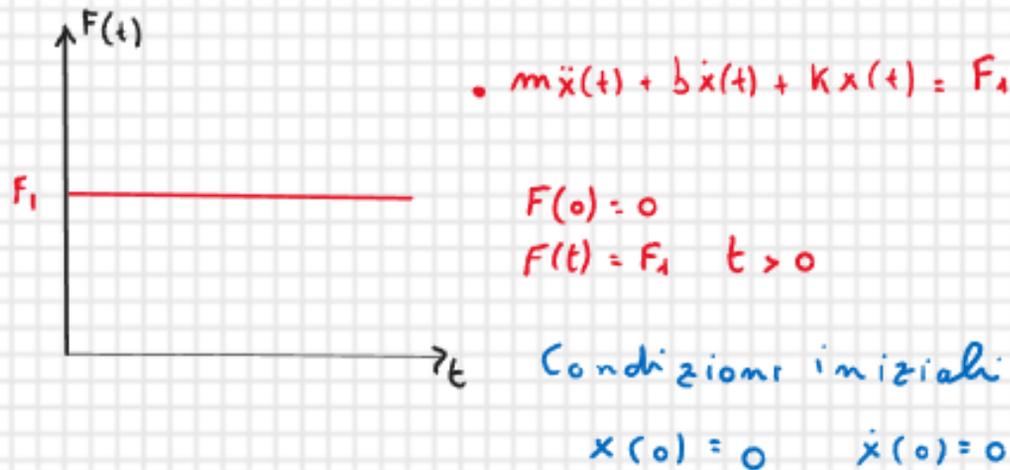


Titolo:

$$m\ddot{x} + b\dot{x} + kx = F(t) \quad \text{Funzione qualsiasi}$$
$$= -m\ddot{y}(t)$$

Integrale di Duhamel



Soluzione eq. moto

$$x(t) = e^{-\nu\omega t} \left[ A \sin \omega t + B \cos \omega t \right] + \frac{F_1}{k}$$

Condizioni al contorno A e B

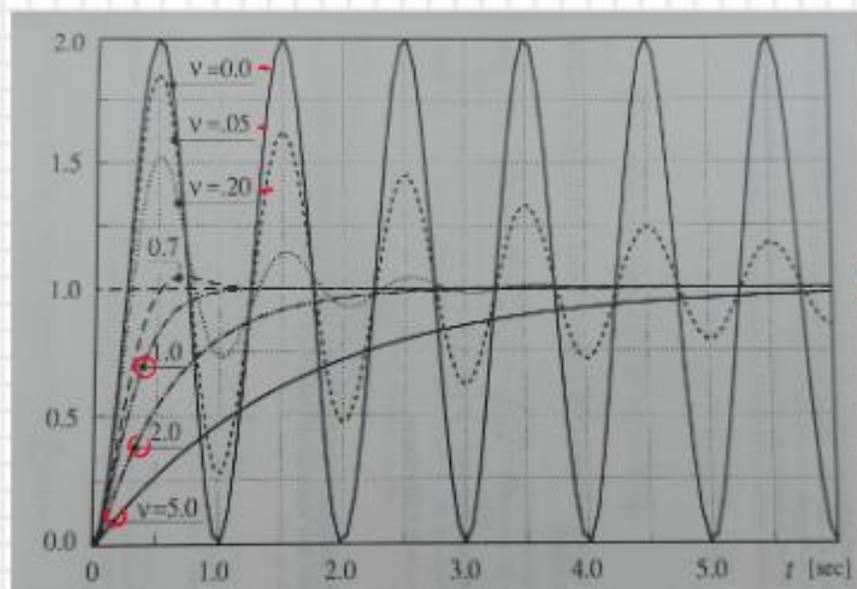
- $x(0) = 0 \Rightarrow B = -\frac{F_1}{k}$
- $\dot{x}(0) = 0 \Rightarrow A = -\frac{F_1}{k} \frac{\nu\omega}{\omega}$

$$\dot{x}(t) = -\nu\omega e^{-\nu\omega t} (A \sin \omega t + B \cos \omega t) + e^{-\nu\omega t} (\omega A \cos \omega t - \omega B \sin \omega t)$$

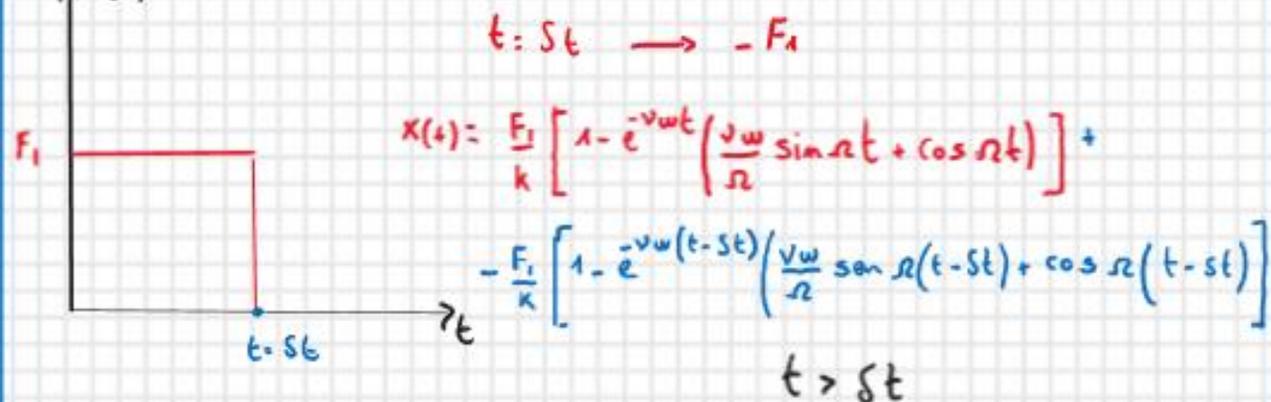
$$x(t) = \frac{F_1}{k} \left[ 1 - e^{-\nu\omega t} \left( \frac{\nu\omega}{\omega} \sin \omega t + \cos \omega t \right) \right] \quad t \rightarrow \infty$$

$x(t) \rightarrow \frac{F_1}{k}$

## Titolo:



$F(t)$



Risposta all'impulso  $F_1 St$

$$g(t) = \frac{F_1}{k} \left[ 1 - e^{-v\omega t} \left( \frac{v\omega}{\Omega} \sin \Omega t + \cos \Omega t \right) \right]$$

$$x(t) = g(t) - g(t - St) \quad St \rightarrow \text{infinitesimo}$$

$$\bullet \quad x(t) = \frac{dg(t)}{dt} St$$

$$\frac{dg(t)}{dt} = -\frac{F_1}{k} \left[ -v\omega e^{-v\omega t} \left( \frac{v\omega}{\Omega} \sin \Omega t + \cos \Omega t \right) + e^{-v\omega t} (v\omega \cos \Omega t - \Omega \sin \Omega t) \right]$$

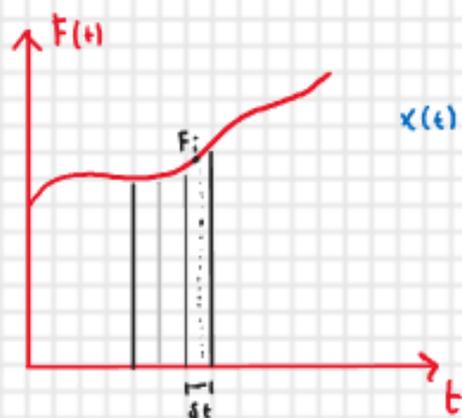
Titolo:

$$X(t) = \frac{F_1}{k} \left( \omega + \frac{v' \omega^2}{\omega} \right) e^{-v \omega t} \sin \omega t \quad \delta t = \frac{F_1}{m \omega} e^{-v \omega t} \sin \omega t \quad \delta t$$

Spostamento dovuto all'impulso applicato a  $t=0$

Impulso agisce a  $t=t_1$

$$x(t) = \frac{F_1}{m \omega} e^{-v \omega (t-t_1)} \sin \omega (t-t_1) \quad t > t_1$$



$$x(t) = \int_0^t \frac{F(t_1)}{m \omega} e^{-v \omega (t-t_1)} \sin \omega (t-t_1) dt_1$$

Integrale di Duhamel

Risposta del nostro sistema ad una forzante  $F(t)$

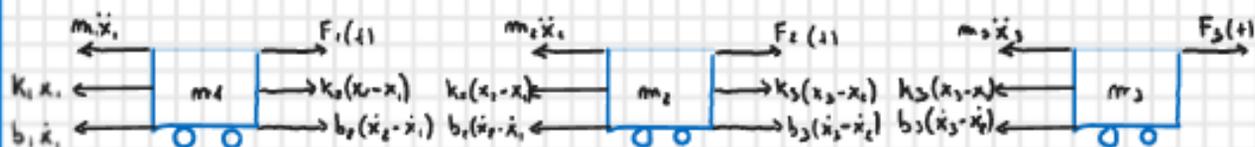
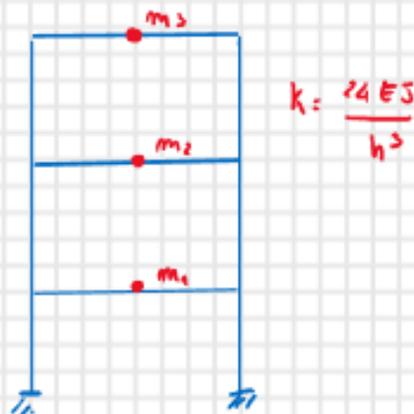
• Esempio semplice  $F(t) = p_0$   $v=0$   $\omega = \omega$

$$x(t) = \frac{1}{m \omega} \int_0^t p_0 \sin \omega (t-t_1) dt_1 = \frac{p_0}{m \omega^2} \cos \omega (t-t_1) \Big|_0^t = \frac{p_0}{m \omega} (1 - \cos \omega t)$$

Integriamo rispetto a  $t_1$

Titolo:

## Sistemi a $n$ gradi di libertà



$$-m_1 \ddot{x}_1 - k_1 x_1 - b_1 \dot{x}_1 + k_2 (x_2 - x_1) + b_2 (\dot{x}_2 - \dot{x}_1) + F_1 = 0$$

$$\bullet m_1 \ddot{x}_1(t) + (b_1 + b_2) \dot{x}_1(t) + b_2 \dot{x}_2(t) + (k_1 + k_2) x_1(t) - k_2 x_2(t) = F_1(t)$$

$$\bullet m_2 \ddot{x}_2(t) - b_2 \dot{x}_1(t) + (b_2 + b_3) \dot{x}_2(t) - b_3 \dot{x}_3(t) - k_2 x_1 + (k_2 + k_3) x_2 - k_3 x_3 = F_2(t)$$

$$\bullet m_3 \ddot{x}_3(t) - b_3 \dot{x}_2(t) + b_3 \dot{x}_3 - k_3 x_2 + k_3 x_3 = F_3(t)$$

$$\begin{bmatrix} m_1 & 0 & 0 \\ 0 & m_2 & 0 \\ 0 & 0 & m_3 \end{bmatrix} \begin{Bmatrix} \ddot{x}_1 \\ \ddot{x}_2 \\ \ddot{x}_3 \end{Bmatrix} + \begin{bmatrix} b_1 + b_2 & -b_2 & 0 \\ -b_2 & b_2 + b_3 & -b_3 \\ 0 & -b_3 & b_3 \end{bmatrix} \begin{Bmatrix} \dot{x}_1 \\ \dot{x}_2 \\ \dot{x}_3 \end{Bmatrix} + \begin{bmatrix} k_1 + k_2 & -k_2 & 0 \\ -k_2 & k_2 + k_3 & -k_3 \\ 0 & -k_3 & k_3 \end{bmatrix} \begin{Bmatrix} x_1 \\ x_2 \\ x_3 \end{Bmatrix} = \begin{Bmatrix} F_1 \\ F_2 \\ F_3 \end{Bmatrix}$$

$$[M] \{\ddot{x}(t)\} + [b] \{\dot{x}(t)\} + [k] \{x(t)\} = \{F(t)\}$$

↑  
Matrice  
massa

↑  
smorzamento

↑  
rigidezza

Titolo:

$$[M] \{\ddot{x}(t)\} = \text{Forze d'inerzia}$$

$$[b] \{\dot{x}(t)\} = \text{Forze viscosse} \quad \{F\} = \text{forzanti}$$

$$[k] \{x(t)\} = \text{Forze elastiche}$$

Studi per oscillazioni libere

$$[M] \{\ddot{x}(t)\} + [k] \{x(t)\} = \{0\}$$

La soluzione

$$\{x(t)\} = \sum_{i=1}^m \{\psi^i\} C_i \sin(\omega_i t + \varphi_i) \cdot$$

$C_i$  e  $\varphi_i \Rightarrow$  condizioni iniziali

$$\{x(t)\} = \{\psi^i\} C_i \sin(\omega_i t + \varphi_i) \cdot$$

$$\{\dot{x}(t)\} = -\{\psi^i\} \omega_i^2 C_i \sin(\omega_i t + \varphi_i) = -\omega_i^2 x(t)$$

$$-\omega_i^2 [M] \{\psi^i\} C_i \sin(\omega_i t + \varphi_i) + [k] \{\psi^i\} C_i \sin(\omega_i t + \varphi_i) = 0$$

$$-\omega_i^2 [M] \{\psi^i\} + [k] \{\psi^i\} = \{0\}$$

Titolo:

$$\bullet \left( -\omega_i^2 [M] + [K] \right) \{ \psi^i \} = \{ 0 \}$$

↑  
incognite

Affinchè si abbiano soluzioni  
non banali del sistema omogeneo

$$\det \left( -\omega_i^2 [M] + [K] \right) = 0 \quad \text{eq. caratteristica}$$

di  $n$   $\omega_i$

$\omega_i^2 \Rightarrow$  reali e positive

$$\omega_1^2 \leq \omega_2^2 < \dots < \omega_n^2$$

$$T_1^2 > T_2^2 > \dots > T_n^2$$

$\omega_1$  (la più piccola)  
sarà la pulsazione  
fondamentale

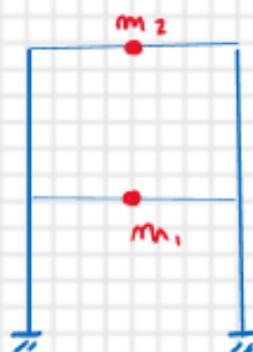
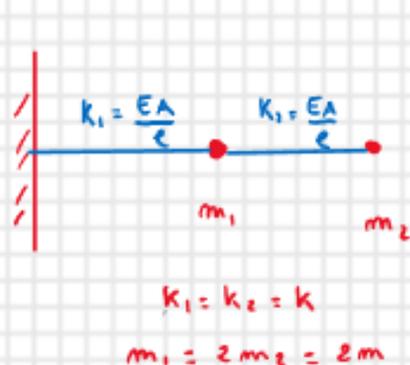
$$\omega_i \Rightarrow \{ \psi^i \}$$

$$\{ x(t) \} = \sum_{i=1}^n \{ \psi^i \} C_i \sin(\omega_i t + \varphi_i)$$

$$\{ x(t) \} = \{ \psi^i \} C_i \sin(\omega_i t + \varphi_i)$$

Se il sistema vibra secondo un modo di  
vibrare i valori max e nulli dello  
spostamento si verificano nello stesso istante

## Titolo:



$$k_i = \frac{24ES}{h^3}$$

- $m_1 \ddot{x}_1(t) + (k_1 + k_2)x_1(t) - k_2 x_2(t) = 0$
- $m_2 \ddot{x}_2(t) - k_2 x_1(t) + k_2 x_2(t) = 0$

$$\begin{bmatrix} m_1 & 0 \\ 0 & m_2 \end{bmatrix} \begin{Bmatrix} \ddot{x}_1 \\ \ddot{x}_2 \end{Bmatrix} + \begin{bmatrix} k_1 + k_2 & -k_2 \\ -k_2 & k_2 \end{bmatrix} \begin{Bmatrix} x_1 \\ x_2 \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0 \end{Bmatrix}$$

$$\left( -\omega_i^2 [M] + [K] \right) \{ \psi^i \} = \{ 0 \}$$

$$\left( -\omega_i^2 m \begin{bmatrix} 2 & 0 \\ 0 & 1 \end{bmatrix} + \begin{bmatrix} 2k & -k \\ -k & k \end{bmatrix} \right) \{ \psi^i \} = \{ 0 \}$$

$$\begin{pmatrix} -2\omega_i^2 m + 2k & -k \\ -k & -\omega_i^2 m + k \end{pmatrix} \{ \psi^i \} = \{ 0 \} \quad \det \begin{pmatrix} -2\omega_i^2 m + 2k & -k \\ -k & -\omega_i^2 m + k \end{pmatrix} = 0$$

$$2\omega_i^4 m^2 - 4\omega_i^2 m k + k^2 = 0 \quad \text{eq. caratteristica}$$

$$\omega_i^2 = \frac{k}{m} \frac{2 \pm \sqrt{2}}{2}$$

$$\underline{\omega_1^2 = \frac{k}{m} \frac{2 - \sqrt{2}}{2}}$$

$$\underline{\omega_2^2 = \frac{k}{m} \frac{2 + \sqrt{2}}{2}}$$

## Titolo:

•  $\omega_1^2 = \frac{k}{m} \frac{2-\sqrt{2}}{2}$

Lo inseriamo nel sistema e calcoliamo gli autovettori

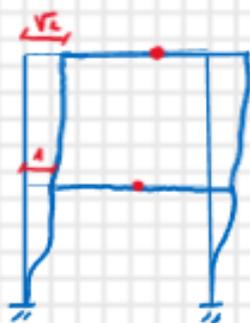
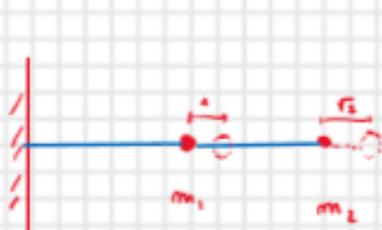
$$\begin{cases} (-2\omega_1^2 m + 2k) \psi_1^{(1)} - k \psi_2^{(1)} \\ -k \psi_1^{(1)} + (-\omega_1^2 m + k) \psi_2^{(1)} = 0 \end{cases} \Rightarrow \psi_1^{(1)} = \frac{\psi_2^{(1)}}{\sqrt{2}}$$

$$\{\psi^{(1)}\} = \begin{Bmatrix} 1 \\ \sqrt{2} \end{Bmatrix} \Rightarrow \begin{Bmatrix} 1/\sqrt{3} \\ \sqrt{2}/3 \end{Bmatrix} \quad \text{1° modo di vibrare}$$

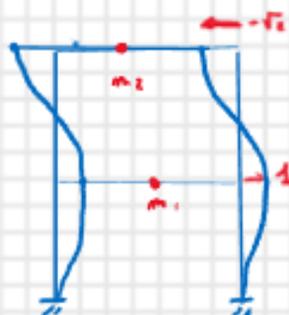
•  $\omega_2^2 = \frac{k}{m} \frac{2+\sqrt{2}}{2} \quad \{\psi^{(2)}\} = \begin{Bmatrix} 1 \\ -\sqrt{2} \end{Bmatrix}$

$$\{x(t)\} = \sum_{i=1}^n \{\psi^{(i)}\} C_i \sin(\omega_i t + \varphi_i)$$

Soluzione (calcolare condizioni iniziali)



1° modo di vibrare  $\begin{Bmatrix} 1 \\ \sqrt{2} \end{Bmatrix}$



2° modo di vibrare  $\begin{Bmatrix} 1 \\ -\sqrt{2} \end{Bmatrix}$