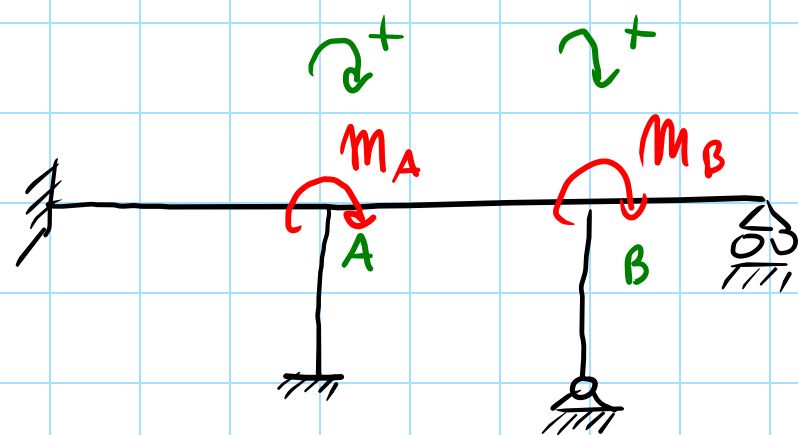


REVIEW DISPLACEMENT METHOD / STIFFNESS MATRIX

1/4/25

D.O.F. : φ_A, φ_B (UNKNOWN)



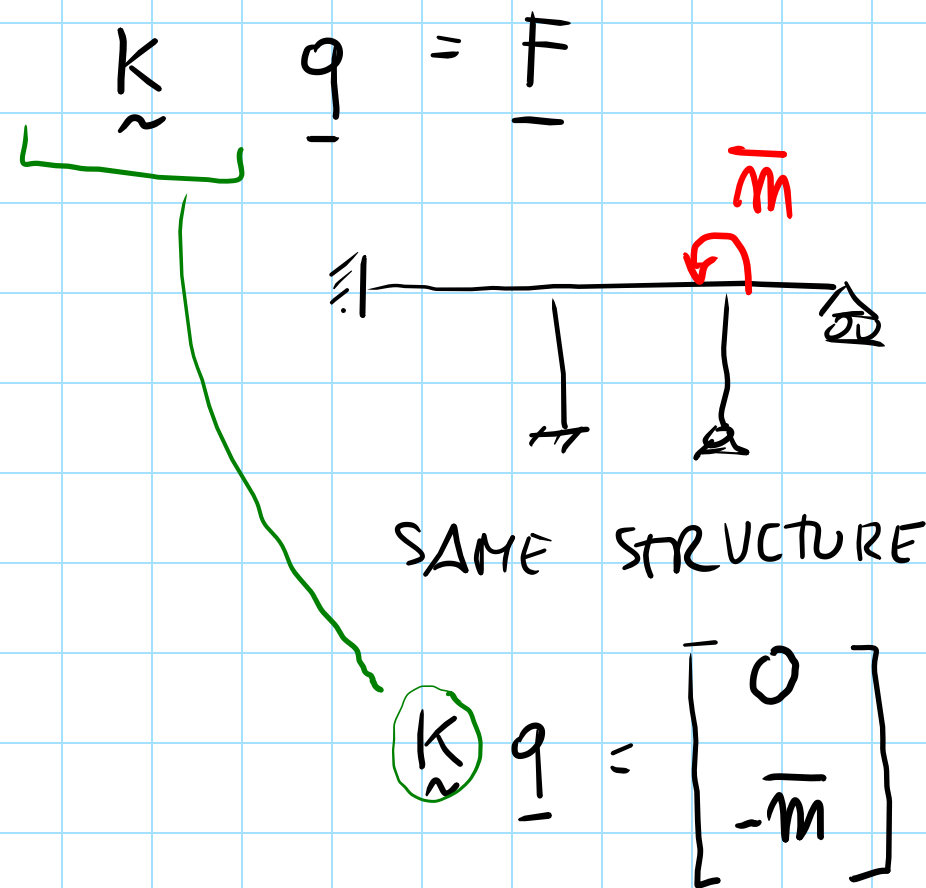
DISPL. METHOD ;
$$\begin{bmatrix} K_{11} & K_{12} \\ K_{21} & K_{22} \end{bmatrix} \begin{bmatrix} \varphi_A \\ \varphi_B \end{bmatrix} = \begin{bmatrix} M_A \\ M_B \end{bmatrix}$$

$\underline{\tilde{K}}$: STIFFNESS MATRIX, DEPENDS ON
GEOMETRY / MECHANICAL PROPERTIES / CONSTRAINTS
OF THE STRUCTURE.

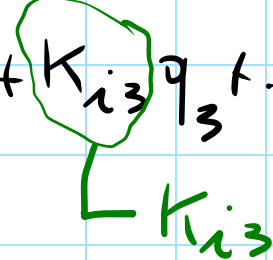
PROPERTIES OF $\underline{\tilde{K}}$:
— SYMMETRIC (TH. BETTI)
— POSITIVE DEFINITE (*)

(*) $\underline{\tilde{K}} \underline{q} \cdot \underline{q} > 0$, $\underline{q} \neq \underline{0}$
 $\underline{\tilde{K}} \underline{q} \cdot \underline{q} = 0$ IF ONLY IF $\underline{q} = \underline{0}$ } DEFINITION

FROM TH. CLARKE & RYAN $\Phi = \frac{1}{2} \underline{F} \cdot \underline{q} > 0 \Rightarrow \frac{1}{2} \underline{\tilde{K}} \underline{q} \cdot \underline{q} > 0$ FOR $\underline{q} \neq \underline{0}$



WE CAN COMPUTE ELEMENTS K_{ij} EITHER BY WRITING DOWN EQUILIBRIUM EQUATIONS AND "EXTRACTING" THE REQUIRED COEFFICIENT:

in the i -th EQUIL EQ $K_{i2}q_2 + K_{i3}q_3 + \dots$


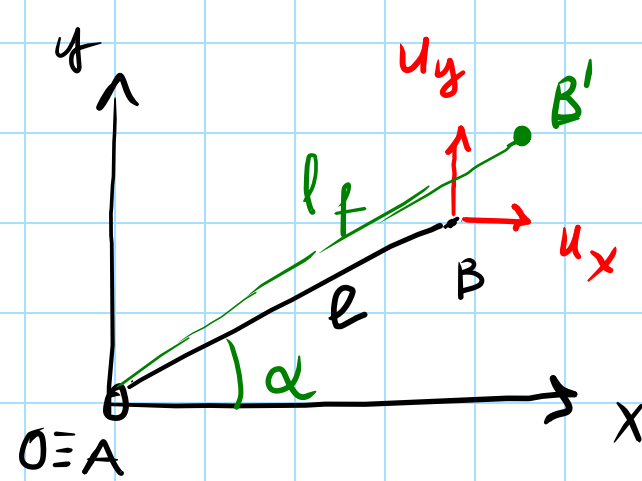
OR USING AN ALTERNATIVE APPROACH:

K_{ij} IS THE "FORCE" RETURNED TO THE i -th D.O.F. WHEN $q_j = 1, q_i = 0$
($i \neq j$)

$$\text{ex: } \left. \begin{array}{l} K_{11}q_1 + K_{12}q_2 = F_1 \\ K_{21}q_1 + K_{22}q_2 = F_2 \end{array} \right\} \Rightarrow \begin{array}{l} q_1 = 1 \\ q_2 = 0 \end{array} \Rightarrow \left. \begin{array}{l} K_{11} \cdot 1 = F_1 \\ K_{21} \cdot 1 = F_2 \end{array} \right\}$$

↳ WITH THIS "SPECIAL" CHOICE $\Rightarrow K_{11} = F_1, K_{21} = F_2$

DISPL. METHOD FOR PLANE TRUSS STRUCTURES



$$l = \sqrt{x_B^2 + y_B^2} \quad ; \quad l_f = \sqrt{(x_B + u_x)^2 + (y_B + u_y)^2} = l_f(u_x, u_y)$$

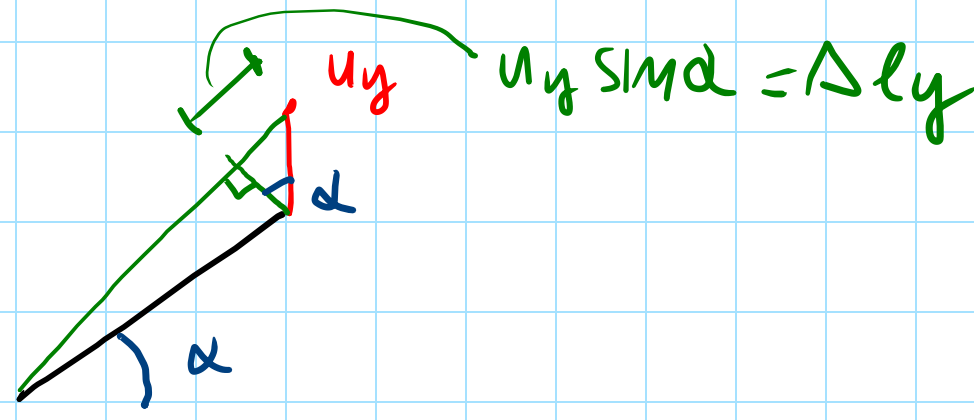
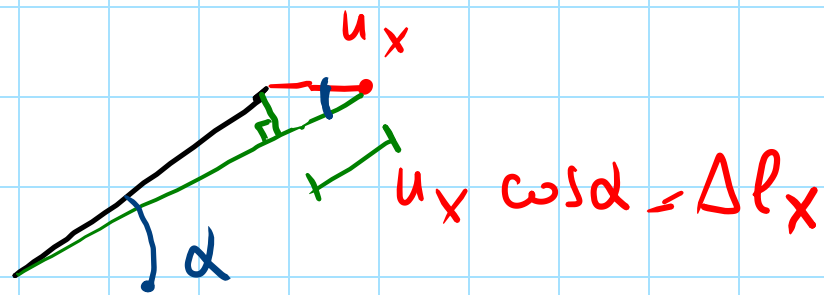
"SMALL DISPLACEMENTS", TAYLOR EXPANSION $l_f(u_x, u_y)$

$$l_f = l_f(0,0) + \left. \frac{\partial l_f}{\partial u_x} \right|_{(0,0)} u_x + \left. \frac{\partial l_f}{\partial u_y} \right|_{(0,0)} u_y + O(\sqrt{u_x^2 + u_y^2})$$

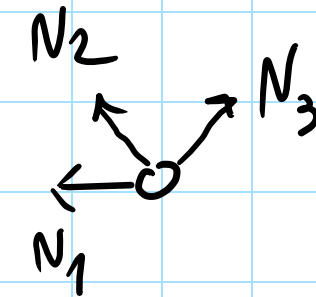
$$l_f \approx l + \left[\frac{1}{2\sqrt{}} \cdot 2(x_B + u_x) \right]_{(0,0)} u_x + \left[\frac{1}{2\sqrt{}} \cdot 2(y_B + u_y) \right]_{(0,0)} u_y$$

$$l_f \approx l + \underbrace{\frac{1}{l} x_B}_{\cos \alpha} u_x + \underbrace{\frac{y_B}{l}}_{\sin \alpha} u_y \Rightarrow l_f - l = \underbrace{\cos \alpha u_x}_{\Delta l_x} + \underbrace{\sin \alpha u_y}_{\Delta l_y}$$

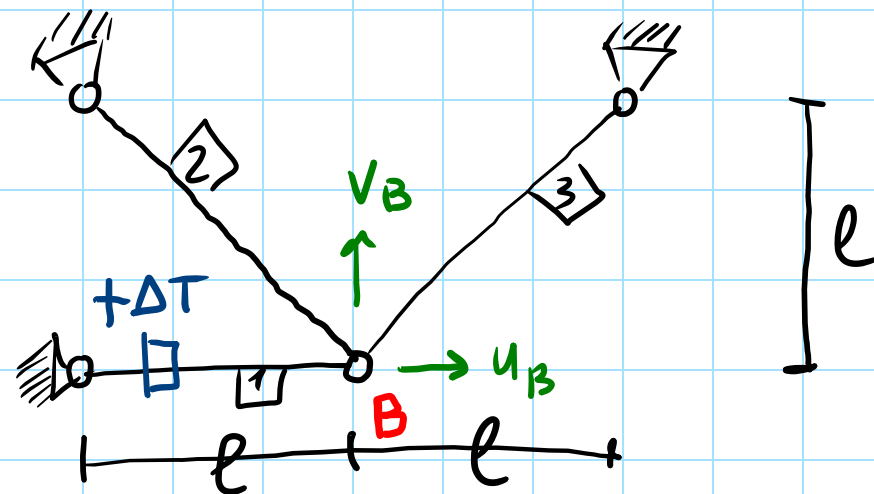
CHANGE IN
LENGTH:
SUM OF 2
CONTRIBUTIONS



$$\Delta l_i = \Delta l_i^{el} + \Delta l_i^{\Delta T}$$



EX



D.O.F.s: u_B, v_B

LOAD: $+\Delta T$ (BAR 1)

NODE B

$$\rightarrow: -N_1 - N_2 \frac{1}{\sqrt{2}} + \frac{N_3}{\sqrt{2}} = 0$$

$$\uparrow: +N_2 \frac{1}{\sqrt{2}} + N_3 \frac{1}{\sqrt{2}} = 0$$

2 DOFS? $\Rightarrow$ 2 EQUIL. EQUATIONS

WHOSE UNKNOWN ARE u_B, v_B

$$A_1 = A ; A_2 = A_3 = \sqrt{2} A$$

$$l_1 = l, l_2 = l_3 = \sqrt{2} l$$

E : const

CONSTITUTIVE LAW ($N_i = f(\Delta l_i)$)

$$N_1 = \frac{EA_1}{l_1} \Delta l_1^{el} = \frac{EA_1}{l_1} (\Delta l_1 - \Delta l_1^{\Delta T})$$

$$= \frac{EA_1}{l_1} (\Delta l_1 - \alpha \Delta T l)$$

$$N_2 = \frac{EA_2}{l_2} \Delta l_2^{(el)}$$

$$; N_3 = \frac{EA_3}{l_3} \Delta l_3^{(el)}$$

$$\Delta l_i = f_i(u_B, v_B)$$

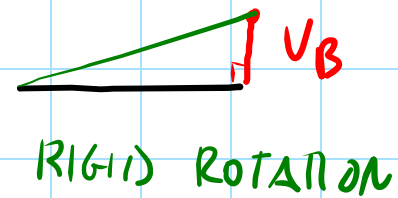
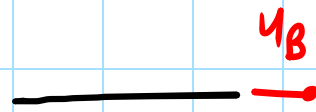
$$u_B, v_B \rightarrow \Delta l_i \rightarrow N_i \rightarrow \text{EQUIL. EQS}$$

u_B

v_B

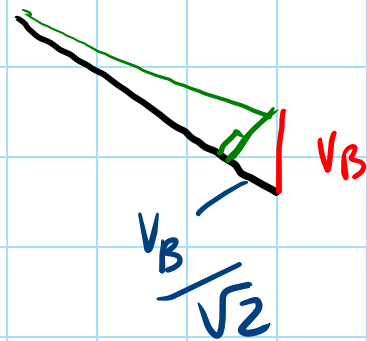
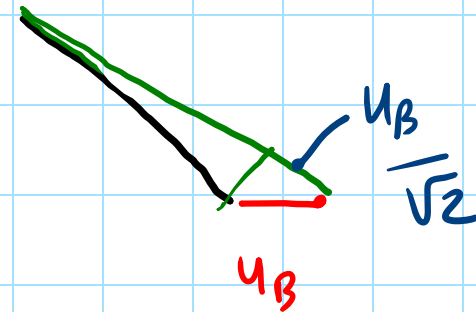
Δl_i

1)



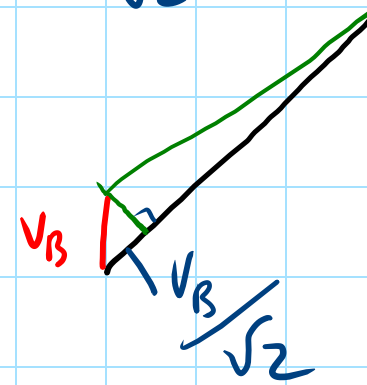
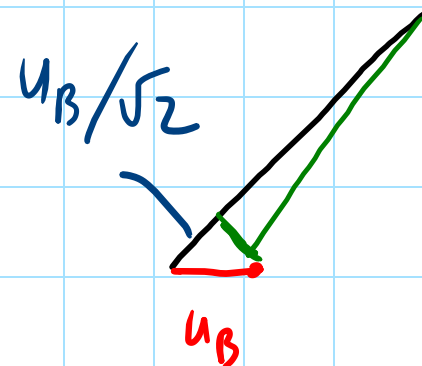
$$\Delta l_1 = +u_B$$

2)



$$\Delta l_2 = +\frac{u_B}{\sqrt{2}} - \frac{v_B}{\sqrt{2}}$$

3)



$$\Delta l_3 = -\frac{u_B}{\sqrt{2}} - \frac{v_B}{\sqrt{2}}$$

NOTE THAT $\frac{EA_1}{l_1} = \frac{EA}{l} = K$; $\frac{EA_2}{l_2} = \frac{EA_3}{l_3} = \frac{EA}{l} = K$

$$1) -k(\Delta l_1 - \alpha \Delta T l) - k \Delta l_2 \frac{1}{\sqrt{2}} + k \Delta l_3 \frac{1}{\sqrt{2}} = 0 \left\{ \right.$$

$$2) +k \Delta l_2 + k \Delta l_3 = 0$$

$$\begin{cases} -k(u_B - \underline{\alpha \Delta T l}) - k \frac{1}{2}(u_B - v_B) - k \frac{1}{2}(u_B + v_B) = 0 \\ k \frac{1}{\sqrt{2}}(u_B - v_B) - k \frac{1}{\sqrt{2}}(u_B + v_B) = 0 \end{cases}$$

LOAD

$$\Rightarrow v_B = 0$$

$$-k u_B - k \frac{1}{2} u_B - k \frac{1}{2} u_B = -k \alpha \Delta T l$$

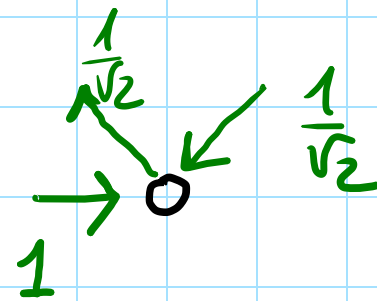
$$-2 u_B = -\alpha \Delta T l$$

$$u_B = \frac{\alpha \Delta T l}{2}$$

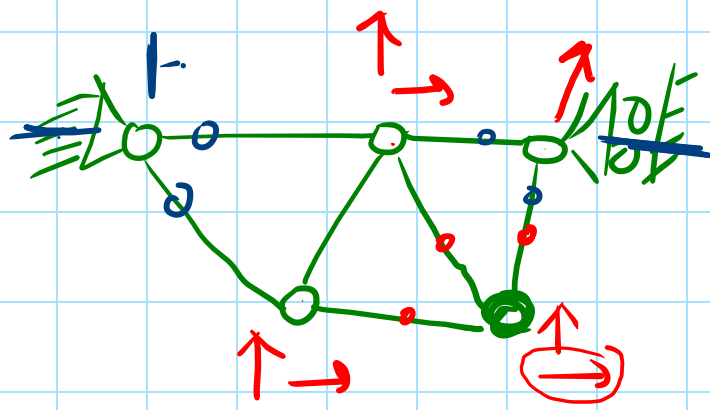
$$N_1 = k(u_B - \alpha \Delta T l) = -k \frac{\alpha \Delta T l}{2}$$

$$N_2 = k \frac{1}{\sqrt{2}} u_B = k \frac{1}{\sqrt{2}} \frac{\alpha \Delta T l}{2} > 0$$

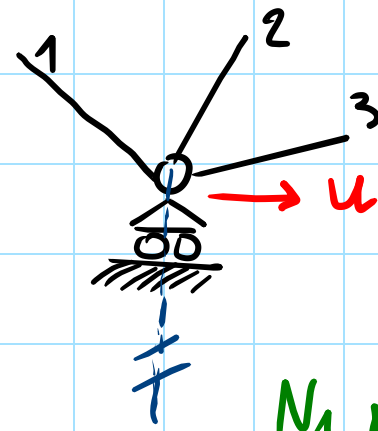
$$N_3 = k \frac{1}{\sqrt{2}} (-u_B) = -\frac{k \alpha \Delta T l}{2\sqrt{2}} < 0$$



$$\left(\times k \frac{\alpha \Delta T l}{2} \right)$$

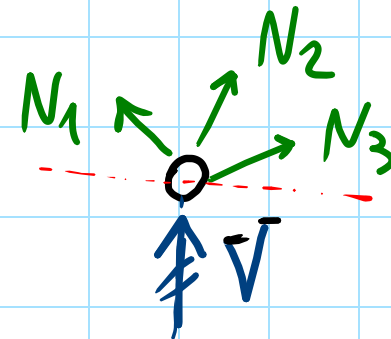


THE CASE OF A ROLLER



ONLY ONE
DOF: u

$\Rightarrow$ ONLY ONE
EQUIL. EQ. (ALONG u)



IN THE EQUIL. EQ
 V DOES NOT APPEAR

V WILL BE "SOLVED" LATER, AFTER HAVING COMPUTED N_1, N_2, N_3 , BY USING THE
'VERTICAL' EQUIL EQ. NOT USED BEFORE.

EX: IN THE STRUCTURE ABOVE THERE ARE $2m_N = 10$ EQUIL EQS

THERE ARE 7 DOFS (IN RED); THEN THE REMAINING 3 EQS ARE EMPLOYED
TO CALCULATE THE 3 REACTIONS

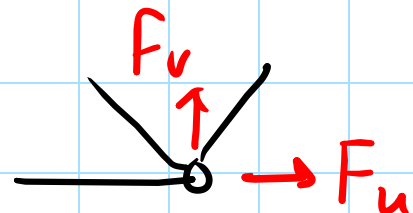
HOW TO OBTAIN $\tilde{K}$: FROM EQUIL. EQS WE SEE THAT THE MATRIX ASSOCIATED WITH THE PROBLEM IS

$$\begin{bmatrix} -2K & 0 \\ 0 & -\frac{2}{\sqrt{2}}K \end{bmatrix} \rightarrow \text{RECALL WE HAVE DROPPED } \frac{1}{\sqrt{2}} \text{ PREVIOUSLY}$$

I GUESS THAT THE "REAL" $\tilde{K}$ IS $\begin{bmatrix} 2K & 0 \\ 0 & K \end{bmatrix}$

NOTE THAT $\tilde{K}$ WILL SOLVE THE PROBLEM

$$\tilde{K} \begin{bmatrix} u_B \\ v_B \end{bmatrix} = \begin{bmatrix} F_u \\ F_v \end{bmatrix}$$



THE CHANGE IN SIGN IS DUE TO THIS FACT: IN BOTH EQS, IMAGINE TO HAVE FORCES F_u, F_v :

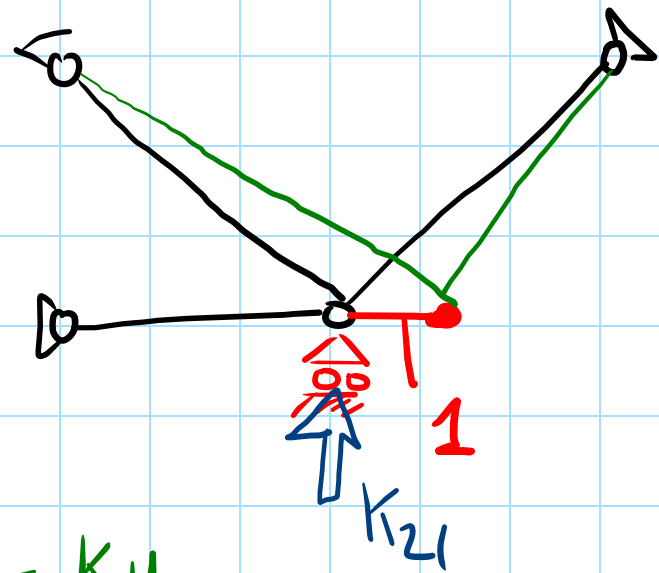
$$1) \left(\begin{smallmatrix} + \\ - \end{smallmatrix} \right) \dots \dots +F_u = 0 \rightsquigarrow \dots = -F_u$$

$$2) \left(\begin{smallmatrix} + \\ + \end{smallmatrix} \right) \dots \dots \frac{1}{\sqrt{2}} \dots \frac{1}{\sqrt{2}} + F_v = 0 \rightsquigarrow \frac{1}{\sqrt{2}} \dots \frac{1}{\sqrt{2}} = -F_v$$

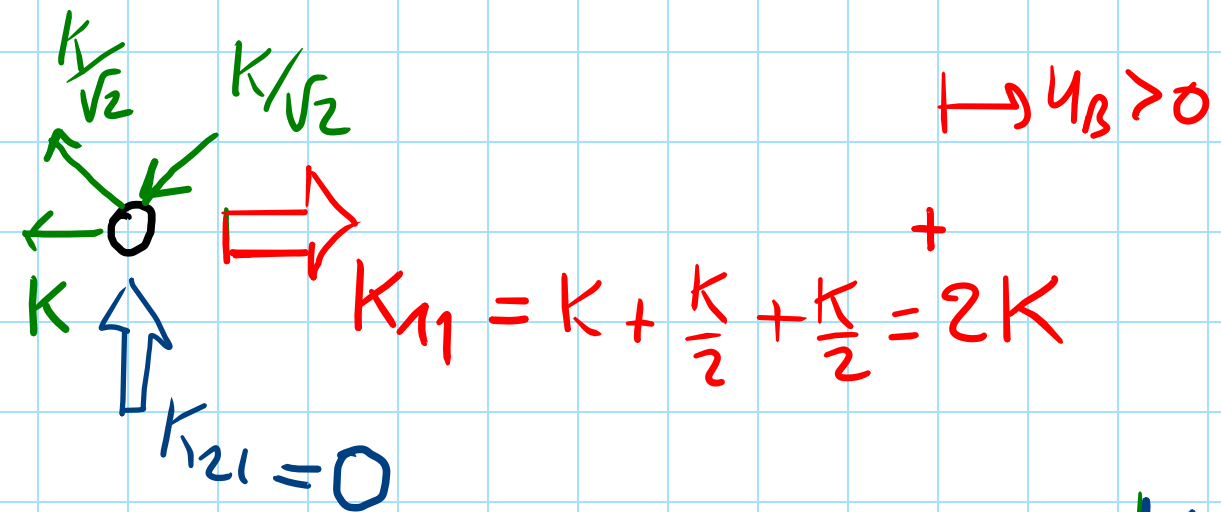
WE NEED TO CHANGE SIGN IN ORDER TO HAVE $+F_u$ AND $+F_v$
WE CANNOT NOW HAVE $\frac{1}{\sqrt{2}}!$

LET US TRY TO COMPUTE K_{11} WITH THE "PRACTICAL" METHOD

$u_B = 1$, $v_B = 0 \Rightarrow$ THE FORCE ALONG DOF u_B



$$N_1 = K u_B$$



BECAUSE THE GREEN FORCES ARE ALREADY IN EQUILIB. ALONG VERTICAL DIRECTION

$$0 = K_{21}$$

STRONG AND WEAK FORMULATIONS OF A DIFFERENTIAL PROBLEMS

CONSIDER THE PROBLEM OF THE ELASTIC LINE:

$$\left. \begin{aligned} (EJ v'')'' &= q \\ + \text{B. CONDITIONS} \end{aligned} \right\} \begin{array}{l} \text{STRONG} \\ \text{FORMULATION} \\ \text{DIFFERENTIAL} \\ \text{EQUATIONS + CONDITIONS} \end{array}$$

OTHER WEAK FORMULATIONS

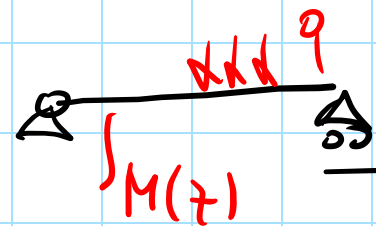
ARE:

STATIONARITY OF TOT. POT. ENERGY

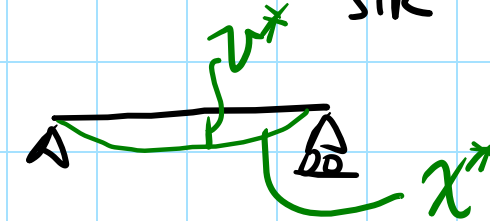
+ OTHER

WE CAN SOLVE ANY STRUCTURAL PROBLEM BY USING TH OF VIRTUAL WORK. FOR ELASTIC LINE:

$$\int_{\text{STR}} M(z) \chi^*(z) dz = \int_{\text{STR}} q(z) v^*(z) dz$$



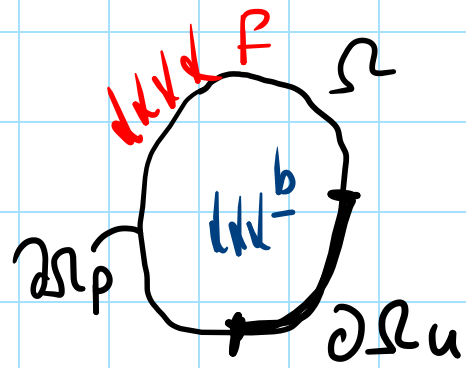
STAT. ADMISS
(REAL) SCH.



KINEM ADMISS SCH.
(FICTITIOUS)

EXAMPLE OF WEAK FORMULATION

IN 3D LINEAR ELASTICITY:



$$\left. \begin{aligned} \operatorname{div} \underline{\underline{\sigma}} + \underline{\underline{b}} &= \underline{\underline{0}} \\ \underline{\underline{\varepsilon}} &= \frac{1}{2} (\nabla \underline{\underline{u}} + \nabla \underline{\underline{u}}^T) \\ \underline{\underline{\sigma}} &= \mathbb{C} \underline{\underline{\varepsilon}} \end{aligned} \right\} \text{in } \Omega$$

$$\left. \begin{aligned} \underline{\underline{u}} &= \underline{\underline{u}}^0 \text{ on } \partial\Omega_u \\ \underline{\underline{\sigma}} \underline{\underline{n}} &= \underline{\underline{p}} \text{ on } \partial\Omega_p \end{aligned} \right\} \text{B.C.'s.}$$

DISPL. FORMULATION: ($\underline{\underline{u}}$ IS THE UNKNOWN)

$$\operatorname{div} \left[\frac{1}{2} \mathbb{C} (\nabla \underline{\underline{u}} + \nabla \underline{\underline{u}}^T) \right] + \underline{\underline{b}} = \underline{\underline{0}} \quad (\Omega)$$

$$\left. \begin{aligned} \underline{\underline{u}} &= \underline{\underline{u}}^0 \\ \frac{1}{2} \mathbb{C} (\nabla \underline{\underline{u}} + \nabla \underline{\underline{u}}^T) \underline{\underline{n}} &= \underline{\underline{p}} \end{aligned} \right\} \text{B.C.'s}$$

STRONG FORM

WEAK FORM

TH. VIRTUAL WORK

$$\begin{aligned} \int_{\Omega} \underline{\underline{\sigma}} \cdot \underline{\underline{\varepsilon}}^* dV &= \int_{\partial\Omega_p} \underline{\underline{p}} \cdot \underline{\underline{u}}^* dS \\ &+ \int_{\Omega} \underline{\underline{b}} \cdot \underline{\underline{u}}^* dV \\ &+ \int_{\partial\Omega_u} \underline{\underline{u}} \cdot \underline{\underline{u}}^0 dS \end{aligned}$$

WHERE $*$ DENOTES AN INDEPENDENT KINEMATIC FIELD.