#### <span id="page-0-0"></span>Introduction to Molecular Structure Molecular Orbital Theory, LCAO, Hydrogen Molecule, Hückel Method

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# <span id="page-3-0"></span>Born-Oppenheimer Approximation

Introduction

Even a simple system like the  ${\sf H}_2^+$  ion has a complex Hamiltonian:

$$
\hat{H} = -\frac{\hbar^2}{2m_e}\nabla^2 - \frac{\hbar^2}{2m_p}\left(\nabla_A^2 + \nabla_B^2\right) + V(\vec{r}; \vec{R}_A, \vec{R}_B)
$$

where:

- $\bullet$   $m_e$  is the electron mass, and  $m_p$  is the proton mass.
- $\nabla^2$  acts on electron's coordinates

• 
$$
\nabla_A^2
$$
 and  $\nabla_B^2$  act on the nuclei

$$
\bullet \;\; V(\vec{r};\vec{R}_A,\vec{R}_B) = -\tfrac{e^2}{4\pi\epsilon_0} \left( \tfrac{1}{|\vec{r}-\vec{R}_A|} + \tfrac{1}{|\vec{r}-\vec{R}_B|} + \tfrac{1}{|\vec{R}_A-\vec{R}_B|} \right)
$$

- **IMPOSS** 10 solve for the wave function if we are to consider all three particles together.
- BO approximation:  $m_p$  is at least 1000 $m_e \rightarrow$  treat the nuclei as fixed
- This *decoupl[es](#page-2-0)* electronic and nuclear degrees [of](#page-4-0) [f](#page-2-0)[re](#page-3-0)[e](#page-4-0)[d](#page-1-0)[o](#page-2-0)[m](#page-5-0)[.](#page-6-0)  $\mathbb{B} \rightarrow \mathbb{R} \oplus \mathbb{R}$

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## <span id="page-4-0"></span>Born-Oppenheimer approximation

Separation of Variables

• Propose a solution

$$
\Psi(\vec{r},\vec{R}_A,\vec{R}_B) = \psi(\vec{r};\vec{R}_A,\vec{R}_B) \chi(\vec{R}_A,\vec{R}_B)
$$

• Substituting this into the full Schrödinger equation yields two equations:

$$
\begin{aligned}\n&\left[-\frac{\hbar^2}{2m_e}\nabla^2 + V(\vec{r}; \vec{R}_A, \vec{R}_B)\right]\psi(\vec{r}; \vec{R}_A, \vec{R}_B) = E_e \psi(\vec{r}; \vec{R}_A, \vec{R}_B) \\
&\left[-\frac{\hbar^2}{2m_p}\left(\nabla_A^2 + \nabla_B^2\right) + E_e\right]\chi(\vec{R}_A, \vec{R}_B) = E\chi(\vec{R}_A, \vec{R}_B)\n\end{aligned}
$$

After we solve the first equation for the fixed nuclei, we can solve the second equation with  $E_e$ .

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#### <span id="page-5-0"></span>Born-Oppenheimer approximation

Potential Energy Surface

- We notice that  $E_e = E_e(\vec{R}_A,\vec{R}_B)$  for both the ground state and excited states.
- Assuming that excited levels do not cross, we can view this function as a *potential energy surface* on which electrons move.



Ground and 1st excited states of the molecule [for](#page-4-0)[ma](#page-6-0)[ld](#page-4-0)[im](#page-5-0)[i](#page-6-0)[n](#page-1-0)[e](#page-2-0) [f](#page-5-0)[o](#page-6-0)[r](#page-1-0) [t](#page-2-0)[w](#page-5-0)o [an](#page-0-0)[gle](#page-34-0)s.

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Linear Combination of Atomic Orbitals (LCAO)

Molecular Orbital Theory (MOT) describes the formation of molecular orbitals (MOs) from atomic orbitals (AOs).



• An *atomic orbital*, in this context, is an orbital that was obtained by solving the Schrödinger equation for an isolated atom, localized on that atom.

<span id="page-8-0"></span>Linear Combination of Atomic Orbitals (LCAO)

• In LCAO, we construct single-electron MO's as superpositions of atomic orbitals on all atoms involved

$$
\Psi(\vec{r}) = \sum_i c_i \phi(\vec{r} - \vec{R}_i)
$$

where *i* runs over all the atoms, and  $\phi$  are the atomic orbitals.

- By placing this wave function into a Schrödinger equation and solving for the coefficients, we obtain *molecular orbitals*.
- This assumes a single orbital per atom  $\rightarrow$  can be extended.

<span id="page-9-0"></span>Linear Combination of Atomic Orbitals (LCAO)

• Begin with the expansion for a molecular orbital  $\psi$  as a linear combination of atomic orbitals (LCAO) from multiple atoms:

$$
\psi = \sum_i c_i \phi_i
$$

- $\bullet \psi$  is the molecular orbital.
- $\phi_i$  is the atomic orbital on atom  $i$  (e.g., hydrogen 1s orbitals).
- $\bullet$   $c_i$  are the coefficients determining the contribution of each atomic orbital.
- $\bullet$  We aim to find the coefficients  $c_i$  by solving the Schrödinger equation:

$$
\hat{H}\psi = E\psi
$$

wh[e](#page-6-0)re  $\hat{H}$  is the molecular Hamiltonian, and [E](#page-8-0) [is](#page-10-0) [it](#page-8-0)[s](#page-9-0) e[n](#page-5-0)e[rg](#page-11-0)[y](#page-12-0)[.](#page-5-0)

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<span id="page-10-0"></span>Linear Combination of Atomic Orbitals (LCAO)

■ Start by substituting the LCAO expansion into the Schrödinger equation:

$$
\hat{H}\left(\sum_{i}c_{i}\phi_{i}\right)=E\left(\sum_{i}c_{i}\phi_{i}\right)
$$

 $\bullet$  Multiply both sides by  $\phi^*_j$  and integrate over all space:

$$
\sum_i c_i \int \phi_j^* \hat{H} \phi_i d\vec{r} = E \sum_i c_i \int \phi_j^* \phi_i d\vec{r}
$$

**3** Define matrix elements:

$$
H_{ij} = \int \phi_j^* \hat{H} \phi_i d\vec{r}, \quad S_{ij} = \int \phi_j^* \phi_i d\vec{r}
$$

- $\bullet$   $H_{ii}$  is the Hamiltonian matrix element (Coulomb and exchange integrals).
- $\bullet$   $S_{ij}$  is the overlap matrix element.

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<span id="page-11-0"></span>Linear Combination of Atomic Orbitals (LCAO)

The generalized secular equation obtained from the LCAO expansion can be written in matrix form:

$$
\vec{H}\vec{c} = E\vec{S}\vec{c}
$$

- $\bullet$  H is the Hamiltonian matrix with elements  $H_{ii}$ .
- $\overrightarrow{S}$  is the overlap matrix with elements  $S_{ii}$ .
- $\vec{c}$  is the vector of coefficients  $c_i$ .
- $\bullet$  Solve the secular equation by finding the eigenvalues  $E$  and the corresponding eigenvectors  $\vec{c}$ :

$$
\left|\vec{H}-E\vec{S}\right|=0
$$

This determinant gives the energy eigenvalues of the molecular orbitals.

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<span id="page-13-0"></span>LCAO Solution for H<sub>2</sub> Molecule

 $\bullet$   $\psi$ <sub>MO</sub> approximated as a linear combination of two H 1s orbitals localized on atoms A and B.

$$
\psi_{\mathsf{MO}} = c_{\mathsf{A}} \phi_{\mathsf{A}} + c_{\mathsf{B}} \phi_{\mathsf{B}}
$$

- Since the problem is very easy, we can construct the solutions by inspection.
- Two types of molecular orbitals arise:
	- Bonding molecular orbital: Constructive interference of the atomic orbitals  $(\psi_{+} = \phi_{A} + \phi_{B})$
	- Anti-bonding molecular orbital: Destructive interference of the atomic orbitals  $(\psi_- = \phi_A - \phi_B)$

<span id="page-14-0"></span>**Solutions** 

• To determine the molecular orbital energies, we solve the Schrödinger equation:

$$
\hat{H}\psi_{\text{MO}}=E\psi_{\text{MO}}
$$

For the LCAO approximation, we expand this equation as:

$$
c_A \hat{H} \phi_A + c_B \hat{H} \phi_B = E(c_A \phi_A + c_B \phi_B)
$$

By multiplying from the left by  $\phi_{\boldsymbol{A}}^*$  and  $\phi_{\boldsymbol{B}}^*$  and integrating over all space, we obtain a system of linear equations:

$$
c_A H_{AA} + c_B H_{AB} = E(c_A S_{AA} + c_B S_{AB})
$$
  

$$
c_A H_{BA} + c_B H_{BB} = E(c_A S_{BA} + c_B S_{BB})
$$

 $\bullet$  Here,  $H_{AA}$ ,  $H_{AB}$ , etc. are the Hamiltonian matrix elements, and  $S_{AA}$ ,  $S_{AB}$ , etc. are the overlap integrals:

$$
H_{AA} = \int \phi_A^* \hat{H} \phi_A d\vec{r}
$$

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**Solutions** 

The system of linear equations can be written in matrix form as:

$$
\begin{pmatrix} H_{AA} - ES_{AA} & H_{AB} - ES_{AB} \ H_{BA} - ES_{BA} & H_{BB} - ES_{BB} \end{pmatrix} \begin{pmatrix} c_A \ c_B \end{pmatrix} = 0
$$

For non-trivial solutions, the determinant of the coefficient matrix must vanish. This leads to the secular determinant:

$$
\begin{vmatrix} H_{AA} - ES_{AA} & H_{AB} - ES_{AB} \\ H_{BA} - ES_{BA} & H_{BB} - ES_{BB} \end{vmatrix} = 0
$$

 $\bullet$  Solving the secular determinant gives two energy eigenvalues  $E_{+}$  and  $E_{-}$ , corresponding to the bonding and anti-bonding molecular orbitals, respectively:

$$
E_{\pm} = \frac{H_{AA} + H_{BB}}{2} \pm \frac{\sqrt{(H_{AA} - H_{BB})^2 + 4H_{AB}^2}}{2}
$$
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Interpretation of Bonding and Anti-Bonding Orbitals



- The bonding molecular orbital  $(\psi_+)$  has lower energy  $(E_+)$  than the anti-bonding molecular orbital  $(E_-\)$ .
- This energy difference leads to the stability of the  $H_2$  molecule in its ground state.

 $\rightarrow$   $\rightarrow$   $\rightarrow$ 

Interpretation of Bonding and Anti-Bonding Orbitals



- The bonding orbital is characterized by a high electron density between the two nuclei, resulting in attractive interaction and bond formation.
- The anti-bonding orbital has a nodal plane between the nuclei, leading to repulsion.

 $\rightarrow$   $\equiv$   $\rightarrow$ 

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## MOT for Diatomic Molecules

Choosing the Atomic Orbitals for LCAO



- Must choose AOs of the same symmetry to avoid zero overlap.
	- *s*,  $p_z$ , and  $d_{z^2} \rightarrow \sigma$  orbitals
	- $p_v$ , and  $p_z \rightarrow \pi$  orbitals
	- $d_{yz}$ , and  $d_{zx} \rightarrow \pi$  orbitals
	- $d_{\mathsf{x}\mathsf{y}}$ , and  $d_{\mathsf{x}^2-\mathsf{y}^2} \to \delta$  orbitals
- The AOs should not be too *diffuse* or too *compact*: e.g.  $(1s, 1s)$  almost does not overlap but (2s, 2s) does.
- The AOs should be close in energy: e.g.  $(2s, 2s)$  and  $(2p, 2p)$  dominate the bonding while (2s, 1s) etc are too small.
- Construct MO diagrams for period 2 homonuclear diatomic molecules:  $N_2$ .  $O<sub>2</sub>$  etc.  $\left\{ \begin{array}{ccc} 1 & 0 & 0 \\ 0 & 1 & 0 \end{array} \right.$ 目

## MOT for Diatomic Molecules

MOT for Period 2 Homonuclear Diatomic Molecules



- $\bullet$  Consider only the valence orbitals  $-$  the core orbitals are far too compact.
- Four σ orbitals from the four 2s– and  $2p<sub>z</sub>$ -orbitals: two antibonding, two bonding.
- Four  $\pi$  orbitals from the four  $p_{\nu}$  and  $p_{\nu}$ orbitals.
- The energy order is difficult to predict.
- Finally we fill the MOs with the number of electrons that we have.
	- HOMO: Highest occupied MO
	- LUMO: Lowest unoccupied MO.

## MOT for Diatomic Molecules

MOT for Heteronuclear Diatomic Molecules

• In the LCAO proposal for the MO

$$
\psi = c_A \phi_A + c_B \phi_B
$$

 $c_A$  and  $c_B$  is not the same anymore.

- $\bullet$  In general, if A is more electronegative than B,  $|c_A|^2 > |c_B|^2$  for bonding and  $|c_A|^2 < |c_B|^2$  for antibonding.
- In the case of heteronuclear molecules, the AOs are different in energy  $\Rightarrow$  Bonding orbitals have lower energies, antibonding orbitals have higher energies due to weaker coupling.



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 $\mathcal{A} \oplus \mathcal{B}$   $\rightarrow$   $\mathcal{A} \oplus \mathcal{B}$   $\rightarrow$   $\mathcal{A} \oplus \mathcal{B}$ 

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Symmetry Adapted Linear Combinations (SALC)

- In forming the MO for polyatomic molecules, choosing symmetry-adapted linear combinations (SALCs) allows us to simplify the problem by using the symmetry properties of the molecule.
- The idea is to combine atomic orbitals in a way that respects the symmetry operations of the molecule.

Mathematical Formulation

$$
\Psi_i = \sum_j c_{ij} \phi_j
$$

where  $\Psi_i$  is the SALC,  $\phi_j$  are the atomic orbitals, and  $c_{ij}$  are the coefficients.

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## MOT for Polyatomic Molecules

Example: H<sub>2</sub>O Molecule



- $H_2O$  has two mirror symmetries and a rotational symmetry by 180 $^o$ , i.e.  $C_{2v}$  symmetry.
- Atomic orbitals of oxygen and hydrogen will combine to form molecular orbitals that respect this symmetry.

Example: H<sub>2</sub>O Molecule



Using H 1s, O1s, O2s and the three O2p orbitals, construct three sets of states that have zero overlap  $\Rightarrow$  smaller matrices to deal with rather than a  $6 \times 6$  matrix.

• Start with O2s, O2 $p_x$ , O2 $p_y$ , O2 $p_z$ , and two hydrogen orbitals:

(a) Symmetry classification of the O and H orbitals (b) Molecular orbitals  $\phi(A_1) = H1_{A} + H1_{S}$  $\phi(B_2) = H1s_A - H1s_B$ 

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Example: H<sub>2</sub>O Molecule



(a) Symmetry classification of the O and H orbitals (b) Molecular orbitals

- Start with O2s, O2 $p_x$ , O2 $p_y$ , O2 $p_z$ , and two hydrogen orbitals:
	- $\phi(A_1) = H1_{A} + H1_{B}$  $\phi(B_2) = H1s_A - H1s_B$
- **The SALCs for water are:**

$$
\psi(A_1) = c_1 \phi(O2s) + c_2 \phi(O2p_z) + c_3 \phi(A_1)
$$
  
\n
$$
\psi(B_1) = \phi(O2p_x)
$$
  
\n
$$
\psi(B_2) = c_4 \phi(O2p_y) + c_5 \phi(B_2)
$$

• Use these SALCs to construct the matrix, and obtain the MOs as shown in the Figure.

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## Conjugated  $\pi$ -Systems

Separation of  $\pi$  and  $\sigma$  systems



- Special class of polyatomic molecules: polyenes and arenes.
- The  $\pi$  and  $\sigma$  systems are completely separate and non-interacting.

 $\leftarrow$   $\Box$ 

 $\rightarrow$   $\equiv$   $\rightarrow$ 

## Conjugated  $\pi$ -Systems

A Famous Example: Graphene

- Graphene consists of a 2D honeycomb lattice of carbon atoms
- $\bullet$  In-plane bonds:  $sp^2$  hybridization

$$
\psi_{sp^2} = \frac{1}{\sqrt{3}} \left( \phi_{2s} + \phi_{2p_x} + \phi_{2p_y} \right)
$$

The  $\pi$  system: Superposition of 2 $p_z$  AOs  $\phi_{2p_z}$ .



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 $\mathcal{A} \oplus \mathcal{B}$   $\mathcal{B}$   $\mathcal{B}$   $\mathcal{B}$   $\mathcal{B}$   $\mathcal{B}$   $\mathcal{B}$   $\mathcal{B}$ 

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Introduction

- $\bullet$  Hückel theory provides a simple approach to study  $\pi$ -electron systems.
- $\bullet$  In Huckel theory, only p-orbitals are considered for forming molecular orbitals.
- The Huckel secular determinant for a conjugated molecule is:

$$
|H - ES| = 0
$$

• The method makes a series of assumptions to simplify the calculations.

Assumptions

- All overlap integrals are set equal to zero unless the orbitals are the same and localized on the same atom:  $S_{ii} = \delta_{ii}$ .
- All diagonal elements of the Hamiltonian matrix are given the same value:  $H_{ii} = \alpha$  with  $\alpha < 0$
- All off-diagonal elements of the hamiltonian are set equal to zero except for those between neighbouring atoms, all of which are set equal to  $\beta$ .

Example:  $\pi$  MOs of 1.3-butadiene C<sub>4</sub>H<sub>6</sub>



 $\left\{ \begin{array}{ccc} 1 & 0 & 0 \\ 0 & 1 & 0 \end{array} \right.$ 

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Example: Butadiene

We begin by constructing the secular determinant for butadiene with four carbon atoms in the  $\pi$ -system:

$$
\det \begin{pmatrix} \alpha - E & \beta & 0 & 0 \\ \beta & \alpha - E & \beta & 0 \\ 0 & \beta & \alpha - E & \beta \\ 0 & 0 & \beta & \alpha - E \end{pmatrix} = 0
$$

where:

- $\alpha$  is the Coulomb integral (on-site energy),
- $\beta$  is the resonance integral (hopping parameter),
- $\bullet$  E is the energy eigenvalue we are solving for.
- Each diagonal element corresponds to  $\alpha E$ , and the off-diagonal elements represent nearest-neighbor interactions between atoms.

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<span id="page-34-0"></span>Example: Butadiene



- The solutions are:  $E=\alpha\pm\beta$ √  $\frac{1}{3\pm\sqrt{5}}$  $\frac{\pm \sqrt{3}}{2}$  $\beta$
- The orbitals are

 $\psi_1 = 0.372\phi_A + 0.602\phi_B + 0.602\phi_C + 0.372\phi_D$  $\psi_2 = -0.602\phi_A - 0.372\phi_B + 0.372\phi_C + 0.602\phi_D$  $\psi_3 = -0.602\phi_A + 0.372\phi_B + 0.372\phi_C - 0.602\phi_D$  $\psi_4 = 0.372\phi_A - 0.602\phi_B + 0.602\phi_C - 0.372\phi_D$ 

where A, B, C, and D refer to the four C atoms in the molecule.

 $\bullet$   $\phi_1$ , and  $\phi_2$  have a higher density in the bonding region and are therefore lower energy while  $\phi_3$ , and  $\phi_4$  have lower density and higher energy.

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