

### **993SM - Laboratory of Computational Physics II week October 4, 2024**

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# **more on the Numerov's algorithm**

- deeper analysis of the Numerov's algorithm
- numerical analysis: search for zeros

### **1D Schroedinger equation** and the Numerov's method

#### The Scroedinger eq.:

$$
-\frac{\hbar^2}{2m}\frac{d^2\psi}{dx^2} + V(x)\psi(x) = E\psi(x)
$$

has the form:

$$
\frac{d^2y}{dx^2} = -g(x)y(x) + s(x)
$$

with:

$$
g(x) = (2m/\hbar^2)[E - V(x)]
$$
 and  $s(x) = 0$ 

### **1D Schroedinger equation for harmonic oscillator** and the Numerov's method

Since  $s(x) = 0$ , the Numerov's formula reduces to:

$$
y_{n+1} \left[ 1 + g_{n+1} \frac{(\Delta x)^2}{12} \right] = 2y_n \left[ 1 - 5g_n \frac{(\Delta x)^2}{12} \right] - y_{n-1} \left[ 1 + g_{n-1} \frac{(\Delta x)^2}{12} \right] + O[(\Delta x)^6]
$$
  
Defining:  $f_n \equiv 1 + g_n \frac{(\Delta x)^2}{12}$  with  $g(x) = 2 \left( \frac{\varepsilon}{\varepsilon} \right) - \frac{x^2}{2}$ 

we rewrite Numerov's formula as

$$
y_{n+1} = \frac{(12-10f_n)y_n - f_{n-1}y_{n-1}}{f_{n+1}}
$$
  
contain  $\varepsilon$ 

### **inside the Numerov's method**

Find a solution  $\psi$  with a given number not nodes and energy En

Possibility 1): we fix En Possibility 2): we do not know  $En$ , but we start from an initial energy guess  $E$ ; the guess is on an energy interval  $[ELow, Eup]$  that we know for sure to contain En

For each value of E: start integrating  $\psi$  from  $x=0$  towards positive values of  $x$ (forward)

During integration, count for the number ncross of changes of sign of ψ: if  $ncross>n \equiv E$  is too high => choose  $[Elow, E]$ if  $ncross < n \implies E$  is too low => choose  $[E, Eup]$ 

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```
search_loop: do
                     set initial lower and upper bounds to the eigenvalue
                   eup=maxval (vpot(:)) : initially = vpot(xmax)
                   elw=minval (vpot(:)) : initially = vpot(0)
                                                                                     Possibility | \cdot |: fix E
                     Set trial energy
                  write(*,"('Trial energy|(0=search with bisection) > ')", advance='no')
                  read (*,*) e
                   if ( e == 0.0 dp ) then
                       ! search eigenvalues with bisection (max 1000 iterations)
                      e = 0.5 dp * (elw + eup)
                      n iter = 1000
                  else
                       ! test a single energy value (no bisection)
                      n iter = 1
                   endif
                   iterate: do k = 1, n_iter | Iterate (k=1 only!) to integrate y
                                outward integration and count number of crossings
                                             ! iterate Numerov's algorithm all over the mesh
                              ncross=0
                              do i =1.mesh-1
                                 y(i+1) = ((12.8_d - d - 18.8_d + 12.8_d + 12.8_d - d - 12.8_d + 12.8_d + 12.8_d - 12.8_d -if (y(i) /= sign(y(i),y(i+1)) ncross=ncross+1
                              end do
                              print *, k, e, ncross, hnodes
                                  …
                                                               half nodes (only for x>0)
                  end do iterate
end do search_loop
```
(ADD!) Calculate and plot eigenfunctions for various values of "**nodes**" (number of **nodes** of the eigenfunction), from nodes=0, 1, 2,… Input also the precise energy value, which is (nodes+1/2).  $\Omega$ .

It may be useful to plot, together with eigenfunctions or eigenfunctions squared, the classical probability density, contained in the fourth column of the output file. It will clearly show the classical inversion points. See harmonic0.gpl on Moodle as an example of a macro for gnuplot.

#### *I/O on the screen:*



Solution for n=1 and energy exact (1.5)



х

# **inside the Numerov's method**

Find a solution  $\psi$  with a given number n of nodes and energy En

Possibility 1): we fix En

Possibility 2): we do not know  $En$ , but we start from an initial energy guess  $E$ ; the guess is on an energy interval  $[ \text{Elow}, \text{Eup}]$  that we know for sure to contain En

For each value of E: start integrating  $\psi$  from  $x=0$  towards positive values of  $x$ (forward)

During integration, count for the number  $ncross$  of changes of sign of  $\psi$ : if  $ncross>n \equiv E$  is too high => choose  $[Elow, E]$ if  $ncross < n \implies E$  is too low => choose  $[E, Eup]$ 

we expect a set of orthogonal eigenvectors with increasing number of zeros io ncreasing energy eigenvalue, more precisely, equal to the energy quantum fumber if, because the number of zeros is equal to the degree of the polynomia We expect a set of orthogonal eigenvectors with increasing number of zeros for increasing energy eigenvalue; more precisely, equal to the energy quantum number n, because the number of zeros is equal to the degree of the polynomial

```
search_loop: do
                set initial lower and upper bounds to the eigenvalue
               eup=maxval (vpot(:)) : initially = vpot(xmax)
               elw=minval (vpot(:)) : initially = vpot(0)
                                                              Possibility 2): iterate on E
               ! Set trial energy
               write(*,"('Trial energy (0=search with bisection) > ')", advance='no')
               read (*,*) e
               if (e == 0.0 dp) then
                  ! search eigenvalues with bisection (max 1000 iterations)
                  e = 0.5 dp * (elw + eup)
                  n iter = 1000
               else
                  ! test a single energy value (no bisection)
                  n_iter = 1
               endif
                                                 Iterate to integrate y
               iterate: do k = 1, n\_iterif (ncross > hnodes) then
                              ! Too many crossings: current energy is too high
                              ! lower the upper bound
                              eup = eelse
                              ! Too few (or correct) number of crossings:
                              ! current energy is too low, raise the lower bound
                              elw = e
                           end if
                            ! New trial value:
                           e = 0.5 dp * (eup+elw)
                                                              iterate on energy
                            ! Convergence criterion:
                           if (eup-elw < 1.d-10) exit iterate
                           \mathbf{I}.
              end do iterate
end do search_loop
```
1. Calculate and plot eigenfunctions for various values of  $n$  (number of nodes of the eigenfunction), from  $n=0, 1, 2,...$  Let the code find the correct energy value (choosing  $"0"$  when the code asks for an energy value). It may be useful to plot, together with eigenfunctions or eigenfunctions squared, the classical probability density, contained in the fourth column of the output file. It will clearly show the classical inversion points. See harmonic0.gpl on Moodle as an example of a macro for gnuplot.

*I/O on the screen:*

Max value for x (typical value: $10$ ) > 5		
Number of grid points (typically a few hundreds) > 500		
output file name > dat		
nodes (type $-1$ to stop) > 1	ncross	hnodes
Trial energy (0=search with bisection) > 0		
6.2500000000000000		
3.1250000000000000		
1.5625000000000000		
0.78125000000000000		
1.1718750000000000 5		<b>UZ C</b> ncross=hnodes is ok,
1.3671875000000000 6		but the stop criterion
. .		is a tolerance on the energy
		(now the uncertainty is $\sim 0.4$ )
1.5000000035797711 34		
1.5000000032159733 35		
1.5000000033978722 36		
1.5000000033069227 37		uncertainty $<$ tolerance = 10^-10

*The results* 

2. Specify an energy value not corresponding to an eigenvalue. Look at the resulting wavefunctions.



2. Specify an energy value not corresponding to an eigenvalue. Look at the resulting wavefunctions.



3. Specify an energy value close to but not exactly corresponding to an eigenvalue. Look at the resulting wavefunctions.



4. Examine the effects of the parameters xmax, mesh. For a given  $\Delta x$ , how large can be the number of nodes? Hint: the nodes are internal to the classical limit range  $\Rightarrow$  $\sqrt{2n+1}$  < xmesh  $\implies$   $n_{max}$  <  $(x$ mesh<sup>2</sup> - 1)/2



# **inside the Numerov's method**

Used to count the number of nodes while "building" the wavefunction:

if (  $y(i)$  /= sign( $y(i)$ ,  $y(i+1)$ ) ) ncross=ncross+1 (ncross means crossing with the x axis, i.e., zeros)

Used to count the number of change of sign of 1-f, determining the classical inversion point

 $f(i)=ddx12*(2.0~dp*(vpot(i)-e))$ if (  $f(i)$  /=  $sign(f(i), f(i-1))$  ) icl=i

**possibility 2): guess on the energy, then refine (iterating, "shooting" method)**

# **inside the Numerov's method:** (rectangular) numerical integration

```
norm = 0.0 dpp(ic): = 0.0 dp
do i=0, icl
   arg = (e - x(i) \times 2/2.0 d\rho)if (arg > 0.0 dp) then
      p(i) = 1.0 dp/sqrt(arg)
   else
      p(i) = 0.0_dp
   end if
   norm = norm + 2.0_dpx+dx*p(i)enddo
```

```
! p allocated from 0:nmesh:
! here p(i)=0 for i>icl, i.e. outside the classical inversion point
I the point x(i) is smaller than the max elongation, dictated by e
! therefore, 1/\sqrt{\sqrt{1+x^2}} = 1/\sqrt{\sqrt{x\sqrt{1+x^2}} - x^2}! the simplest integration (rectangular method)
```
#### Polinomi ortogonali classici rispetto al peso w nell'intervallo I:

∫*I*  $w(x)P_n(x)P_m(x) \propto \delta_{n,m}$ 



$$
\textstyle \int_{-\infty}^\infty dx \ e^{-x^2} H_n(x) H_m(x) = \sqrt{\pi} 2^n n! \delta_{n,m}
$$

$$
H_0 = 1
$$
  
\n
$$
H_1 = 2x
$$
  
\n
$$
H_2 = 4x^2 - 2
$$
  
\n
$$
H_3 = 8x^3 - 12x
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\n
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H_4 = 16x^4 - 48x^2 + 12
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\n
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H_4 = 16x^4 - 48x^2 + 12
$$

# **Part II: search for zeros**

- 1. determine an interval  $[x_L,x_U]$  at whose extremes the function  $y(x)$ has discordant signs (to be sure that it contains  $x_0$ )
- 2. calculate the midpoint of the interval  $[x_L,x_U]: x_M = (x_L + x_U)/2$  and evaluate  $y(x_M)$ ;
- 3. If  $y(x_M)=0$  then  $x_M = x_0$  and the search ends.
- 4. Otherwise, take as the new interval the one at whose extremities the function has discordant signs (depending on the case it will be necessary to redefine  $x_L=x_M$  or  $x_U=x_M$ ): it contains  $x_0$
- 5. iterate points 2 4 until:
- a) the uncertainty on the location of  $x_0$  decreases below a preestablished an absolute threshold  $(|x_U-x_L| \le \epsilon)$ , or a relative threshold  $(|x_U-x_L| \le \epsilon |x_L|)$  or  $\le \epsilon |x_U|$ , where  $x_0$  is approximated by  $x<sub>L</sub>$  or  $x<sub>U</sub>$ ); or
- b)  $|y(x_M)| \leq \varepsilon$ '; or
- c) a maximum number of iterations is exceeded.

1. determine an interval  $[x_L,x_U]$  at whose extremes the function  $y(x)$  has discordant signs (to be sure that it contains  $x_0$ )

Implementation:

 $y(xL) * y(xU) < 0$ 

 $y(xL)$  /= sign(y(xL),  $y(xU)$ ) or  $y(xU)$  /= sign(y(xU),  $y(xL)$ )

4. take as the new interval the one at whose extremities the function has discordant signs (depending on the case it will be necessary to redefine  $x_L = x_M$  or  $x_U = x_M$ ): it contains  $x_0$ 

```
if (y(xL) * y(xM) < 0) then
    xU = xMelse if (y(xU) * y(xM) < 0) then
    xL = xMend if
Implementation:
```
- 5. iterate points 2 4 until:
- a) the uncertainty on the location of  $x_0$  decreases below a preestablished an absolute threshold ( $|x_U-x_L| < \varepsilon$ ), or a relative threshold  $(|x_U-x_L| \le \varepsilon |x_L| \text{ or } \le \varepsilon |x_U|$ , where  $x_0$  is approximated by  $x_L$  or  $x_U$ )

Implementation - which is the best criterion?

- use an absolute threshold: possible problems (for roundoff errors) if  $x_0$  is large
- use a relative threshold: possible problems if  $x_0$  is small
- other possible problems if  $y(x)$  is too flat close to  $x_0$

### **Exercise:**

# **implement this algorithm** if you never did, or if you have forgotten…