## Modes of mass transport in electrochemistry

Mass transport from one region in solution to another, arising from:

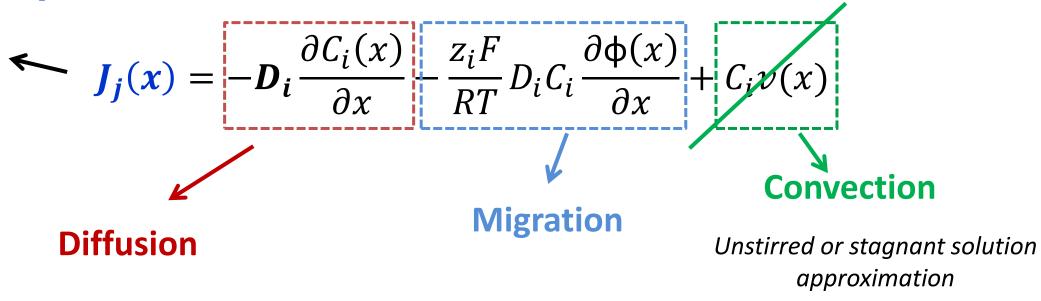
- Convection (net transport of a volume element of solution, e.g. stirring or hydrodynamics), which can be *natural* (presence of <u>density gradient in</u> <u>solution</u>) or *forced* (hydrodynamic techniques)
- **Migration**, which describes the transport of charges in solution under the effect of an electrical field (gradient of electrical potential)
- Diffusion, which describes the transport of a species in solution under the effect of a gradient of chemical potential (or concentration)
- Migration and diffusion arise from gradients of electrochemical potential  $\overline{\mu}_i$

## Mass transport in solution: the flux (J)

The presence of gradients in solution is the driving force for mass transfer, which is described in terms of a **flux**, *J*, *which operates to alleviate the magnitude of the gradients* 

Nernst-Planck Equation (for linear mass transfer)

 $[mol cm^{-2} s^{-1}]$ 



# The effect of gradients in electrochemical potential

The flux is proportional to the gradient of electrochemical potential (*i* = charged species)

$$J_i \propto \nabla \overline{\mu_i} \qquad \nabla = \frac{\partial}{\partial x} \qquad Vector \ operator \ (gradient) \\ for \ linear \ mass \ transfer$$

- The flux is equivalent to a current density (*i* = charged species)
- <u>Large concentration of supporting electrolyte</u>: the effect of migration on the electroactive species *i* can be neglected

$$\boldsymbol{J_i(x)} = -\boldsymbol{D_i} \frac{\partial C_i(x)}{\partial x} - \frac{z_i F}{RT} D_i C_t \frac{\partial \phi(x)}{\partial x}$$



Fick's laws: Differential equations describing the flux and the concentration gradient of a species as functions of time (t) and position (x)

 $\frac{\partial C_i(x,t)}{\partial x}$ 

Assumption: linear (1D) diffusion

Fick's 1<sup>st</sup> law

A. Fick (1829 - 1901)

Net diffusive mass transport of a species i, described as number of moles passing through a unit area per unit time (mol s<sup>-1</sup> cm<sup>-2</sup>) Diffusion coefficient Constant term [cm<sup>2</sup>/s]

 $-J_i(x,t) = \boldsymbol{D_i}$ 

*Concentration gradient The flux is proportional to the concentration gradient* 



*dx* = *step length* Number of moles  $J_{i}(x,t) = \frac{1}{A} \cdot \frac{\frac{N_{i}(x)}{2} - \frac{N_{i}(x+dx)}{2}}{dt} \cdot \frac{dx^{2}}{dx^{2}}$  $N_{i}(x) \qquad N_{i}(x + dx)$   $\leftarrow \qquad N_{i}(x + dx)$   $\frac{N_{i}(x + dx)}{2}$   $\frac{N_{i}(x)}{2}$  $J_i(x,t) = \frac{dx^2}{2dt} \cdot \frac{N_i(x) - N_i(x + dx)}{Adx} dx$ X **Concentration**  $J_i(x,t) = \frac{dx^2}{2dt} \frac{C_i(x) - C_i(x+dx)}{dx}$ 1° Fick's law

constant term

**Diffusion coefficient D**<sub>i</sub>



A. Fick (1829 - 1901)

- The diffusion coefficient D reflects the intrinsic tendency of a molecule to diffuse and depends on the molecule size (bigger molecules have typically smaller D)
- Molecular species typically have  $D \approx 10^{-6} 10^{-5} \text{ cm}^2/\text{s}$
- Determines how far a species diffuses in a certain time

$$\sqrt{\langle x^2 \rangle} = \sqrt{2Dt}$$

root-mean-square displacement at time t

• Provides an estimation for the **thickness of the diffusion layer** 

How thick is a diffusion layer?



A. Fick (1829 - 1901) Area (A)  $J_i(x + dx, t)$  $J_i(x,t)$ x+dx

Fick's 2<sup>nd</sup> law

$$\frac{\partial C_i(x,t)}{\partial t} = \boldsymbol{D}_i \left( \frac{\partial^2 C_i(x,t)}{\partial x^2} \right)$$

Gradient of concentration of the species i with time

$$\frac{\partial C_i(x,t)}{\partial t} = \frac{J_i(x,t) - J_i(x+dx,t)}{dx}$$
$$J_i(x+dx,t) = J_i(x,t) + \left(\frac{\partial J_i(x,t)}{\partial x}\right) dx$$
$$\frac{\partial C_i(x,t)}{\partial t} = -\left(\frac{\partial J_i(x,t)}{\partial x}\right)$$
from Fick's 1<sup>st</sup> law... 
$$\frac{\partial C_i(x,t)}{\partial t} = \left(\frac{\partial}{\partial x}\right) \left[ D_i \left(\frac{\partial C_i(x,t)}{\partial x}\right) \right]$$

 $\partial x$ 

∂t

$$\frac{\partial C_i(x,t)}{\partial t} = D_i \left( \frac{\partial^2 C_i(x,t)}{\partial x^2} \right)$$

Linear diffusion equation

$$\frac{\partial C_i}{\partial t} = D_i \nabla^2 C_i$$

General Fick's 2<sup>nd</sup> law for any geometry

Laplacian operator

Туре	Variables	$\nabla^2$	Example
Linear	x	$\partial^2/\partial x^2$	Shielded disk electrode
Spherical	r	$\partial^2/\partial r^2 + (2/r)(\partial/\partial r)$	Hanging drop electrode
Cylindrical (axial)	r	$\partial^2/\partial r^2 + (1/r)(\partial/\partial r)$	Wire electrode
Disk	r, z	$\partial^2/\partial r^2 + (1/r)(\partial/\partial r) + \partial^2/\partial z^2$	Inlaid disk ultramicroelectrode <sup>b</sup>
Band	<i>x</i> , <i>z</i>	$\partial^2/\partial x^2 + \partial^2/\partial z^2$	Inlaid band electrode <sup>c</sup>

 TABLE 4.4.2
 Forms of the Laplacian Operator for Different Geometries<sup>a</sup>

Bard & Faulkner, Electrochemical Methods, 2<sup>nd</sup> Ed, 2001

## The electrochemical problem

Purely diffusive electroactive species

$$Ox + ne^- \rightleftharpoons Red$$

$$J_{Ox}(0,t) = -J_{Red}(0,t) \qquad -J_{Ox}(0,t) = \frac{i}{nFA} = D_{Ox} \left( \frac{\partial C_{Ox}(x,t)}{\partial x} \right)_{x=0} \qquad \begin{array}{c} \text{Concentration profile} \\ \text{at the electrode} \\ \text{surface} \end{array}$$

$$\frac{i}{FA} = -\sum_{i} n_i \cdot J_i(0, t) = \sum_{i} n_i \cdot D_i \left[ \frac{\partial C_i(x, t)}{\partial x} \right]_{x=0}$$

#### **Electrochemical problem**

Solving the mass-transfer equations (diffusion)



Relationship describing the **concentrations** of the species as functions of the variables **x** and **t** 

## How to solve partial differential equations?

$$\frac{\partial C(x,t)}{\partial t} = D\left(\frac{\partial^2 C(x,t)}{\partial x^2}\right)$$

#### Partial differential equation (PDE)

- Concentration of the species are **functions of both distance (x) and time (t)**
- Ordinary differential equations (ODEs) contain derivatives of functions of a single variable only
- PDEs are often characterized by **multiple solutions**
- The specific solution of a PDE requires the definition of the **boundary conditions** for <u>each</u> <u>species</u> involved in the redox process
- Different boundary conditions usually lead to different solutions of PDE

# The boundary conditions

**1. Initial conditions** (concentration profile at **t** = **0**)

 $C_{Ox}(x,0) = f(x)$  Es.  $C_{Ox}(x,0) = C_{Ox}^*$   $C_{Red}(x,0) = 0$ 

**2. Semi-infinite boundary conditions** (approximation related to the volume of the cell reactors compared to the size of the diffusion layer)

 $\lim_{x \to \infty} C_{Ox}(x,t) = C^*_{Ox} \qquad \qquad \lim_{x \to \infty} C_{Red}(x,t) = 0$ 

electrochemical cells are usually (not always!) much bigger than the diffusion layer

**3. Electrode surface boundary conditions** (related to the concentrations of the species (or their gradients) at the electrode surface, x = 0)

This condition is specific for the particular electrochemical experiment considered

Potential-controlled experiments

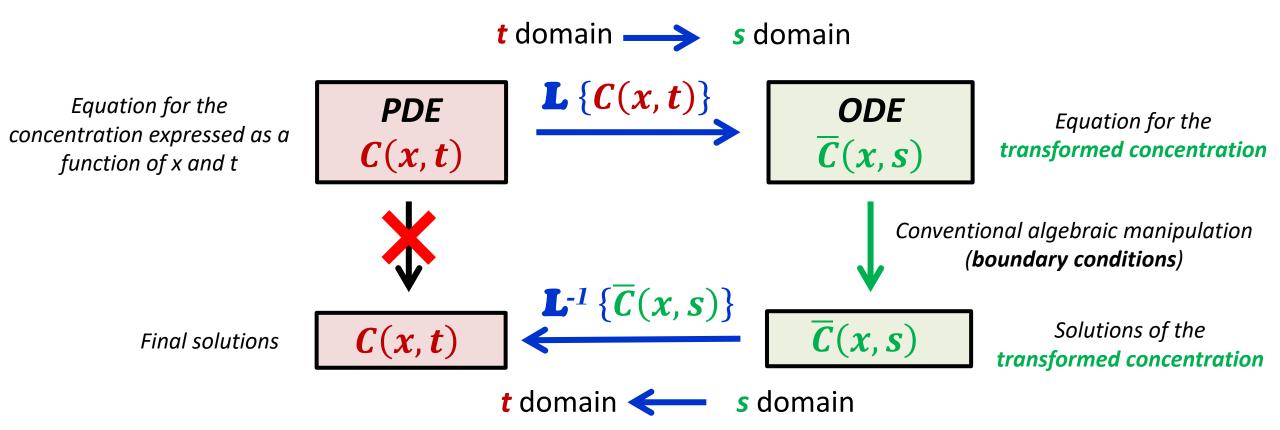
$$C_{Ox}(0,t) = f(E)$$

$$\frac{C_{Ox}(0,t)}{C_{Red}(0,t)} = f(E)$$

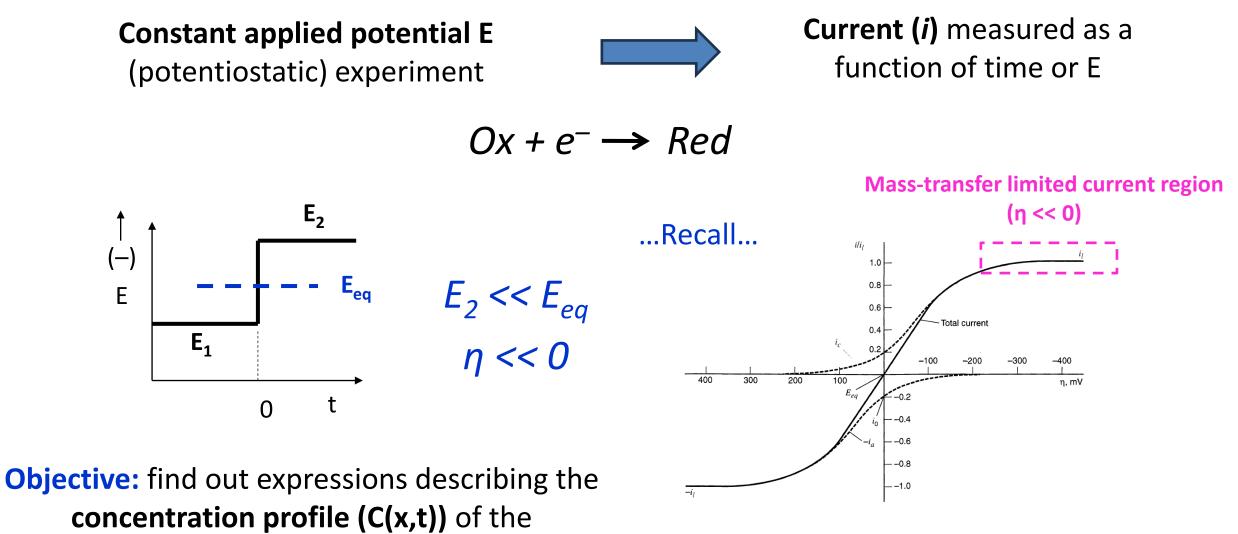
## The Laplace transformation

$$\mathbf{L}\{\mathbf{F}(t)\} = \overline{\mathbf{F}}(s) \equiv \int_0^\infty e^{-st} F(t) dt$$

Laplace transformation is a mathematical tool, enabling a conversion of the problem into a domain where simpler mathematical manipulation is possible



#### **Step-potential experiment**



electroactive species and the current vs. time

Bard & Faulkner, 2<sup>nd</sup> Ed., Wiley, 2001

## Step-potential experiment

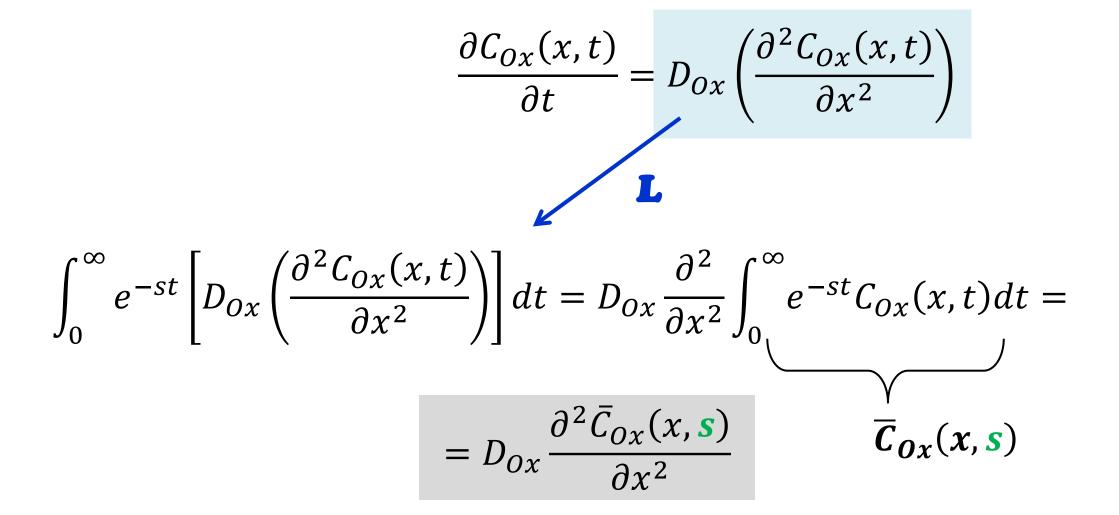
- Planar electrode
- Unstirred solution
- Only diffusive mass transport (linear diffusion)
- <u>Initial condition</u>  $C_{Ox}(x,0) = C_{Ox}^*$  for t = 0
- Microelectrode (reaction volume is assumed to be so high and the area of the electrode so small that the current flowing does not significantly affect the bulk concentration of the electroactive species)
- Semi-infinite condition

$$\lim_{x\to\infty} C_{Ox}(x,t) = C_{Ox}^*$$

<u>Surface concentration condition</u> (very large η, diffusion-limited current regardless of the ET kinetics)

3<sup>rd</sup> boundary condition (specific step-E experiment)

$$C_{Ox}(0,t) = 0 \quad \text{for } t > 0$$



$$\frac{\partial C_{Ox}(x,t)}{\partial t} = D_{Ox} \left( \frac{\partial^2 C_{Ox}(x,t)}{\partial x^2} \right)$$

$$\int_{0}^{\infty} e^{-st} \frac{\partial C_{Ox}(x,t)}{\partial t} dt = [e^{-st} C_{Ox}(x,t)]_{0}^{\infty} - \int_{0}^{\infty} C_{Ox}(x,t)(-se^{-st}) dt$$
Integration by parts
$$= -C_{Ox}(x,0) + s \int_{0}^{\infty} C_{Ox}(x,t)(e^{-st}) dt$$

$$= -C_{Ox}(x,0) + s \int_{0}^{\infty} C_{Ox}(x,t)(e^{-st}) dt$$

$$\int_{a}^{b} g(x)f'(x) dx = [g(x)f(x)]_{a}^{b} - \int_{a}^{b} f(x)g'(x) dx$$
Intial boundary
Intial boundary

condition

$$-C_{Ox}^* + s\overline{C}_{Ox}(x, s) = D_{Ox}\left(\frac{\partial^2 \overline{C}_{Ox}(x, s)}{\partial x^2}\right)$$

$$\frac{\partial^2 \overline{C}_{Ox}(x, s)}{\partial x^2} - \frac{s}{D_{Ox}} \overline{C}_{Ox}(x, s) + \frac{C_{Ox}^*}{D_{Ox}} = 0$$

$$\overline{C}_{Ox}(x,s) = \frac{C_{Ox}^*}{s} + A'(s)e^{\left[-\sqrt{\frac{s}{D_{Ox}}x}\right]} + B'(s)e^{\left[\sqrt{\frac{s}{D_{Ox}}x}\right]}$$
General solution

$$\overline{\boldsymbol{C}}_{\boldsymbol{O}\boldsymbol{X}}(\boldsymbol{x},\boldsymbol{s}) = \frac{C_{\boldsymbol{O}\boldsymbol{X}}^*}{s} + A'(s)e^{\left[-\sqrt{\frac{s}{D_{\boldsymbol{O}\boldsymbol{X}}}}\boldsymbol{x}\right]} + B'(s)e^{\left[\sqrt{\frac{s}{D_{\boldsymbol{O}\boldsymbol{X}}}}\boldsymbol{x}\right]}$$

$$\lim_{x\to\infty} C_{Ox}(x,t) = C_{Ox}^*$$

1

$$\lim_{x\to\infty} \overline{C}_{ox}(x,s) = \frac{C_{ox}^*}{s}$$

Semi-infinite boundary

condition

B'(s) should be equal to 0

$$\overline{\boldsymbol{C}}_{\boldsymbol{O}\boldsymbol{X}}(\boldsymbol{x},\boldsymbol{s}) = \frac{C_{\boldsymbol{O}\boldsymbol{X}}^*}{s} + A'(s)e^{\left[-\sqrt{\frac{s}{D_{\boldsymbol{O}\boldsymbol{X}}}}\boldsymbol{x}\right]}$$

$$\overline{C}_{Ox}(x, s) = \frac{C_{Ox}^*}{s} + A'(s)e^{\left[-\sqrt{\frac{s}{D_{Ox}}}x\right]}$$

$$= 0 \quad \text{for } \begin{array}{c} x = 0 \\ t > 0 \end{array} \quad \overrightarrow{C}_{Ox}(0, s) = 0 \end{array}$$

$$C_{Ox}(0,t) = 0 \quad \text{fo}$$

3<sup>rd</sup> boundary condition

(surface concentration)

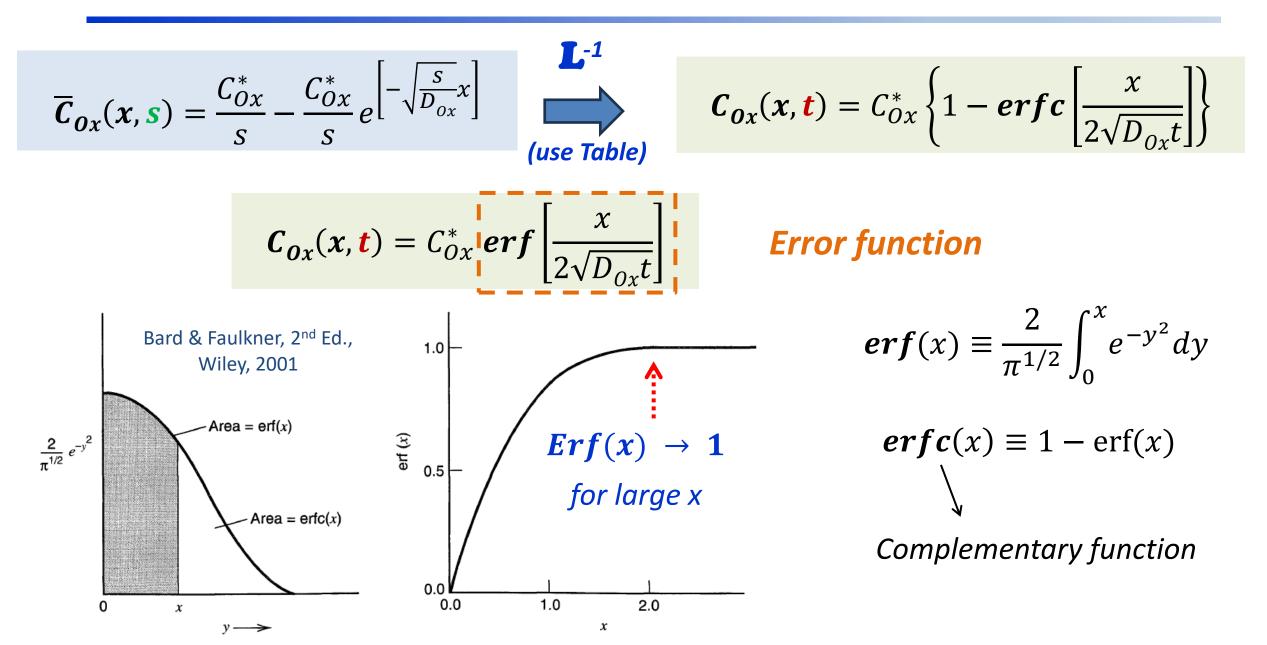
$$O = \frac{C_{Ox}^*}{s} + A'(s)e^{\left[-\sqrt{\frac{s}{D_{Ox}}}x\right]} \qquad A'(s) = -\frac{C_{Ox}^*}{s}$$

$$x = 0$$

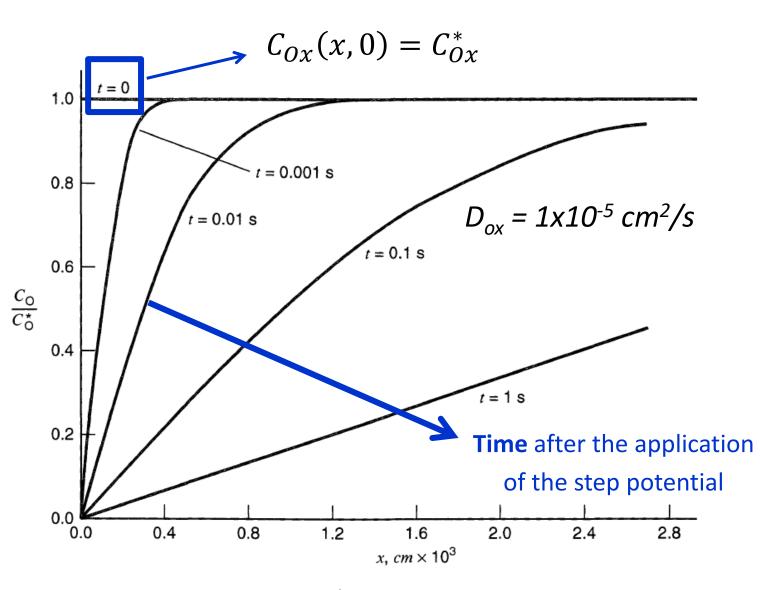
 $\overline{\boldsymbol{C}}_{\boldsymbol{O}\boldsymbol{X}}(\boldsymbol{0},\boldsymbol{s})=0$ 

$$\overline{C}_{Ox}(x,s) = \frac{C_{Ox}^*}{s} - \frac{C_{Ox}^*}{s} e^{\left[-\sqrt{\frac{s}{D_{Ox}}}x\right]}$$

#### The concentration profiles



## The concentration profiles

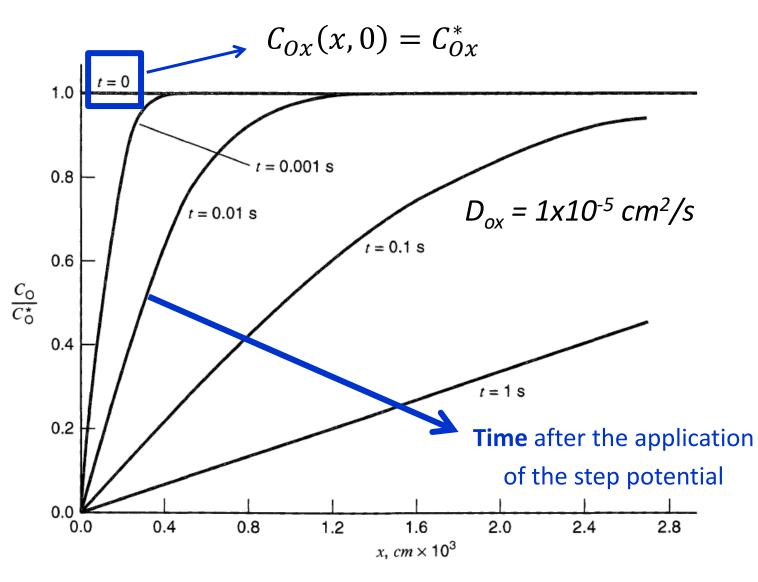


 $\boldsymbol{C}_{\boldsymbol{O}\boldsymbol{X}}(\boldsymbol{x},\boldsymbol{t}) = C_{\boldsymbol{O}\boldsymbol{X}}^* \boldsymbol{erf} \left| \frac{\boldsymbol{x}}{2\sqrt{D_{\boldsymbol{O}\boldsymbol{x}}t}} \right|$ 

- Time-dependent concentration profiles near the electrode surface (as a function of the distance)
- Ox depletion at the surface with increasing times
- The **slope** of each concentration profile indicates the Ox **concentration gradient**  $(\partial C_{Ox}(x,t)/\partial x)$  at various t

Bard & Faulkner, 2<sup>nd</sup> Ed., Wiley, 2001

## The concentration profiles



Bard & Faulkner, 2<sup>nd</sup> Ed., Wiley, 2001

$$\boldsymbol{C}_{\boldsymbol{0}\boldsymbol{x}}(\boldsymbol{x},\boldsymbol{t}) = C_{\boldsymbol{0}\boldsymbol{x}}^* \, \boldsymbol{erf} \left[ \frac{\boldsymbol{x}}{2\sqrt{D_{\boldsymbol{0}\boldsymbol{x}}\boldsymbol{t}}} \right]$$

Profiles asymptotically tend to
 C<sub>ox</sub>\* value

$$\frac{\partial C_{Ox}(x,t)}{\partial x} = 0 \quad \Longrightarrow \quad C_{Ox}(x,t) = C_{Ox}^*$$

• Time-dependent thickness of the diffusion layer, which can be defined in terms of  $\sqrt{D_{0x}t}$  ( $\approx 6\sqrt{D_{0x}t}$ )

## The Cottrell Equation

$$-J_{0x}(\mathbf{0},t) = \frac{i}{nFA} = D_{0x} \left( \frac{\partial C_{0x}(x,t)}{\partial x} \right)_{x=\mathbf{0}} \qquad \overline{C}_{0x}(x,s) = \frac{C_{0x}^*}{s} - \frac{C_{0x}^*}{s} e^{\left[ -\sqrt{D_{0x}}x \right]}$$
$$\overline{L} \qquad \overline{L} \qquad$$

$$\frac{\boldsymbol{u}(\boldsymbol{s})}{nFA} = D_{OX} \left( \frac{\partial \boldsymbol{\mathcal{L}}_{OX}(\boldsymbol{x}, \boldsymbol{s})}{\partial \boldsymbol{x}} \right)_{\boldsymbol{x}=\boldsymbol{0}} = D_{OX} \left[ -\frac{\mathcal{L}_{OX}}{s} \left( -\frac{s^{1/2}}{D_{OX}^{1/2}} \right) e^{\left[ -\sqrt{D_{OX}} \boldsymbol{x} \right]} \right]_{\boldsymbol{x}=\boldsymbol{0}} = D_{OX} \left( \frac{\mathcal{L}_{OX}}{s^{1/2} D_{OX}^{1/2}} \right)$$

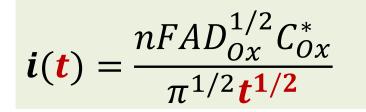
$$\bar{\iota}(s) = \frac{nFAD_{0x}^{1/2}C_{0x}^*}{s^{1/2}} \longrightarrow i(t) = \frac{nFAD_{0x}^{1/2}C_{0x}^*}{\pi^{1/2}t^{1/2}}$$

#### **Cottrell Equation**

## The Cottrell Equation



F. G. Cottrell (1877 - 1948)



- Effect of depletion of Ox near the surface
- Faradaic current inversely proportional to t<sup>1/2</sup> (i decrease with time)
- The linear dependence of i vs. 1/t<sup>1/2</sup> is diagnostic test for diffusion-controlled electrochemical process
- Cottrell behaviour not observed at very short times due to capacitive or nonfaradaic current ("charging" current),

 $i_C \propto 1/\exp(t)$ 

#### Chronoamperometry

 $i(t) = \frac{nFAD_{0x}^{1/2}C_{0x}^*}{\pi^{1/2}t^{1/2}}$ 

Diffusion-limited current at E<sub>2</sub>

