

**II partial written test**  
**academic year 2011/2012**  
**January 16, 2012**

**Exercise 2:** *Effective masses and density of levels*

Consider bulk Silicon, whose conduction band minima  $E_c$  are near the Brillouin zone boundary along  $\langle 100 \rangle$  directions. Assume a parabolic conduction band with ellipsoidal constant energy surfaces around the minima, described by:

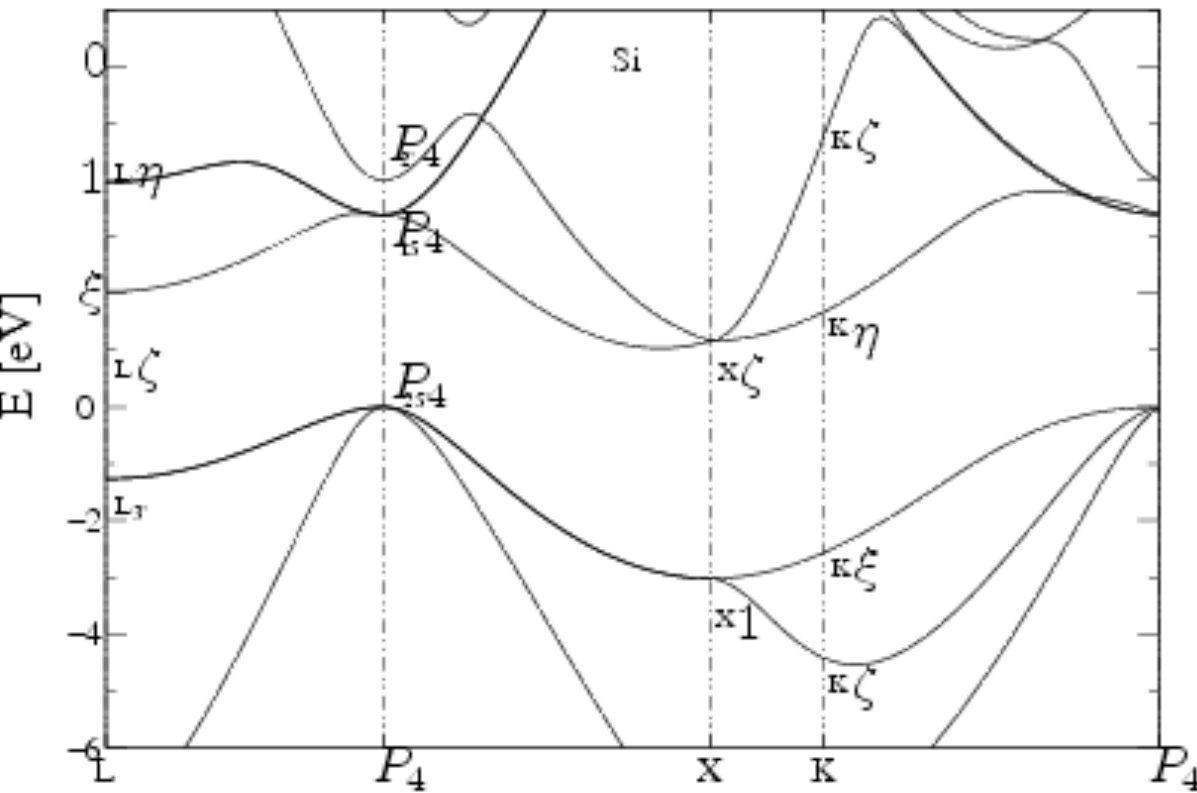
$$E(\mathbf{k}) = E_c + \frac{\hbar^2}{2} \left( \frac{k_\ell^2}{m_\ell^*} + \frac{k_t^2}{m_t^*} + \frac{k_t^2}{m_t^*} \right)$$

( $t$ =transverse,  $\ell$ =longitudinal, with  $m_\ell^* = 0.98m_0$ ,  $m_t^* = 0.19m_0$ , where  $m_0$  is the free electron mass;  $\vec{k}$  is with respect to the location of the minima).

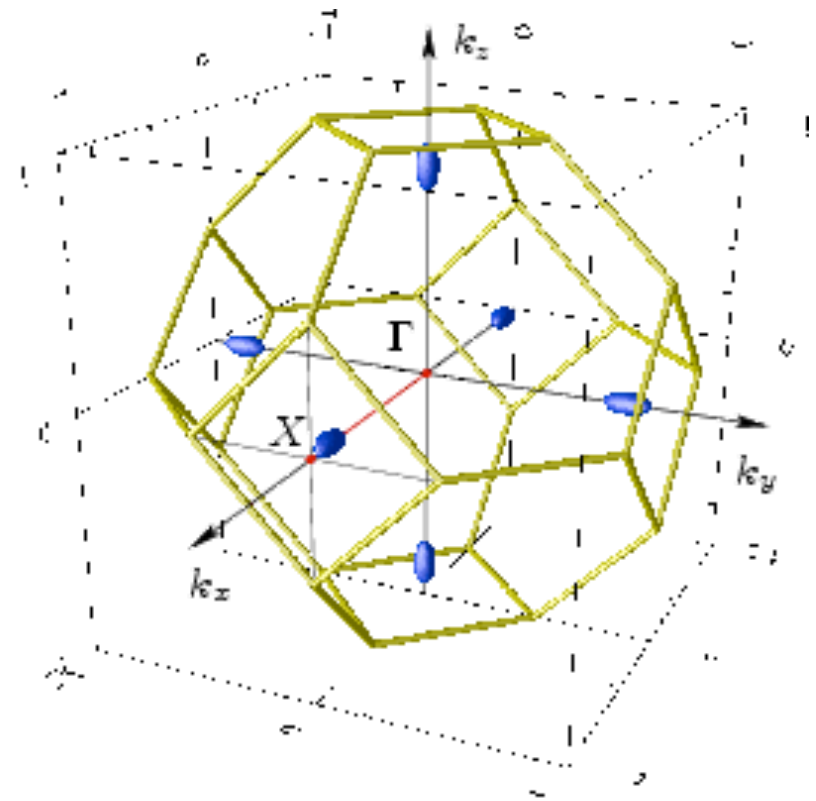
1. How many equivalent minima there are?
2. Write the expression of the density of states  $g(E)$  around one of the conduction band minima, in terms of  $E_c, m_\ell^*, m_t^*$ .
3. Calculate the number of states per unit energy for an energy 100 meV above the conduction band bottom, in a  $100 \times 100 \times 10$  nm piece of silicon. Write the result in units of  $eV^{-1}$ .

# Band structure of selected semiconductors

## Si



(a) Band diagram of silicon.



(b) First conduction band valleys.

indirect gap!

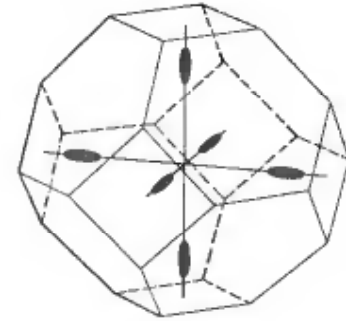
# Band structure of selected semiconductors

16/01/2012, II test  
Ex. 2

*Silicon* The crystal has the diamond structure, so the first Brillouin zone is the truncated octahedron appropriate to a face-centered cubic Bravais lattice. The conduction band has six symmetry-related minima at points in the  $\langle 100 \rangle$  directions, about 80 percent of the way to the zone boundary (Figure 28.5). . . .

Figure 28.5

Constant-energy surfaces near the conduction band minima in silicon. There are six symmetry-related ellipsoidal pockets. The long axes are directed along  $\langle 100 \rangle$  directions.



from Ashcroft-Mermin,  
Ch 28, Typical semiconductor band structure, p 569

From the Sommerfeld model:

$$g(E) = \frac{m}{\hbar^2 \pi^2} \sqrt{\frac{2mE}{\hbar^2}}, \quad E > 0$$

we obtain:

$$g(E) = \frac{\sqrt{m_e m_t^2}}{\hbar^2 \pi^2} \sqrt{\frac{2(E - E_c)}{\hbar^2}}, \quad E > 0$$

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$$g(E) \cdot V = ? \text{ where } V = 10^5 \text{ mm}^3 = 10^5 (10^{-7} \text{ cm})^3 \\ = 10^5 \cdot 10^{-21} \text{ cm}^3 = 10^{-16} \text{ cm}^3$$

$$\text{and } E = E_{\text{e}} + 100 \text{ meV}$$

$$\text{and } m_e = 0.98 m_0, m_t = 0.19 m_0$$

$$g(E) = \frac{\sqrt{0.98 \cdot 0.19^2} m_0^{3/2} \sqrt{2 \cdot 0.1 \text{ eV}}}{h^2 \pi^2} =$$

$$= \frac{\sqrt{0.98 \cdot 0.19^2} (9.11 \cdot 10^{-28} \text{ g})^{3/2} (2 \cdot 0.1 \cdot 1.6 \cdot 10^{-12} \text{ erg})^{1/2}}{\pi^2 (1.05 \cdot 10^{-27} \text{ erg} \cdot \text{s})^3}$$

$$= \frac{0.188 (9.11)^{3/2} (3.2)^{1/2} 10^{-\frac{14}{2} - \frac{28 \cdot 3}{2}} \cdot 10^{-\frac{12 \cdot 1}{2}}}{\pi^2 1.05^3 10^{-27 \cdot 3 - 81}} = 10^{33}$$

$$= 0.256 \cdot 10^{33} \text{ cm}^{-3} \text{ erg}^{-1} \Rightarrow g(E) \cdot V = 0.256 \cdot 10^{17} \text{ erg}^{-1}$$

$$[1 \text{ eV} = 1.6 \cdot 10^{-12} \text{ erg}] \rightarrow \frac{1}{\text{erg}} = \frac{1.6 \cdot 10^{-12}}{\text{eV}}$$

$$= 0.256 \cdot 1.6 \cdot 10^5 \text{ eV}^{-1}$$

$$= 0.4 \cdot 10^5 \text{ eV}^{-1}$$

$$[h] = [g][t] \quad \frac{[m]^{3/2} [g]^{1/2}}{[g]^2 [t]^2 [g][t]} = \frac{1}{[g]} \frac{1}{[L]^3}$$

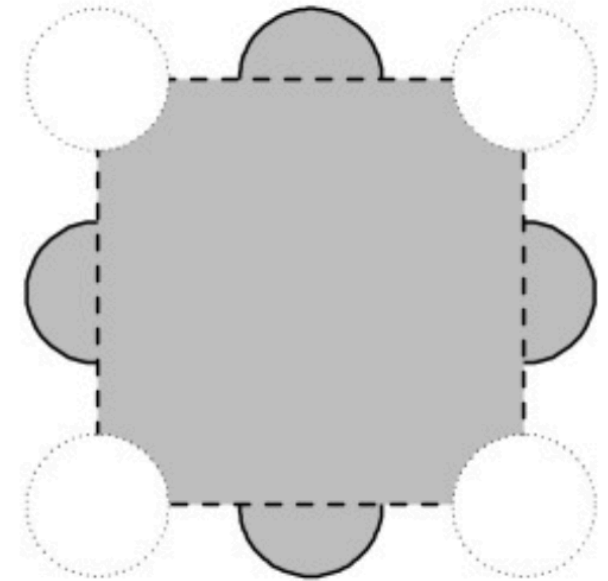
$$[g] = [F \cdot l] = [m a \cdot l] = [m] [l]^2 [t]^{-2}$$

$$\frac{[m]^{3/2}}{[g]^{3/2} [t]^3} = \frac{[m]^{3/2}}{[m]^{3/2} [L]^{3 \cdot \frac{3}{2}} [t]^{-2 \cdot \frac{3}{2}} [t]^3} \quad \text{Ok!}$$

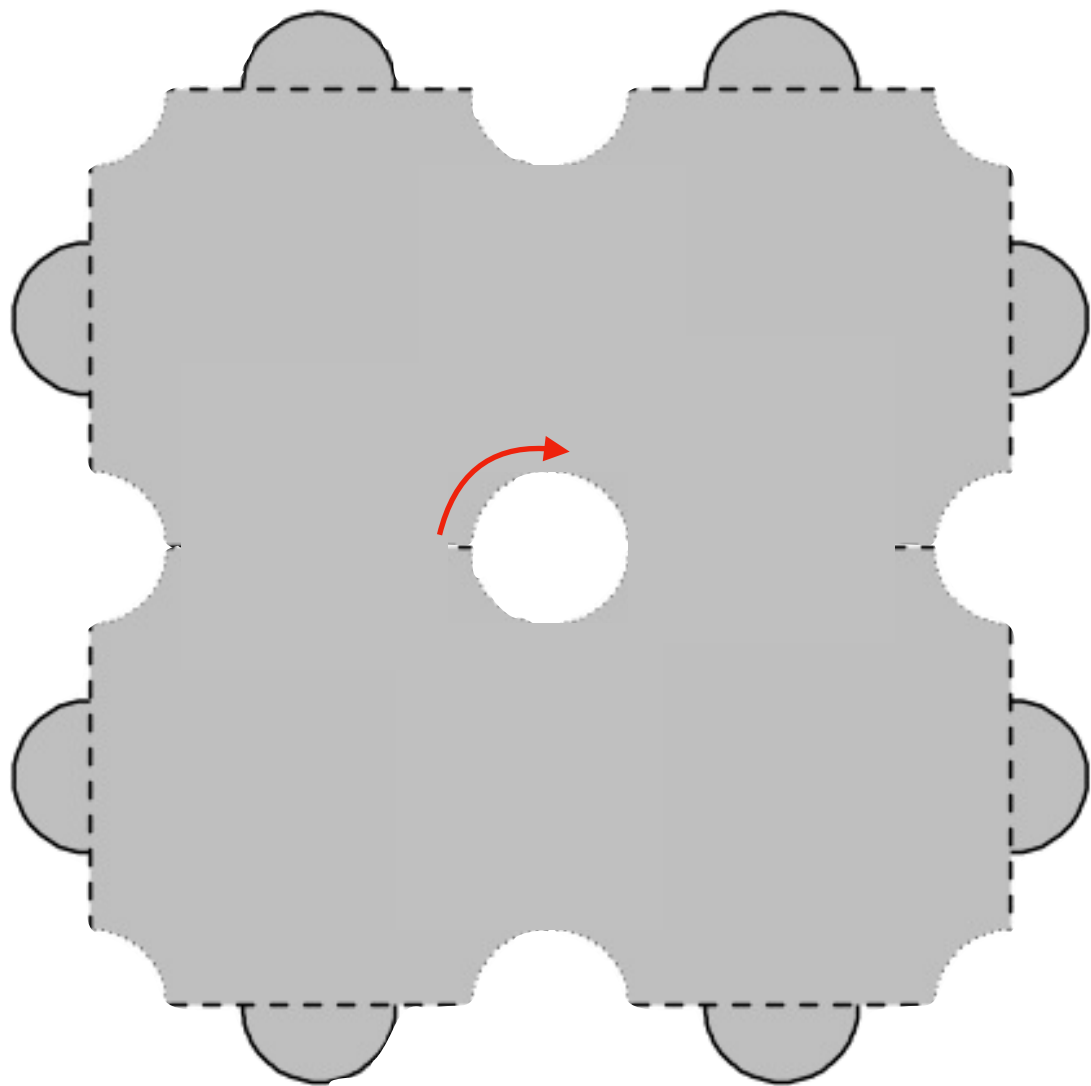
### Exercise 3: *Semi metals*

Bismuth is a "semi metal"; it has the second lowest thermal conductivity (after mercury) and the highest Hall coefficient, a high electrical resistance (or low electrical conductivity) (look for instance at Tab 1.2 and 1.6 of A&M book!).

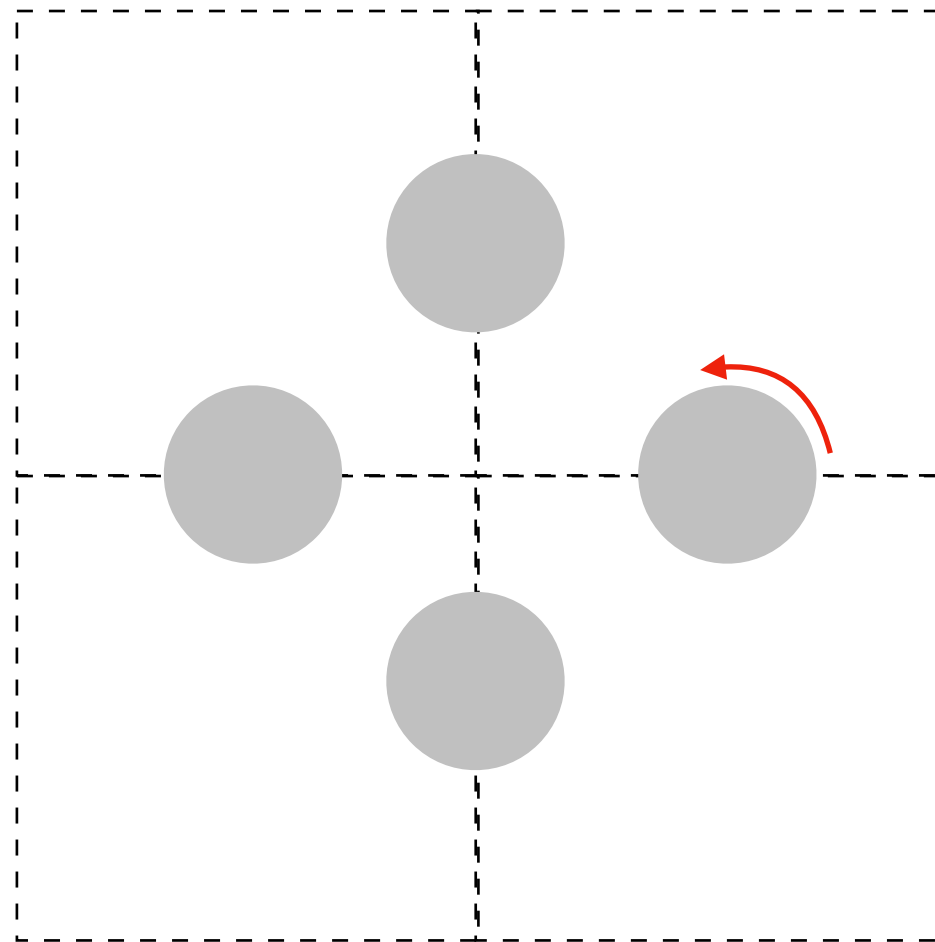
The unit cell is rhombohedral with two atoms (see Tab 7.5 A&M book), so it *could* be an insulator. However, there is a little band overlap that makes the situation similar to the case of a divalent metal with simple cubic lattice, whose Fermi-surfaces in  $(k_x, k_y)$  plane is shown in the figure. We refer therefore for simplicity to this case.



1. Make the same picture using the repeated zone scheme. Which part of the Fermi-surface can be described as electron-like and which as hole-like?
2. By which factor is the specific heat of the electrons at low temperatures ( $k_B T \ll E_F$ ) smaller than the electronic specific heat in the model of free electrons? For numerical estimations use the following data: the radius of the electronic Fermi-sphere is  $k_e = 0.1G_0$ , where  $G_0$  is the shortest reciprocal lattice vector; the bands in the vicinity of  $E_F$  are parabolic with effective masses  $m_e = m_h = 0.1m_0$ .
3. Why is the conductivity of such metals smaller than in the model of free electrons?



**h-like**



**e-like**

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Specific heat of free  $e^-$  :

$$c_v = \frac{\pi^2}{2} \left( \frac{k_B T}{\mathcal{E}_F} \right) m k_B \quad (2.81)$$

here  $\mathcal{E}_F = \frac{\hbar^2 k_F^2}{2m^*}$  with  $k_e = 0.1 G_0$   
 $m^* = 0.1 m_0$

$$G_0 = \frac{2\pi}{a} ; m_0 = \frac{2}{a^3} \quad \text{but } k_F^3 = 3\pi^2 m_0 = \frac{6\pi^2}{a^3}$$

$$\left( \frac{m}{\mathcal{E}_F} \right)^{3D} \propto \frac{\rho^3}{k_F^2} m_0 = m_0 k_F$$

$$\left[ \frac{c_v^*}{c_v^{\text{free}}} = \frac{k_e m^*}{k_F^{\text{free}} m_0} = 0.1 \cdot \frac{0.1 \cdot 2\pi/a}{(6\pi^2)^{1/3}/a} = 0.1 \cdot \frac{0.628}{3.9} \approx 0.016 \right]$$