



**993SM - Laboratory of
Computational Physics
VIII week
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Numerical integration - I

M. Peressi - UniTS - Laurea Magistrale in Physics
Laboratory of Computational Physics - Unit VII

- **deterministic methods in ID**
equispaced points (trapezoidal, Simpson...),
others...
- **Monte Carlo methods**
(acceptance-rejection, sample mean,
importance sampling...)

Error handling:

sample mean

block average

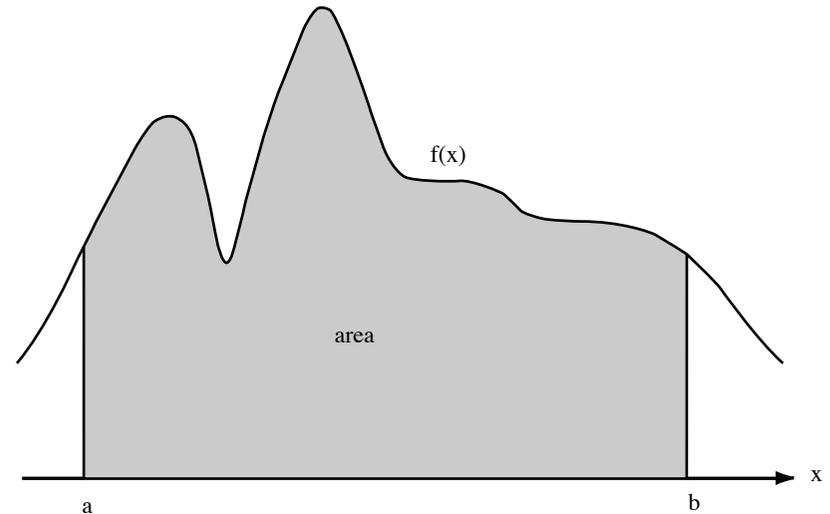
reduction of the variance

Deterministic methods

Deterministic methods

Start from the geometrical interpretation of a definite integral:

$$F = \int_a^b f(x) dx$$



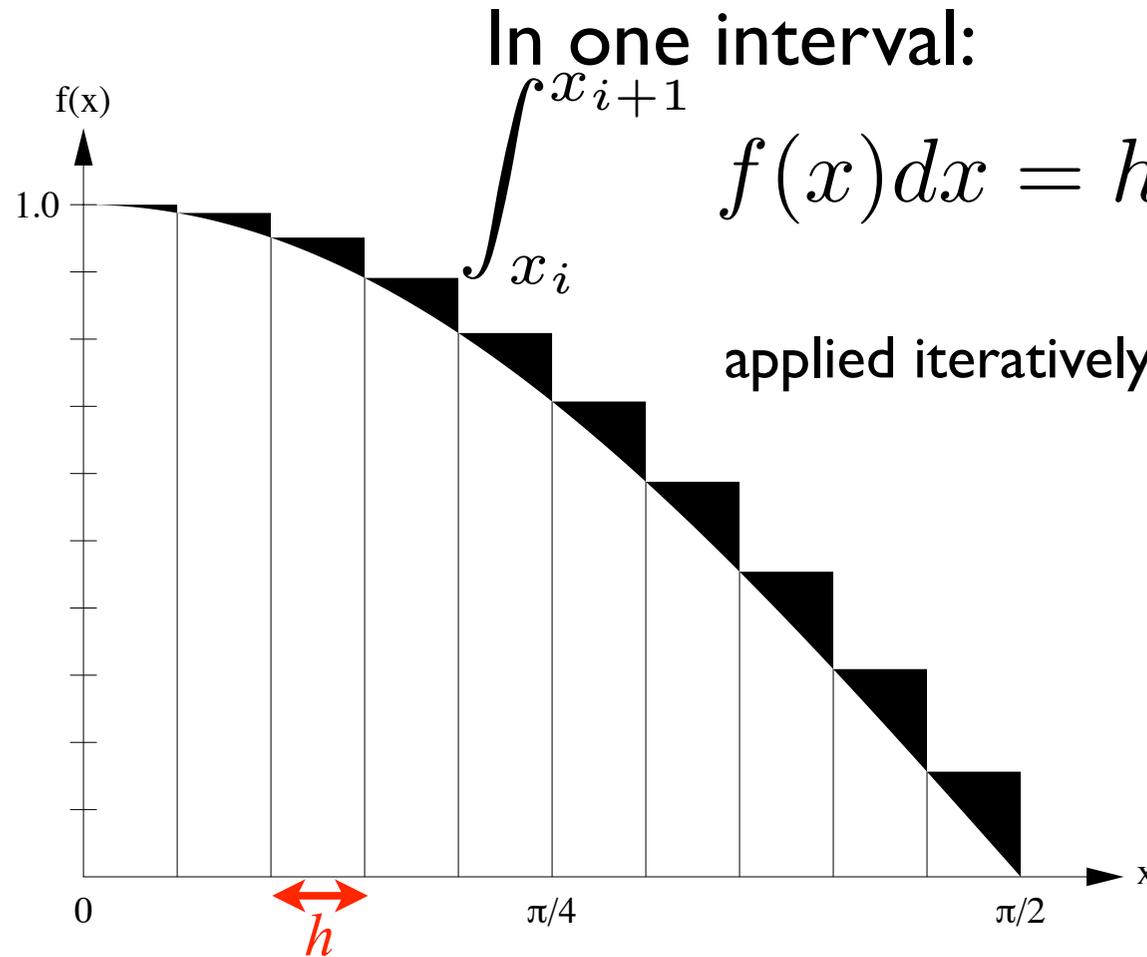
Divide the integration interval into “small” intervals:

$$\Delta x = \frac{b - a}{n},$$

$$x_n = x_0 + n \Delta x.$$

(Note: n intervals $\Leftrightarrow n + 1$ points)

Deterministic methods: rectangular rule



In one interval:

$$\int_{x_i}^{x_{i+1}} f(x) dx = h f_i$$

with error:

$$\mathcal{O}(h^2 f'), \propto 1/n^2$$

applied iteratively over consecutive intervals:

$$F_n = \sum_{i=0}^{n-1} f(x_i) \Delta x.$$

with a total error:

$$\mathcal{O}(h f'), \propto 1/n$$

: The rectangular approximation for $f(x) = \cos x$ for $0 \leq x \leq \pi/2$.

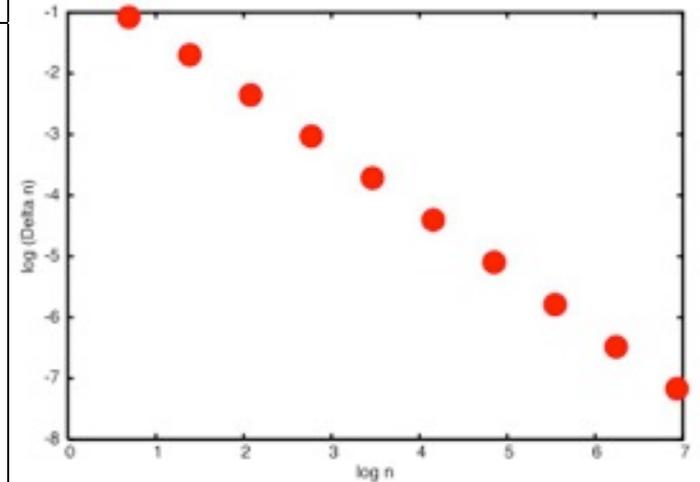
Deterministic methods: rectangular rule - error

$$I = \int_0^{\pi/2} \cos(x) dx = 1$$

$$F_n = \frac{\pi}{2n} \sum_0^{n-1} \cos x_i; \quad x_i = i \frac{\pi}{2n}$$

$$\Delta_n = F_n - I$$

n	F_n	Δ_n
2	1.34076	0.34076
4	1.18347	0.18347
8	1.09496	0.09496
16	1.04828	0.04828
32	1.02434	0.02434
64	1.01222	0.01222
128	1.00612	0.00612
256	1.00306	0.00306
512	1.00153	0.00153
1024	1.00077	0.00077



Rectangular approximations of the integral of $\cos x$ from $x = 0$ to $x = \pi/2$ as a function of n , the number of intervals. The error Δ_n is the difference between the rectangular approximation and the exact result of unity. Note that the error Δ_n decreases approximately as n^{-1} , that is, if n is increased by a factor of 2, Δ_n decreases by a factor 2.

Deterministic methods: generalities

- sum values of $f(x_i)$ with $x_i \in [a, b]$
- we want to have $F = \int_a^b f(x)dx$
as accurate as possible but with the
minimum number of calculations of $f(x_i)$

OK simple algorithms, but if the number of calculations is too high, improve the algorithm!

Deterministic methods: trapezoidal rule

In one interval:

with error:

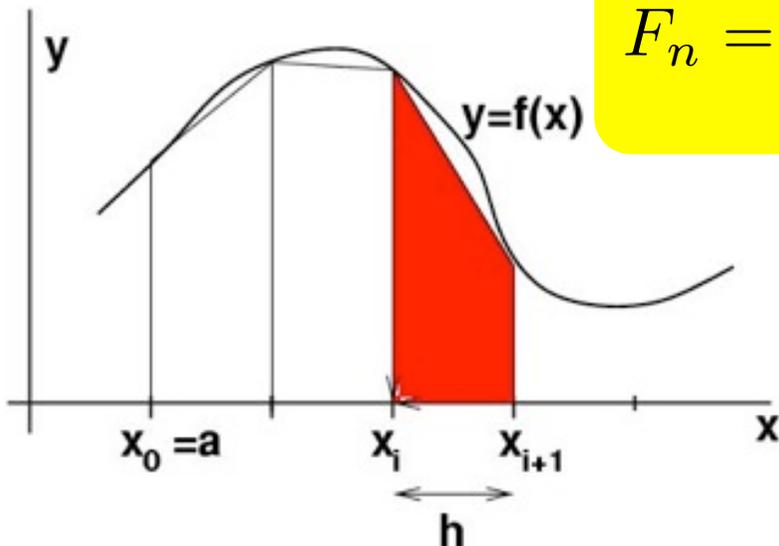
$$\int_{x_i}^{x_{i+1}} f(x) dx = h \left[\frac{1}{2} f_i + \frac{1}{2} f_{i+1} \right] \quad \mathcal{O}(h^3 f'''), \propto 1/n^3$$

Applied iteratively over consecutive intervals:

$$F_n = \left[\frac{1}{2} f(x_0) + \sum_{i=1}^{n-1} f(x_i) + \frac{1}{2} f(x_n) \right] \Delta x.$$

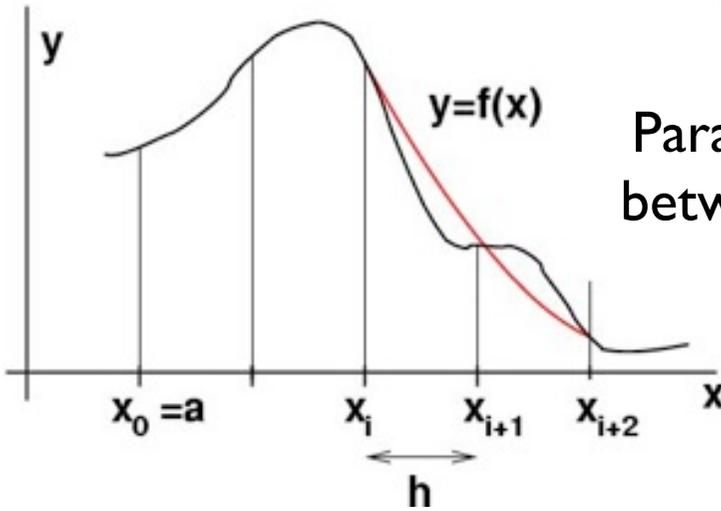
with a total error:

$$\mathcal{O}(h^2 f'''), \propto 1/n^2$$



Deterministic methods:

Simpson's rule



Parabolic interpolation procedure
between triplets of adjacent points

In one interval:

$$\int_{x_i}^{x_{i+2}} f(x) dx = h \left[\frac{1}{3} f_i + \frac{4}{3} f_{i+1} + \frac{1}{3} f_{i+2} \right] + \mathcal{O}(h^5 f^{IV}) \quad (\text{error} \propto 1/n^5)$$

Iteratively applied to the whole interval of integration (**odd** number of points!):

$$\int_{x_0}^{x_n} f(x) dx = h \left[\frac{1}{3} f_0 + \frac{4}{3} f_1 + \frac{2}{3} f_2 + \frac{4}{3} f_3 + \dots + \frac{2}{3} f_{n-2} + \frac{4}{3} f_{n-1} + \frac{1}{3} f_n \right] + \mathcal{O}(h^4 f^{IV}) \quad (\text{error} \propto 1/n^4)$$

Errors in deterministic methods

Error estimate for numerical integration with deterministic methods

$$\int f(x)dx = F_n + error$$

How to evaluate the error? Consider the Taylor expansion of the integrand function and then integrate:

$$f(x) = f(x_i) + f'(x_i)(x - x_i) + \frac{1}{2}f''(x_i)(x - x_i)^2 + \dots, (*)$$

$$\int_{x_i}^{x_{i+1}} f(x) dx = f(x_i)\Delta x + \frac{1}{2}f'(x_i)(\Delta x)^2 + \frac{1}{6}f''(x_i)(\Delta x)^3 + \dots(**)$$

$$\Delta x \equiv x_{i+1} - x_i$$

Error estimate for numerical integration: Rectangular approximation

$$\int_{x_i}^{x_{i+1}} f(x) dx \approx f(x_i) \Delta x$$

Compare with (**):

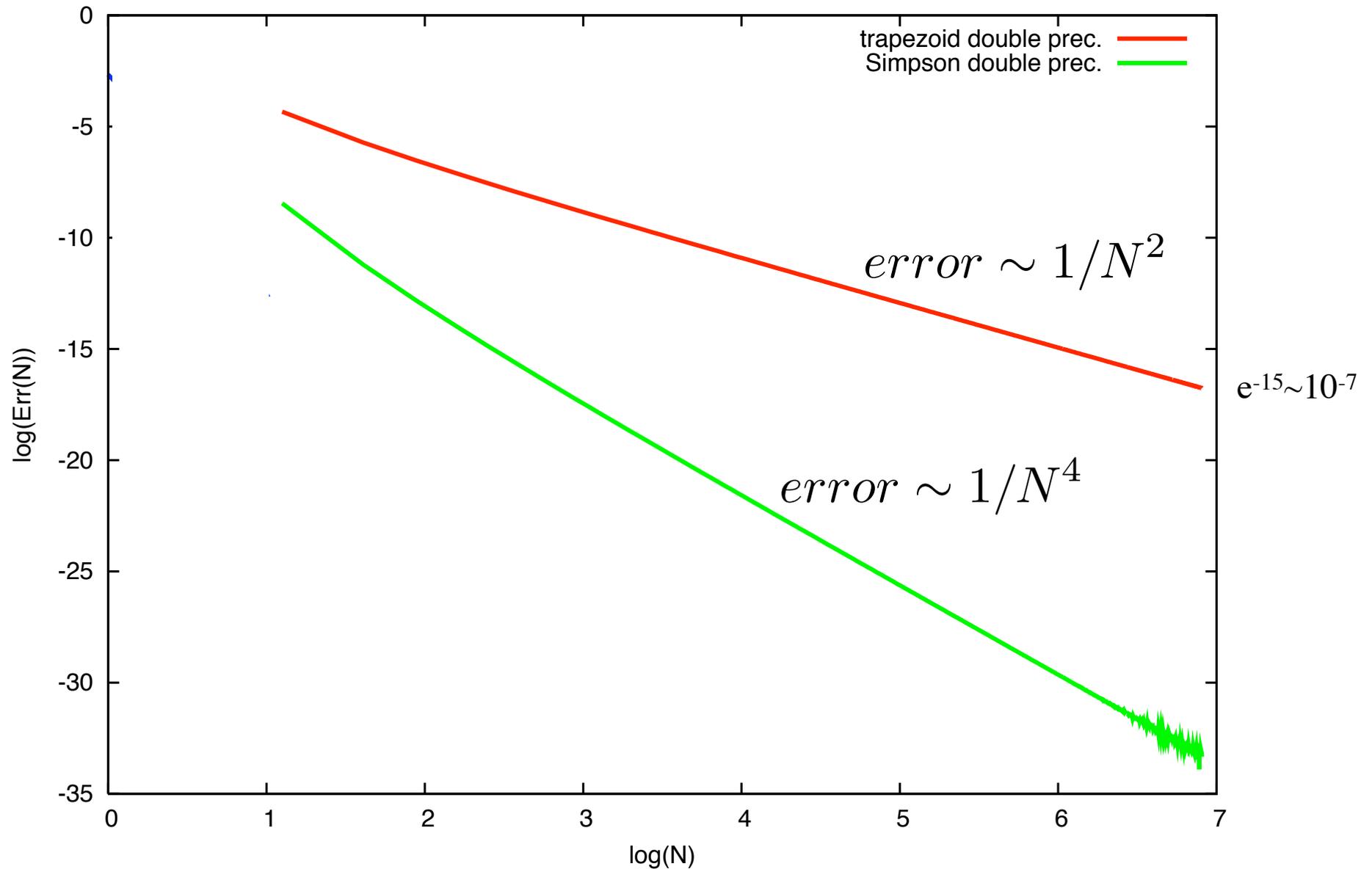
$$\int_{x_i}^{x_{i+1}} f(x) dx = f(x_i) \Delta x + \frac{1}{2} f'(x_i) (\Delta x)^2 + \frac{1}{6} f''(x_i) (\Delta x)^3 + \dots$$

error

(leading order in Δx)

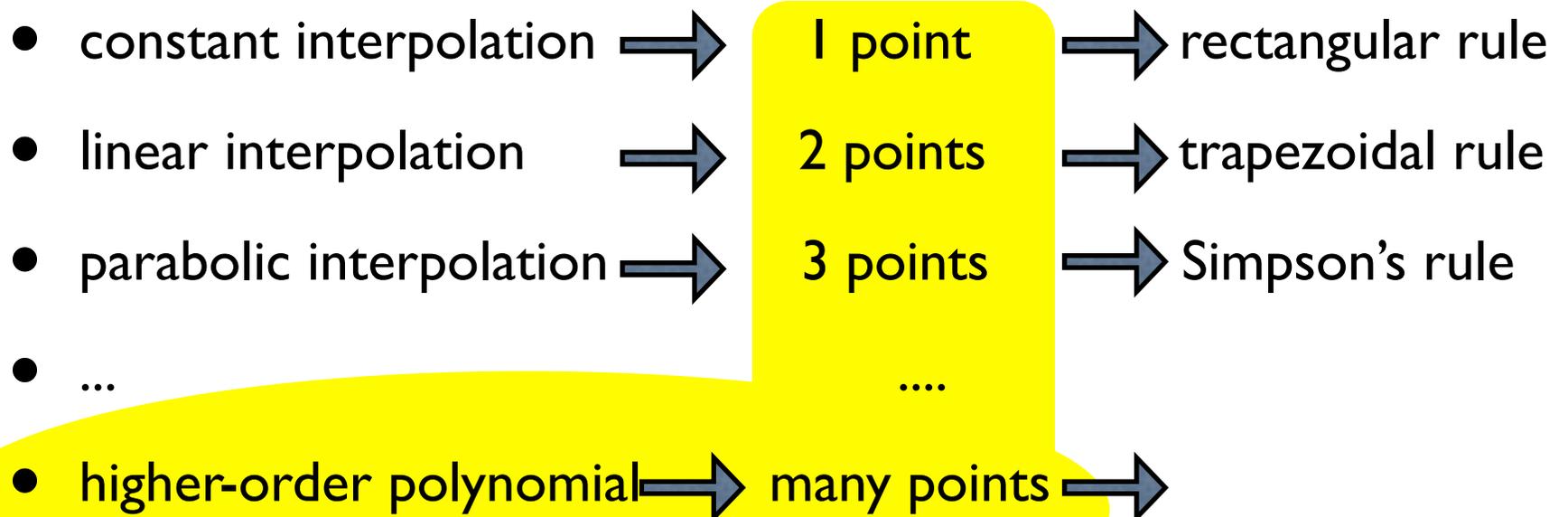
For n intervals ($\Delta x = (b - a)/n$): error is $n(\Delta x)^2 \sim 1/n$

Numerical integration - deterministic methods: comparison of errors in 1D



Deterministic methods - I

We use a piecewise polynomial interpolation:

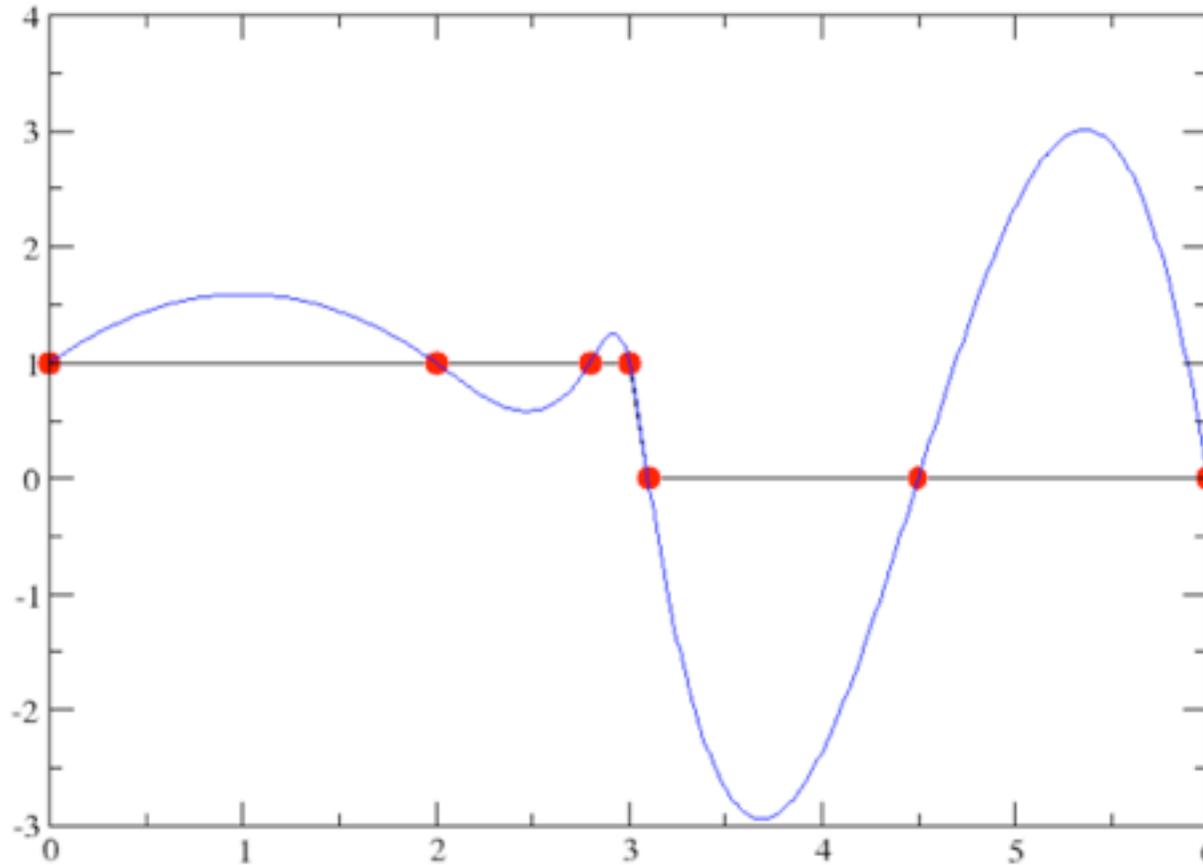


NOT CONVENIENT!

Warning: using higher degrees does not always improve accuracy!

(see also: Runge phenomenon (polynomial interpolation, oscillation at the edges of an interval), Gibbs phenomenon ...)

Deterministic methods - II



• $(x_i, f(x_i))$

Warning:

using high-order piecewise polynomial interpolation: possible strong oscillations between consecutive $(x_i, f(x_i))$, giving a bad interpolation of $f(x)$.

Here: $f(x)$ step function; - linear interp.; - cubic spline

Other deterministic methods

Numerical integration; other deterministic methods:

- in the simplest **equally-spaced-point** methods, we choose **weights** to calculate the average of the function:

$$\int_a^b f(x)dx \approx F_N = \sum_{i=1}^N v_i f(x_i)$$

rectangular rule: $x_i = a + \frac{b-a}{N}i, \quad v_i = \frac{b-a}{N} \quad \forall i = 1, \dots, N-1$

trapezoidal rule: $x_i = a + \frac{b-a}{N}i, \quad v_i = \frac{b-a}{N} \quad \forall i \neq 1, N; \quad v_1 = v_N = \frac{b-a}{2N}$

...

(at variance with these methods, in **MC methods** such as the ‘importance sampling’, we **choose only points, not weights**)

Numerical integration; other deterministic methods:

- in the simplest **equally-spaced-point** methods, we choose **weights** to calculate the average of the function:

$$\int_a^b f(x)dx \approx F_N = \sum_{i=1}^N v_i f(x_i)$$

idea: choose **not only weights but also points:**
more degrees of freedom!

$$x_i = ?, \quad v_i = ?$$

Another deterministic method:

Gaussian quadrature - I

Consider $\int_a^b f(x)dx$ and a function $W(x)$ defined on $[a, b]$

We can always formally write:

$$\int f(x)dx = \int W(x)F(x)dx \approx \sum_{j=1}^N w_j F(x_j) = \sum_{j=1}^N w_j \frac{f(x_j)}{W(x_j)} = \sum_{j=1}^N v_j f(x_j)$$

with $v_j = w_j/W(x_j)$

to be determined, depending on $W(x)$ (*)

(This will be convenient in particular if the resulting $F(x)$ is smooth,
but not necessarily)

(*) in general: $w_j \neq W(x_j)$

Another deterministic method:

Gaussian quadrature - II

Consider $\int_a^b f(x)dx$ and a function $W(x)$ defined on $[a, b]$

We can always formally write:

$$\int f(x)dx = \int W(x)F(x)dx \stackrel{(*)}{\approx} \sum_{j=1}^N w_j F(x_j) = \sum_{j=1}^N w_j \frac{f(x_j)}{W(x_j)} = \sum_{j=1}^N v_j f(x_j)$$

For a given $W(x)$, the N points and weights $\{x_j\}, \{w_j\}$ can be chosen to make the approximate relationship $(*)$ an exact equality if $F(x)$ is a $2N-1$ degree polynomial.

Another deterministic method: Gaussian quadrature - III

Consider
$$\int_{x_1}^{x_2} W(x)F(x)dx = \sum_{j=1}^N w_j F(x_j)$$

$F(x)$ a $2N - 1$ degree polynomial. Which are the N $\{x_j\}$, $\{w_j\}$?

If there is a set of polynomials $\{p_N(x)\}$ which are orthogonal in the same interval and for the same weight function $W(x)$:

$$\langle p_N | p_{N'} \rangle_W = \delta_{N,N'} , \text{ i.e., } \int_{x_1}^{x_2} W(x)p_N(x)p_{N'}(x)dx = \delta_{N,N'}$$

the points $\{x_i\}$ are exactly the roots of the $p_N(x)$ polynomials. The weights $\{w_j\}$ are related to them, but in general $w_j \neq W(x_j)$.

Gauss-Legendre quadrature

Consider $\int_{x_1}^{x_2} W(x)F(x)dx = \sum_{j=1}^N w_j F(x_j)$

with $F(x)$ a $2N - 1$ degree polynomial.

If: $W(x) = 1$ and $x_1 = -1, x_2 = 1$,

the Legendre polynomials $\{P_N(x)\}$ defined by:

$$(j + 1)P_{j+1} = (2j + 1)xP_j - jP_{j-1}$$

are orthogonal in $[-1, 1]$ with $W(x)=1$;

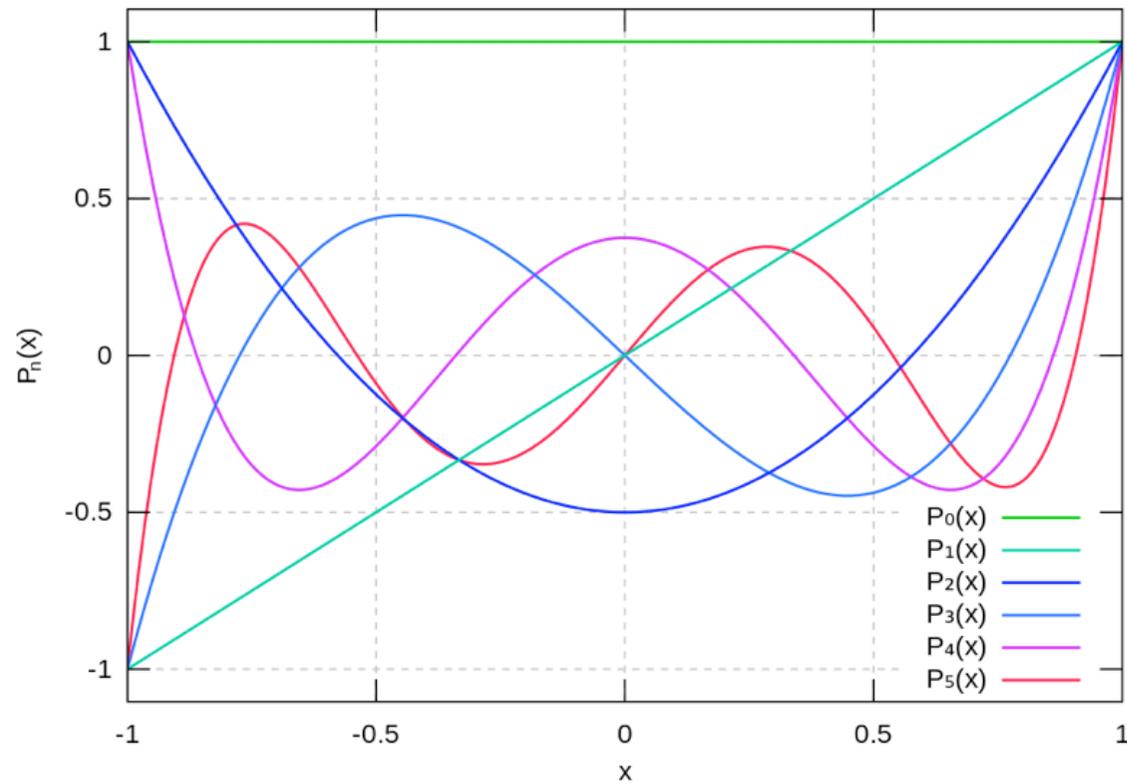
$\{x_j\}, \{w_j\}$ are such that $P_N(x_i) = 0$ and

$$w_i = \frac{2}{(1 - x_i^2)[P'_N(x_i)]^2}$$

The first few Legendre polynomials are:

n	$P_n(x)$
0	1
1	x
2	$\frac{1}{2}(3x^2 - 1)$
3	$\frac{1}{2}(5x^3 - 3x)$
4	$\frac{1}{8}(35x^4 - 30x^2 + 3)$
5	$\frac{1}{8}(63x^5 - 70x^3 + 15x)$

legendre polynomials



polynomials are odd or even in $x \Rightarrow$ roots are even

Legendre polynomials in Physics: examples of applications

1) For a polynomial expansion of a gravitational or coulombic potential:

$$\frac{1}{|\mathbf{x} - \mathbf{x}'|} = \frac{1}{\sqrt{r^2 + r'^2 - 2rr' \cos \gamma}} = \sum_{\ell=0}^{\infty} \frac{r'^{\ell}}{r^{\ell+1}} P_{\ell}(\cos \gamma)$$

where r and r' are the lengths of the vectors \mathbf{x} and \mathbf{x}' respectively and γ is the angle between those two vectors.

2) solution of Laplace's equation of the static potential, $\nabla^2 \Phi(\mathbf{x}) = 0$, in a charge-free region of space, if the boundary conditions have axial symmetry :

θ is the angle between the position of the observer and the $\hat{\mathbf{z}}$ axis (the zenith angle) ;

the solution for the potential will be

$$\Phi(r, \theta) = \sum_{\ell=0}^{\infty} \left[A_{\ell} r^{\ell} + B_{\ell} r^{-(\ell+1)} \right] P_{\ell}(\cos \theta).$$

3) solving Schrödinger equation in three dimensions for a central force :

the **associated Legendre polynomials** are derivatives of ordinary Legendre polynomials ($m \geq 0$)

$$P_{\ell}^m(x) = (-1)^m (1 - x^2)^{m/2} \frac{d^m}{dx^m} (P_{\ell}(x))$$

The Legendre polynomials are closely related to the spherical harmonics

$$Y_{\ell,m}(\theta, \phi) = \sqrt{\frac{(2\ell + 1)(\ell - m)!}{4\pi(\ell + m)!}} P_{\ell}^m(\cos \theta) e^{im\phi} \quad -\ell \leq m \leq \ell.$$

Gaussian quadrature

In practice, we choose $W(x)$ and N and use the set of N points and weights $\{x_j\}, \{w_j\}$ for the approximate integration:

$$\int f(x)dx = \int W(x)F(x)dx \approx \sum_{j=1}^N w_j F(x_j) = \sum_{j=1}^N w_j \frac{f(x_j)}{W(x_j)} = \sum_{j=1}^N v_j f(x_j)$$

Gauss-Legendre quadrature

For: $x_1 = -1$, $x_2 = 1$

(see slide 10 for the list of the first polynomials)

N	i	x_i	w_i	degree
1	1	0	2	1
2	1	-0.577350269189626	1	3
	2	0.577350269189626	1	
3	1	-0.774596669241483	0.555555555555556	5
	2	0	0.888888888888889	
	3	0.774596669241483	0.555555555555556	
4	1	-0.861136311594053	0.347854845137454	7
	2	-0.339981043584856	0.652145154862546	
	3	0.339981043584856	0.652145154862546	
	4	0.861136311594053	0.347854845137454	

degree of the polynomial exactly integrable

The integration in an interval $[a,b]$ different from $[-1,1]$ (“old”) can be easily done performing the scaling:

$$x_{new} = \frac{b-a}{2}x_{old} + \frac{b+a}{2} \quad \text{and} \quad w_{new} = \frac{b-a}{2}w_{old}$$

Gauss-Legendre quadrature

In case of classical, well known, orthogonal polynomials, ready-to-use subroutines exist for the computation of Abscissas and Weights $\{x_j\}, \{w_j\}$

e.g. GAULEG(x_1, x_2, x, w, n) of *Numerical Recipes* which, given x_1, x_2, n , provides as output the arrays $x(n), w(n)$

Some programs:

on <https://moodle2.units.it/>

int.f90 (trapeziodal and Simpson integration)

gauleg-llorder.f90

gauleg-others.f90 (generation of points up to 15 points in $[-1, 1]$ using GAULEG adapted from “Numerical Recipes” (self-contained) and some tests for easy-to-integrate functions)

In the subdirectory: **gauss-nr90/**

find the original routine from “Numerical Recipes” and related external routines/modules/interfaces and a main program for test (see following slide)

Gauss-Legendre from Numerical Recipes

Use of GAULEG:

In order to use the routines of Numerical Recipes, you have to compile and link the main program with:

- the subroutine **gauleg.f90** which gives points and abscissas
- **nrtype.f90** containing type declarations;
- **nrutil.f90** containing **moduli** and utilities;
- **nr.f90** containing (through the interfaces) the conventions to call the subroutines with the main program

i) You must first **compile** these files **with the option “-c”**:
this **produces .mod and .o (the objects)**.

ii) In a second step compile the main program.

iii) Finally you link all the files *.o and produce the executable:

```
gfortran -c nrtype.f90 nrutil.f90 nr.f90 gauleg.f90
```

```
gfortran -c gauleg_nr_test.f90
```

```
gfortran -o a.out gauleg_nr_test.o nrtype.o nrutil.o nr.o gauleg.o
```

gauleg.f90 from Numerical Recipes

```
SUBROUTINE gauleg(x1,x2,x,w)
  USE nrtype; USE nrutil, ONLY : arth,assert_eq,nrerror
  IMPLICIT NONE
  REAL(SP), INTENT(IN) :: x1,x2
  REAL(SP), DIMENSION(:), INTENT(OUT) :: x,w
  REAL(DP), PARAMETER :: EPS=3.0e-14_dp
  INTEGER(I4B) :: its,j,m,n
  INTEGER(I4B), PARAMETER :: MAXIT=10
  REAL(DP) :: x1,xm
  REAL(DP), DIMENSION((size(x)+1)/2) :: p1,p2,p3,pp,z,z1
  LOGICAL(LGT), DIMENSION((size(x)+1)/2) :: unfinished
  n=assert_eq(size(x),size(w),'gauleg')
  m=(n+1)/2
  xm=0.5_dp*(x2+x1)
  x1=0.5_dp*(x2-x1)
  z=cos(PI_D*(arth(1,1,m)-0.25_dp)/(n+0.5_dp))
  ...
  ...
  ...
  x(1:m)=xm-x1*z
  x(n:n-m+1:-1)=xm+x1*z
  w(1:m)=2.0_dp*x1/((1.0_dp-z**2)*pp**2)
  w(n:n-m+1:-1)=w(1:m)
END SUBROUTINE gauleg
```

nrtype.f90 from Numerical recipes

```
MODULE nrtype
  INTEGER, PARAMETER :: I4B = SELECTED_INT_KIND(9)
  INTEGER, PARAMETER :: I2B = SELECTED_INT_KIND(4)
  INTEGER, PARAMETER :: I1B = SELECTED_INT_KIND(2)
  INTEGER, PARAMETER :: SP = KIND(1.0)
  INTEGER, PARAMETER :: DP = KIND(1.0D0)
  INTEGER, PARAMETER :: SPC = KIND((1.0,1.0))
  INTEGER, PARAMETER :: DPC = KIND((1.0D0,1.0D0))
  INTEGER, PARAMETER :: LGT = KIND(.true.)
  REAL(SP), PARAMETER :: PI=3.141592653589793238462643383279502884197_sp
  REAL(SP), PARAMETER :: PI02=1.57079632679489661923132169163975144209858_sp
  REAL(SP), PARAMETER :: TWOPI=6.283185307179586476925286766559005768394_sp
  REAL(SP), PARAMETER :: SQRT2=1.41421356237309504880168872420969807856967_sp
  REAL(SP), PARAMETER :: EULER=0.5772156649015328606065120900824024310422_sp
  REAL(DP), PARAMETER :: PI_D=3.141592653589793238462643383279502884197_dp
  REAL(DP), PARAMETER :: PI02_D=1.57079632679489661923132169163975144209858_dp
  REAL(DP), PARAMETER :: TWOPI_D=6.283185307179586476925286766559005768394_dp
  TYPE sprs2_sp
    INTEGER(I4B) :: n,len
    REAL(SP), DIMENSION(:), POINTER :: val
    INTEGER(I4B), DIMENSION(:), POINTER :: irow
    INTEGER(I4B), DIMENSION(:), POINTER :: jcol
  END TYPE sprs2_sp
  TYPE sprs2_dp
    INTEGER(I4B) :: n,len
    REAL(DP), DIMENSION(:), POINTER :: val
    INTEGER(I4B), DIMENSION(:), POINTER :: irow
    INTEGER(I4B), DIMENSION(:), POINTER :: jcol
  END TYPE sprs2_dp
END MODULE nrtype
```

nr.f90 from Numerical Recipes

```
MODULE nr
  INTERFACE
    SUBROUTINE gauleg(x1,x2,x,w)
      USE nrtype
      REAL(SP), INTENT(IN) :: x1,x2
      REAL(SP), DIMENSION(:), INTENT(OUT) :: x,w
    END SUBROUTINE gauleg
  END INTERFACE
  ! ... the original file contains several other INTERFACES ...
END MODULE nr
```

nrutil.f90 (Here only for: array_copy, arth, assert_eq, nrerror)

```
MODULE nrutil
  USE nrtype
  IMPLICIT NONE
  INTEGER(I4B), PARAMETER :: NPAR_ARTH=16,NPAR2_ARTH=8

  ...

  ...

  ...

  -----,-----,-----,-----,-----
  INTERFACE array_copy
    MODULE PROCEDURE array_copy_r, array_copy_d, array_copy_i
  END INTERFACE

  ...

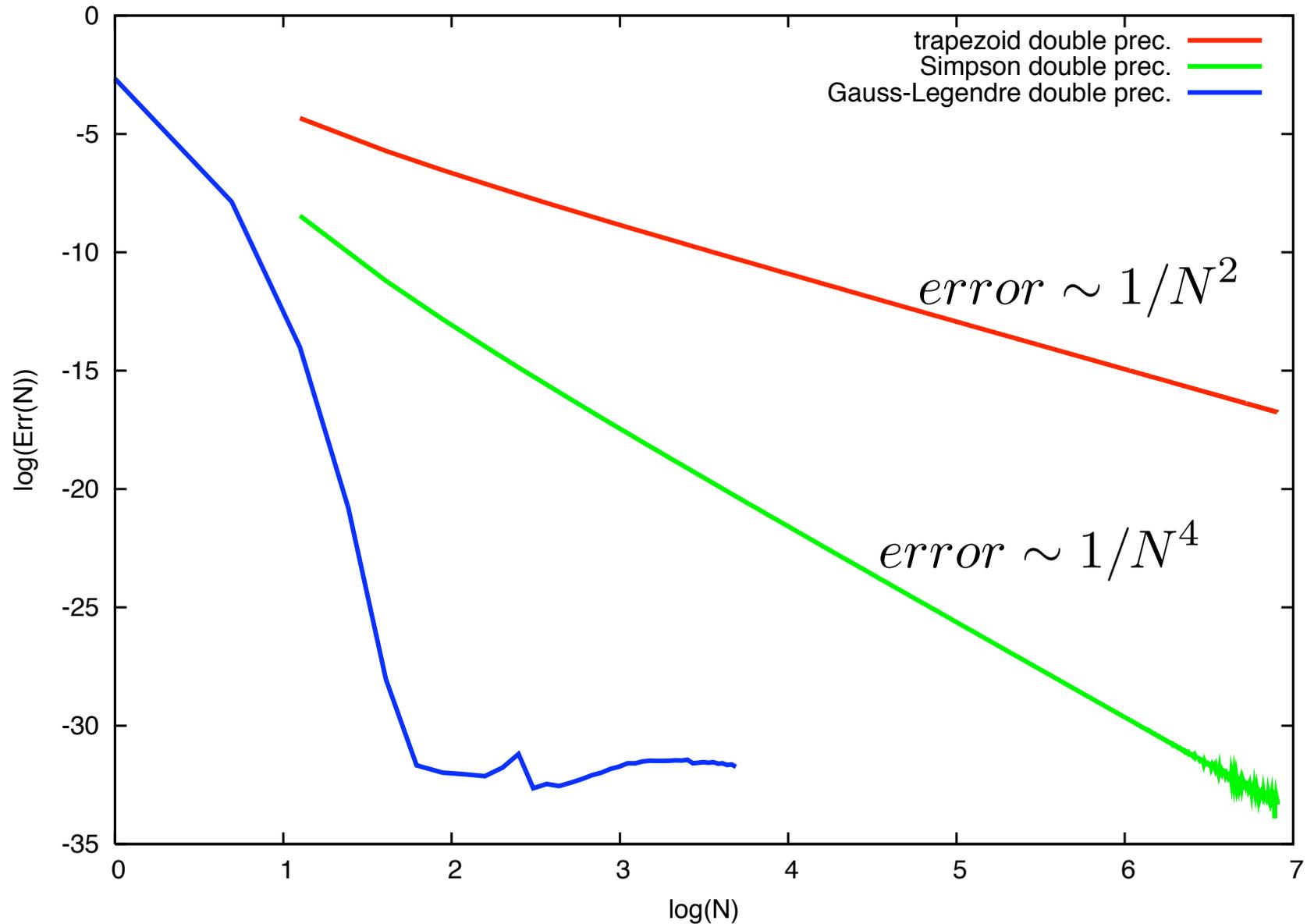
  ! ... l'originale contiene ancora molte altre INTERFACES....
CONTAINS

  SUBROUTINE array_copy_r(src,dest,n_copied,n_not_copied)
    REAL(SP), DIMENSION(:), INTENT(IN) :: src

    ...

    ! .... and many other FUNCTIONS and SUBROUTINES ....
END MODULE nrutil
```

Numerical integration, deterministic methods: comparison of errors in 1D



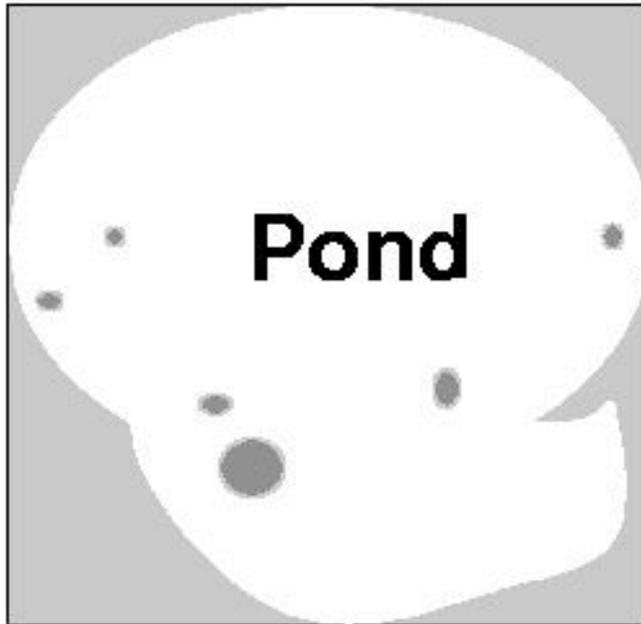
(double precision needed to appreciate the convergence of Gauss-Legendre numerical estimate)

Monte Carlo methods

Monte Carlo methods:

“acceptance-rejection” or “hit or miss”
(to calculate areas)

which is A_{pond} ?



- enclose the pond in a box of Area A_{box}
- throw pebbles uniformly and randomly in the box
- count the number of pebbles felt in the pond with respect to the number felt in the box
- Assuming a uniform distribution, the number of pebbles falling into the ponds is proportional to the area of the pond:

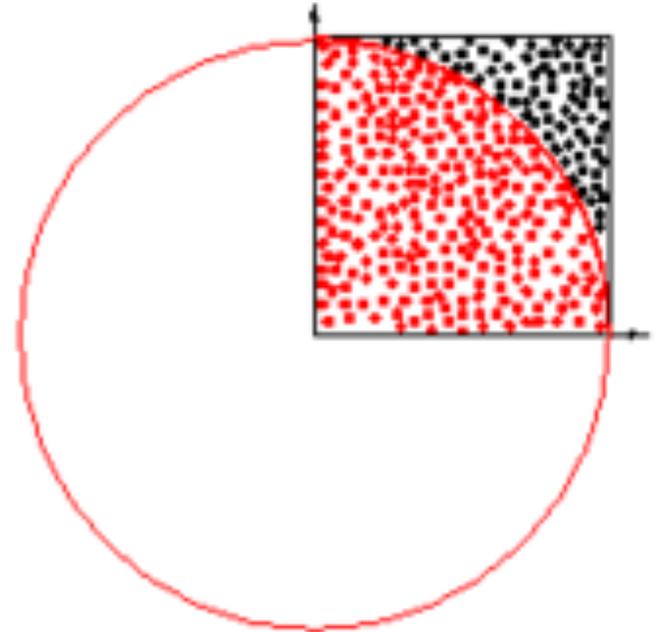
$$\frac{N_{\text{pond}}}{N_{\text{pond}} + N_{\text{box}}} = \frac{A_{\text{pond}}}{A_{\text{box}}}$$
$$\Rightarrow A_{\text{pond}} = \frac{N_{\text{pond}}}{N_{\text{pond}} + N_{\text{box}}} A_{\text{box}}$$

Monte Carlo methods:

“acceptance-rejection” or “hit or miss”
(to calculate areas)

$$\pi = ???$$

N random points in the unit square
- coordinates x_i, y_i -



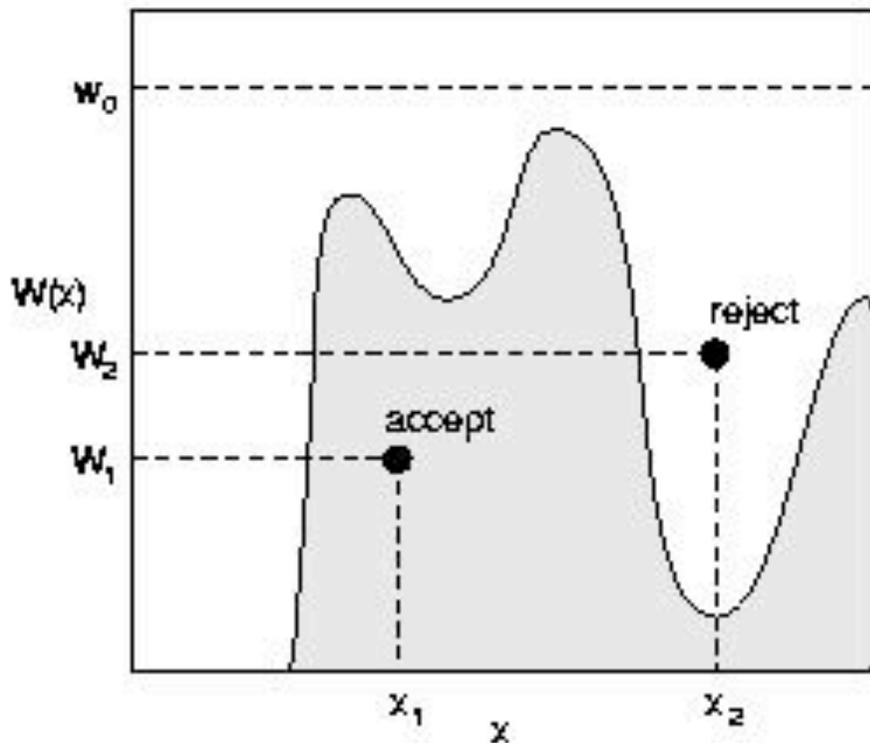
Then, the number of points N_c lying within the quarter circle (i.e. fulfilling the relation $x^2 + y^2 \leq 1$) is compared to the total number N of points and the fraction will give us an approximate value of π :

$$\pi(N) = 4 \frac{N_c(N)}{N}$$

Monte Carlo methods:

“acceptance-rejection” or “hit or miss”
(to calculate definite integrals)

$$\int W(x) dx = ?$$



For $W(x)$ positive in the integration interval, the value of the area under $W(x)$ can be obtained by producing random points (i.e. (x,y) random pairs) uniformly distributed in a rectangle containing $W(x)$.

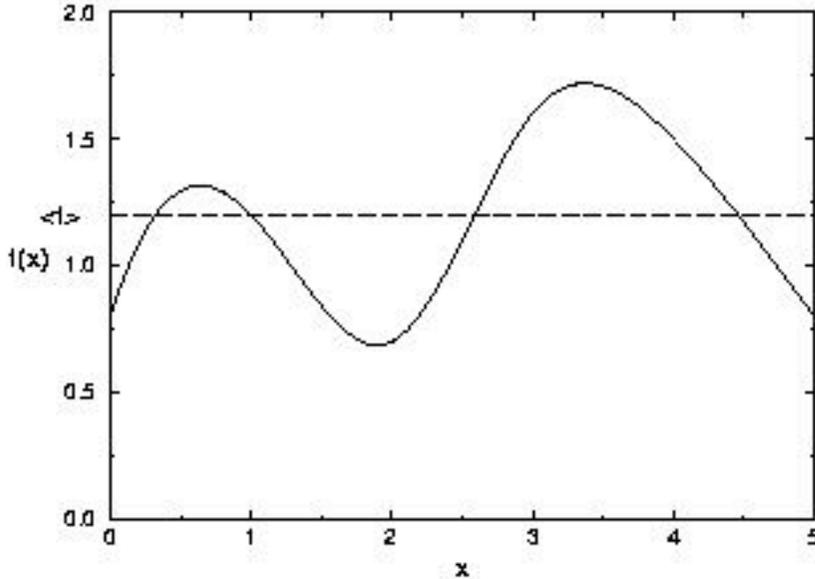
For each point (x,y) compare y with $W(x)$: if $y < W(x)$, the point is accepted. The area under $W(x)$ is the number of points accepted divided by the total number of points generated and multiplied by the area of the rectangle.

(remember: also used to generate random numbers x_i distributed according $W(x)$)

Other simple Monte Carlo methods

We can always write:

$$I = \int_a^b f(x) dx = (b - a) \langle f \rangle$$

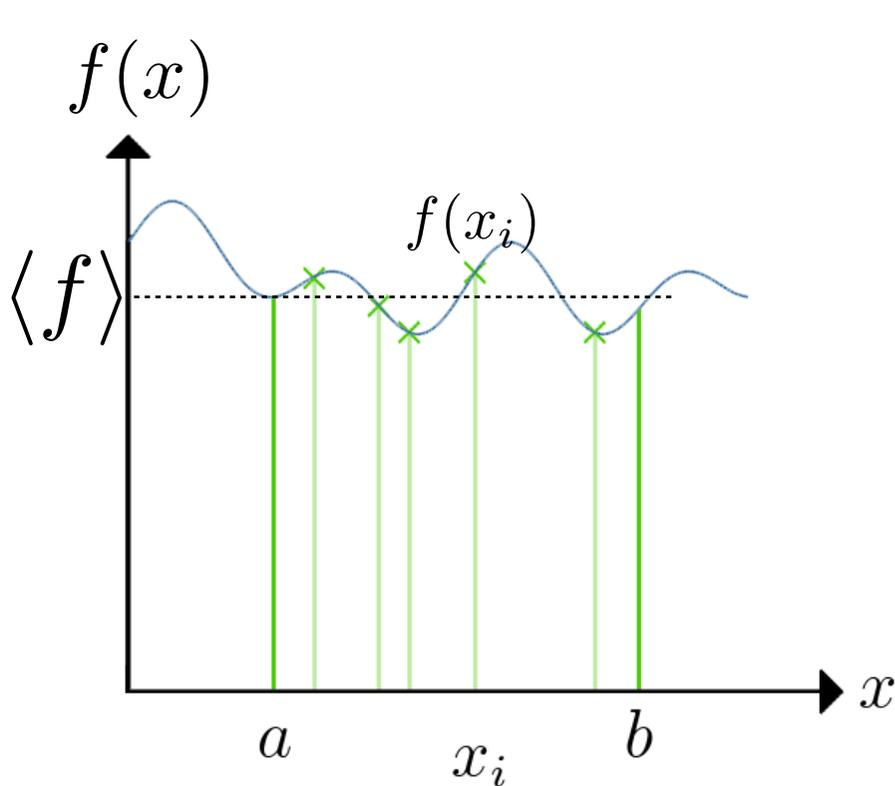


i.e., the value of the integral of $f(x)$ between a and b equals the length of the interval $(b-a)$ times the average value of the function $\langle f \rangle$ over the same interval.

(If $f:[a,b] \rightarrow \mathbb{R}$ is a continuous function, then there exists a number c in $[a,b]$ such that $f(c) = \langle f \rangle$ (mean value theorem for integration))

how to estimate $\langle f \rangle$ efficiently and accurately?

A simple Monte Carlo method: “sample mean”



$$I = \int_a^b f(x) dx = (b - a) \langle f \rangle$$

The **sample mean** can be calculated by sampling the function (*if smooth enough...*) with a sequence of N uniform random numbers in [a,b]:

$$\langle f \rangle \approx \frac{1}{N} \sum_{i=1}^N f(x_i)$$

$$\int_a^b f(x) dx \approx (b - a) \frac{1}{N} \sum_{i=1}^N f(x_i) = (b - a) \langle f \rangle$$

Monte Carlo methods: error estimate

Example: MC estimate of π (exact value known)

We can use either acceptance-rejection or sample mean method: $I = 4 \int_0^1 \sqrt{1-x^2} = \pi = 3.1416\dots$

Since we know the “exact” result I , we can calculate the **error** in two ways:

1) the **actual error** from the difference with respect to the exact value:

$$\Delta_n = |F_n - I| \quad \text{with} \quad F_n = (b-a) \frac{1}{n} \sum_{i=1}^n f(x_i), \quad x_i \text{ random}$$

2) the numerical error from the **variance of the data** $\{f(x_i)\}$:

$$\sigma^2 = \langle f^2 \rangle - \langle f \rangle^2,$$

where

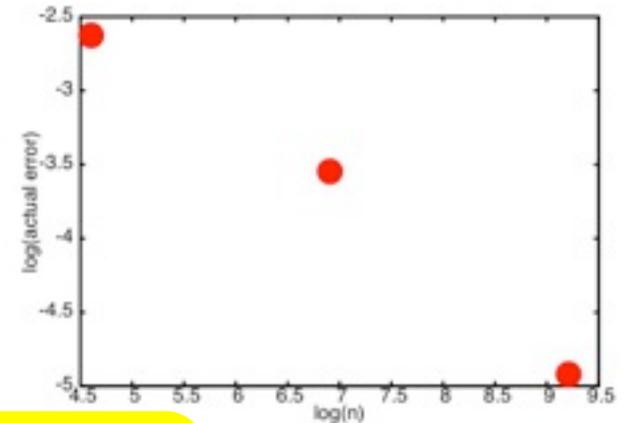
$$\langle f \rangle = \frac{1}{n} \sum_{i=1}^n f(x_i) \quad \text{and} \quad \langle f^2 \rangle = \frac{1}{n} \sum_{i=1}^n f(x_i)^2$$

Monte Carlo methods: error estimate

Results:

$$I = 4 \int_0^1 \sqrt{1-x^2} = \pi = 3.1416\dots$$

n	F_n	actual error	σ_n
10^2	3.0692	0.0724	0.85550
10^3	3.1704	0.0288	0.8790
10^4	3.1489	0.0073	0.8850



1) the **actual error** Δ_n decreases as $1/n^{1/2}$

2) the numerical error from the **variance of the data**, σ_n , is roughly constant and is much larger than the actual error

what is the correct error estimate?

Monte Carlo methods: error estimate

...typically you do not know which is the “actual error” (you do not know the “true” value and you cannot compare your result with that!)...
but **we would like to give an error to our numerical estimate...**
(to which extent is our numerical estimate reliable?)

Two methods to estimate the error numerically
from the variance of the data
(**“reduction of variance”**):

I) average of the averages

II) block average

MC error handling: method I

“average of the averages”

make additional runs of n trials each.

Let M_α be the average of each run :

run α	M_α	actual error
1	3.1489	0.0073
2	3.1326	0.0090
3	3.1404	0.0012
4	3.1460	0.0044
5	3.1526	0.0110
6	3.1397	0.0019
7	3.1311	0.0105
8	3.1358	0.0058
9	3.1344	0.0072
10	3.1405	0.0011

one run $\equiv n = 10^4$ trials each

Examples of Monte Carlo measurements of the mean value of $f(x) = 4\sqrt{1-x^2}$ in the interval $[0, 1]$. A total of 10 measurements of $n = 10^4$ trials each were made. The mean value M_α and the actual error $|M_\alpha - \pi|$ for each measurement are shown.

$$\text{Calculate: } \sigma_m^2 = \langle M^2 \rangle - \langle M \rangle^2 \quad \text{with } \langle M \rangle = \frac{1}{m} \sum_{\alpha=1}^m M_\alpha, \quad \langle M^2 \rangle = \frac{1}{m} \sum_{\alpha=1}^m M_\alpha^2$$
$$\implies \sigma_m = 0.0068$$

σ_m is consistent with the results for the actual errors

MC error handling: method II

“block averages”

Instead of doing additional measurements, divide them into “s SUBSETS” and let S_k be the average within each subset :

subset k	S_k
1	3.14326
2	3.15633
3	3.10940
4	3.15337
5	3.15352
6	3.11506
7	3.17989
8	3.12398
9	3.17565
10	3.17878

The variance associated to the average of the subsets $\sigma_s^2 = \langle S^2 \rangle - \langle S \rangle^2$ gives $\sigma_s = 0.025$, but

σ_s/\sqrt{s} , which for our example is approximately $0.025/\sqrt{(10)} \approx 0.008$.

is consistent with the actual error

Monte Carlo methods:

error estimate - variance reduction

summary

$$\sigma_n / \sqrt{n} \approx \sigma_m \approx \sigma_s / \sqrt{s}$$

from the variance of
the whole set of data

Note: for
uncorrelated data !

(proof)

the variance
of the

average of
the averages

from the variance
of the

block averages

the most convenient!
but: change block size
and check that
it does not change

Monte Carlo methods: summary

We have introduced :

* “acceptance-rejection”

* “sample mean” to estimate $\langle f \rangle \approx \frac{1}{N} \sum_{i=1}^N f(x_i)$

both OK for smoothly varying functions, but
not very efficient for rapidly varying functions

How to improve the efficiency of MC integration?

A trick for numerical integration: “reduction of variance”

(Note: same word, but different meaning w.r.t. previous slides on error handling)

Given a function $f(x)$ to integrate, suppose that $g(x)$ exists, whose integral is known and such that:

$$|f(x) - g(x)| \ll \varepsilon$$

Therefore:

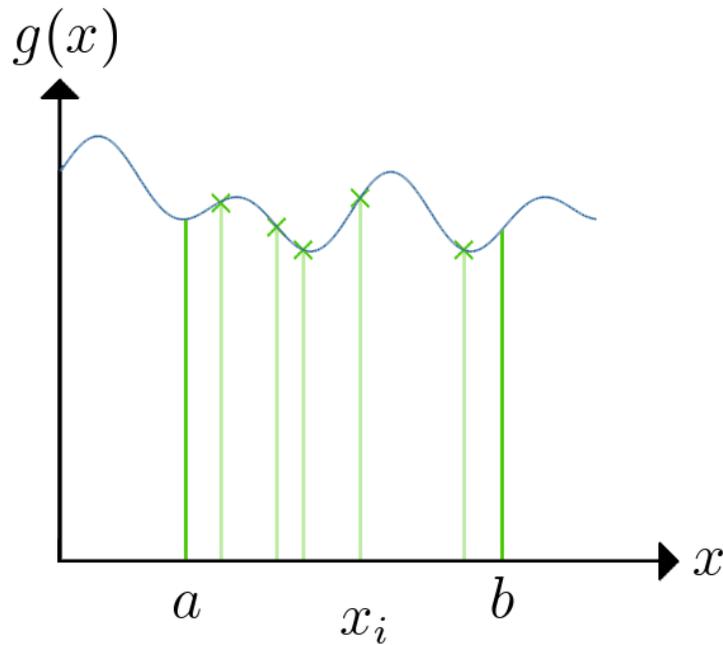
$$F = \int_a^b f(x) dx = \int_a^b ((f(x) - g(x)) + g(x)) dx = \int (f(x) - g(x)) dx + \int g(x) dx$$



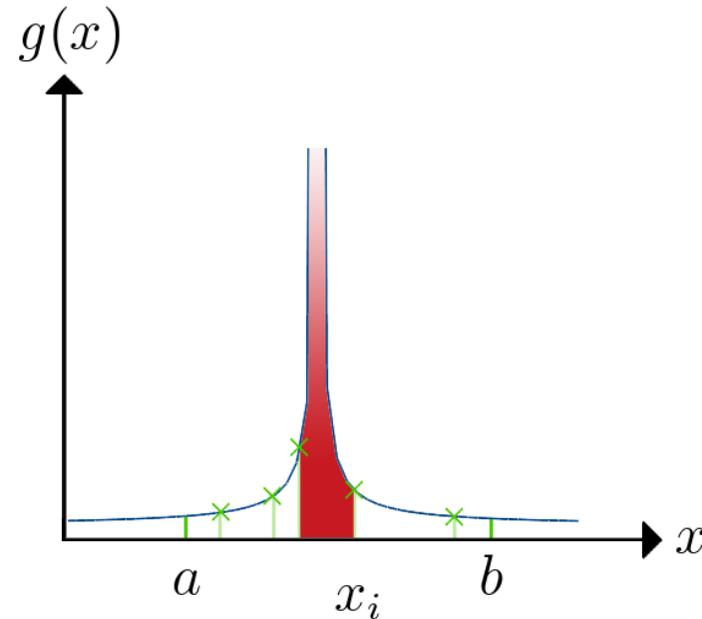
easy to calculate

Another simple Monte Carlo method: “importance sampling”

Mean value: easy to calculate for smoothly varying functions.
But not for functions rapidly varying.



smooth function



function with singularity

How to manage such cases?

Another simple Monte Carlo method: “importance sampling”

Mean value: easy to calculate for smoothly varying functions.

Idea: in order to calculate:

$$\langle f \rangle \approx \frac{1}{N} \sum_{i=1}^N f(x_i)$$

consider a **distribution function** $p(x)$ **easy to integrate analytically and close to** $f(x)$:

$$F = \int_a^b f(x) dx = \int_a^b \left[\frac{f(x)}{p(x)} \right] p(x) dx = \left\langle \frac{f(x)}{p(x)} \right\rangle \int_a^b p(x) dx$$

where $\left\langle \frac{f(x)}{p(x)} \right\rangle \approx \frac{1}{N} \sum_{i=1}^N \left[\frac{f(x_i)}{p(x_i)} \right]$

(particular case:
uniform distrib.
 $p(x)=1/(b-a)$...)

with $\{x_i\}$ distributed according to $p(x)$

Monte Carlo methods: “importance sampling”

Calculate:

$$F = \int_0^1 e^{-x^2} dx.$$

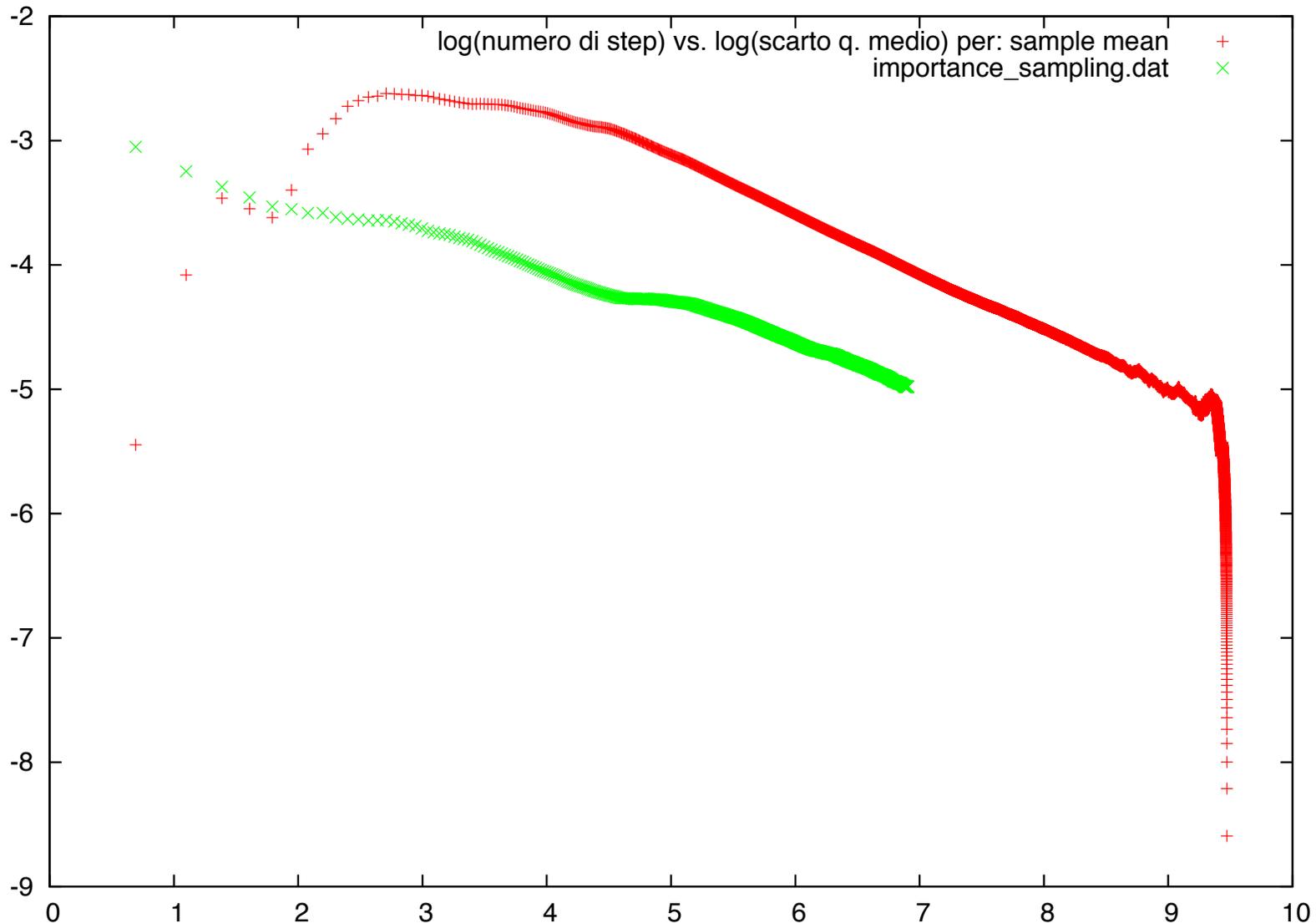
with “sample mean” with random numbers with uniform distribution or using the “importance sampling” with $p(x) = e^{-x}$

	$p(x) = 1$	$p(x) = Ae^{-x}$
n (trials)	4×10^5	8×10^3
F_n	0.7471	0.7469
σ	0.2010	0.0550
σ/\sqrt{n}	3×10^{-4}	6×10^{-4}
Total CPU time (s)	35	1.35
CPU time per trial (s)	10^{-4}	2×10^{-4}

← efficient !

(pay attention to the normalization of $p(x)$...)

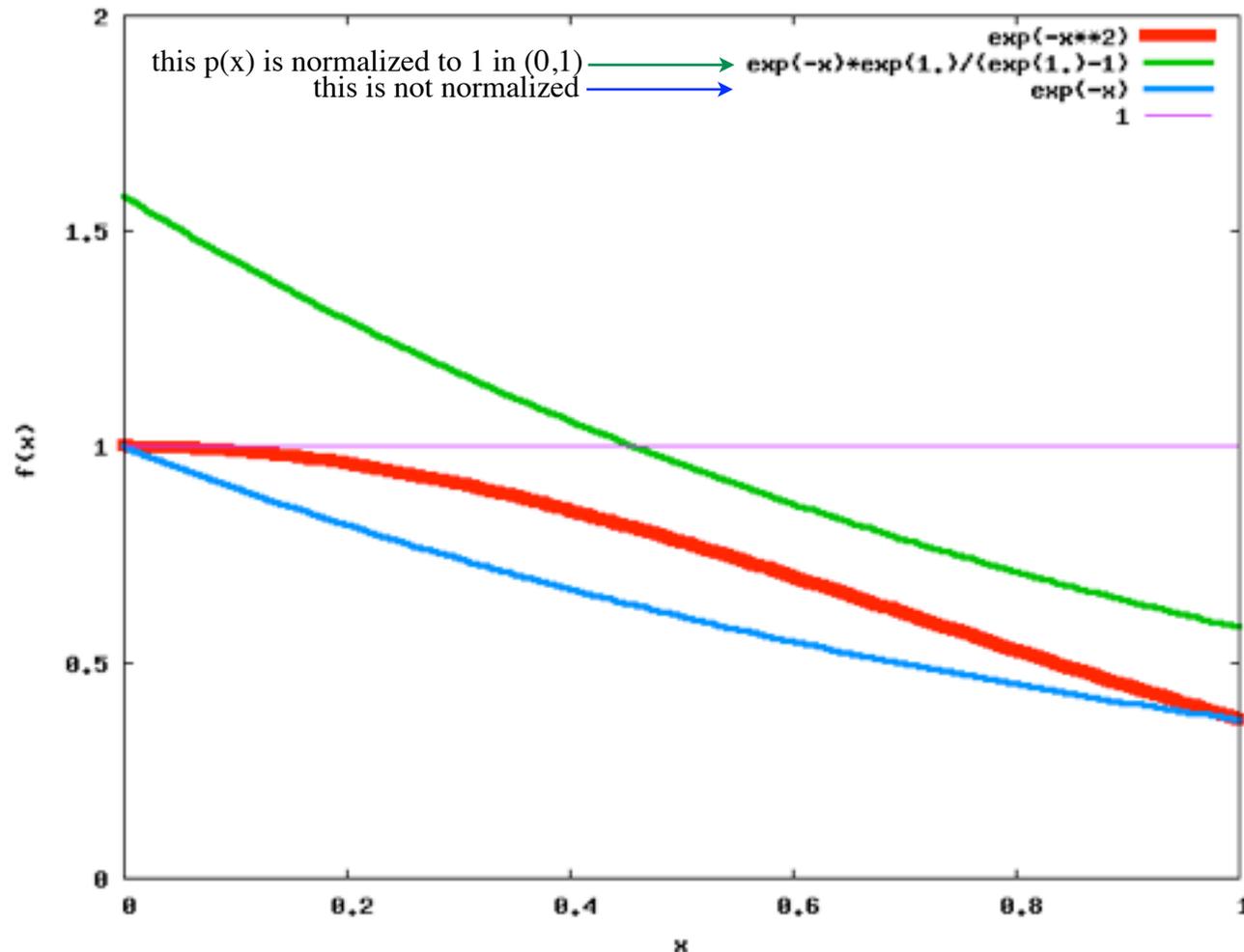
$\text{error}(\text{MC}) \sim I/\sqrt{N} \Rightarrow$ see $\log(\text{error})$ vs. $\log(N)$
but with different prefactors
for sample means vs importance sampling



Choice of the importance sampling function

$$F = \int_0^1 e^{-x^2} dx.$$

$$F = \int_a^b f(x) dx = \int_a^b \left[\frac{f(x)}{p(x)} \right] p(x) dx = \left\langle \frac{f(x)}{p(x)} \right\rangle \int_a^b p(x) dx$$



(pay attention to the normalization of p(x)..)

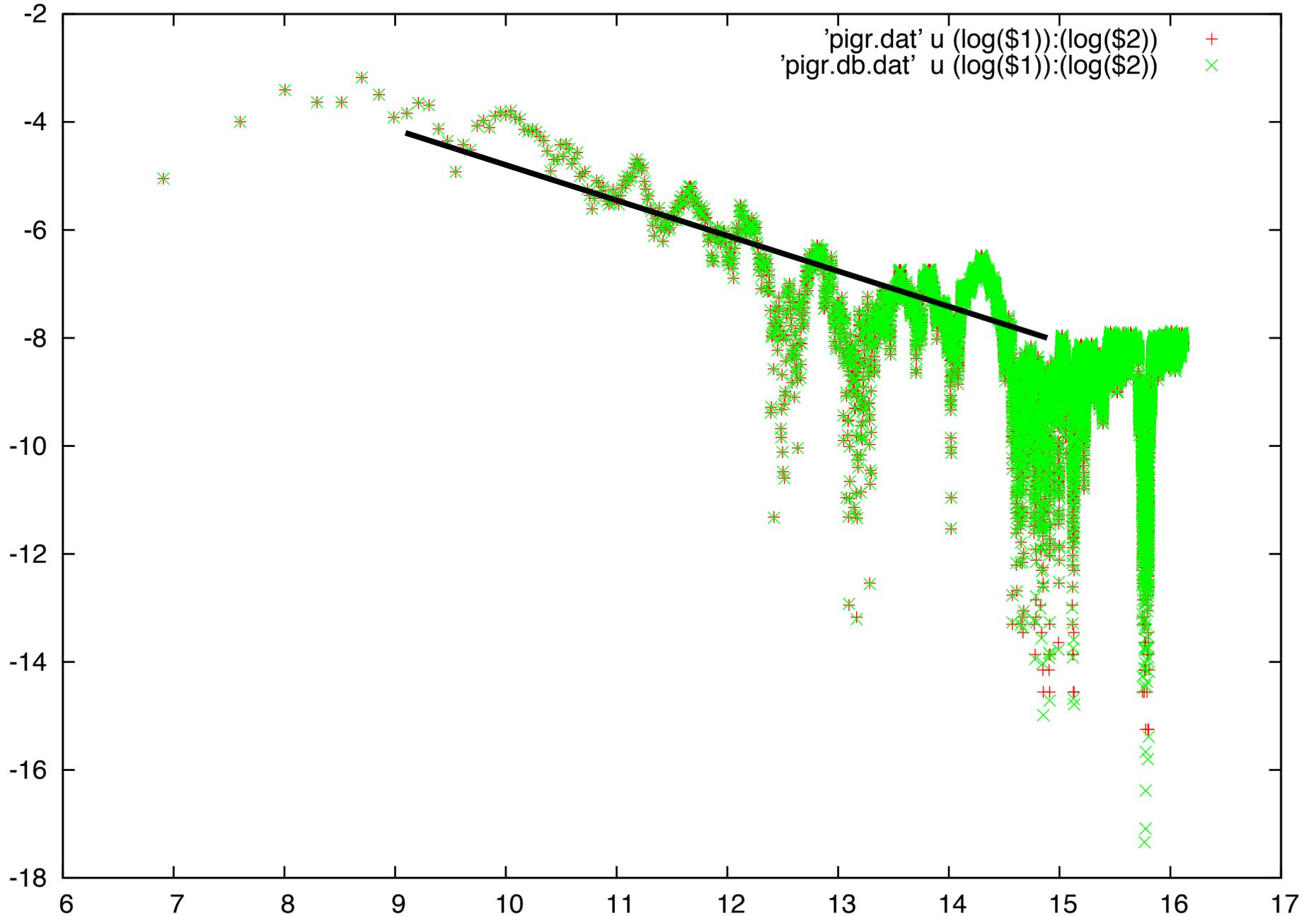
Some programs

on <https://moodle2.units.it/>

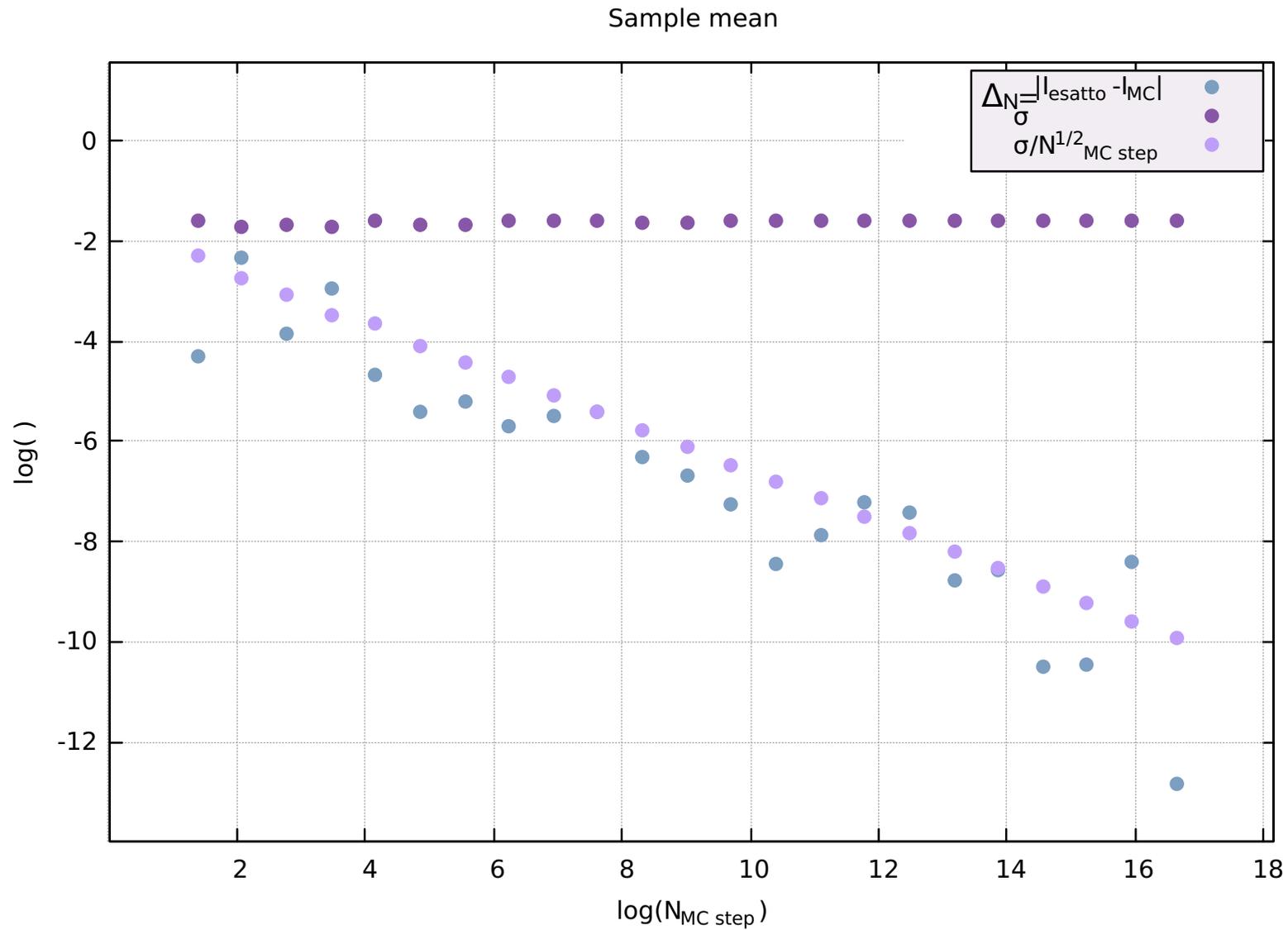
pi.f90 Monte Carlo integration for the calculation of π

For the other exercises: write yourself the codes

$\text{error}(\text{MC}) \sim 1/\sqrt{N} \Rightarrow \text{see } \log(\text{error}) \text{ vs. } \log(N)$



error(MC) $\sim 1/\sqrt{N}$: “true” error and statistical error



(credits: G. Lautizi, a.y. 2019-20)

Summary of numerical integration (MC and deterministic) methods

MC sample mean

$$\int_a^b f(x)dx = (b-a) \langle f \rangle \approx (b-a) \frac{1}{N} \sum_{i=1}^N f(x_i) \quad \text{with } \{x_i\} \text{ randomly uniformly distributed in } [a, b]$$

(it can be considered as Importance sampling with $p(x) = \frac{1}{b-a}$ in $[a, b]$)

MC importance sampling

$$\int_a^b f(x)dx = \int_a^b \frac{f(x)}{p(x)} p(x)dx = \left\langle \frac{f(x)}{p(x)} \right\rangle \int_a^b p(x)dx \approx \frac{1}{N} \sum_{i=1}^N \frac{f(x_i)}{p(x_i)} \int_a^b p(x)dx$$

with $\{x_i\}$ randomly distributed according $p(x)$

Deterministic, equispaced points

$$\int_a^b f(x)dx \approx \sum_{i=1}^N v_i f(x_i) \quad \text{with } x_i = a + \frac{b-a}{N}i, \quad v_i \text{ to be determined}$$

Deterministic, non equispaced points

$$\int_a^b f(x)dx \approx \sum_{i=1}^N v_i f(x_i) \quad \text{with } \{x_i\}, \{v_i\} \text{ to be determined}$$

Error estimate:
comparison between
deterministic and MC
methods
in d -dimension

Error estimate for numerical integration with deterministic methods

(Reminder from previous slides)

$$\int f(x)dx = F_n + error$$

How to evaluate the error? Consider the Taylor expansion of the integrand function and then integrate:

$$f(x) = f(x_i) + f'(x_i)(x - x_i) + \frac{1}{2}f''(x_i)(x - x_i)^2 + \dots ,$$

$$\int_{x_i}^{x_{i+1}} f(x) dx = f(x_i)\Delta x + \frac{1}{2}f'(x_i)(\Delta x)^2 + \frac{1}{6}f''(x_i)(\Delta x)^3 + \dots (*)$$

$$\Delta x \equiv x_{i+1} - x_i$$

Error estimate for numerical integration: Rectangular approximation

(Reminder from
previous slides)

$$\int_{x_i}^{x_{i+1}} f(x) dx \approx f(x_i) \Delta x$$

Compare \square with (*):

$$\int_{x_i}^{x_{i+1}} f(x) dx = f(x_i) \Delta x + \frac{1}{2} f'(x_i) (\Delta x)^2 + \frac{1}{6} f''(x_i) (\Delta x)^3 + \dots$$

error

(leading order in Δx)

For n intervals ($\Delta x = (b - a)/n$): error is $n(\Delta x)^2 \sim 1/n$

(...and similarly for
higher-order approximations)

Numerical integration: multidimensional integrals

$$F = \int_R f(x, y) dx dy$$

The rectangular approximation gives $\Delta x \Delta y \sim (\Delta x)^2 \sim 1/n$, being n the number of parts (or pairs of points) of the integration domain:

$$\int_{x_i}^{x_{i+1}} \int_{y_i}^{y_{i+1}} f(x, y) dx dy \approx f(x_i, y_i) \Delta x \Delta y \quad (*)$$

The Taylor expansion of the integrand function gives:

$$f(x, y) = f(x_i, y_i) + f'_x(x_i, y_i)(x - x_i) + f'_y(x_i, y_i)(y - y_i) + \dots$$

$$\int_{x_i}^{x_{i+1}} \int_{y_i}^{y_{i+1}} f(x, y) dx dy = f(x_i, y_i) \Delta x \Delta y + f'_x(x_i, y_i) \frac{(\Delta x)^2}{2} \Delta y + f'_y(x_i, y_i) \Delta x \frac{(\Delta y)^2}{2} + \dots (**)$$

(*) against (**) => **error**
(leading order in Δx)

For n intervals: error is $n(\Delta x)^3 \sim 1/n^{1/2}$

Numerical integration: multidimensional integrals

Therefore for rectangular approx.:

$$d=1: \text{error} \sim 1/n \qquad d=2: \text{error} \sim 1/n^{1/2}$$

In general:

if the error decreases as n^{-a} for $d = 1$, then the error decreases as $n^{-a/d}$ in d dimensions.

**Classical formulas with equispaced points:
slowly decreasing error for multidimensional integration !**

Numerical integration: error in MC methods

$$\sigma_n / \sqrt{n} \approx \sigma_m \approx \sigma_s / \sqrt{s}$$

(σ_n is roughly constant with n ; for **uncorrelated points**,
the variance of the averages goes like $\sim 1/n^{1/2}$)

- The average function value

$$\langle f \rangle \equiv \frac{1}{N} \sum_{i=1}^N f(x_i)$$

- The average squared function value

$$\langle f^2 \rangle \equiv \frac{1}{N} \sum_{i=1}^N f^2(x_i)$$

- Estimate of the integrand (+/- standard error)

$$\int f dV \approx V \langle f \rangle \pm V \sqrt{\frac{\langle f^2 \rangle - \langle f \rangle^2}{N}}$$

Numerical integration: errors in multidimensional integrals

d	Rect.	Trap.	Simps.	MC
1	$1/n$	$1/n^2$	$1/n^4$	$1/n^{1/2}$
2	$1/n^{1/2}$	$1/n$	$1/n^2$	$1/n^{1/2}$
4	$1/n^{1/4}$	$1/n^{1/2}$	$1/n$	$1/n^{1/2}$
...

if the error decreases as n^{-a} for $d = 1$, then the error decreases as $n^{-a/d}$ in d dimensions.

the error for all Monte Carlo integration methods decreases as $n^{-1/2}$ independently of the integral.

Monte Carlo convenient for multidimensional integration !

Summary:

advantages of MC integration methods

- convergence as $\sim N^{1/2}$ in any dimension **regardless of the smoothness of the integrand**
- **simplicity**: only two simple steps required (namely, producing a set of sampling points and evaluating the integrand function over such points)
- **generality**: sampling can be used even on domains that do not have a natural correspondence with the 'standard' domain $[0, 1]^d$ and thus are not well-suited to numerical quadrature
- better suited than quadrature for **integrands with singularities** (importance sampling can handle this problem)
- **flexibility**: easy to add more points as needed (in the Gaussian quadrature, increasing the accuracy implies doing calculations from scratch)