SZ 2500) cluster surveys. The green shaded area marks the mass and redshift range to be covered by the Euclid cluster sample.

Photometric catalogs capable of detecting system down to group mass scale but have a much less cleaner selection function which hamper they cosmological exploitation

MASS MEASUREMENTS FROM X-ray DATA

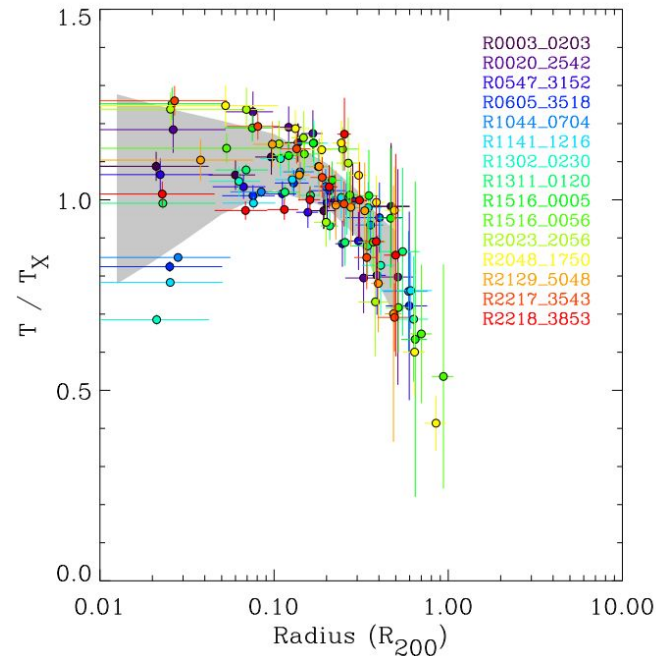
- From hydrostatic equilibrium:

$$M(< r) = -\frac{r k_B T(r)}{G \mu m_p} \left(\frac{d \ln \rho_{\text{gas}}(r)}{d \ln r} + \frac{d \ln T(r)}{d \ln r} \right)$$

Assumptions:

- Hydrostatic equilibrium (Negligible non-thermal pressure support)
- Spherical symmetry

Temperature profiles from *XMM-Newton* observations (Pratt+06)



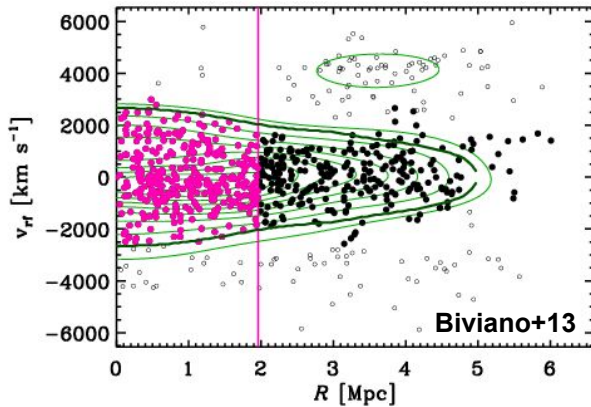
MASS MEASUREMENTS FROM SPECTROSCOPIC DATA

- Dynamical mass estimates (Jeans equation):

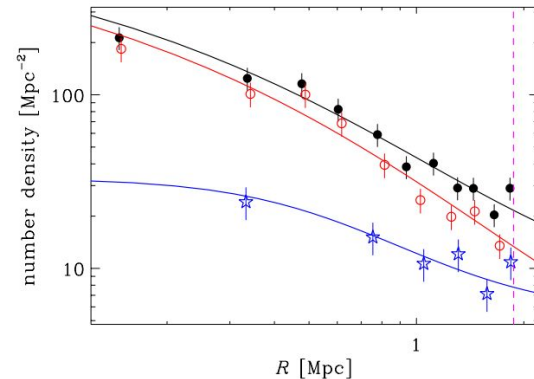
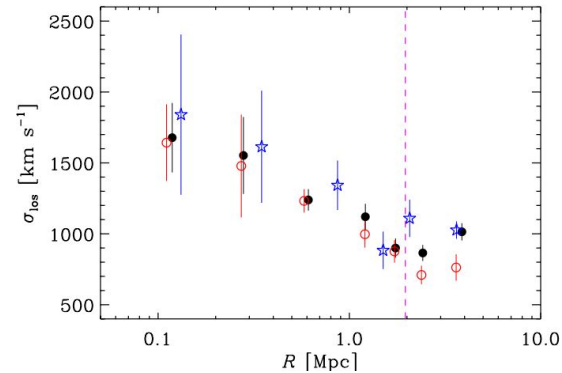
$$M(< r) = -\frac{r\sigma_r^2}{G} \left(\frac{d \ln \sigma_r^2}{d \ln r} + \frac{d \ln n_{\text{glx}}}{d \ln r} + 2\beta \right)$$

Assumptions:
Spherical symmetry
Dynamical equilibrium

- Caustic method (projected phase-space distribution):



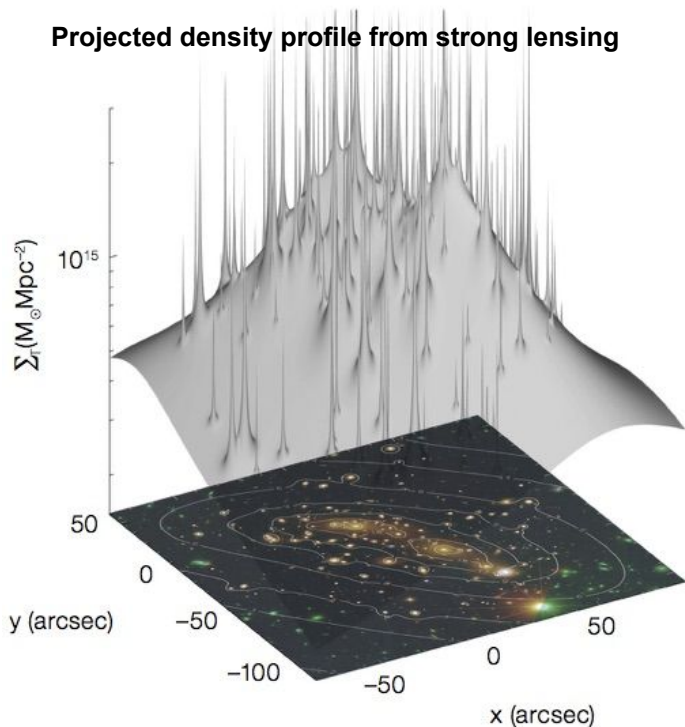
L.o.s. velocity dispersion and member galaxy density profiles from VLT/VIMOS (Biviano+13)



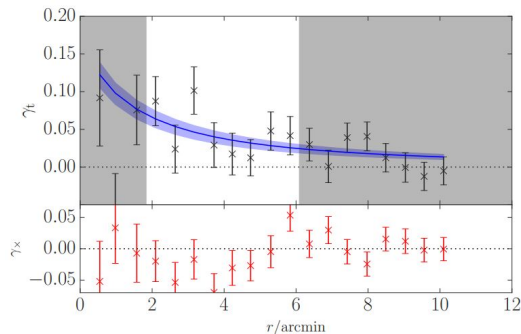
MASS MEASUREMENTS FROM IMAGING

- **Strong and Weak Lensing mass measurements:**

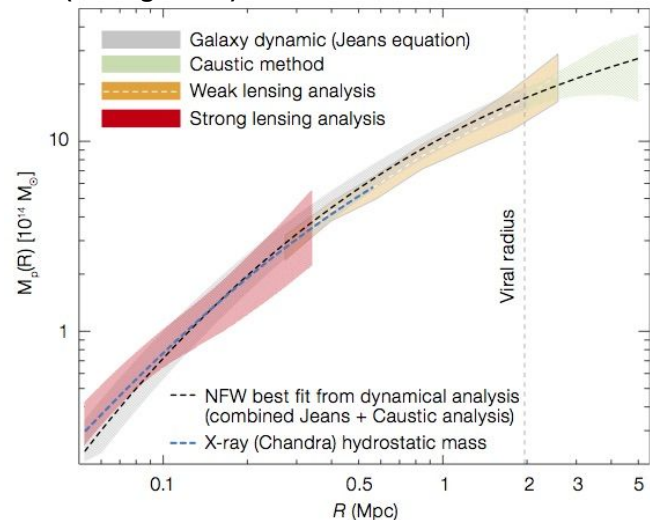
Projected density profile from strong lensing



Tangential shear profile from WL (Dietrich+18)



Cluster mass profile from different techniques (Battaglia+16)



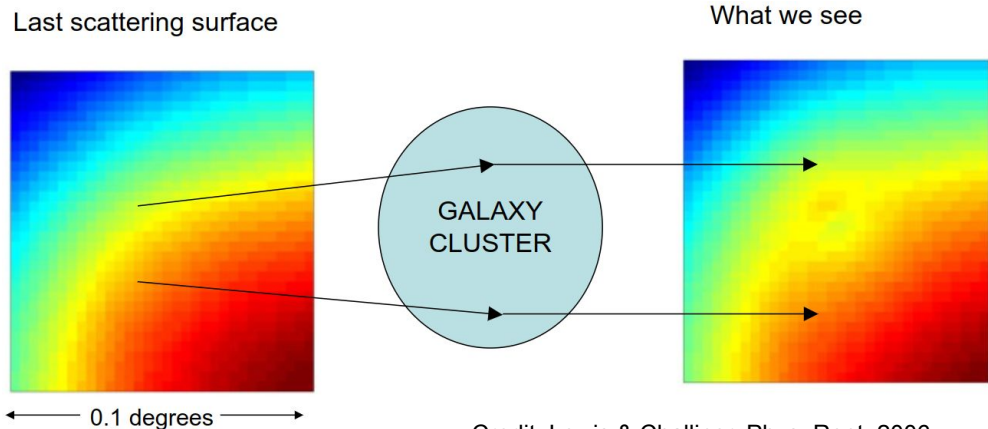
$$\gamma_t(n_s(z), \Sigma(R)) \Rightarrow \Sigma(R) = \int_{-\infty}^{\infty} d\chi \Delta\rho\left(\sqrt{R^2 + \chi^2}\right)$$

Assumption:

- **Parametric form for the halo density profile (e.g. NFW, Einasto profiles; Navarro+97, Einasto 1965) and correlated structures (2-halo term)**

CMB CLUSTER LENSING

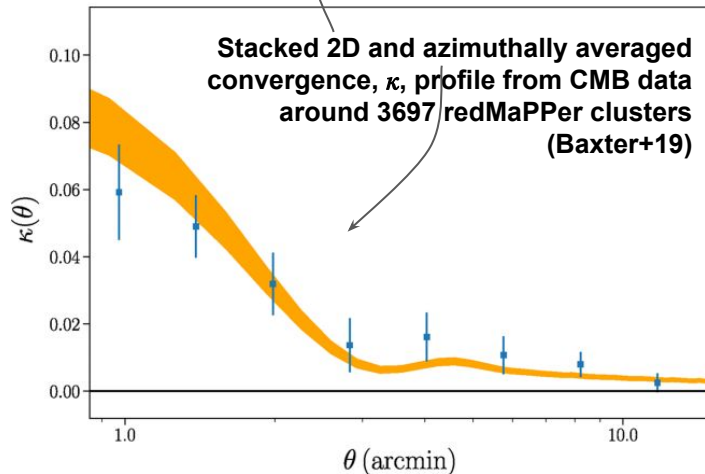
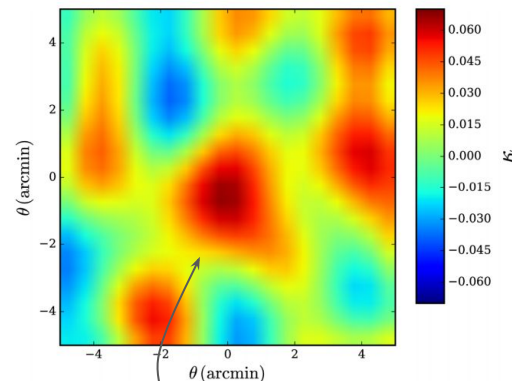
- **Lensing by GC induces a dipole-like distortion in the CMB:**



Credit: Lewis & Challinor, Phys. Rept. 2006

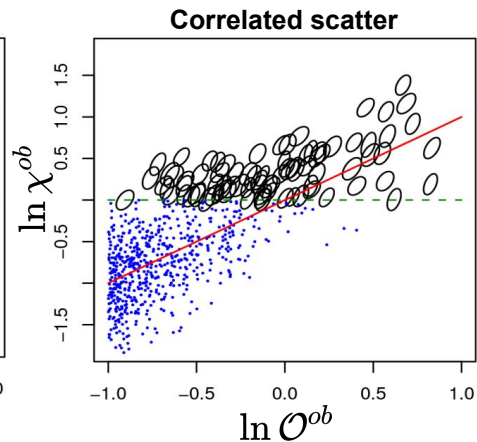
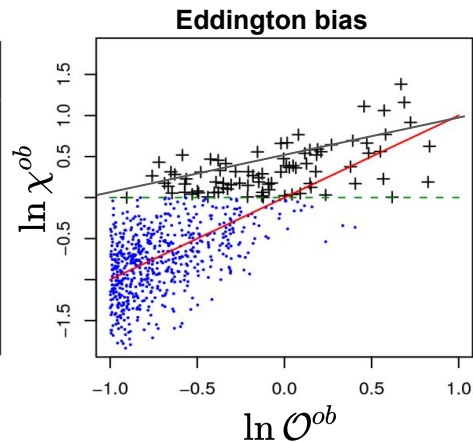
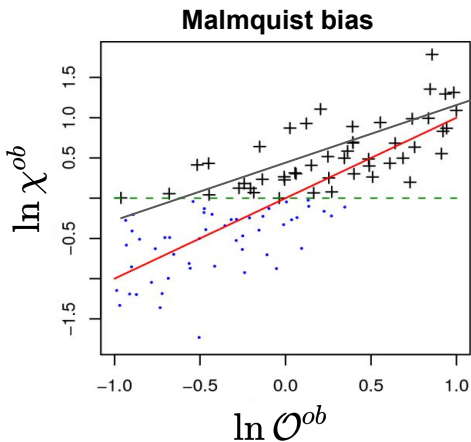
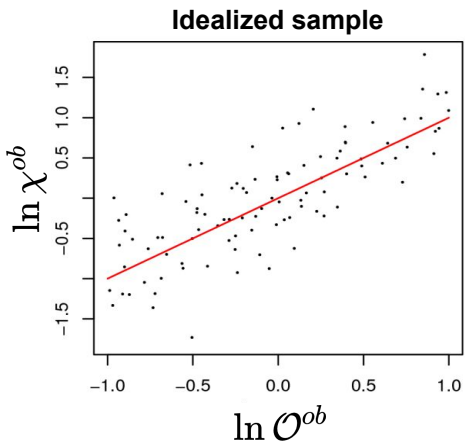
The distortion is quite small ($\sim 10\mu\text{K}$ for $10^{15}M_{\odot}$ halo) but can be used to calibrate the mass of high redshift clusters.

The lensing signal can also be detected in polarization data (see e.g. Raghunathan+19).



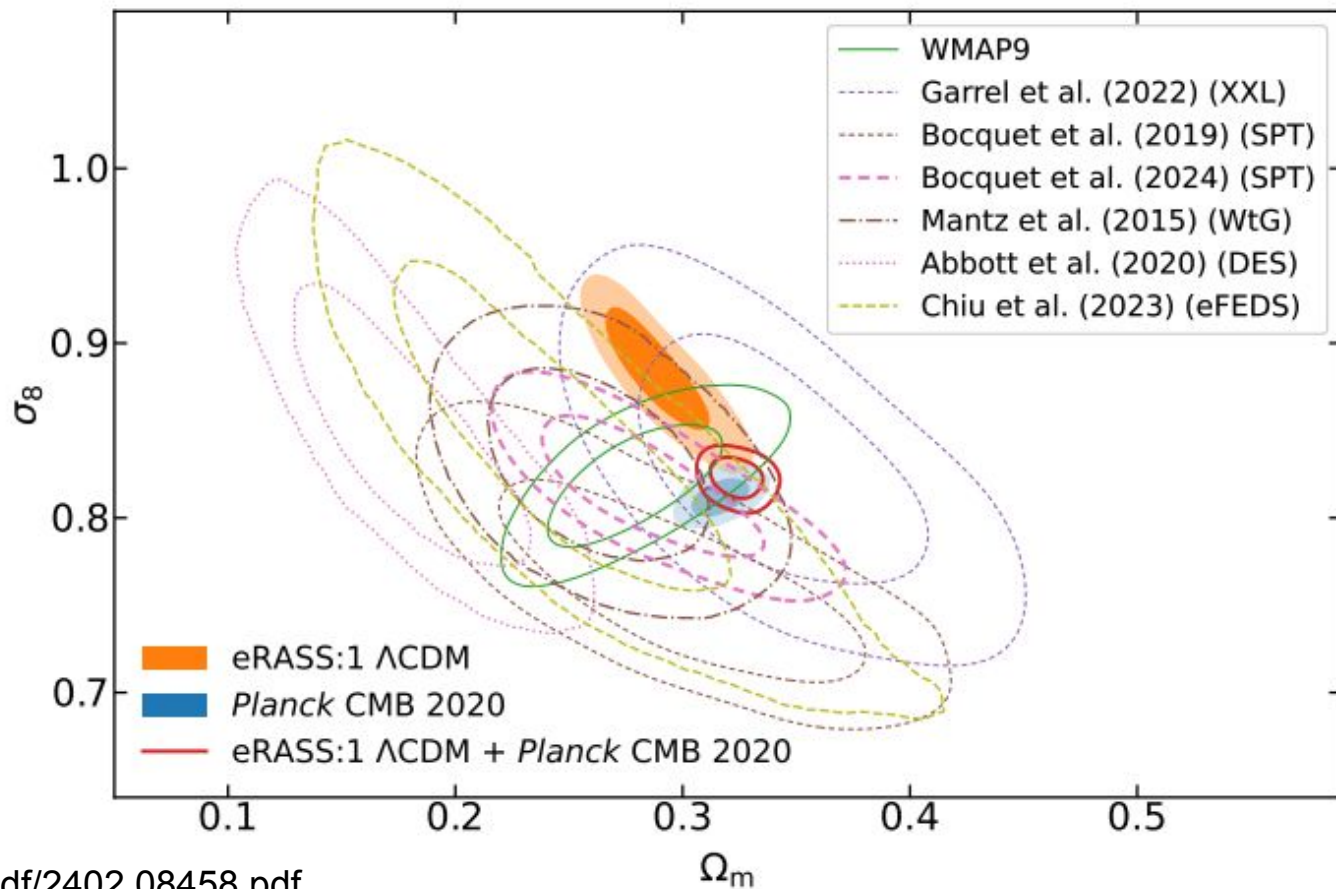
FROM THEORY TO OBSERVATION: SCALING RELATIONS

- Scaling relation(s) calibration:

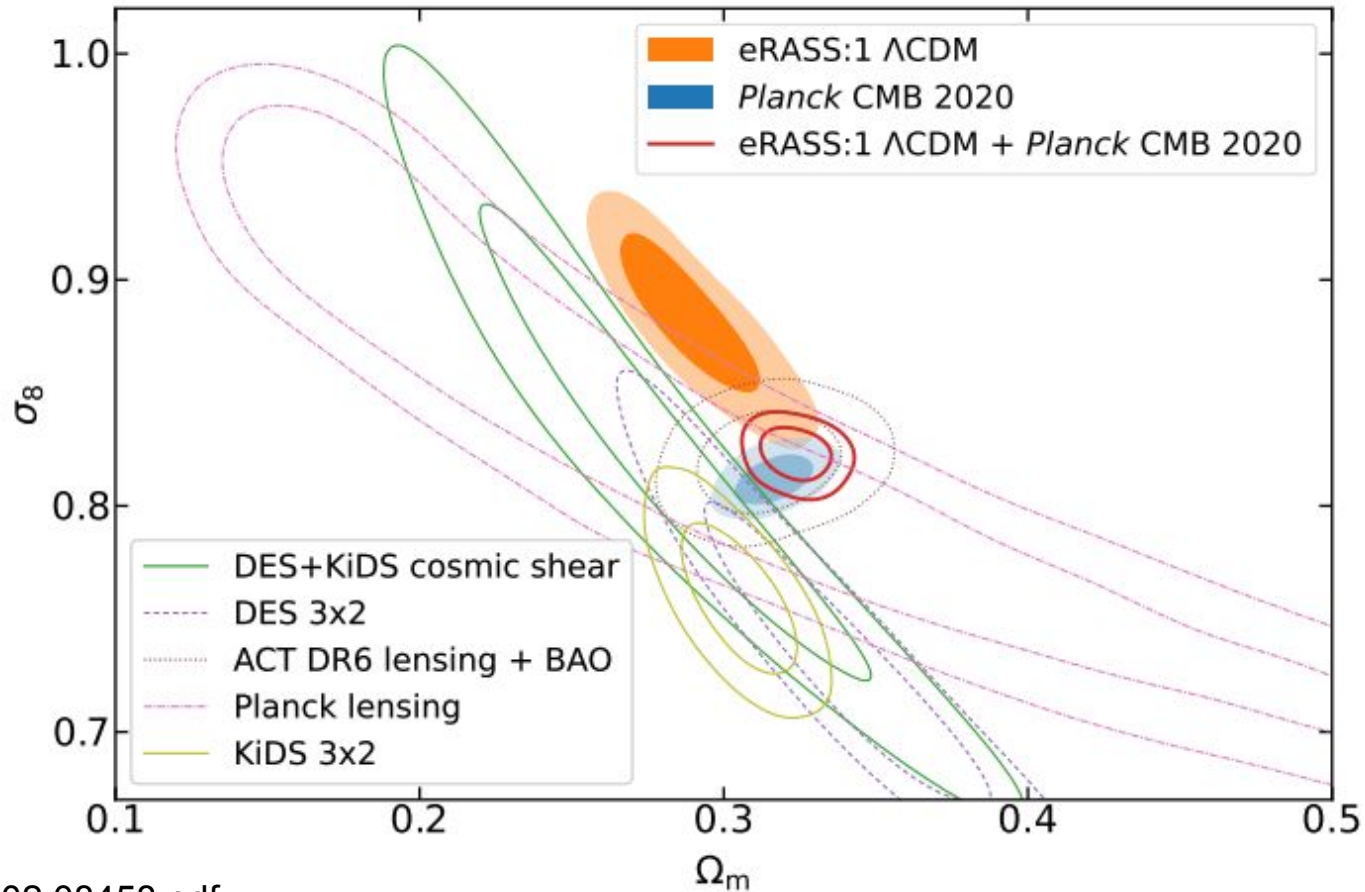


Credit A. Mantz

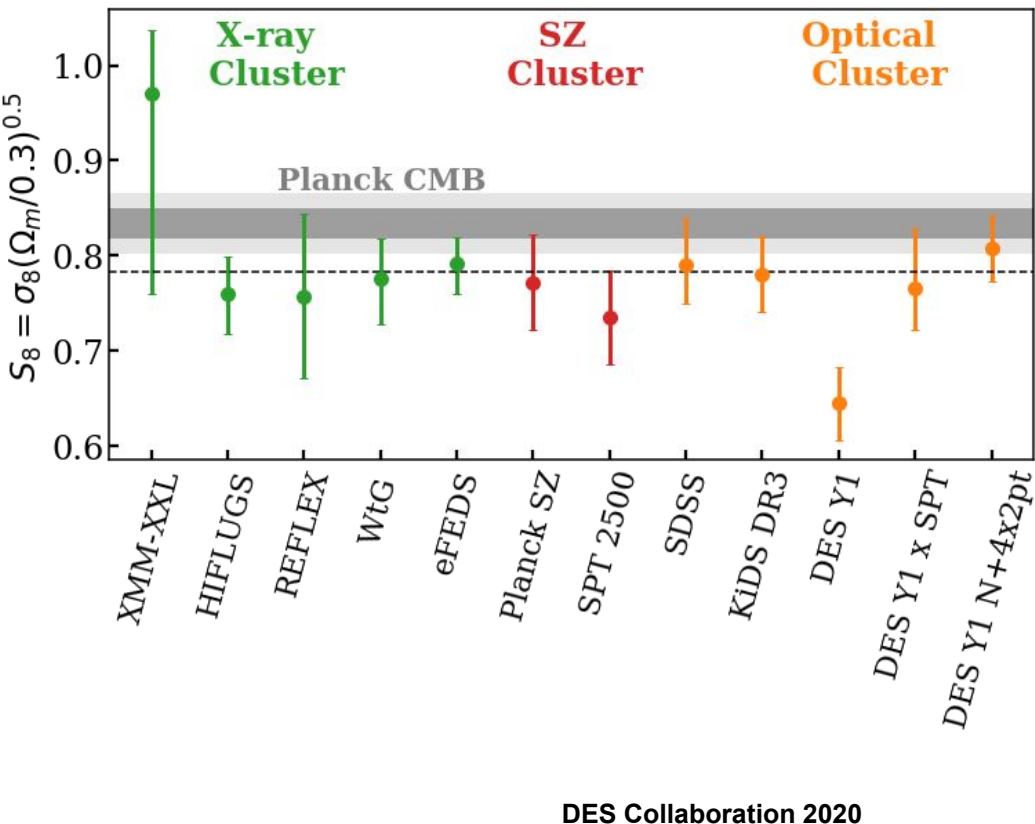
RECENT CONSTRAINTS FROM CLUSTER NC



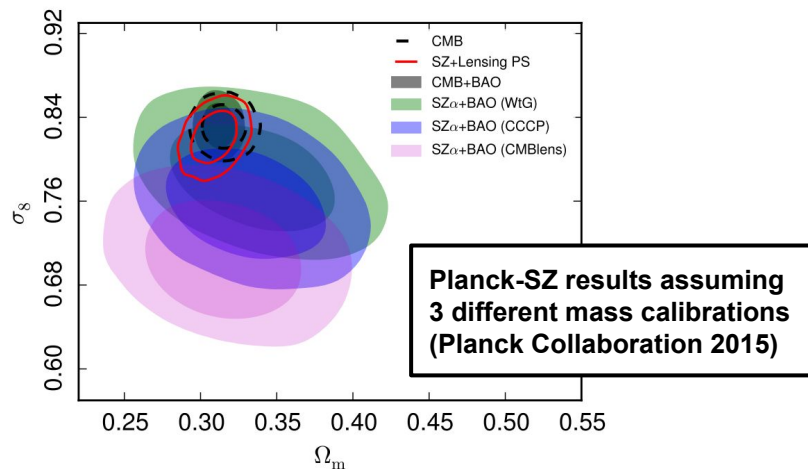
RECENT CONSTRAINTS FROM CLUSTER NC



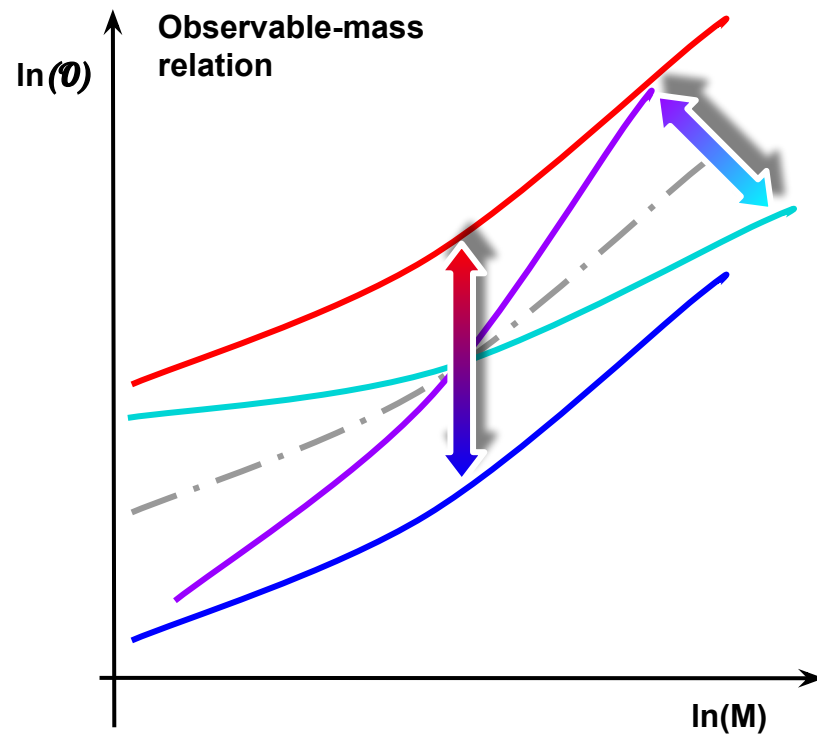
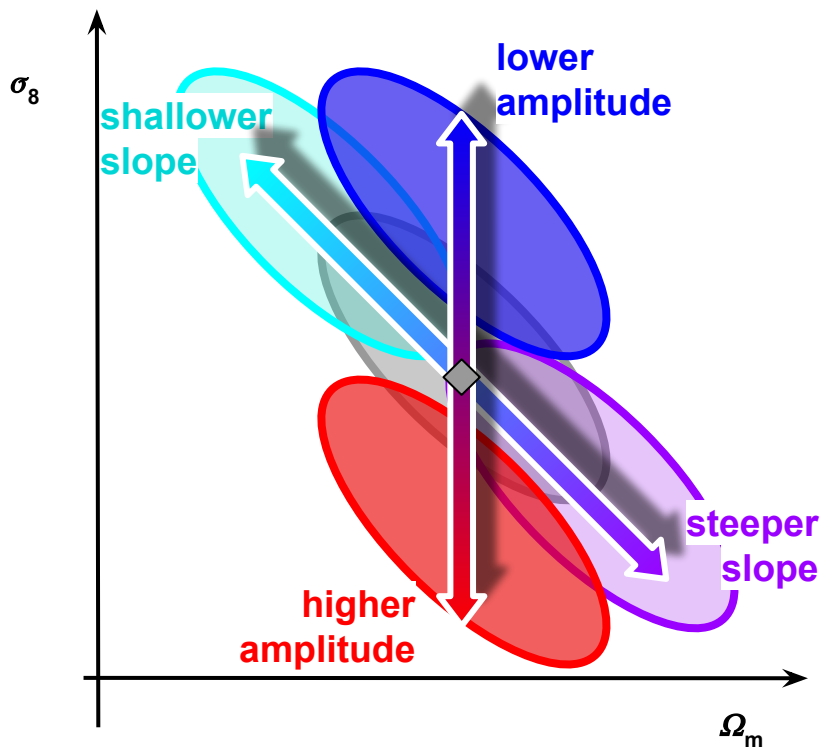
LIMITATIONS FOR CLUSTER COSMOLOGY STUDIES



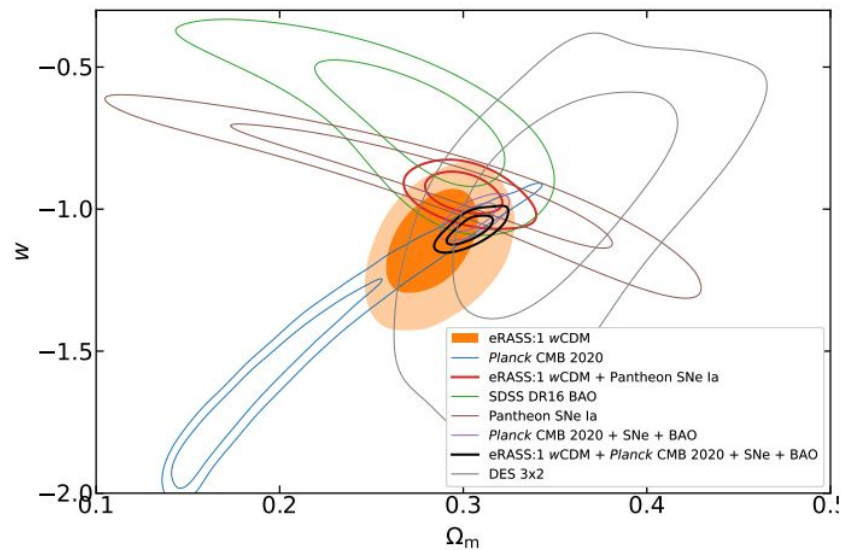
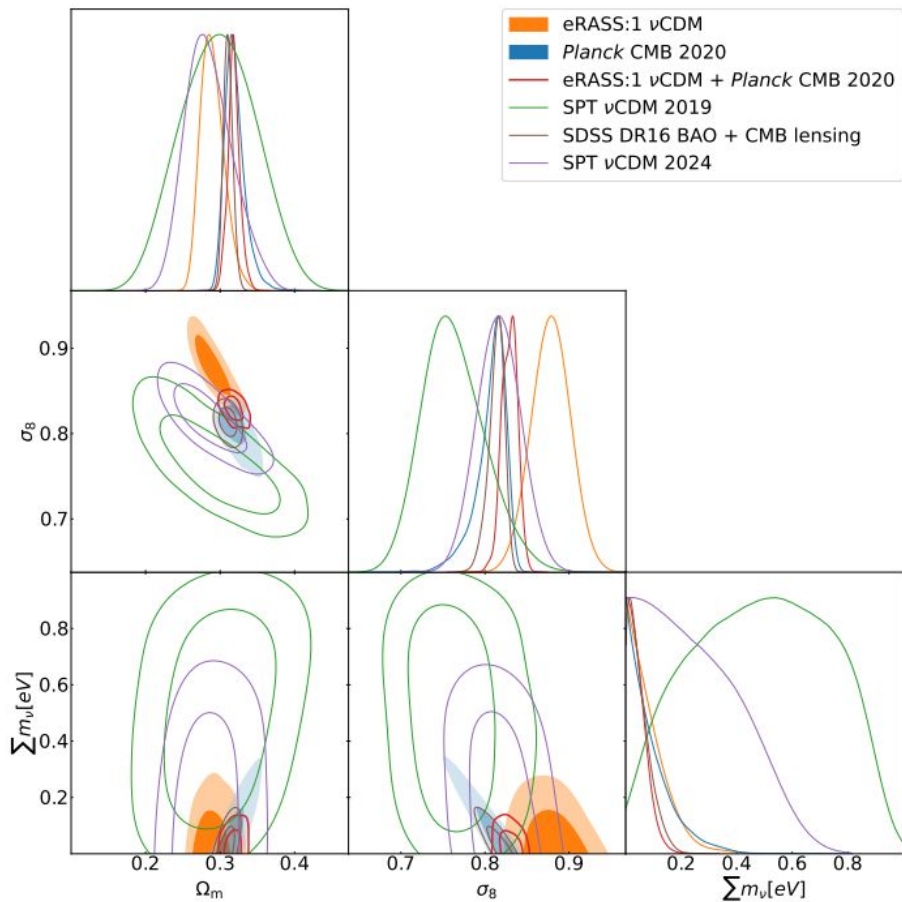
- Cosmological constraints independent and competitive with other cosmological probes
- Slight to moderate tension between different cluster studies
- Currently limited by the mass (i.e. scaling relation) calibration



LIMITATIONS FOR CLUSTER COSMOLOGY STUDIES



RECENT CONSTRAINTS FROM CLUSTER NC

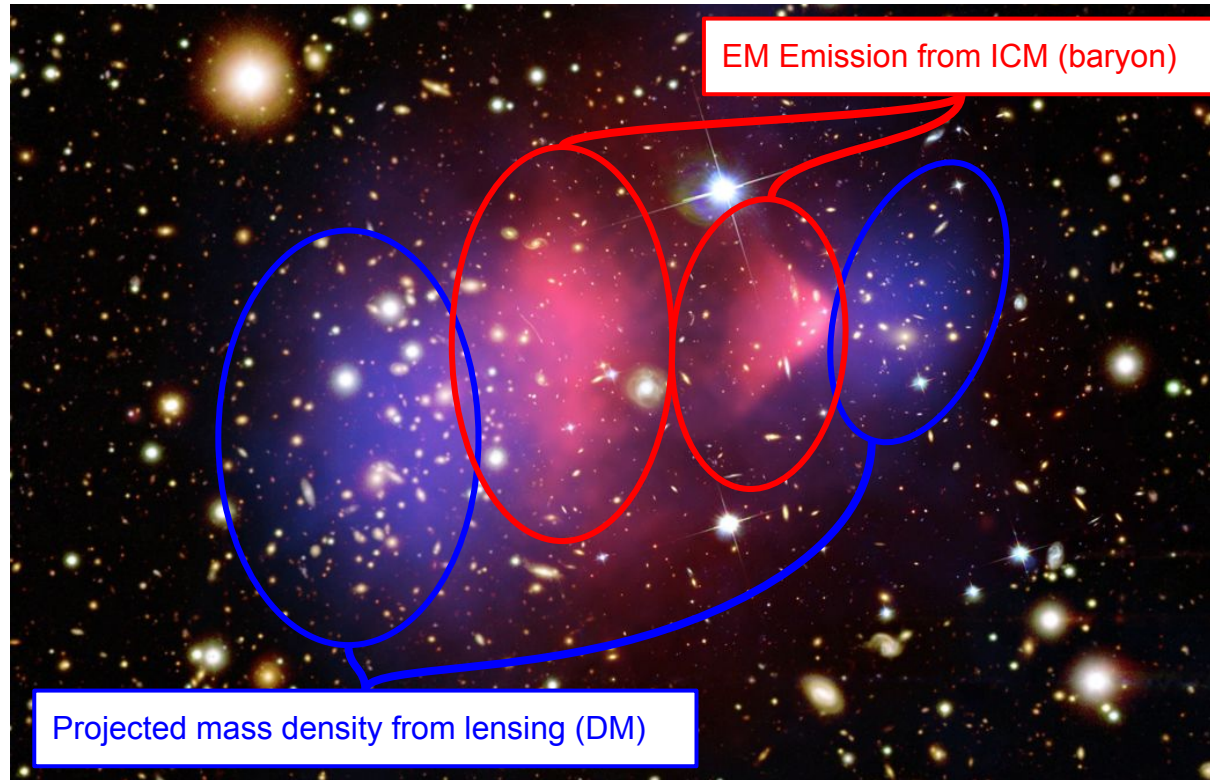


OTHER COSMOLOGICAL TESTS WITH GALAXY CLUSTERS

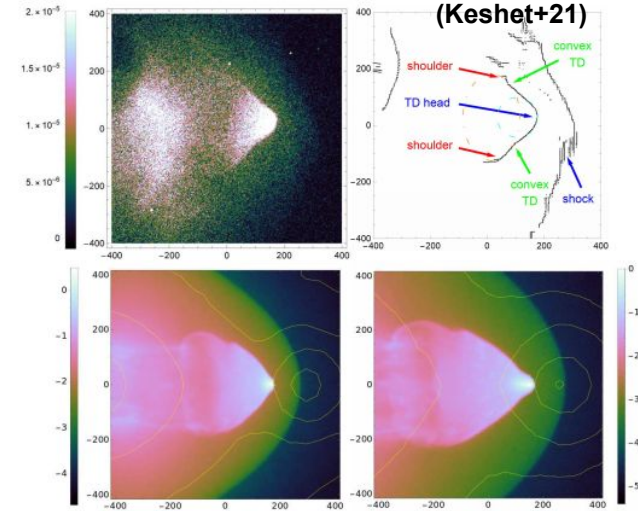
For a review: [Allen+2011](#) or [Kravtsov+2012](#)

OTHER COSMOLOGICAL TESTS WITH GALAXY CLUSTERS

- The Bullet Cluster (DM nature)



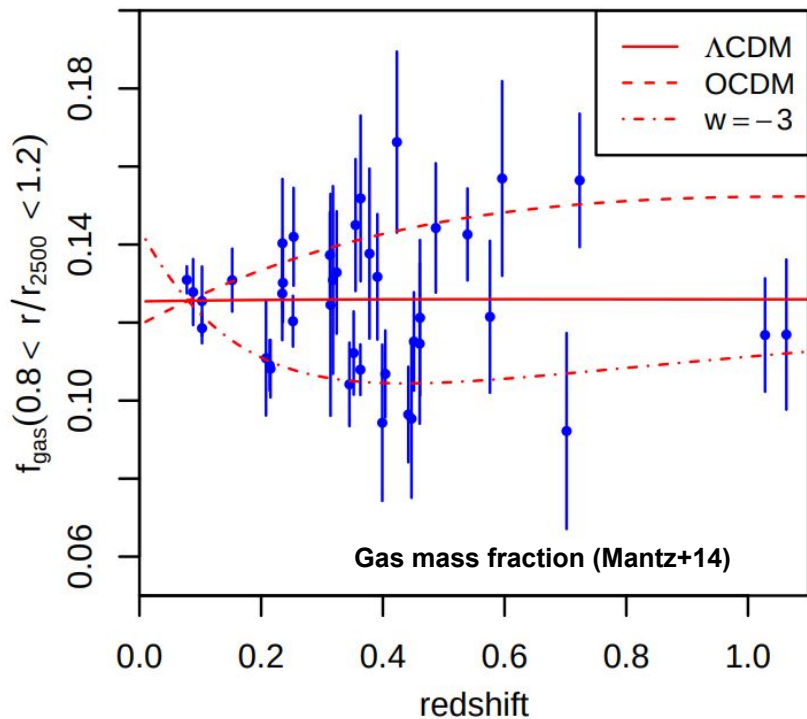
Real vs simulated Bullet-like shock
(Keshet+21)



The offset between the EM and WL signal peaks, along with the shape of the shock wave, provide compelling evidence for the presence of dark matter; moreover it allows to place constraints on the dark matter cross-section

OTHER COSMOLOGICAL TESTS WITH GALAXY CLUSTERS

- Gas mass fraction ($\Omega_m, \Omega_\Lambda, w$):



$$f_{\text{gas}} = \frac{M_{\text{gas}}}{M_{\text{tot}}} \propto \left(\frac{\Omega_b}{\Omega_m} \right) \left[\frac{d^{\text{ref}}(z)}{d(z)} \right]^{3/2}$$

With priors on

- $\Omega_b h^2$ (important)
- h (less important),

the low- z data constrain Ω_m :

$$f_{\text{gas}}(z \lesssim 0.15) \propto \frac{\Omega_b}{\Omega_m} h^{3/2}$$

Apparent evolution constrains dark energy:

$$f_{\text{gas}}(z) \propto d(z)^{-3/2}$$

Credit A. Mantz

OTHER COSMOLOGICAL TESTS WITH GALAXY CLUSTERS

- H_0 from X-ray and SZ distance measurements:

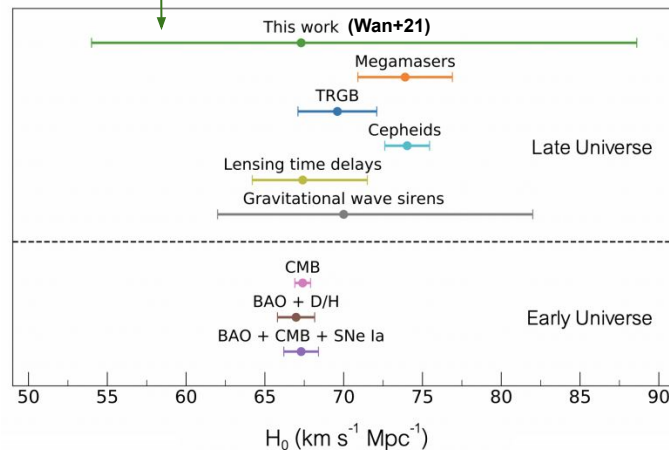
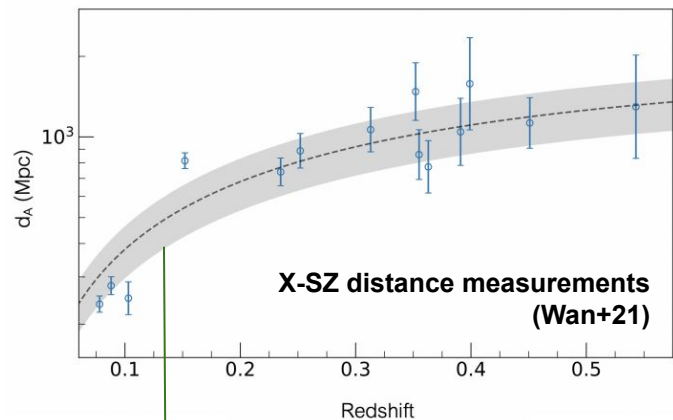
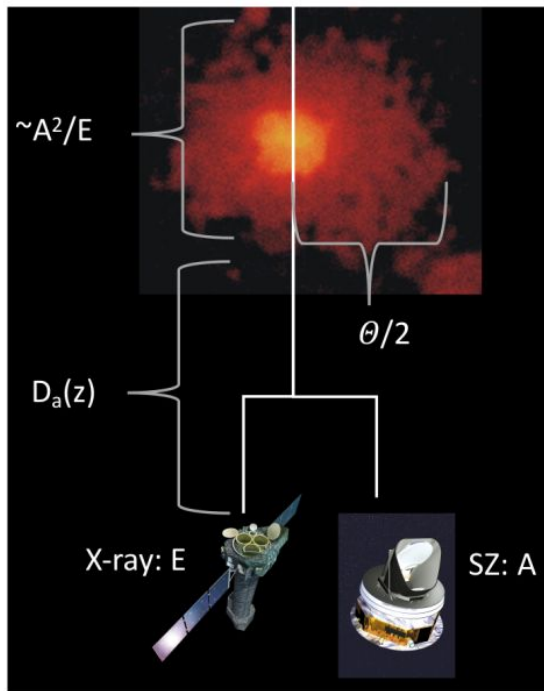
Based on a distance measuring techniques that depend on a comparison of 2 observables (Cavaliere+77):

$$E \propto \int n_e^2 dl$$

$$A \propto \int n_e dl$$

If the structure of the gas is known, given the angular size ϑ of the system, the angular diameter distance is given by:

$$D_A(z) = A^2 / (E\theta)$$



HALO PROFILE

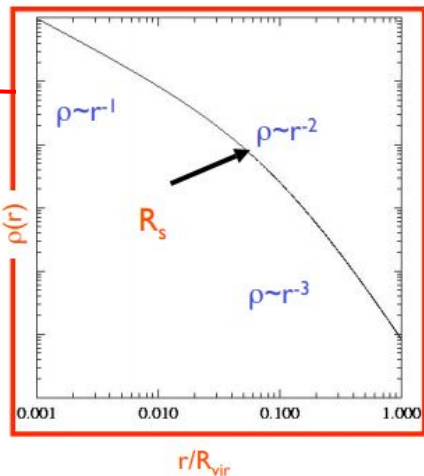
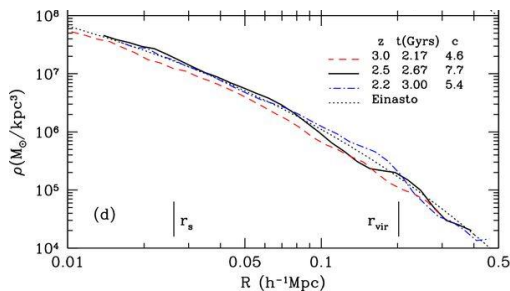
From n-body/hydro simulations we can predict the dark matter/gas halo profiles. For LCDM models E.g. Navarro+97 and Einasto 1965:

$$\rho_{\text{NFW}} = \frac{\rho_s}{(r/r_s)(1+r/r_s)^2}$$

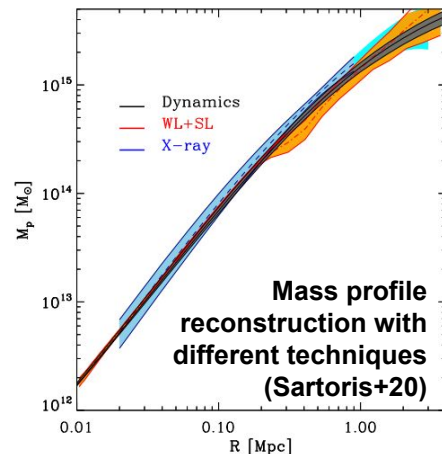
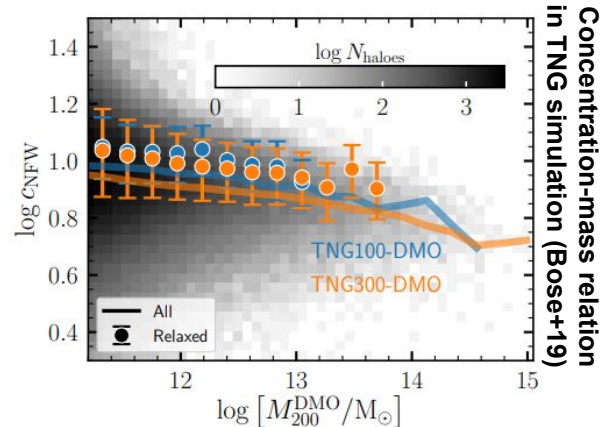
$$\rho_{\text{Ein}} = \rho_{-2} \exp \left\{ -\frac{2}{\alpha} \left[\left(\frac{r}{r_{-2}} \right)^\alpha - 1 \right] \right\}$$

Observationally, cluster profiles can be inferred from strong and weak lensing, galaxy dynamics, and ICM (X-ray, SZ) measurements

DM halo profiles for halos of different mass and redshift



$$c_{\text{NFW}} = r_{200}/r_s$$

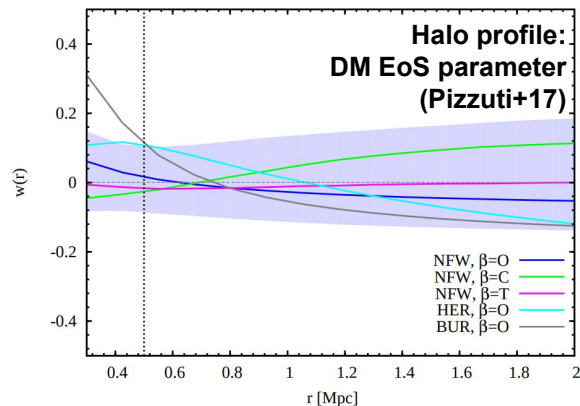
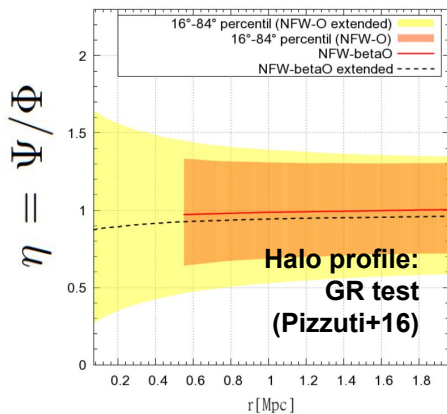


Concentration-mass relation in TNG simulation (Bose+19)

OTHER COSMOLOGICAL TESTS WITH GALAXY CLUSTERS

- Galaxy cluster mass profile:

The shape/slope of the halo profile, especially in the inner regions, can be used to test several fundamental physics model, such as the nature of dark matter (e.g. warm vs cold, interacting DM) or GR test.



Halo profile: Interacting DM vs DM (Vega-Ferrero+20)

