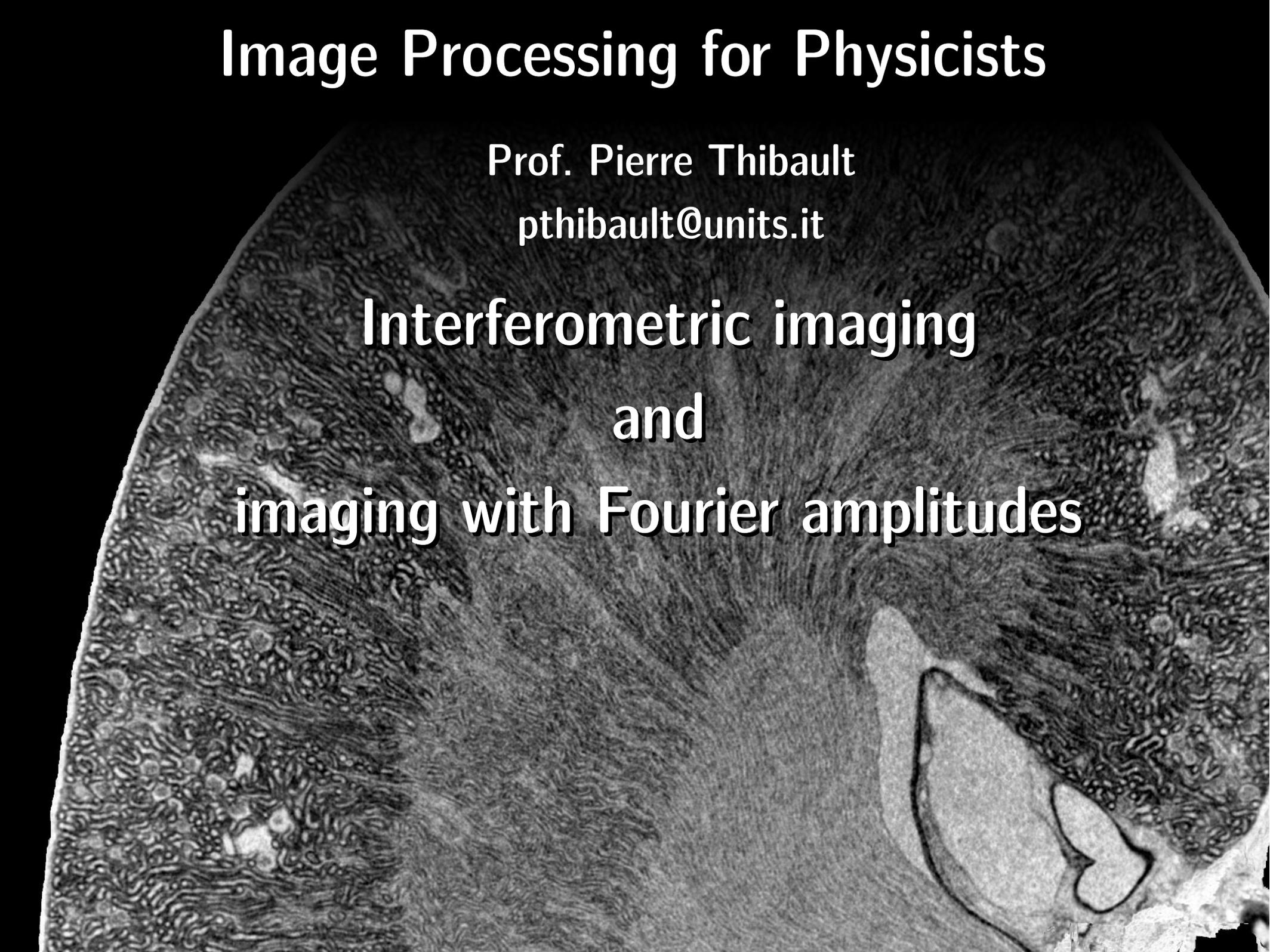


# Image Processing for Physicists

Prof. Pierre Thibault

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Interferometric imaging  
and  
imaging with Fourier amplitudes



# Overview

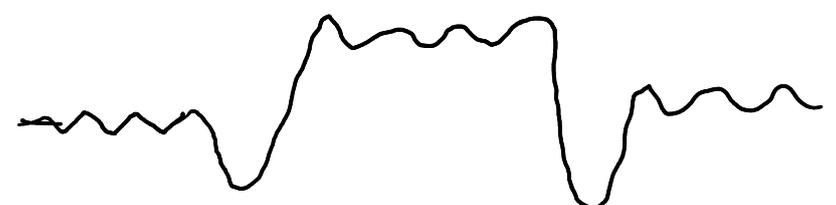
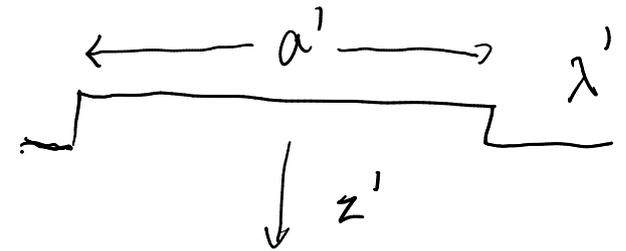
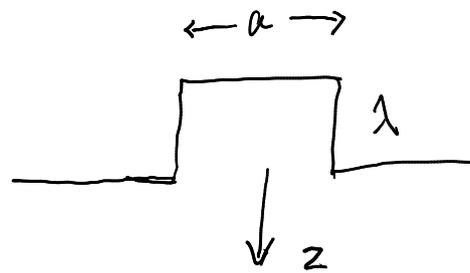
- The phase problem
- Holography: on/off-axis
- Grating interferometric imaging
- Imaging using far-field amplitude measurements
  - Fourier transform holography
  - Coherent diffraction imaging
  - Ptychography

# Wave propagation



Reminder: angular spectrum formulation

$$\psi(r; z) = \mathcal{F}^{-1} \left\{ \mathcal{F} \{ \psi(r; z=0) \} \exp(i\pi \underbrace{u^2 \lambda z}_{\text{unitless number}}) \right\}$$



identical if  $\frac{1}{a^2} \lambda z = \frac{1}{a'^2} \lambda' z'$

$\sqrt{\lambda z}$  = characteristic length

$$\frac{a^2}{\lambda z} = f$$

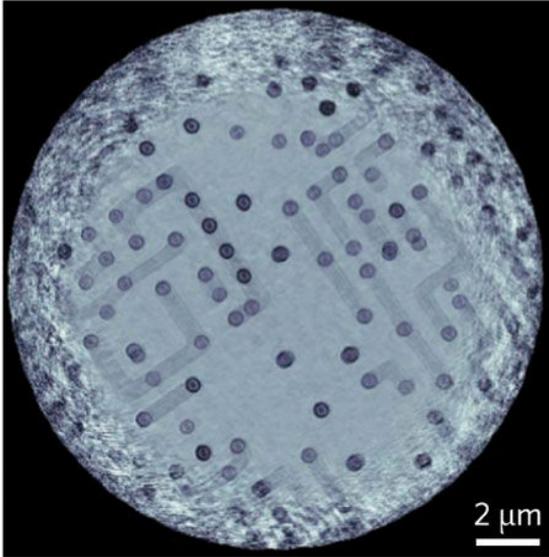
Fresnel number

$f \ll 1$ ; far-field

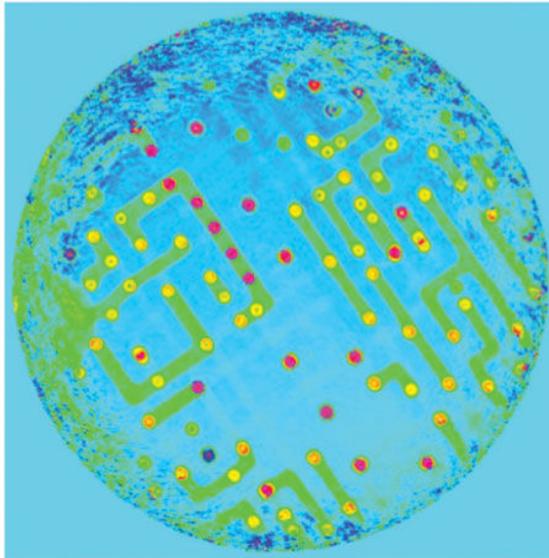
$f \gg 1$ ; near-field

# Complex-valued images

X-ray transmission image



Amplitude  
attenuation  
of the wave

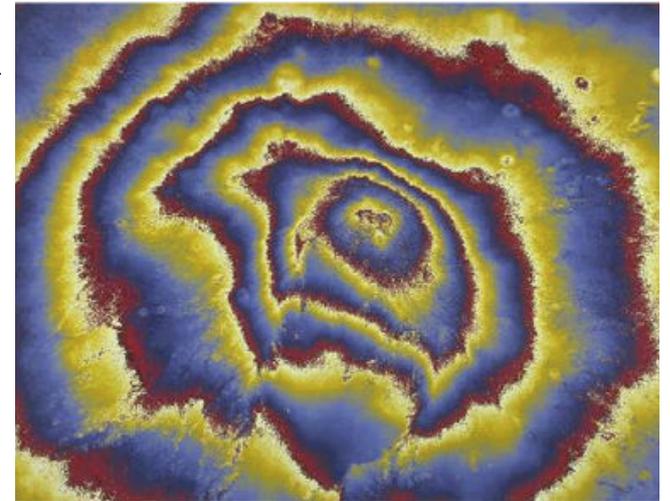


phase  
delay in  
the wavefield  
(refraction)

synthetic aperture radar



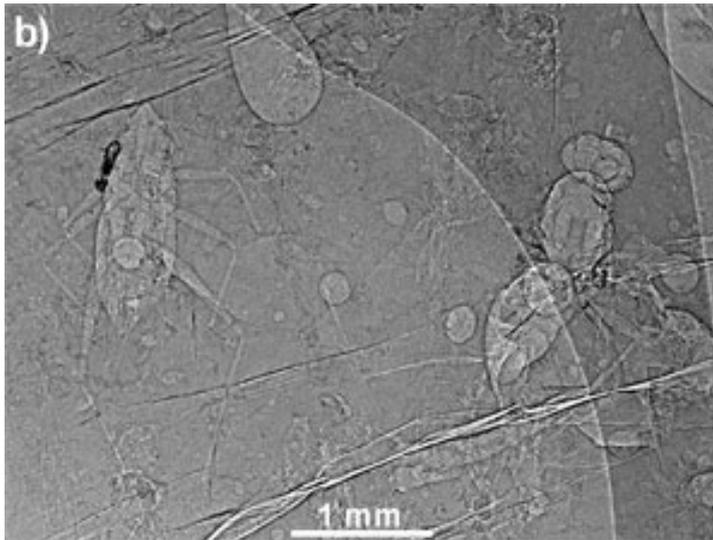
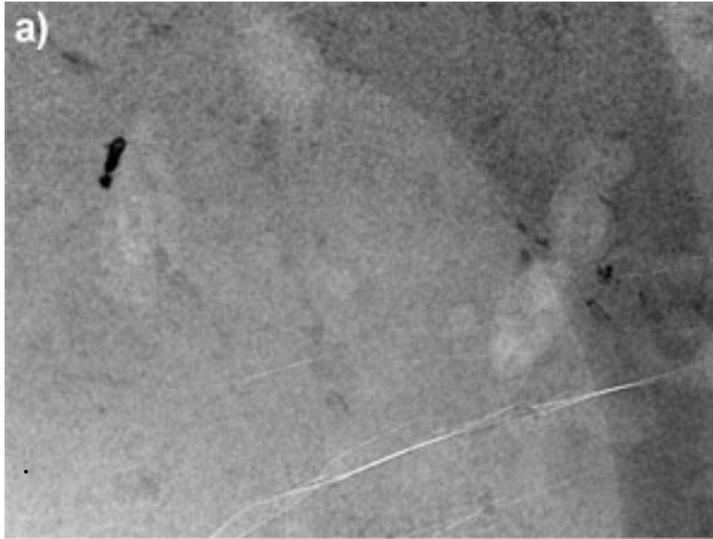
phase  
unwrapping →



phase  
image →

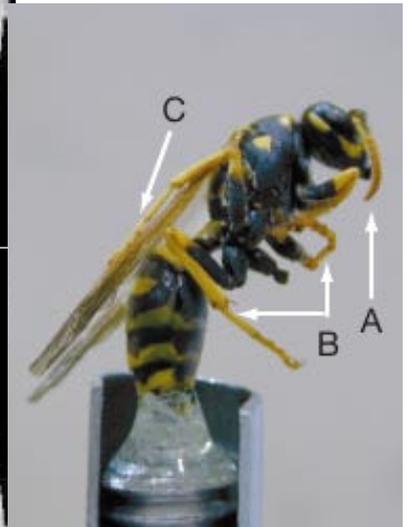
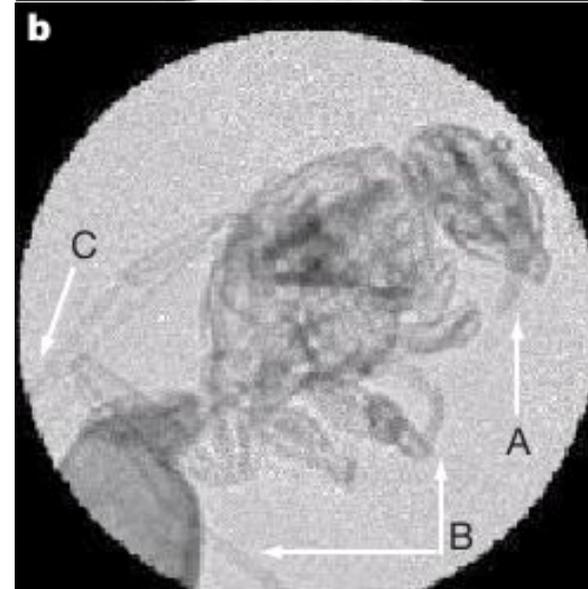
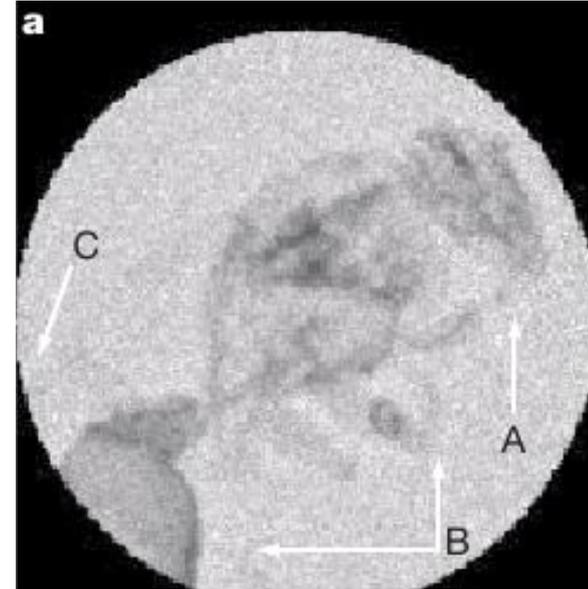
# Phase-contrast

Hard X-ray propagation-based  
phase contrast



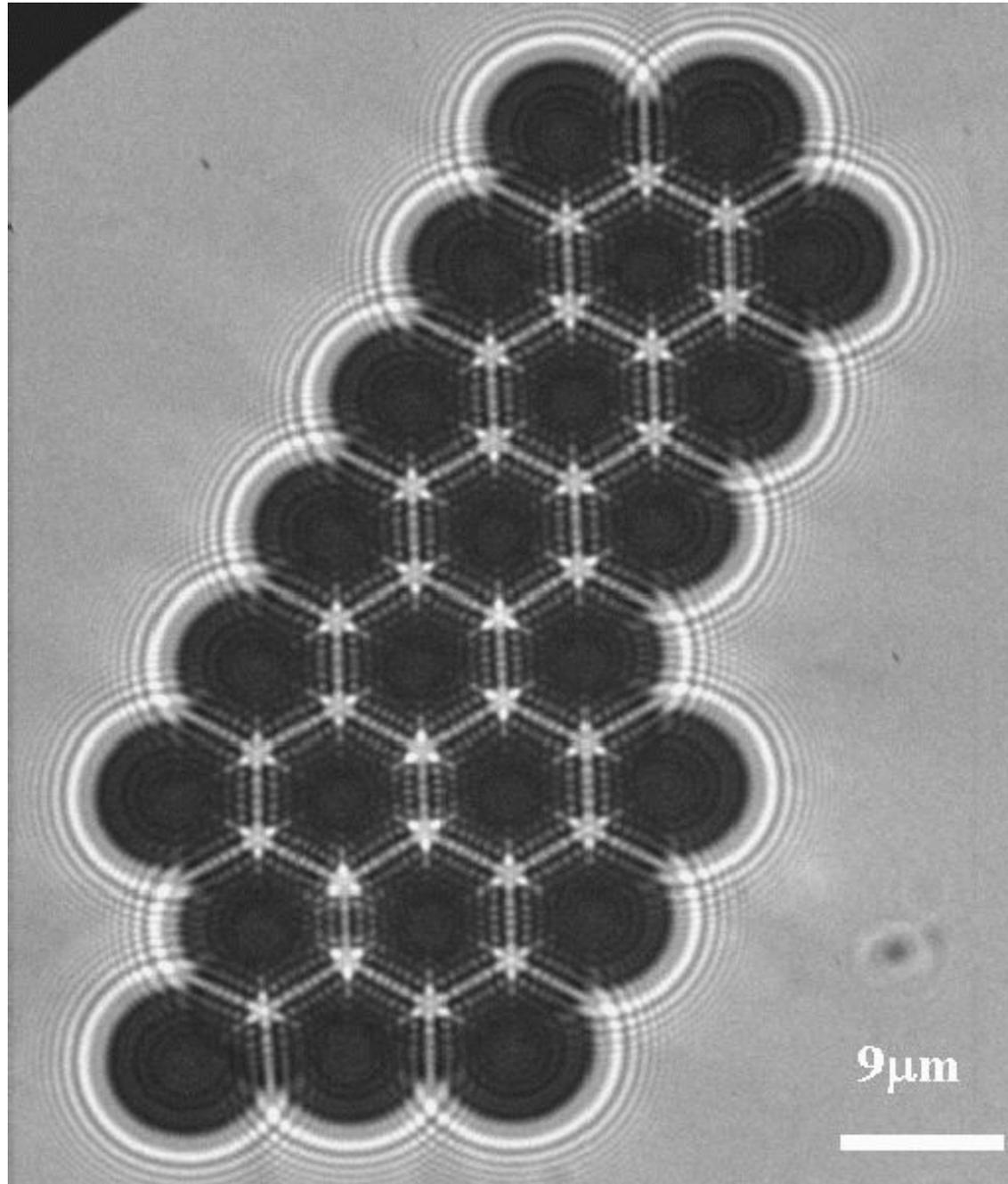
Source:  
[www.esrf.eu/news/general/amber/amber/](http://www.esrf.eu/news/general/amber/amber/)

Neutron phase contrast



Source: Allman et al. Nature **408** (2000).

# Inline holography



Source: Mayo et al. Opt Express 11 (2003).

# Inline holography

Measure  $I(\vec{r}) = |\psi(r; z)|^2$

\* plane monochromatic wave

\* weak transmission of imaged object

object plane  $\psi(r; z=0) = A (1 + \epsilon(r; z=0))$

constant (plane wave)  $\epsilon(r; z=0)$  small perturbation

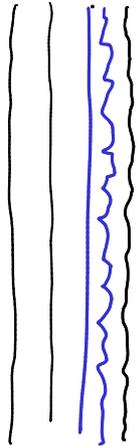
$$I(\vec{r}, z) = |A (1 + \epsilon(r; z))|^2 = |A|^2 (1 + \epsilon(r; z) + \epsilon^*(r; z) + |\epsilon(r; z)|^2)$$

$$\approx |A|^2 (1 + \underbrace{\epsilon(\vec{r}; z)}_{\text{propagated by } z} + \underbrace{\epsilon^*(r; z)}_{\text{equivalent to propagation by } -z})$$

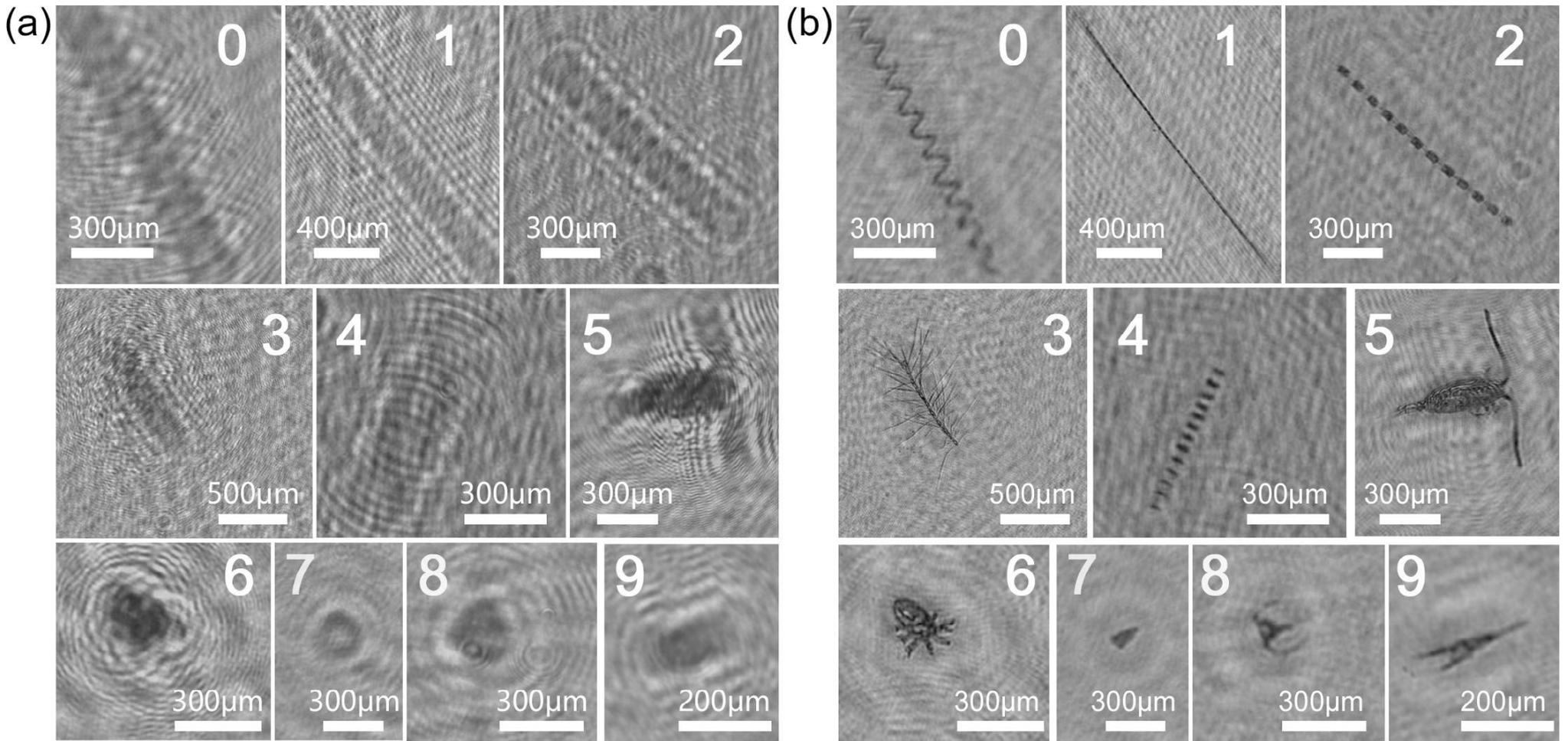
propagated by  $z$

equivalent to propagation by  $-z$

"twin image" problem



# Digital inline holography



trick: apply angular spectrum propagator

# The phase problem

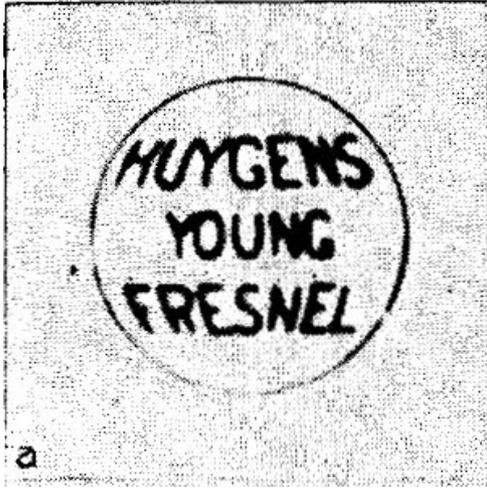
We always measure  $|\psi|^2$ . phases are lost.

We need to recover the phase part of the wavefield

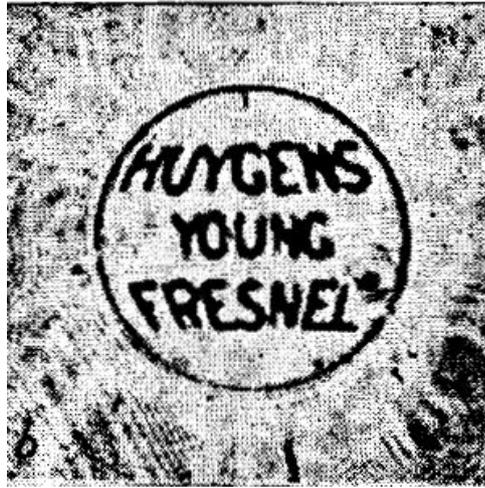
\* sometimes the phase is interesting

\* most often: the phase is an auxiliary quantity for proper interpretation of the wavefield.

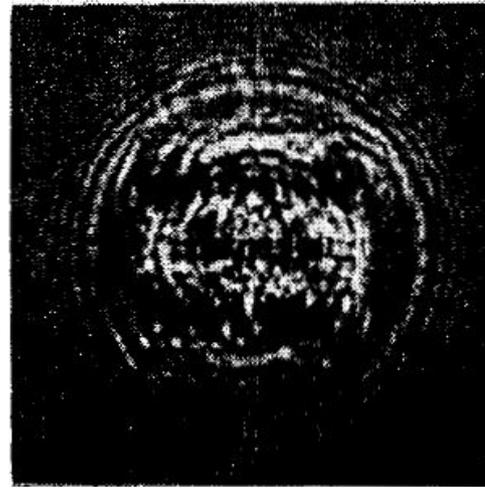
# In-line holography



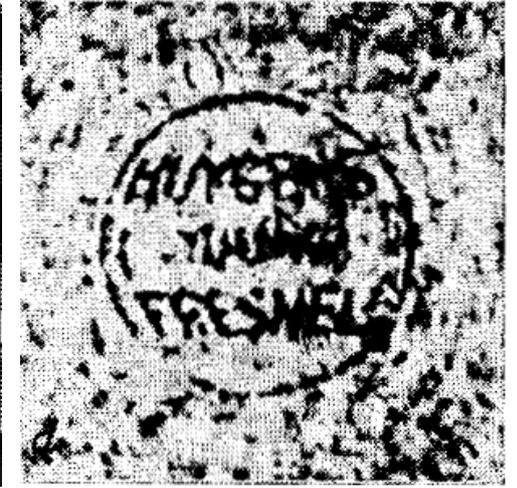
↑  
mask



↑  
"in focus"



↑  
after propagation



↑  
D. Gabor, Nature **161**, 777-778 (1948).  
↓  
propagated

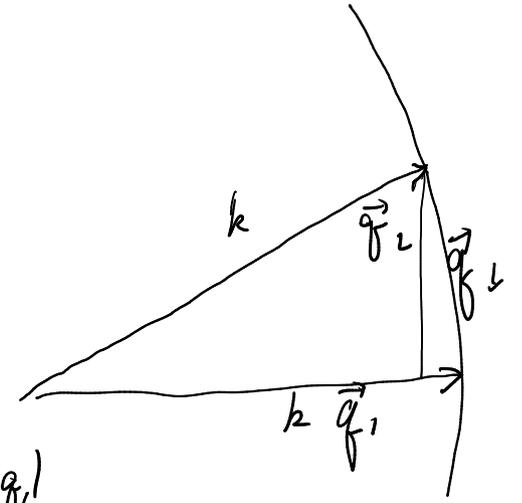
↙  
twin image  
degrades quality

# Tilted reference

Two tilted plane waves:  

$$e^{i\vec{q}_1 \cdot \vec{r}} + e^{i\vec{q}_2 \cdot \vec{r}} = e^{ikz} \left( 1 + e^{i\vec{q}_1 \cdot \vec{r}_1} \right)$$

propagates at  
 angle  $\sin\theta = \frac{|\vec{q}_1|}{k}$



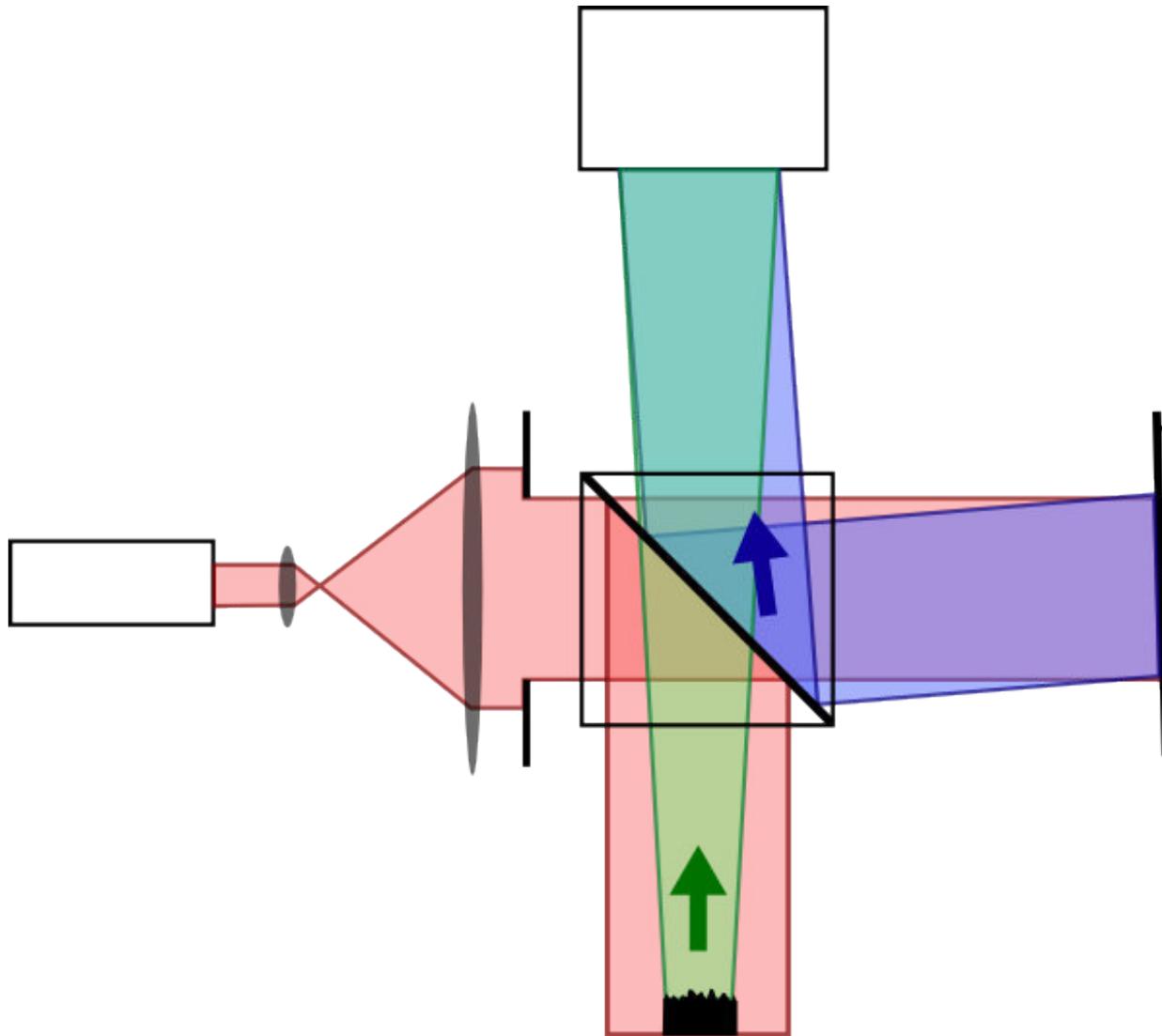
$$I = \left| 1 + e^{i\vec{q}_1 \cdot \vec{r}} \right|^2 = \left( 1 + e^{i\vec{q}_1 \cdot \vec{r}} + e^{-i\vec{q}_1 \cdot \vec{r}} + 1 \right) = 2 + 2 \cos(\vec{q}_1 \cdot \vec{r})$$

oscillations ("fringes") with spatial frequency

$$\vec{u} = \frac{\vec{q}_1}{2\pi}$$

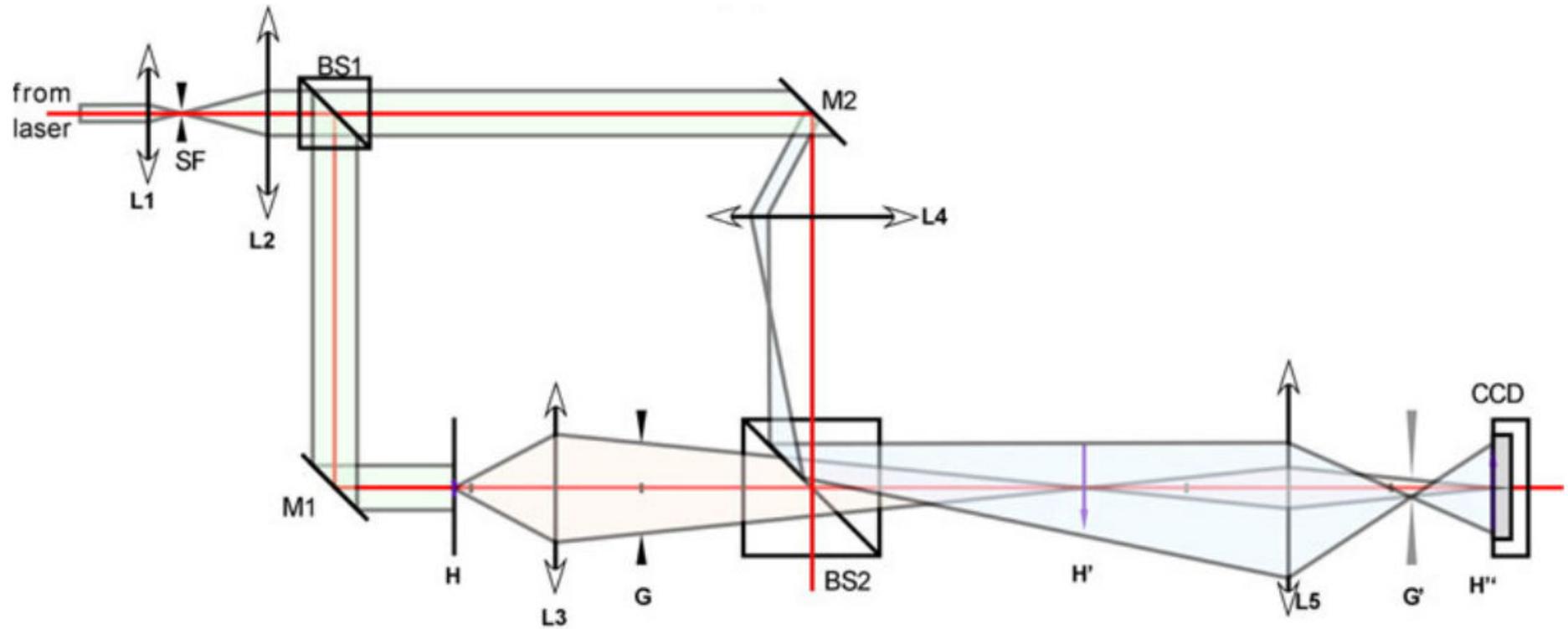
with a sample:  $\left| a(\vec{r}) + e^{i\vec{q}_1 \cdot \vec{r}} \right|^2$

# Fringe interferometry



Twyman-Green interferometer

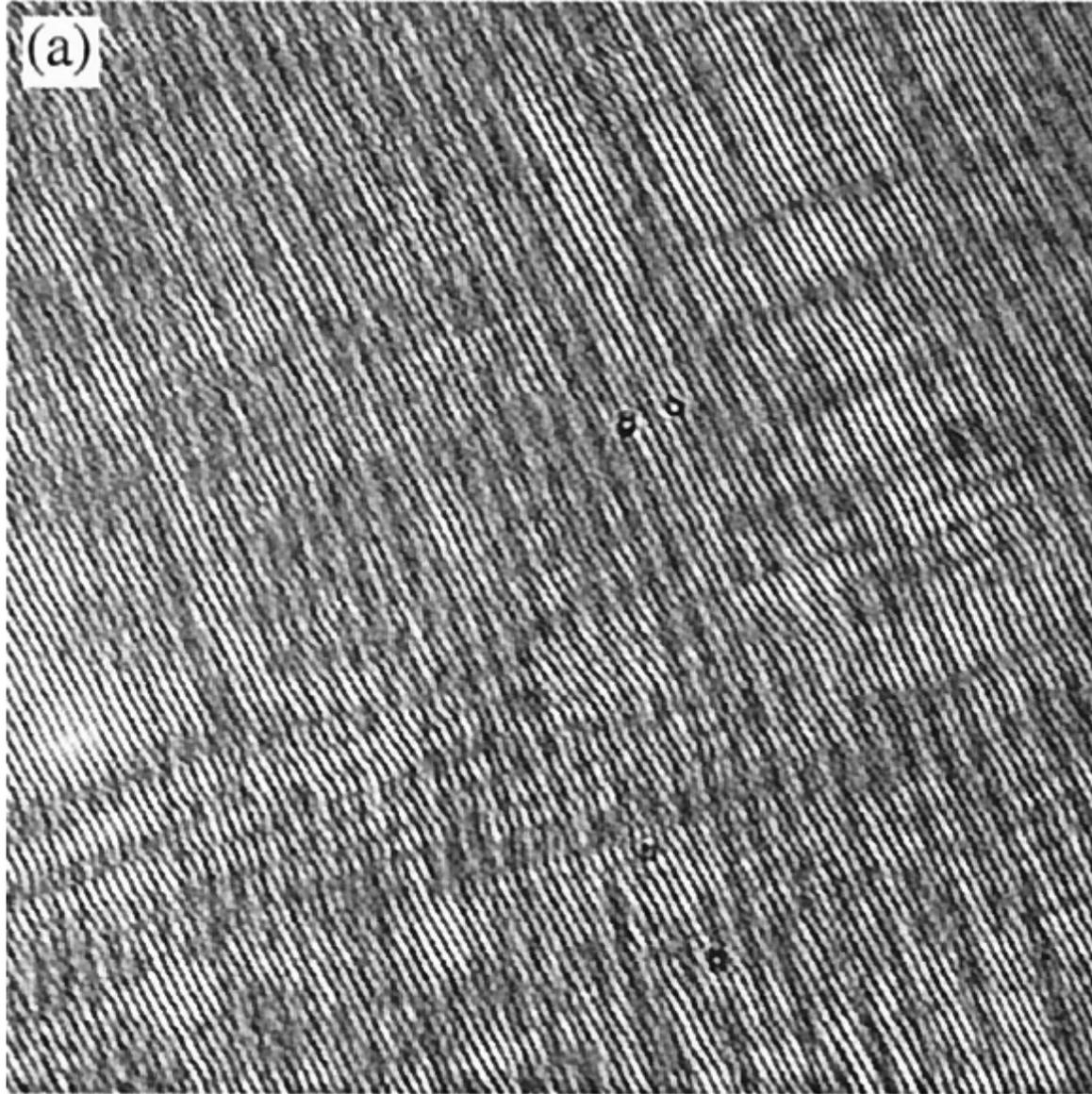
# Visible light interferometer



## Mach-Zehnder interferometer

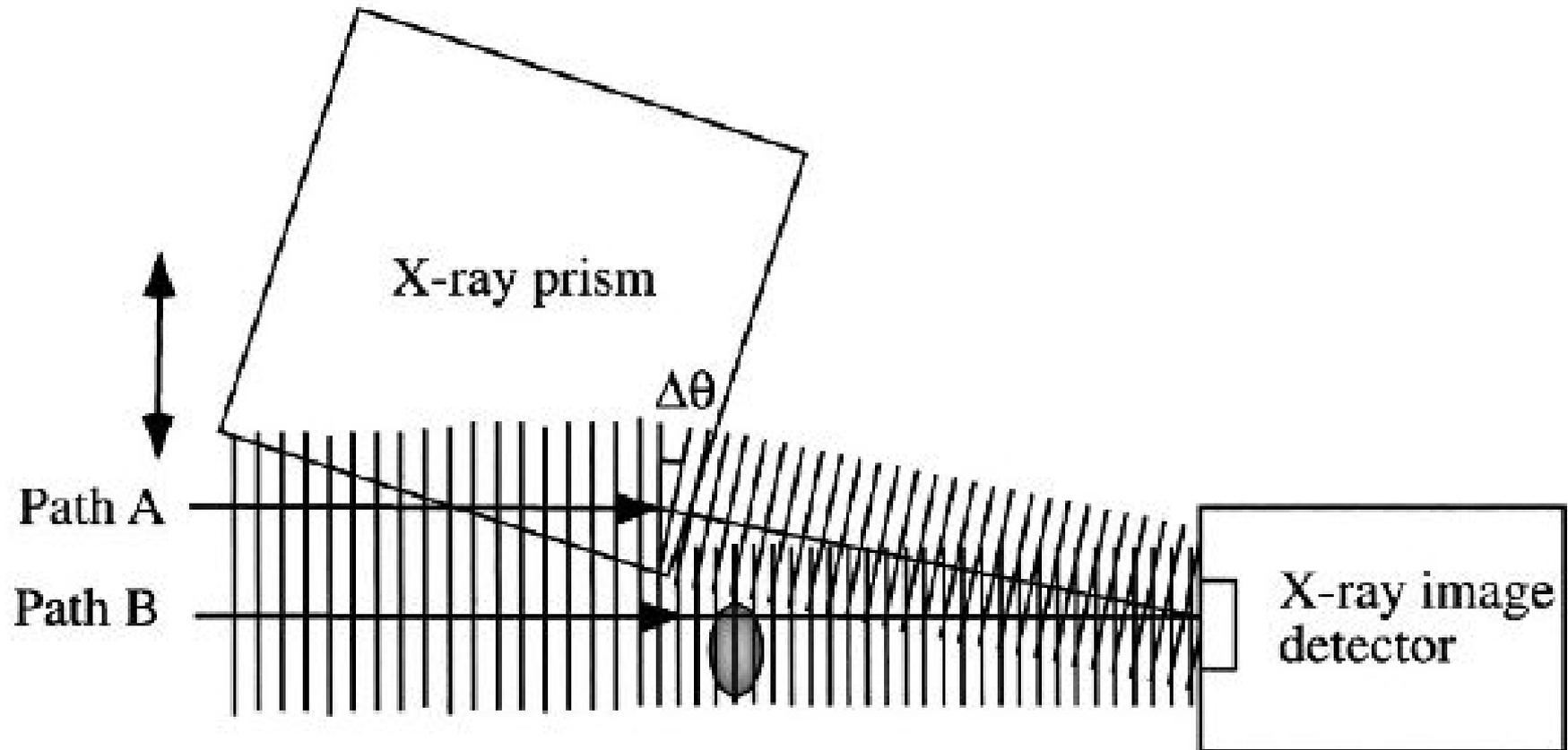
Source: M. K. Kim, SPIE Rev. 1, 018005 (2010).

# Fringe interferometry



Source: Cucho et al. *Appl. Opt.* **39**, 4070 (2000)

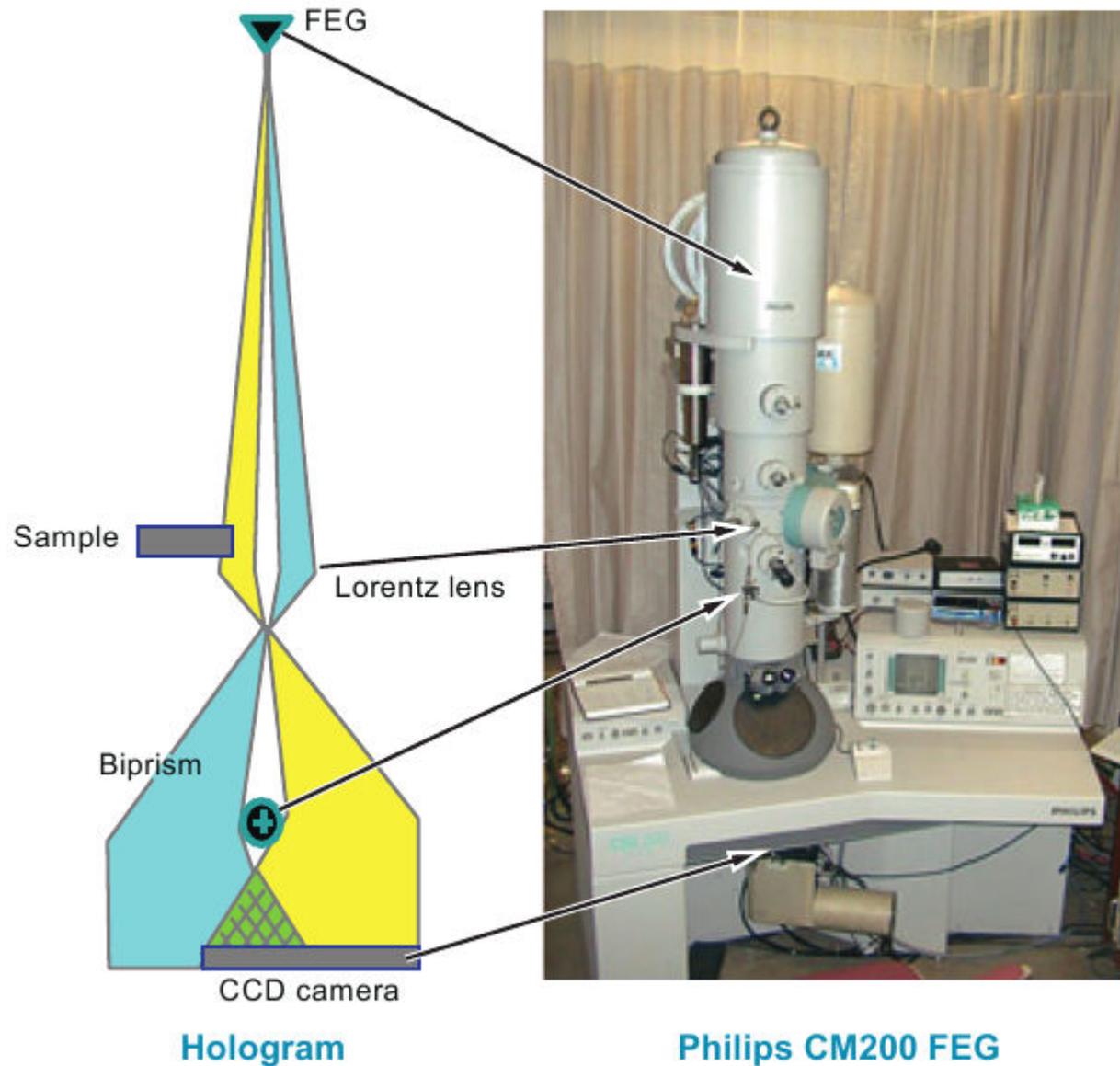
# Off-axis X-ray holography



Source: Y. Kohmura, J. Appl. Phys. **96**, 1781-1784 (2004)

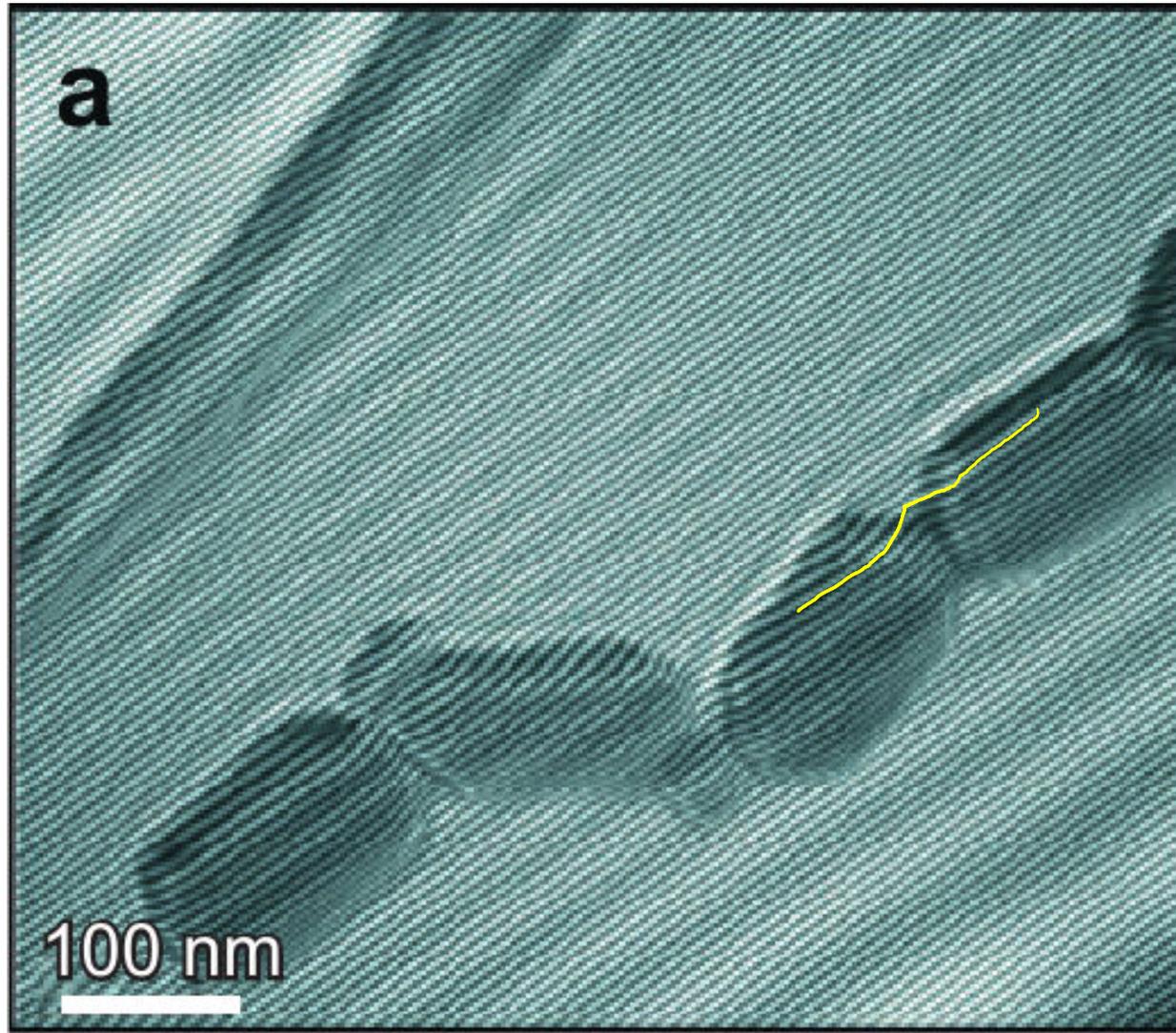
# Off-axis electron holography

## Electron microscopy



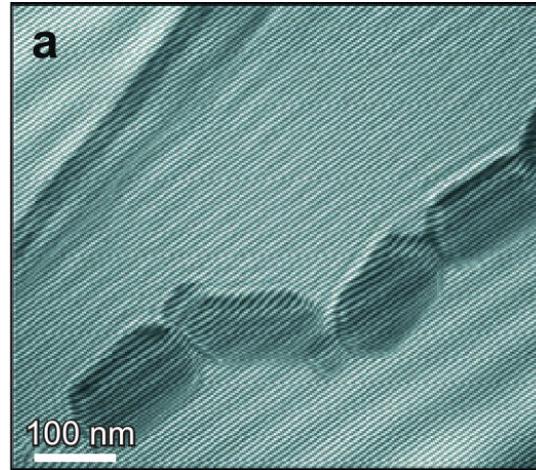
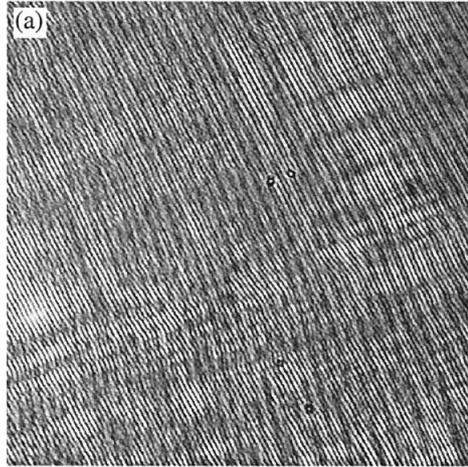
Source: M. R. McCartney, *Ann. Rev. Mat. Sci.* **37** 729-767 (2007)

# Off-axis electron holography



Source: M. R. McCartney, *Annu. Rev. Mat. Sci.* **37** 729-767 (2007)

# Fringe interferometry



$a \rightarrow$  attenuation  
 $\varphi \rightarrow$  phase shift  
 (refraction)

$$\psi = \psi_o + \psi_r$$

object      reference

$$\psi_r(\vec{r}) = A e^{i\vec{q}_\perp \cdot \vec{r}}$$

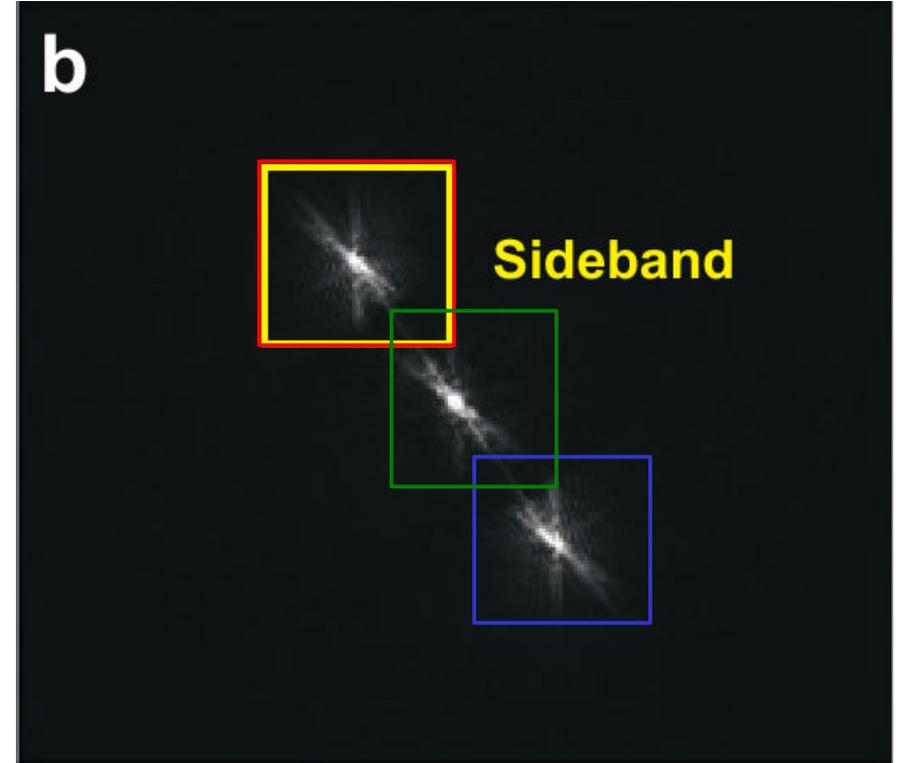
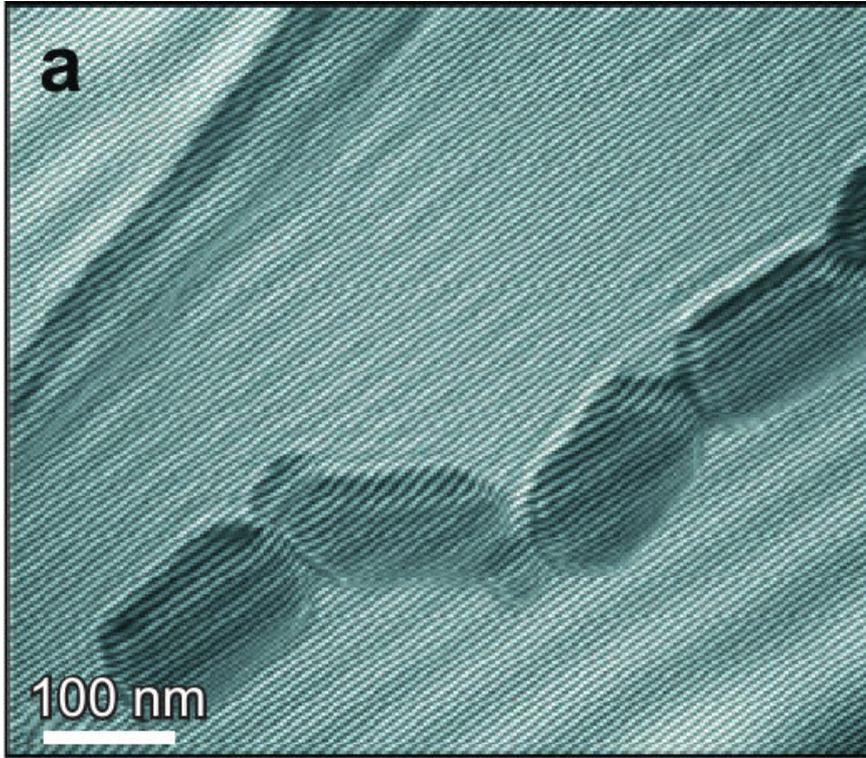
$$\psi_o(\vec{r}) = A a(\vec{r}) e^{i\varphi(\vec{r})}$$

Measurement:

$$|\psi|^2 = (\psi_o + \psi_r)(\psi_o + \psi_r)^*$$

$$= |A|^2 \left( 1 + a^2(\vec{r}) + \underbrace{2a(\vec{r}) \cos(\vec{q}_\perp \cdot \vec{r} - \varphi)}_{\substack{a(\vec{r}) e^{i(\varphi - \vec{q}_\perp \cdot \vec{r})} \\ + a(\vec{r}) e^{-i(\varphi - \vec{q}_\perp \cdot \vec{r})}} \right)$$

# Off-axis holography



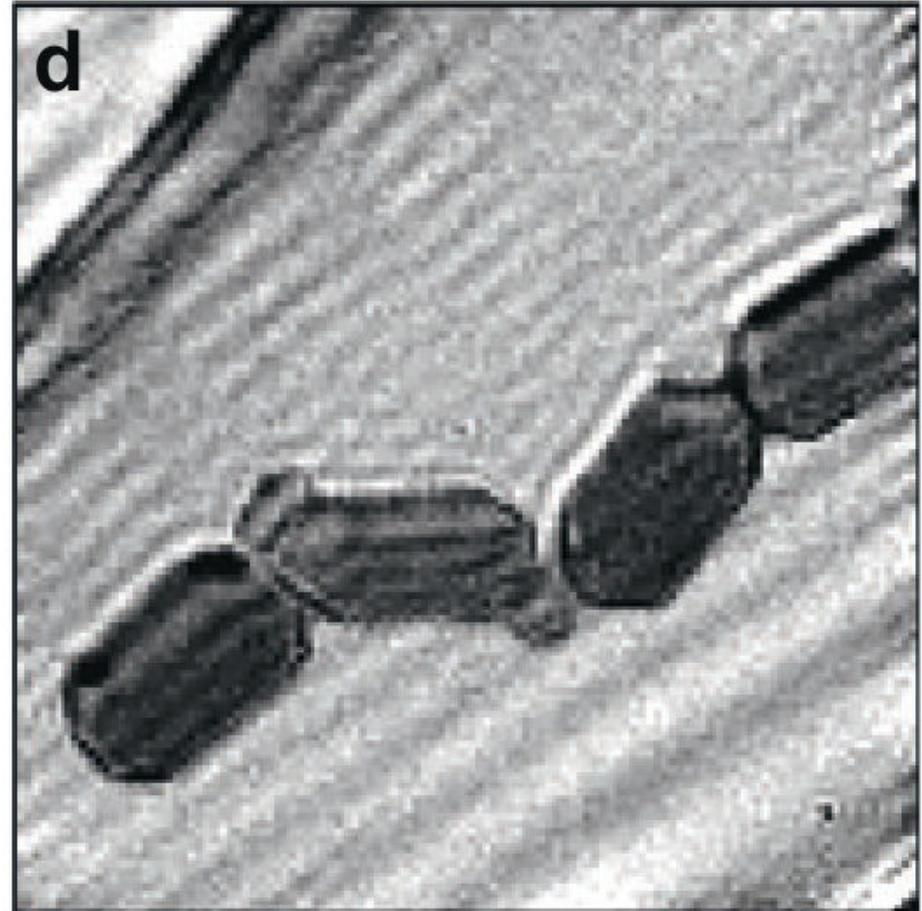
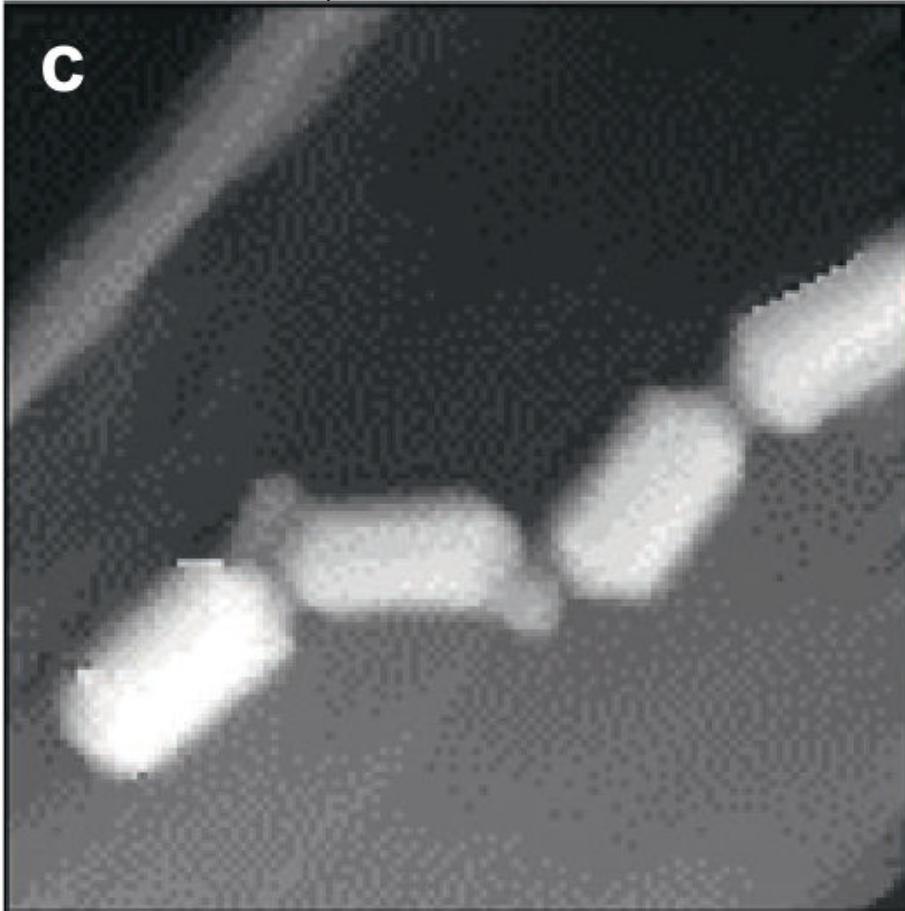
$$\mathcal{F}\{| \psi |^2\} = |A^2| \left[ \mathcal{F}\{a^2(r) + 1\} + \mathcal{F}\{\psi_0\} \left( \vec{u} + \frac{\vec{q}_\perp}{2\pi} \right) + \mathcal{F}\{\psi_0^*\} \left( \vec{u} - \frac{\vec{q}_\perp}{2\pi} \right) \right]$$

Source: M. R. McCartney, Annu. Rev. Mat. Sci. **37** 729-767 (2007)

# Off-axis holography

phase  $\varphi$

attenuation ( $a$ )



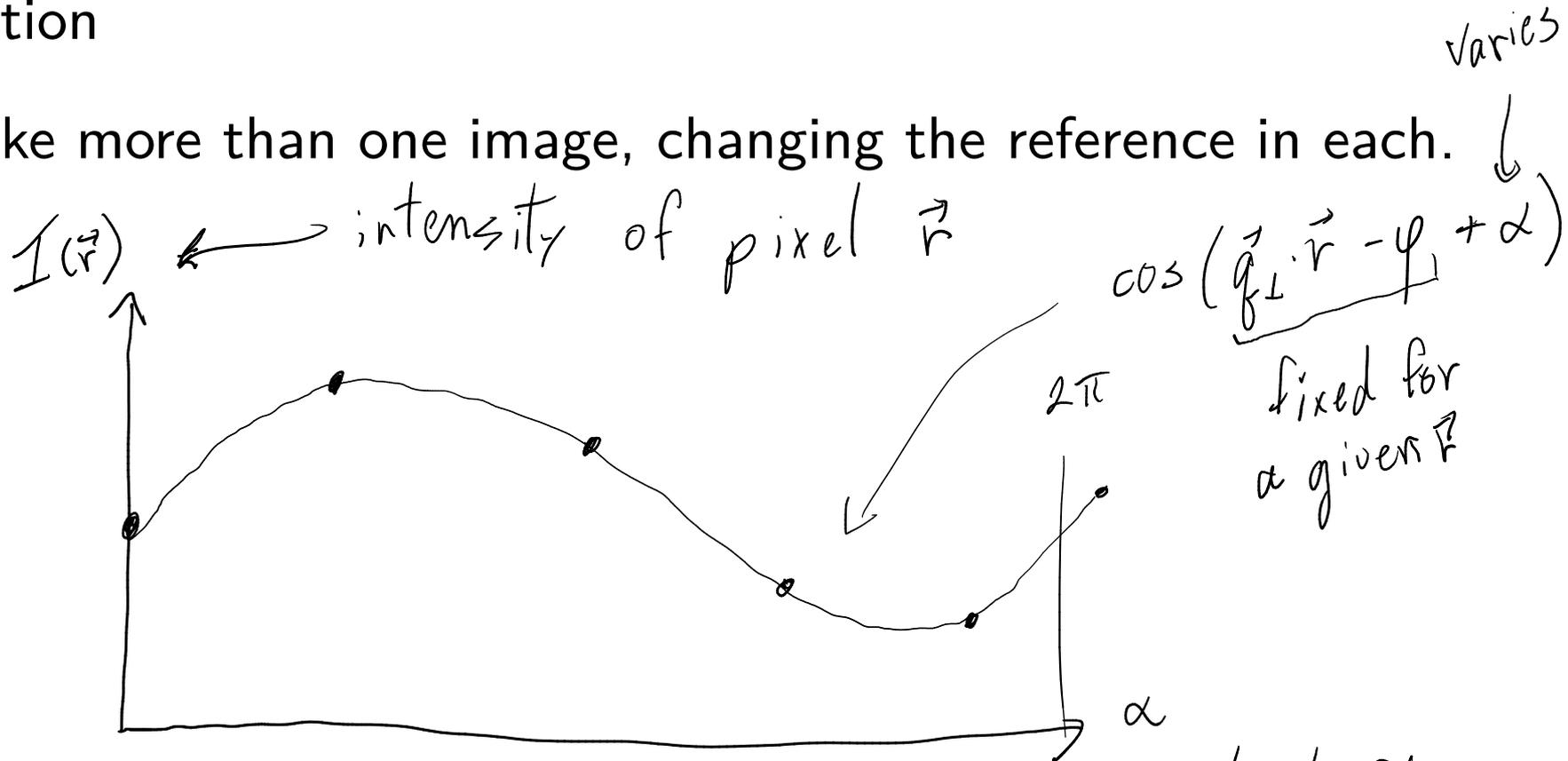
Price paid: loss in resolution

Source: M. R. McCartney, Annu. Rev. Mat. Sci. **37** 729-767 (2007)

# Phase stepping

- Encoding phase **and** amplitude in a single image has a price: resolution

→ Take more than one image, changing the reference in each.

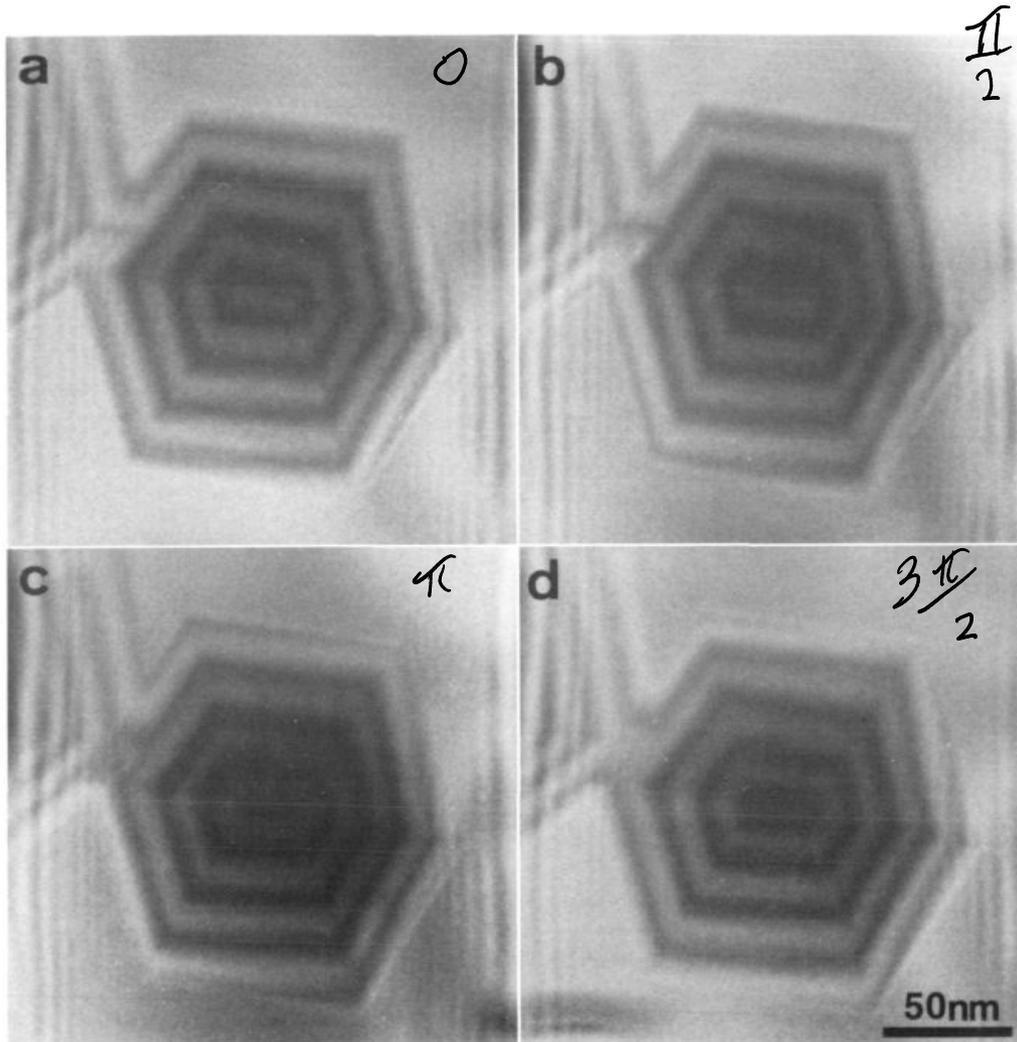


$$\psi_r = A e^{i(\vec{q}_\perp \cdot \vec{r} + \alpha)} = A e^{i(\vec{q}_\perp \cdot (\vec{r} + \vec{s}))}$$

physical shift

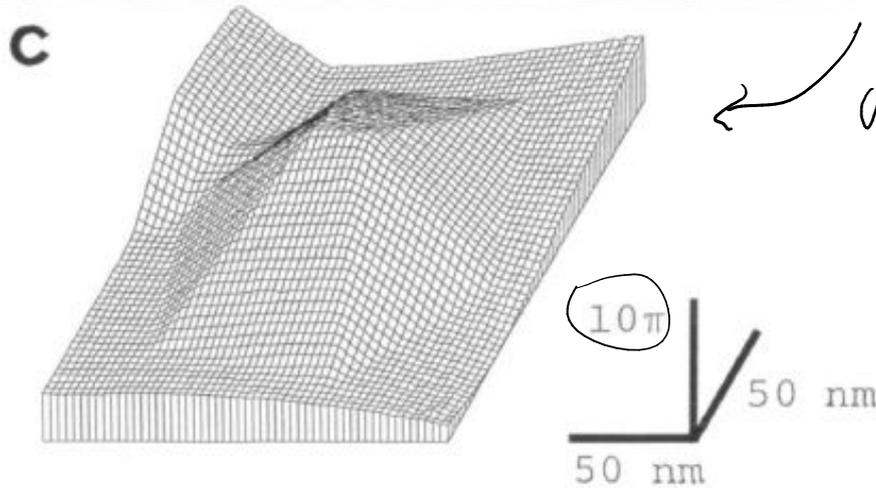
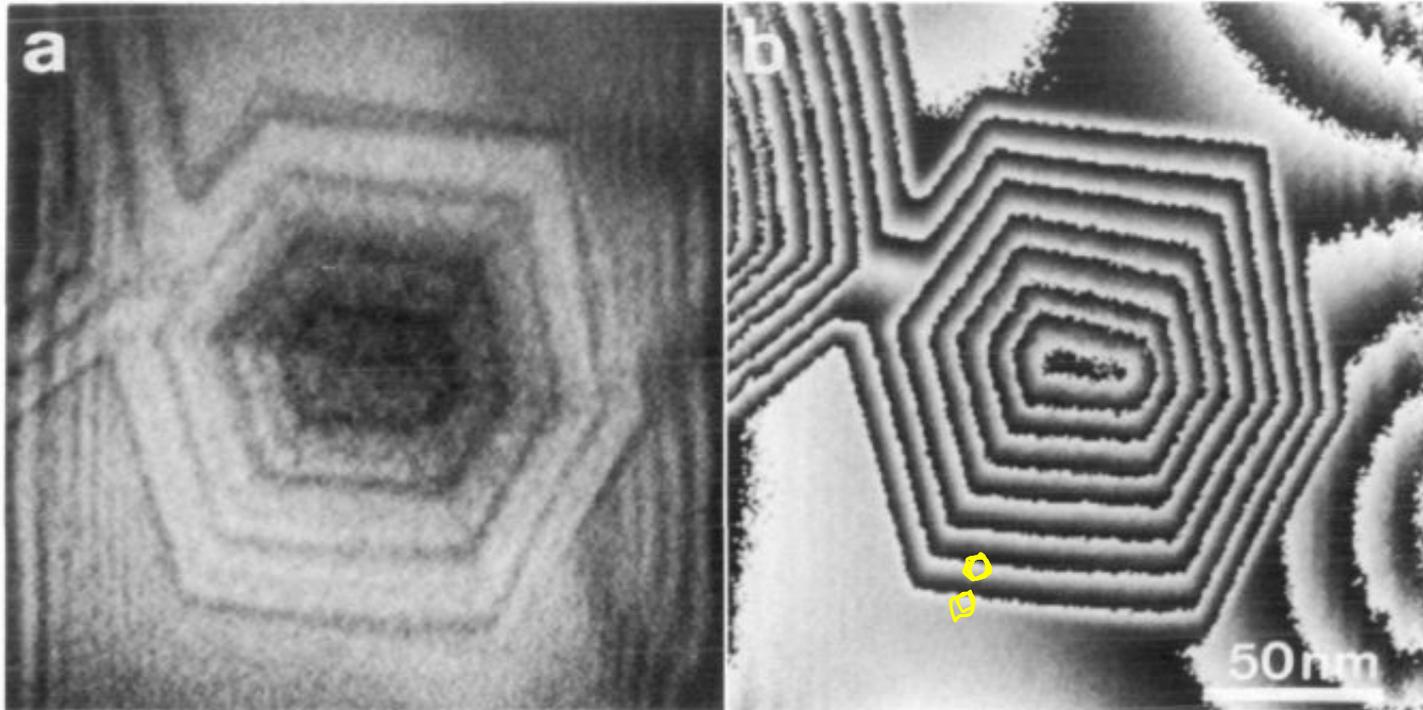
# Fringe scanning

Transmission  
Electron  
Microscopy



Source: K. Harada, J. Electron Microsc. **39** 470-476 (1990)

# Fringe scanning

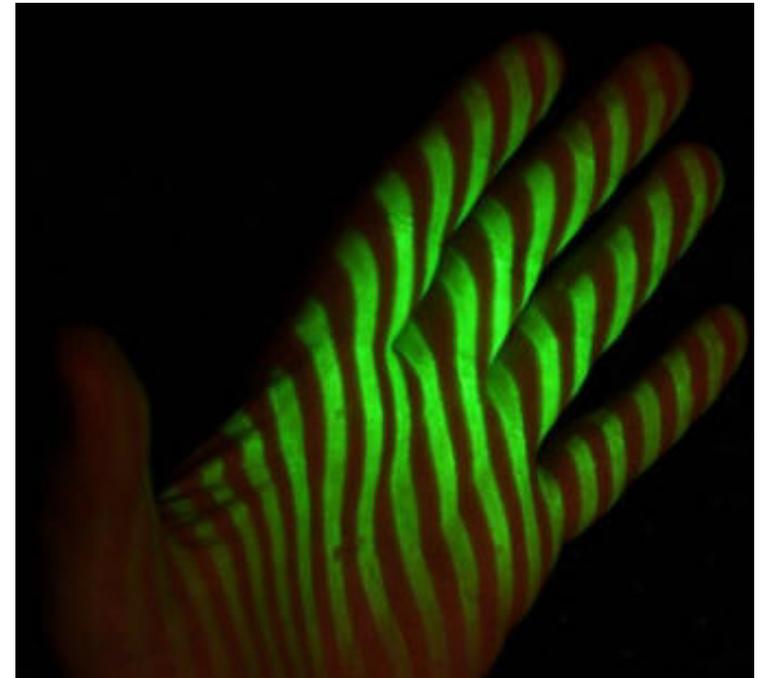
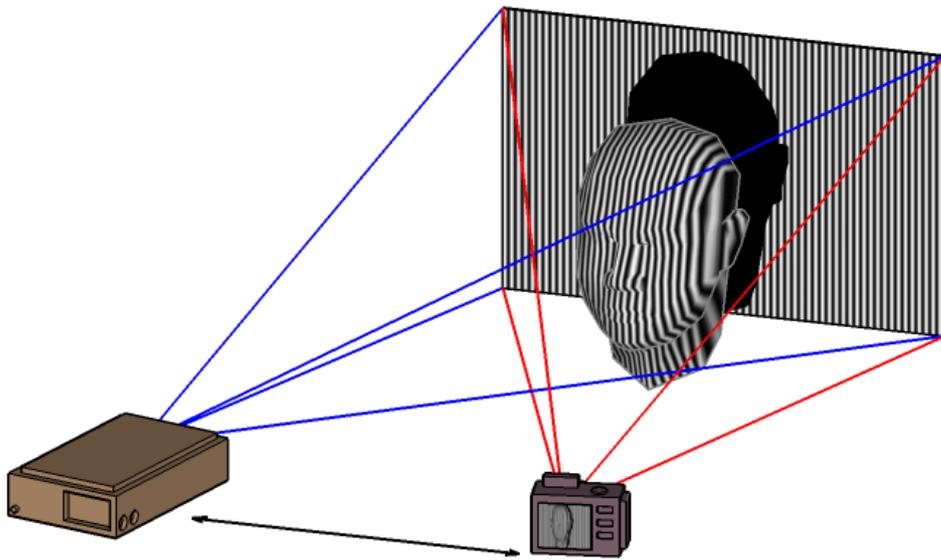


native resolution ✓  
price paid: more measurements required

Source: K. Harada, J. Electron Microsc. **39** 470-476 (1990)

# Structured light sensing

- Project a structured light pattern onto sample
- Distortions of light pattern allow reconstruction of sample shape



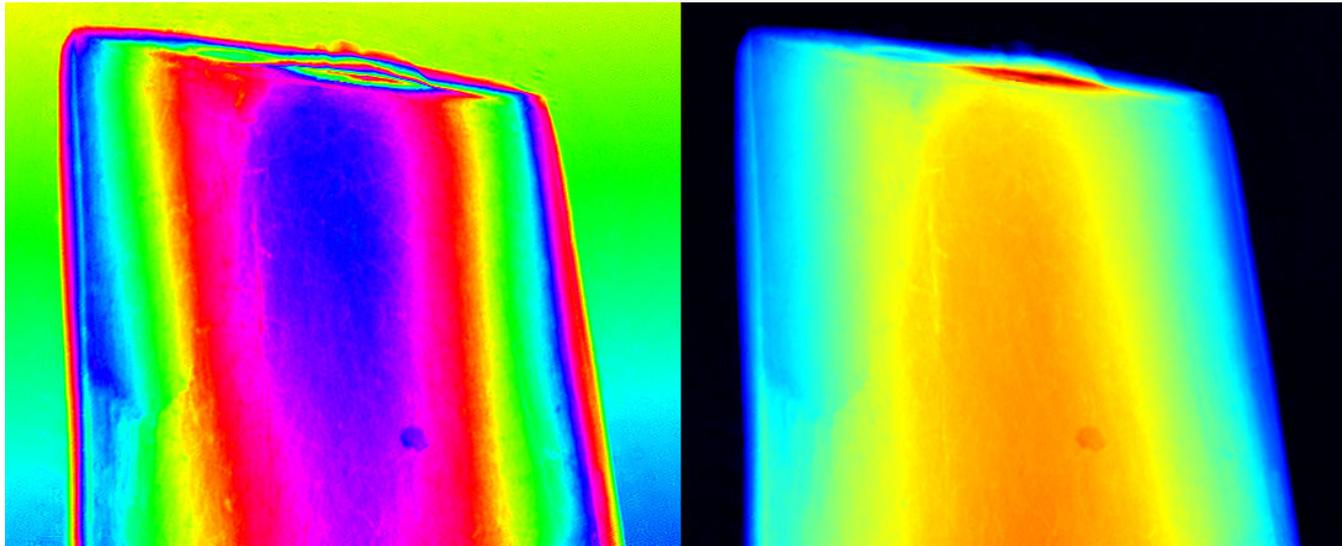
# Phase unwrapping

- Phase is measured only in the interval  $[0, 2\pi)$
- Physical phase shifts (which can be larger) are wrapped on this interval
  - Any multiple of  $2\pi$  is possible
- Unwrapping: use correlations in the image to guess the total phase shift.
- Main difficulties:
  - aliasing: phase shifts are too rapid for the image sampling
  - noise: produces local singularities (vortices)
- Many strategies exist

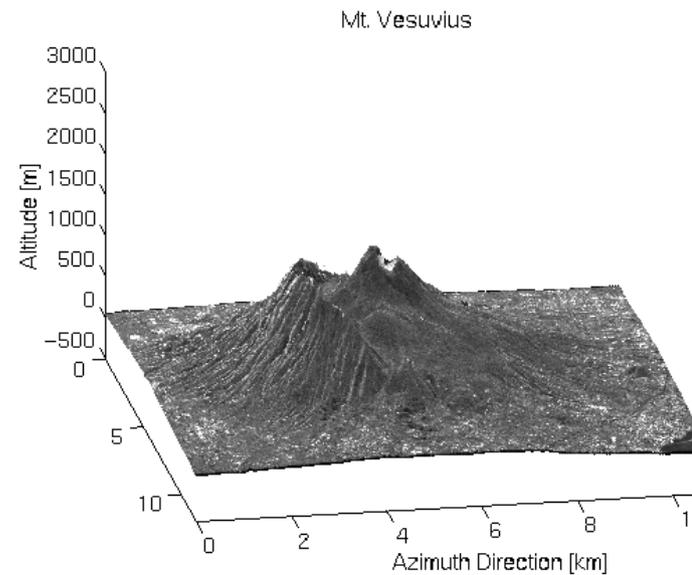
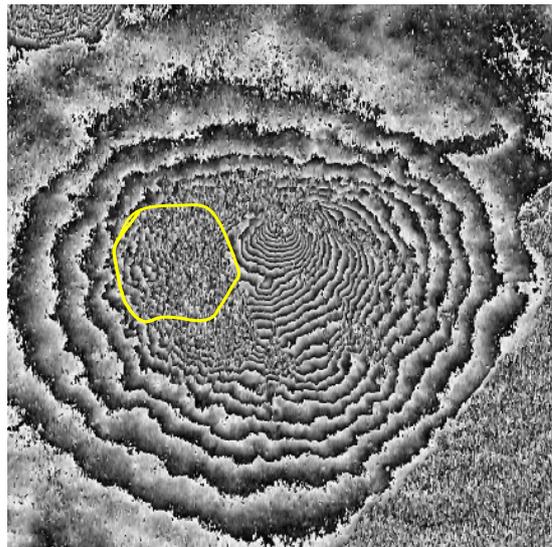
# Complex-valued images

## Phase unwrapping

"easy"



aliasing

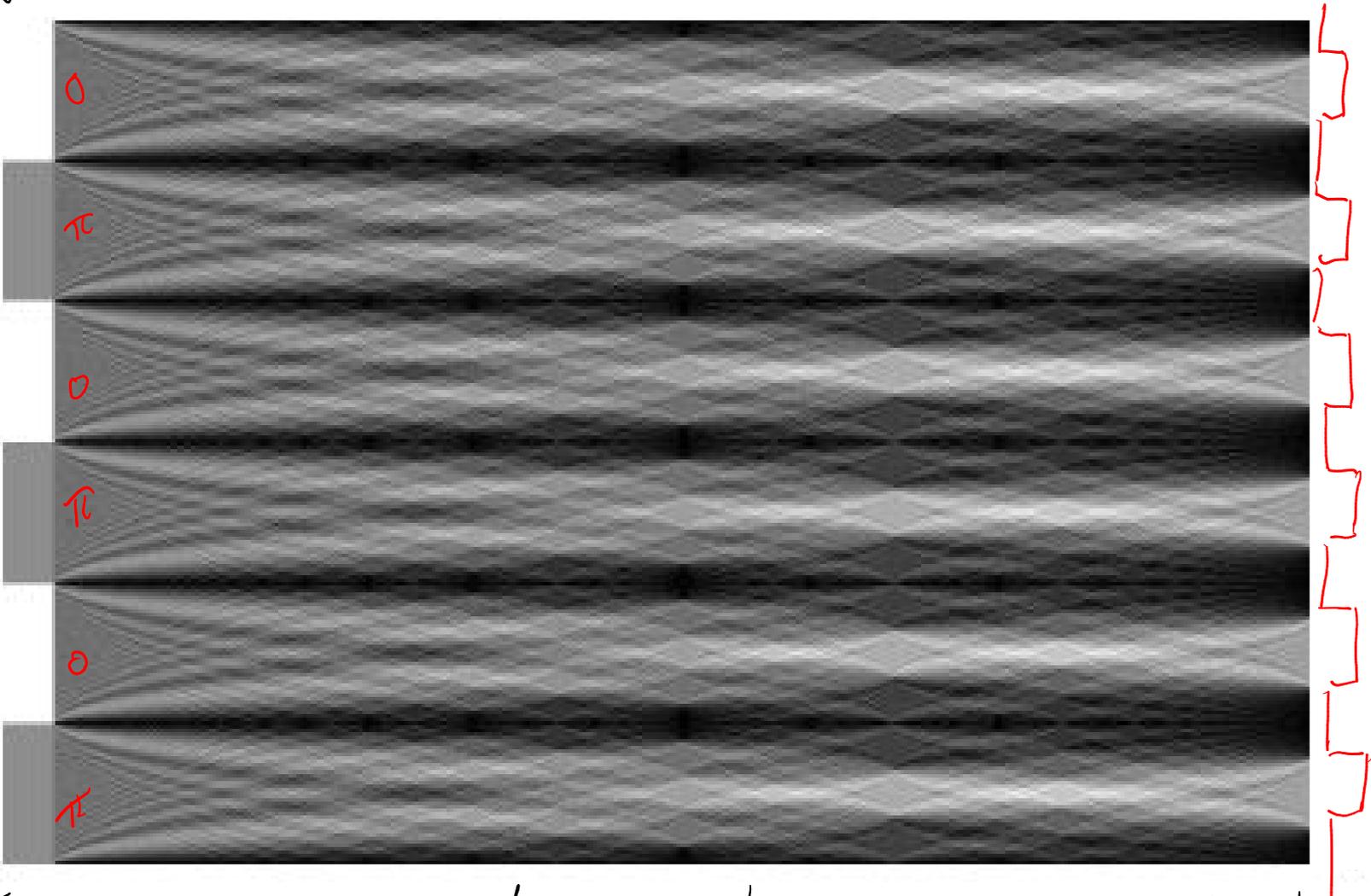


Source: <http://earth.esa.int/workshops/ers97/program-details/speeches/rocca-et-al/>

# Grating interferometry

## Diffraction from a grating

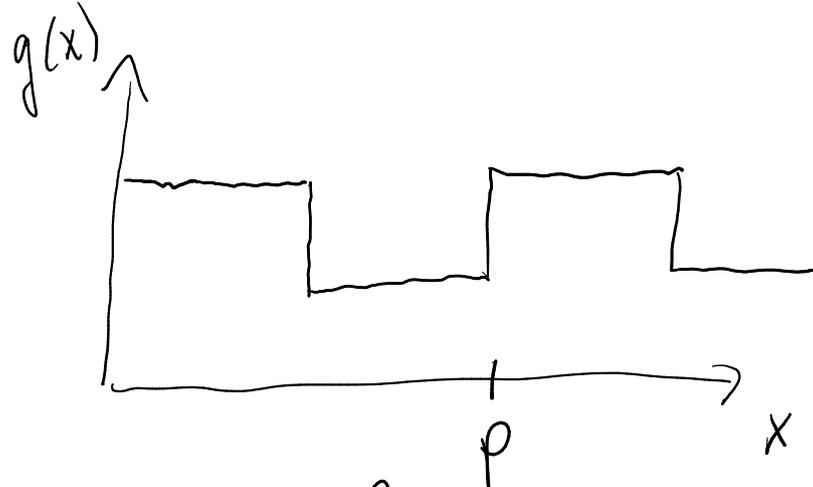
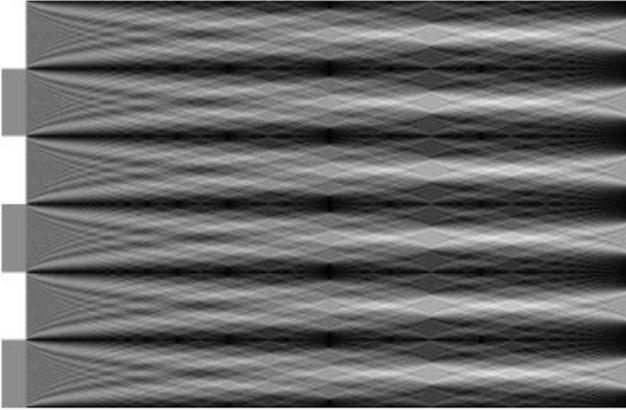
Phase grating



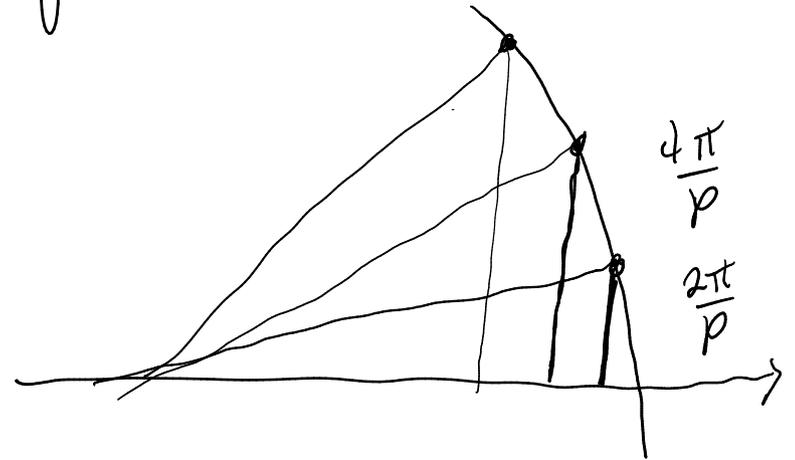
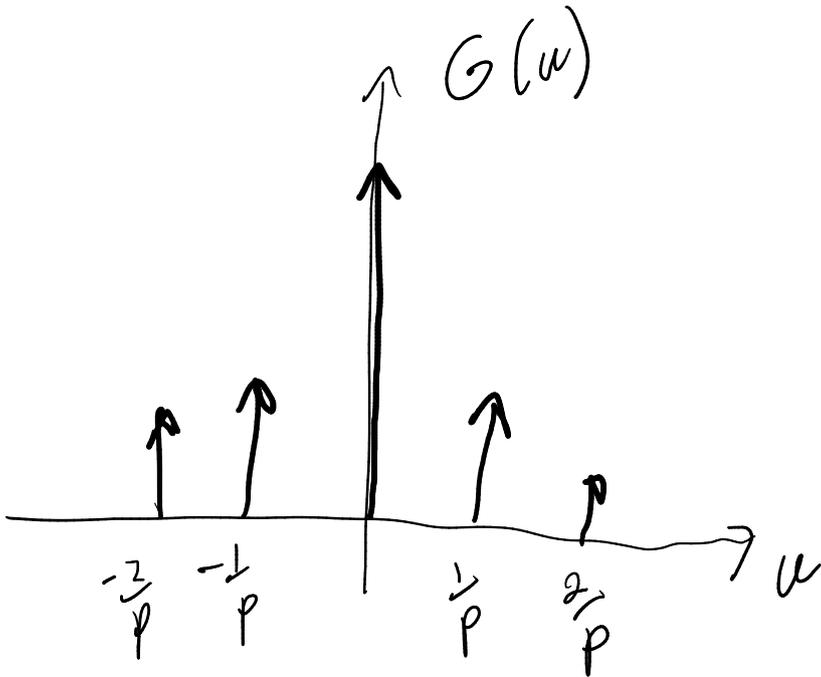
Talbot effect: periodic structure reforms periodically through propagation

# Grating interferometry

## Diffraction from a grating

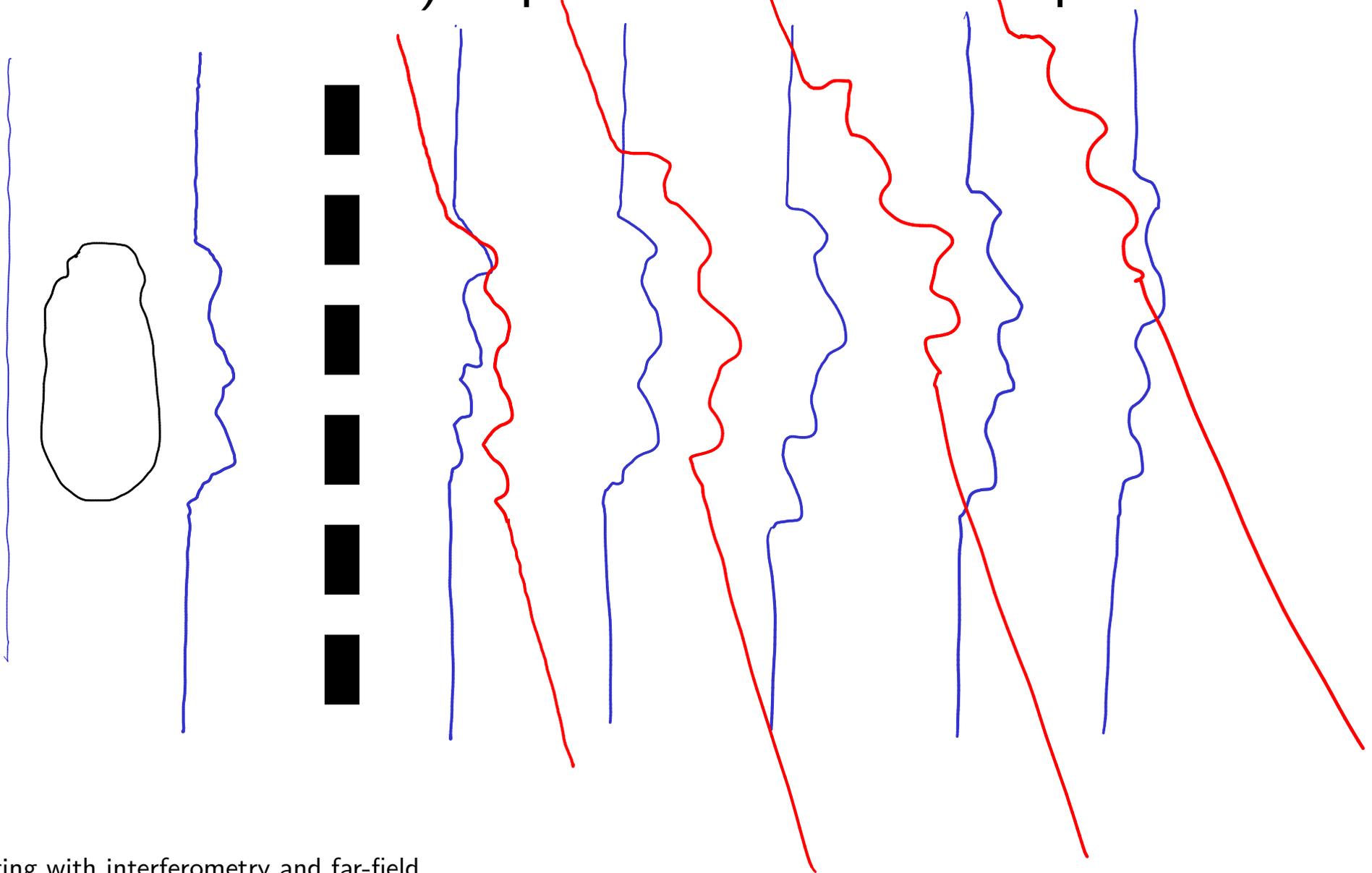


$$g(x) = \sum_{-\infty}^{\infty} g_n e^{2\pi i x n / p}$$



# Grating interferometry

Observing the interference between two (slightly offset) copies of the same sample.



# Grating interferometry

Observing the interference between two (slightly offset) copies of the same sample.

phase grating, as an example, has dominant orders  $\pm 1$

$$\psi(\vec{r}; z=0) = \psi_0(\vec{r}) * g(\vec{r})$$

$$\psi(\vec{r}; z) = \mathcal{F}^{-1} \left\{ \mathcal{F} \left\{ \psi_0(\vec{r}) g(\vec{r}) \right\} e^{-i\pi u^2 \lambda z} \right\}$$

$$= \mathcal{F} \left\{ \tilde{\psi}_0(\vec{u}) * G(\vec{u}) e^{-i\pi u^2 \lambda z} \right\}$$

$$= \mathcal{F}^{-1} \left\{ \tilde{\psi}_0(\vec{u}) * \left( g_1 \delta\left(\vec{u} - \frac{\hat{x}}{p}\right) + g_{-1} \delta\left(\vec{u} + \frac{\hat{x}}{p}\right) \right) e^{-i\pi u^2 \lambda z} \right\}$$

$$= \mathcal{F}^{-1} \left\{ g_1 \tilde{\psi}_0\left(\vec{u} - \frac{\hat{x}}{p}\right) + g_{-1} \tilde{\psi}_0\left(\vec{u} + \frac{\hat{x}}{p}\right) e^{-i\pi u^2 \lambda z} \right\}$$

$$\mathcal{F}^{-1} \left\{ \tilde{\Psi}_0 \left( u - \frac{\hat{x}}{p} \right) e^{-\pi i u^2 \lambda z} \right\} = \mathcal{F}^{-1} \left\{ \tilde{\Psi}_0(\vec{u}') e^{-i\pi \lambda z \left( u'^2 + \frac{c}{p^2} + \frac{2u'_x}{p} \right)} \right\}$$

$$= \int_{-\infty}^{\infty} \tilde{\Psi}_0 \left( \vec{u} - \frac{\hat{x}}{p} \right) e^{-\pi i u^2 \lambda z} e^{2\pi i \vec{u} \cdot \vec{r}} du$$

$$\begin{aligned} \vec{u}' &= \vec{u} - \frac{\hat{x}}{p} & u'^2 &= \left( \vec{u}' + \frac{\hat{x}}{p} \right)^2 \\ & & &= u'^2 + \frac{2u'_x}{p} + \frac{1}{p^2} \end{aligned}$$

$$= e^{-\frac{i\pi \lambda z}{p^2}} \int_{-\infty}^{\infty} d^2 \vec{u}' \underbrace{\tilde{\Psi}_0(\vec{u}') e^{-i\pi \lambda z u'^2}}_{\text{usual argument for propagation}} \underbrace{e^{-\frac{2\pi i \lambda z}{p} \vec{u}' \cdot \hat{x}} e^{2\pi i \left( \vec{u}' + \frac{\hat{x}}{p} \right) \cdot \vec{r}}}_{1}$$

$$= e^{-\frac{i\pi \lambda z}{p^2}} \Psi_0 \left( \vec{r} - \frac{\lambda z}{p} \hat{x}; z \right) e^{\frac{2\pi i x}{p}}$$

⋮

$$\Psi(\vec{r}; z) = e^{-\frac{i\pi\lambda z}{p^2}} \left( g_1 e^{\frac{2\pi i x}{p}} \psi_0\left(\vec{r} - \frac{\lambda z \hat{x}}{p}; z\right) + g_{-1} e^{-\frac{2\pi i x}{p}} \psi_0\left(\vec{r} + \frac{\lambda z \hat{x}}{p}; z\right) \right)$$

$$\begin{aligned} I(\vec{r}; z) &= |\Psi(\vec{r}; z)|^2 \\ &= \Psi(\vec{r}; z) \Psi^*(\vec{r}; z) \\ &= |g_1|^2 \left| \psi_0\left(\vec{r} - \frac{\lambda z \hat{x}}{p}; z\right) \right|^2 + |g_{-1}|^2 \left| \psi_0\left(\vec{r} + \frac{\lambda z \hat{x}}{p}; z\right) \right|^2 \\ &\quad + g_1 g_{-1}^* e^{\frac{4\pi i x}{p}} \psi_0\left(\vec{r} - \frac{\lambda z \hat{x}}{p}; z\right) \psi_0^*\left(\vec{r} + \frac{\lambda z \hat{x}}{p}; z\right) + \text{c.c.} \end{aligned}$$

$$\psi_0 = a(r) e^{i\varphi(r)}$$

$$g_1 = g_{-1}^*$$

$$I(\vec{r}; z) = |g_1|^2 \left( |a(\vec{r} - \frac{\lambda z}{p} \hat{x})|^2 + |a(\vec{r} + \frac{\lambda z}{p} \hat{x})|^2 \right)$$

$$+ |g_1|^2 e^{\frac{4\pi i x}{p}} a(\vec{r} - \frac{\lambda z}{p} \hat{x}) a(\vec{r} + \frac{\lambda z}{p} \hat{x}) e^{i\varphi(\vec{r} - \frac{\lambda z}{p} \hat{x}) - i\varphi(\vec{r} + \frac{\lambda z}{p} \hat{x})} + \text{c.c.}$$

$\frac{\lambda z}{p}$  assumed to be small compared to resolution of the system

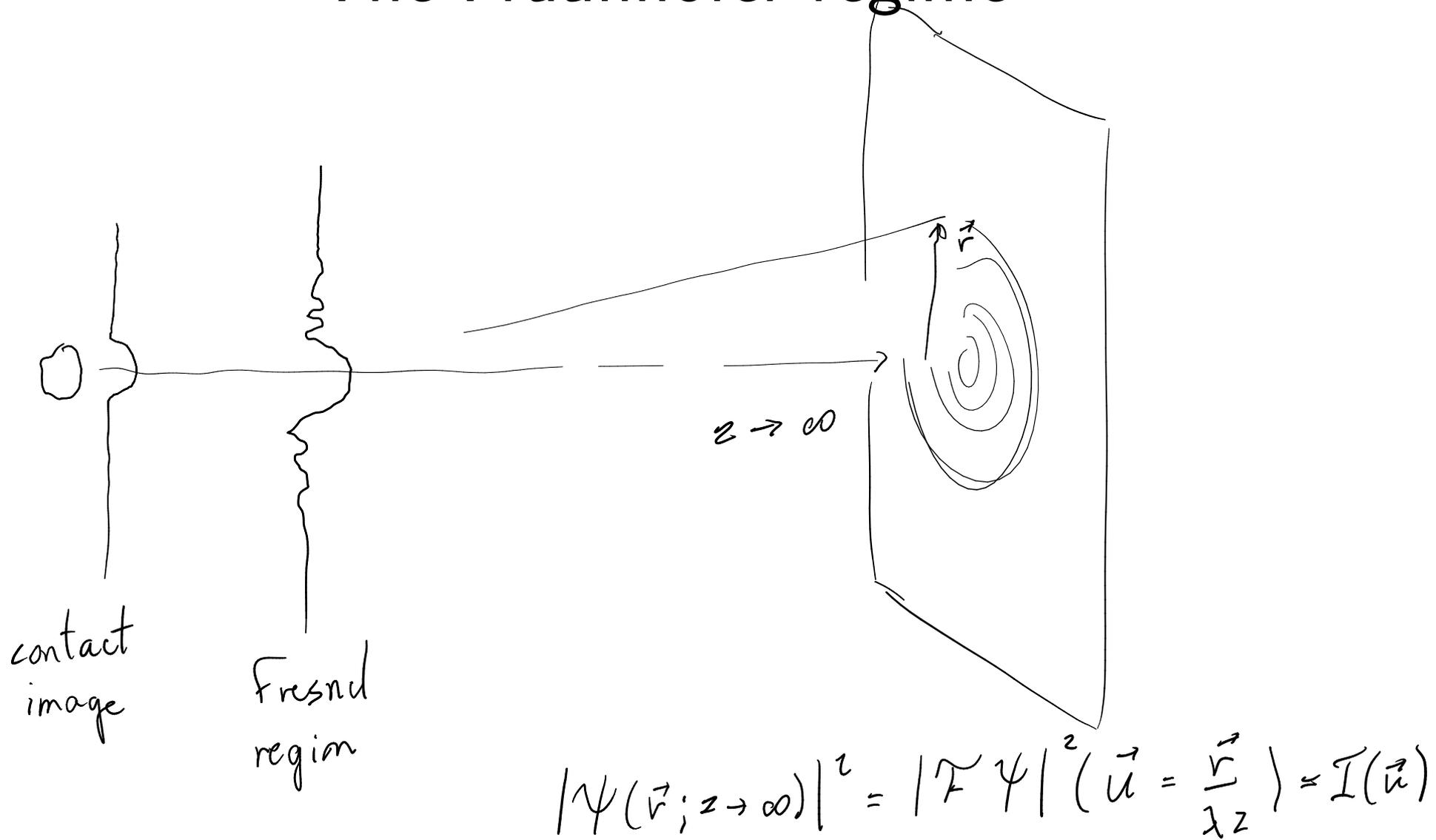
$$I(\vec{r}; z) \approx 2|g_1|^2 |a(\vec{r})|^2$$

$$+ 2|g_1|^2 |a(\vec{r})|^2 \cos\left(\frac{4\pi x}{p} - 2\frac{\lambda z}{p} \frac{\partial \varphi}{\partial x}\right)$$

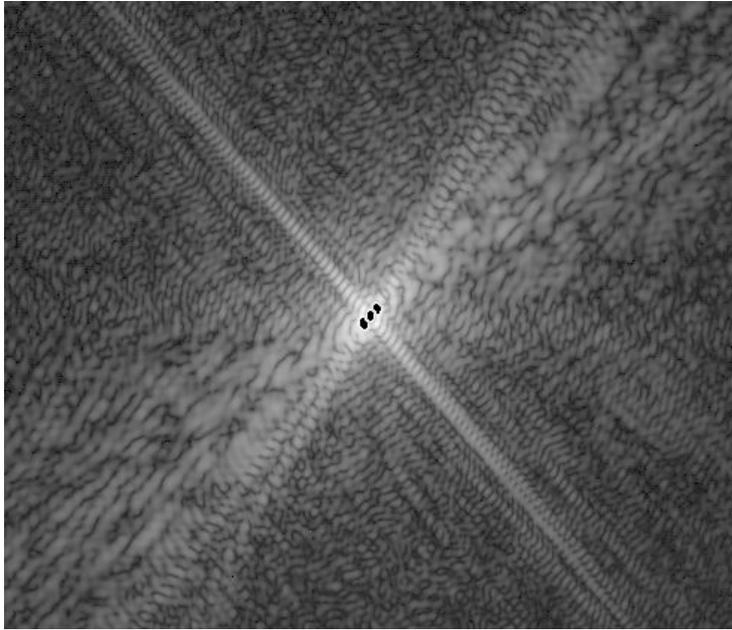
$$\varphi(\vec{r} - \frac{\lambda z}{p} \hat{x}) - \varphi(\vec{r} + \frac{\lambda z}{p} \hat{x}) \approx -2\frac{\lambda z}{p} \frac{\partial \varphi}{\partial x}$$

# Far-field diffraction

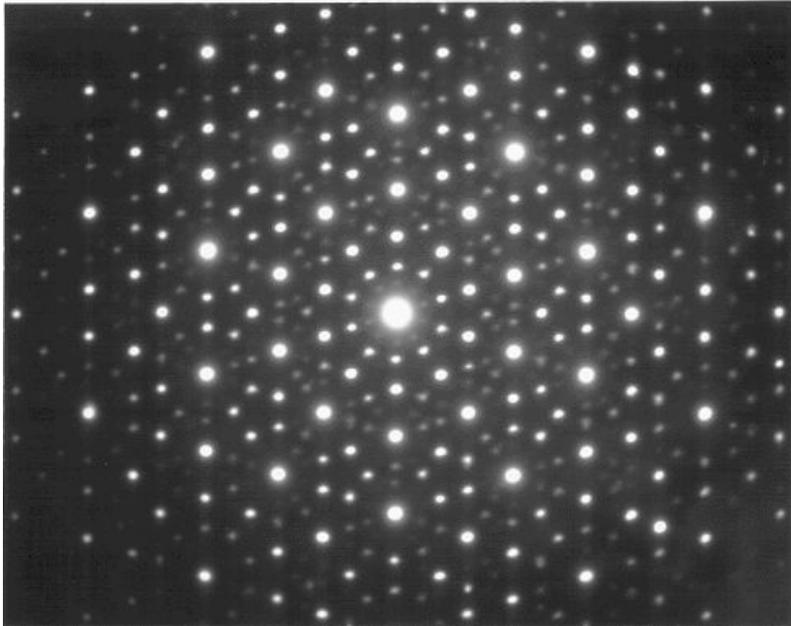
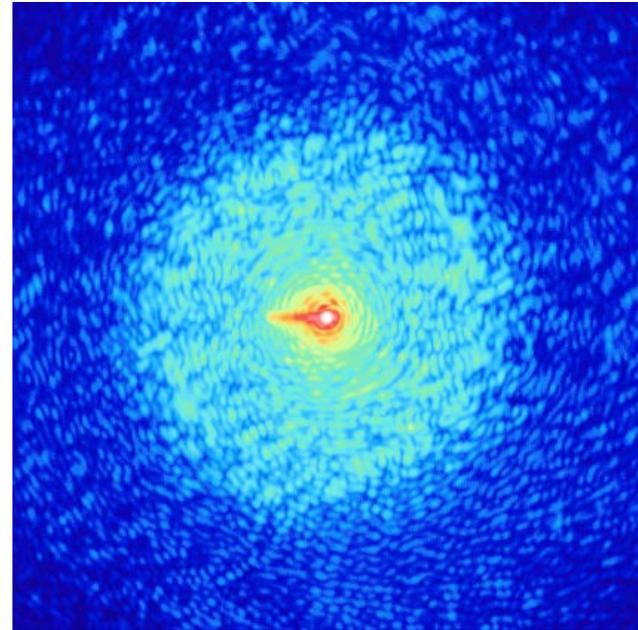
## The Fraunhofer regime



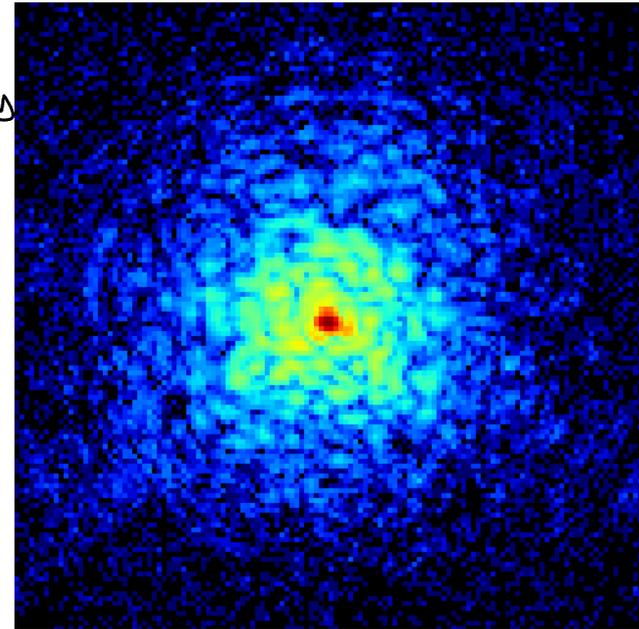
# Diffraction patterns



speckles



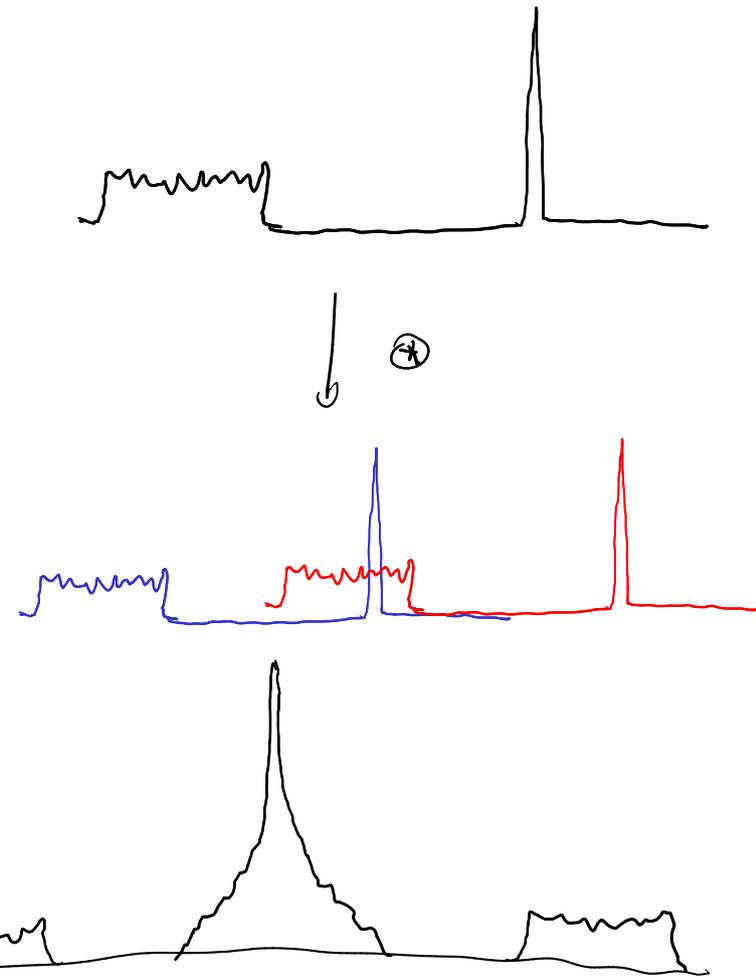
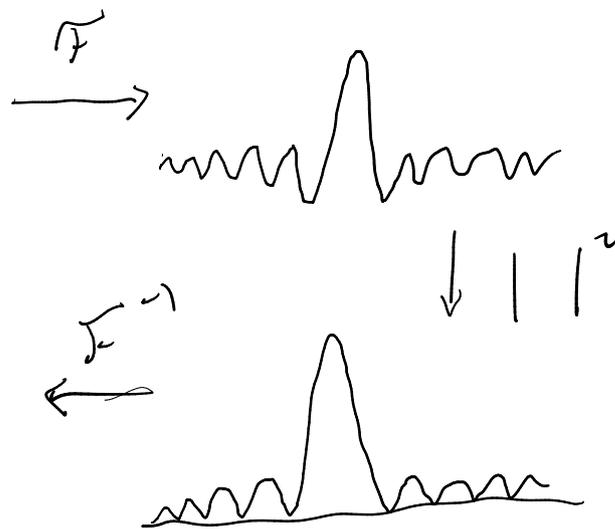
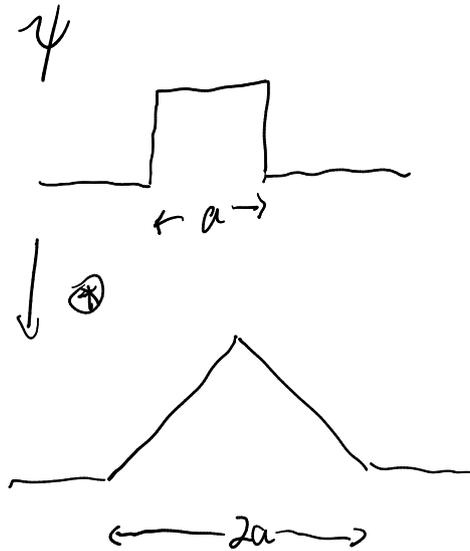
Bragg peaks



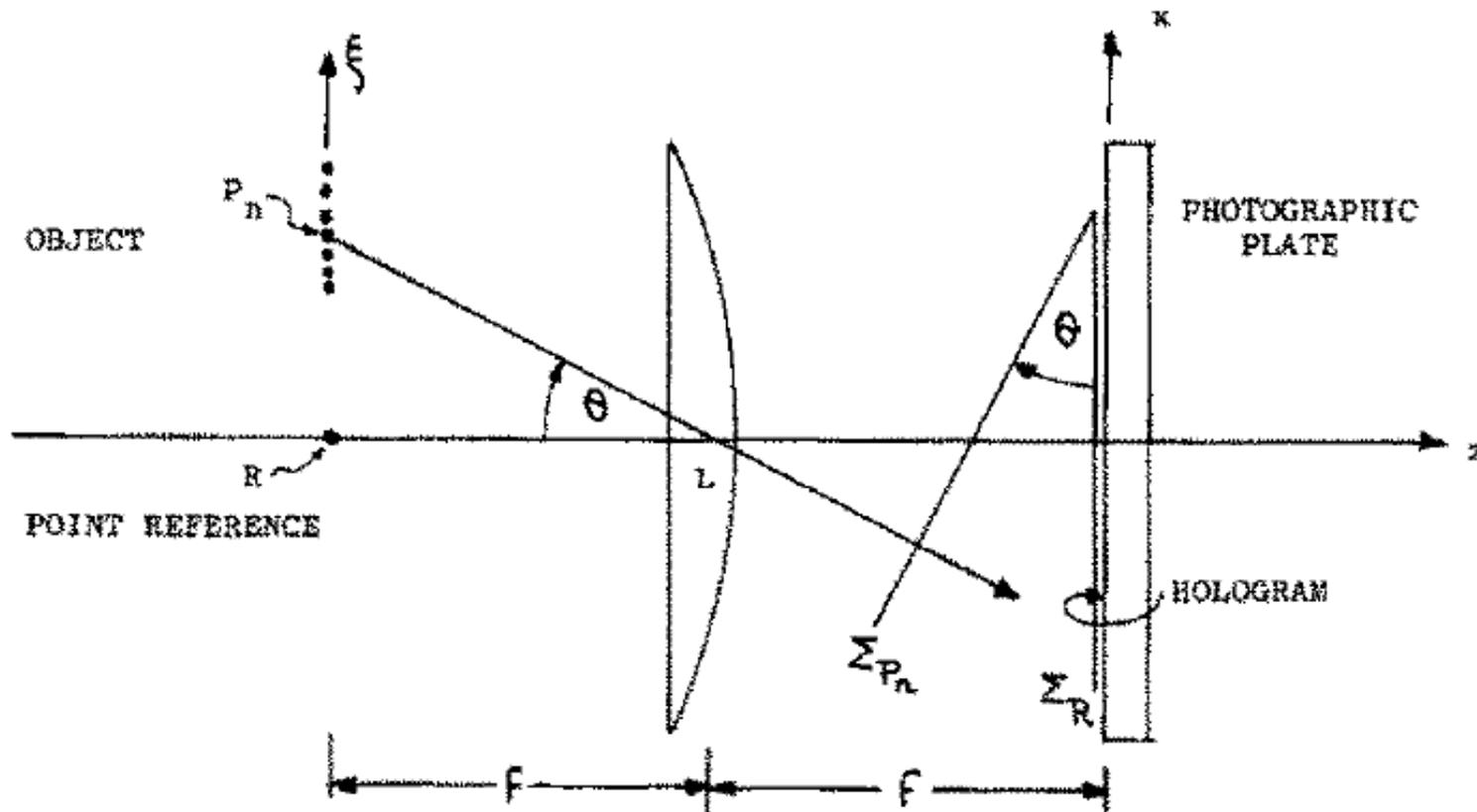
# Diffraction and autocorrelation

$$\mathcal{F}^{-1}\{I(\vec{u})\} = \mathcal{F}^{-1}\{\psi(u)\psi^*(u)\} = \psi(\vec{r}) \otimes \psi^*(\vec{r})$$

autocorrelation



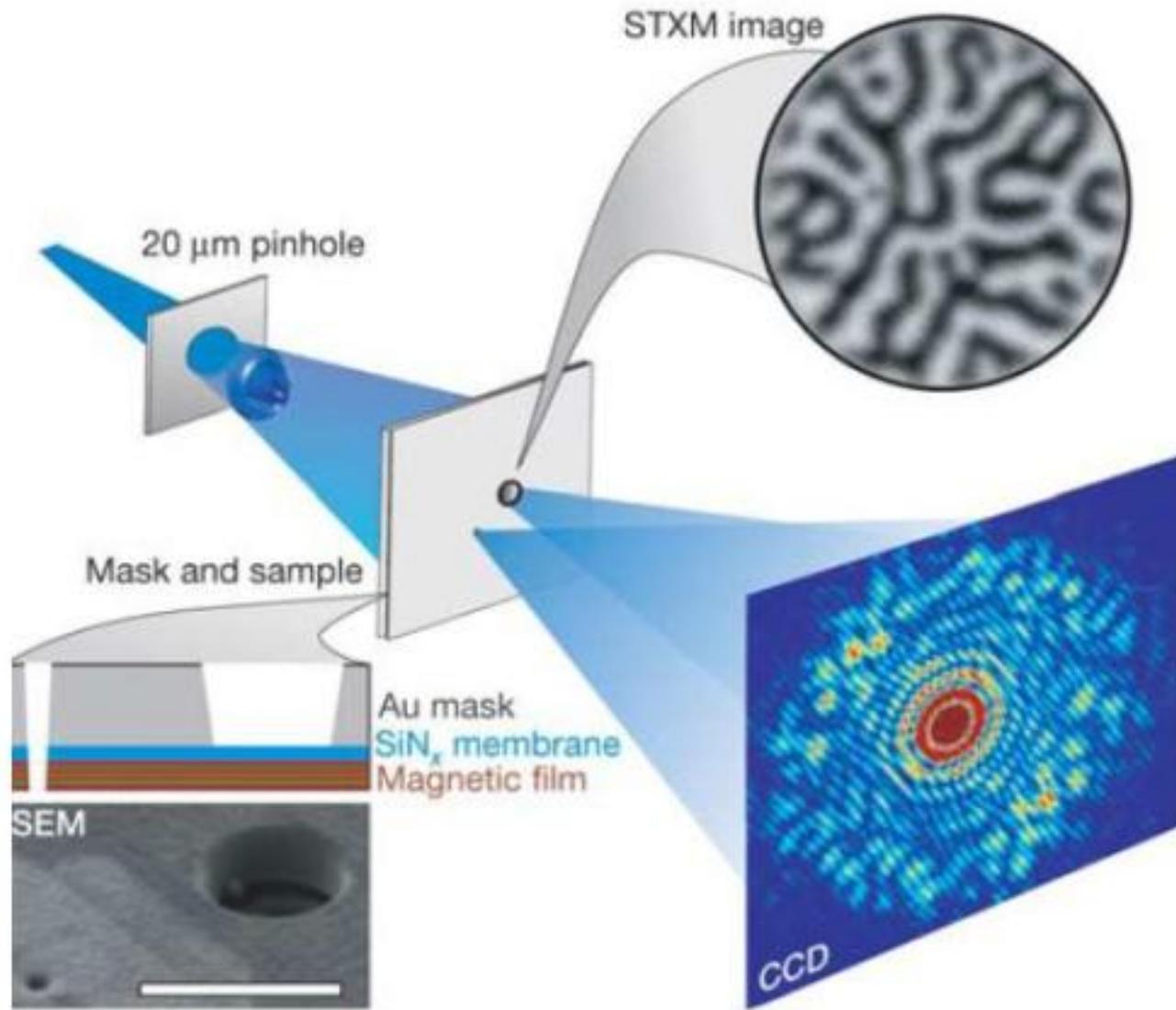
# Fourier transform holography



**Fig. 1. Recording of a Fourier-transform hologram with a lens  $L$ .  $\Sigma_R$  = reference wavefront.**

Source: G. Stroke, Appl. Phys. Lett. **6**, 201-203 (1965).

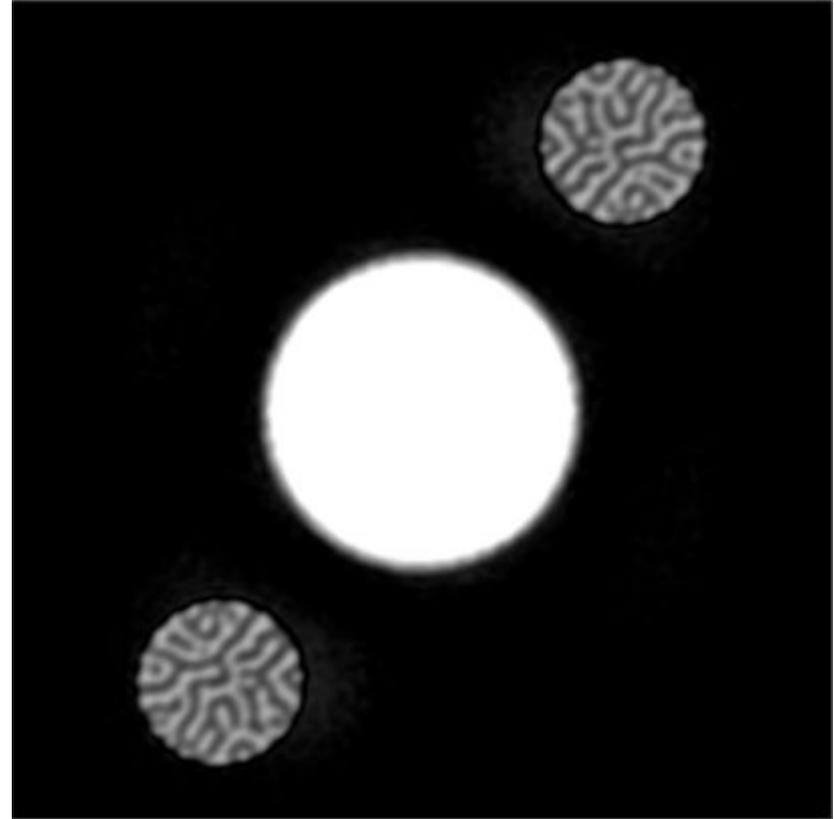
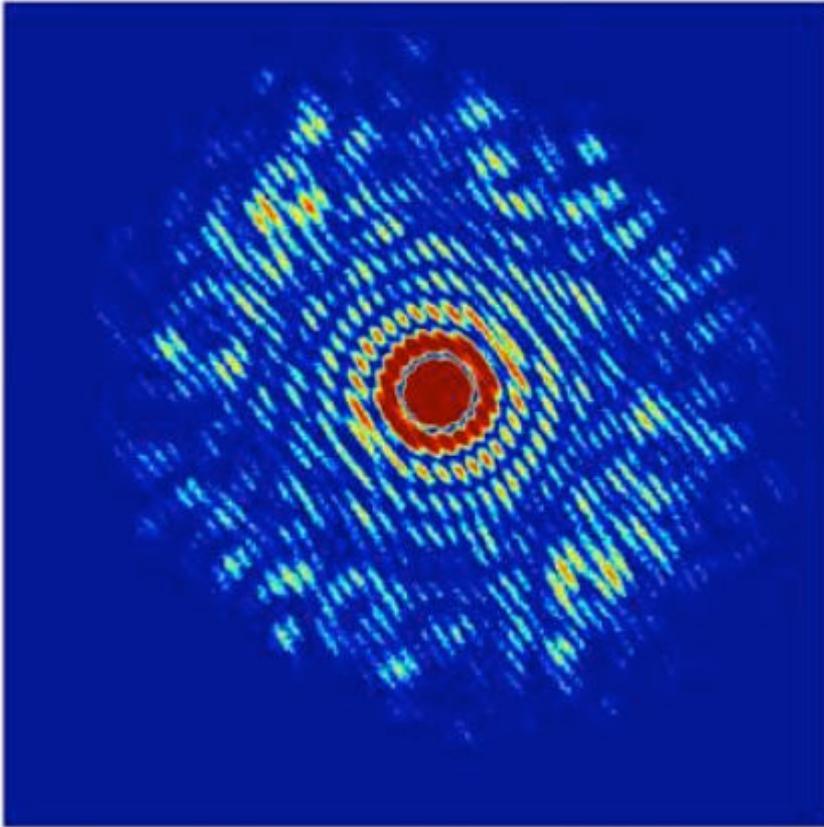
# Fourier transform holography



Source: S. Eisebitt et al., Nature **432**, 885-888 (2004).

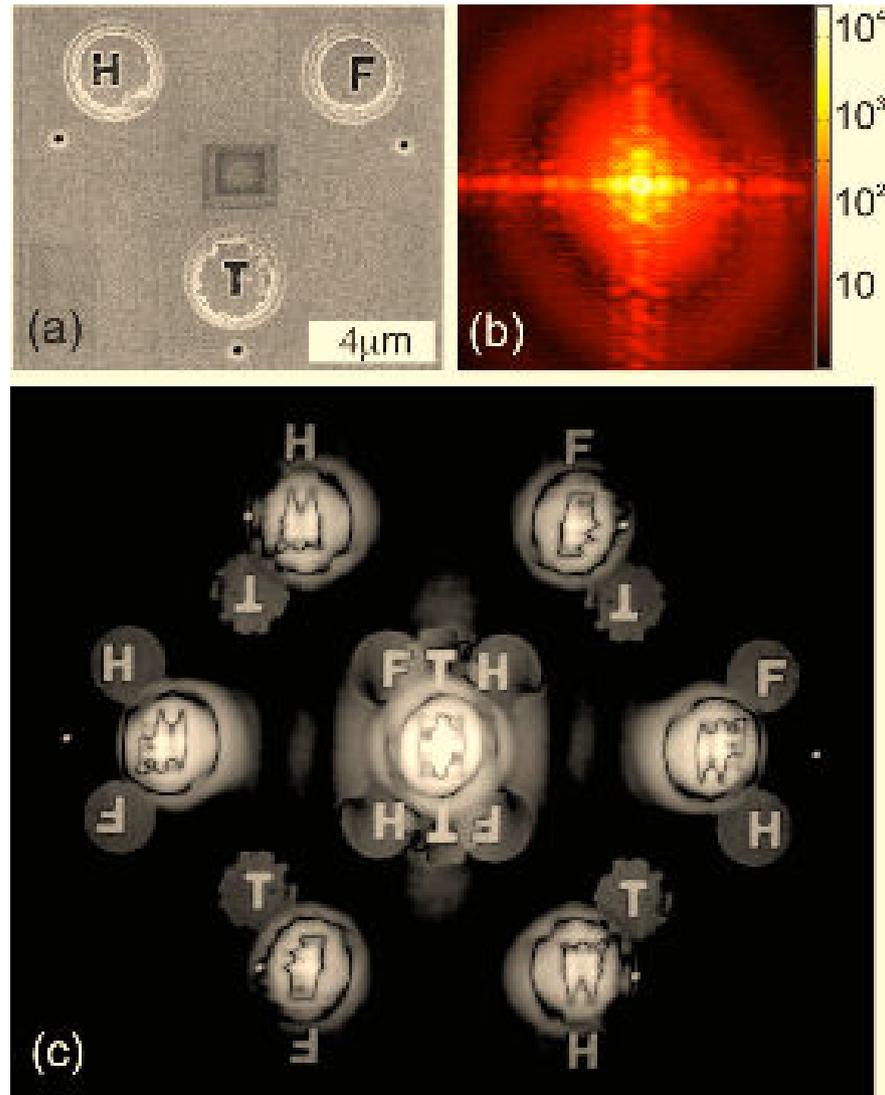
# Fourier transform holography

# Fourier transform holography



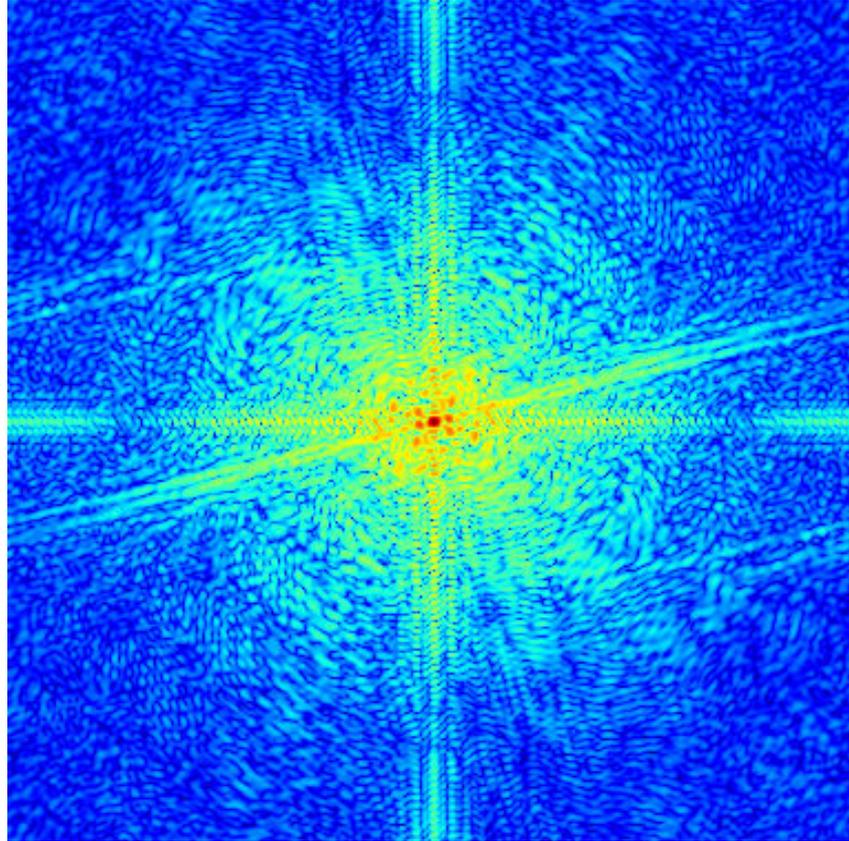
# Fourier transform holography

## Multiple references



Source: W. Schlotter et al., Opt. Lett. **21**, 3110-3112 (2006).

# Coherent diffractive imaging

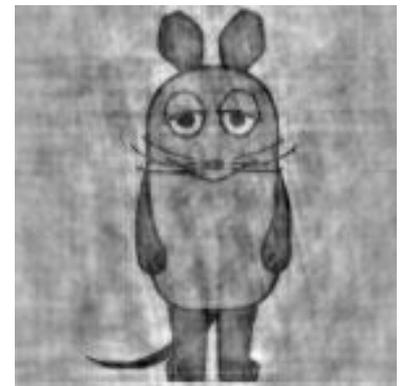
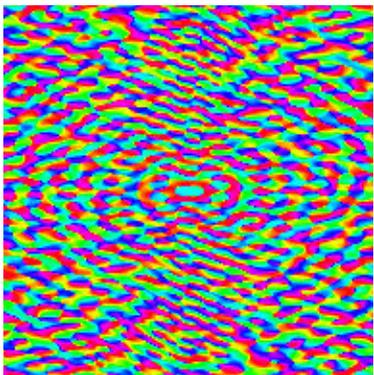
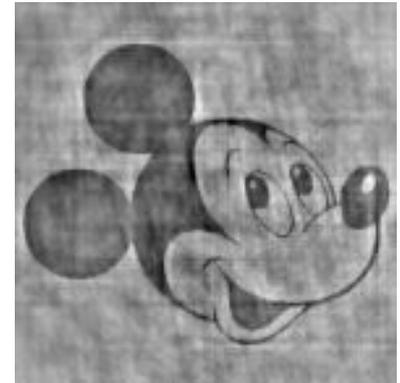
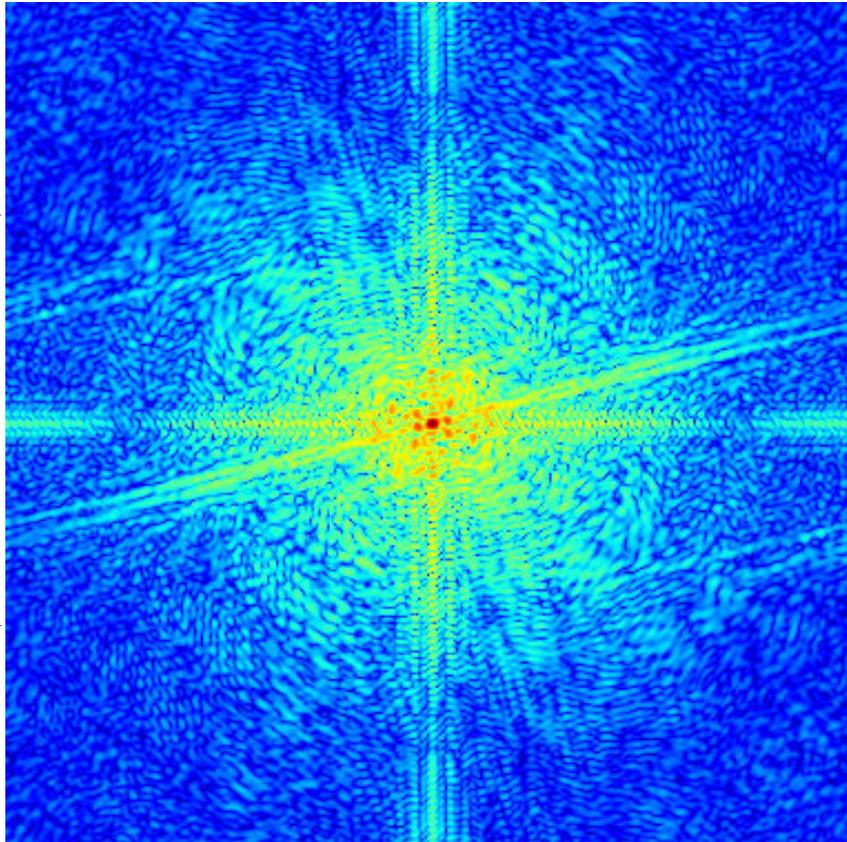
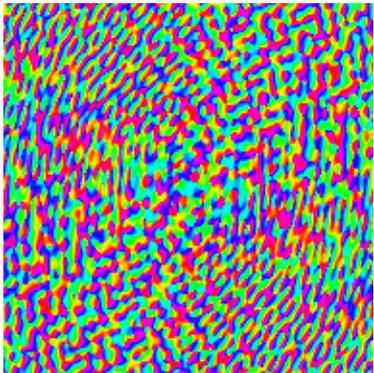


# The phase problem

$$e^{i\varphi(\omega)}$$

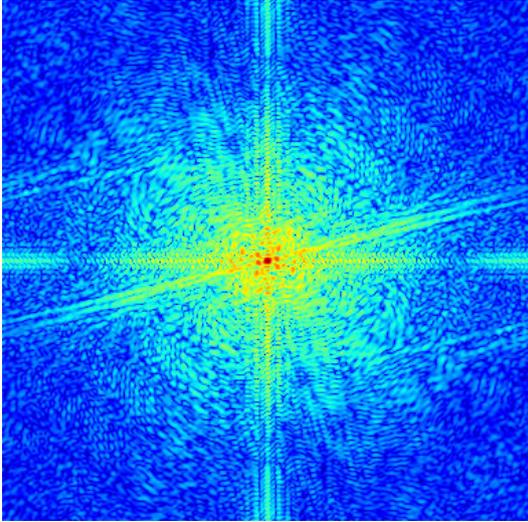
$$|\mathcal{F}\{\varphi\}|^2(\vec{u})$$

*phases carry  
very important  
information*

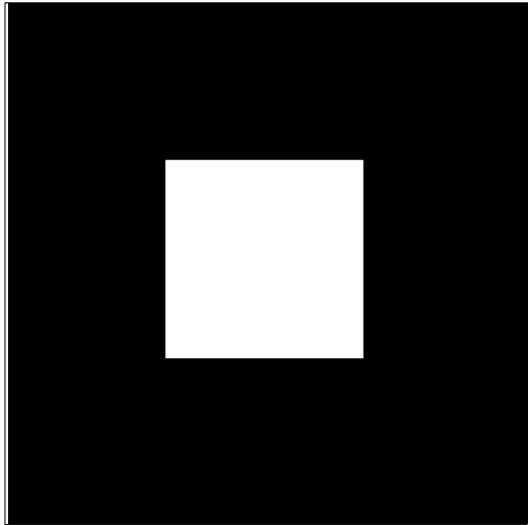


# Coherent diffractive imaging

*Two constraints*

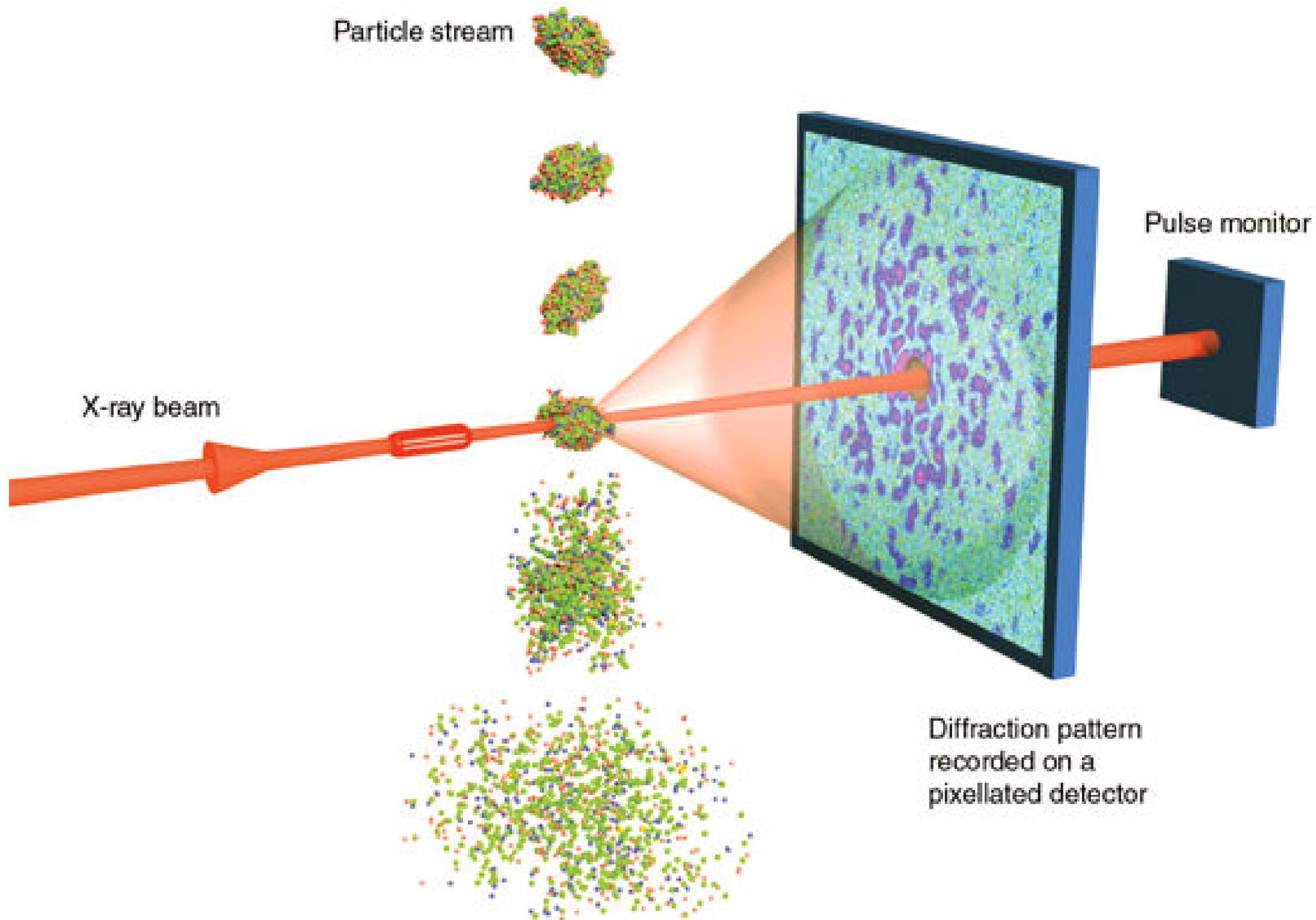


1. Solution is consistent with measured Fourier amplitude



2. Solution is isolated in space

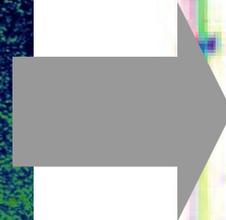
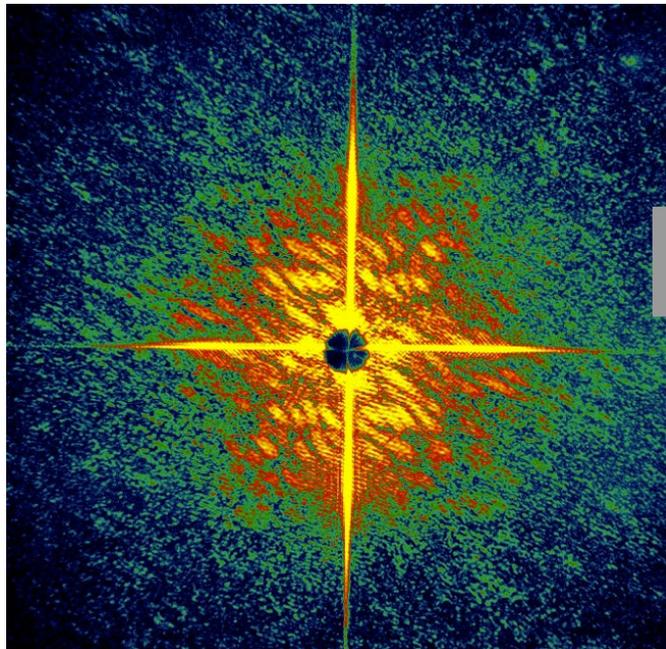
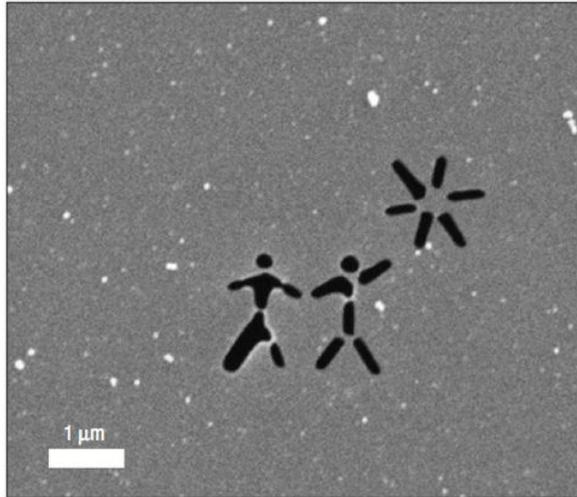
# Radiation damage limits on radiation



R. Neutze *et al*, Nature **406**, 752 (2000)

K. J. Gaffney *et al*, Science **316**, 1444 (2007)

# “Diffraction before destruction”

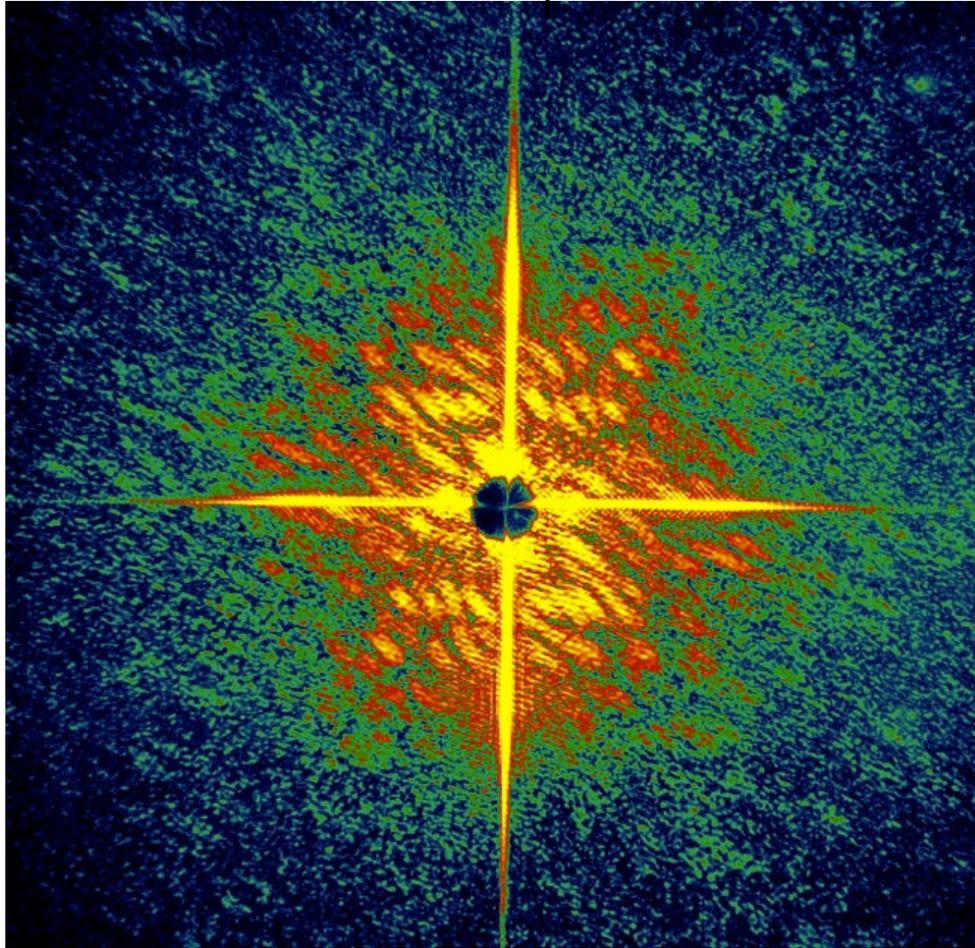


H. N. Chapman *et al*, Nat. Phys. **2**, 839 (2006)

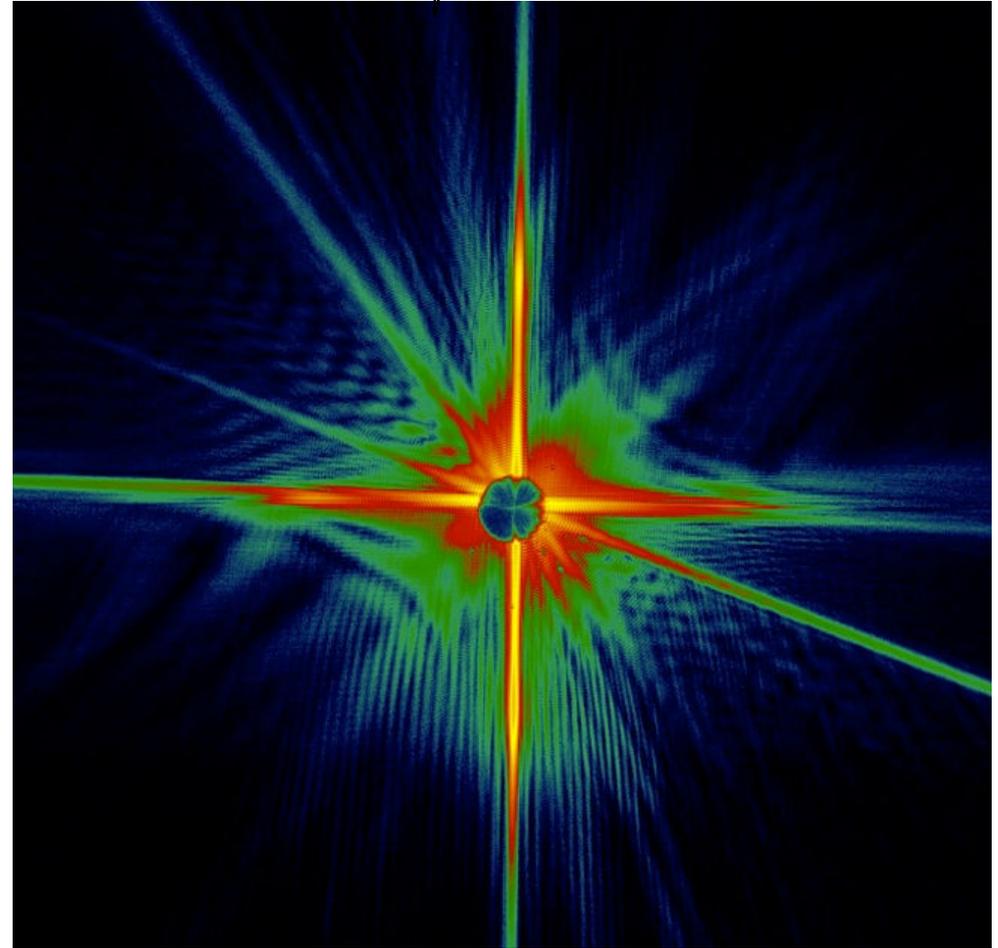
# “Diffraction before destruction”

The imaging pulse vaporized the sample

1<sup>st</sup> pulse



2<sup>nd</sup> pulse (no more sample)



H. N. Chapman *et al*, Nat. Phys. **2**, 839 (2006)

# Ptychography

- Scanning an isolated illumination on an extended specimen
- Measure full coherent diffraction pattern at each scan point
- Combine everything to get a reconstruction

## **Dynamische Theorie der Kristallstrukturanalyse durch Elektronenbeugung im inhomogenen Primärstrahlwellenfeld**

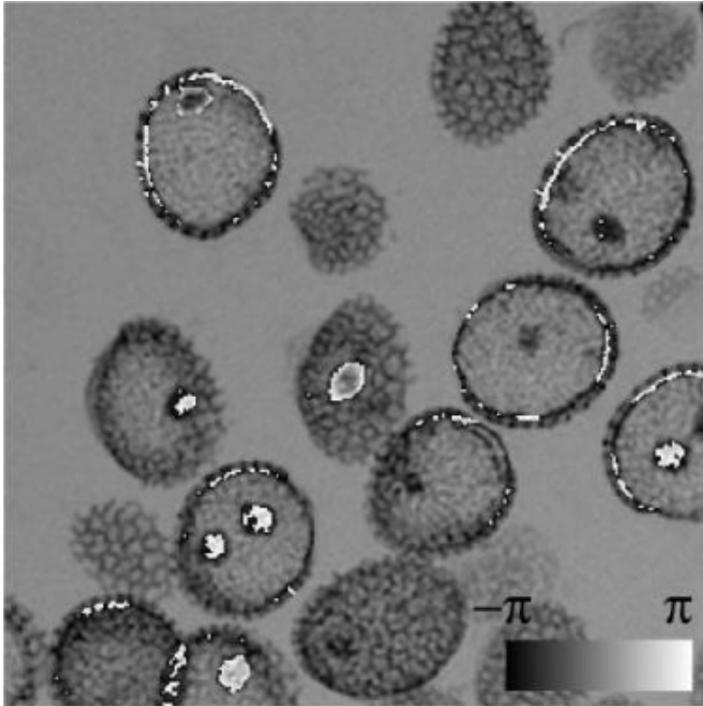
Von R. Hegerl und W. Hoppe

Some time ago a new principle was proposed for the registration of the complete information (amplitudes and phases) in a diffraction diagram, which does not – as does Holography – require the interference of the scattered waves with a single reference wave. The basis of the principle lies in the interference of neighbouring scattered waves which result when the object function  $g(x, y)$  is multiplied by a generalized primary wave function  $p(x, y)$  in Fourier space (diffraction diagram) this is a convolution of the Fourier transforms of these functions. The above mentioned interferences necessary for the phase determination can be obtained by suitable choice of the shape of  $p(x, y)$ . To distinguish it from holography this procedure is designated “ptychography” ( $\pi\tau v\zeta = \text{fold}$ ). The procedure is applicable to periodic and aperiodic structures. The relationships are simplest for plane lattices. In this paper the theory is extended to space lattices both with and without consideration of the dynamic theory. The resulting effects are demonstrated using a practical example.

# Ptychography

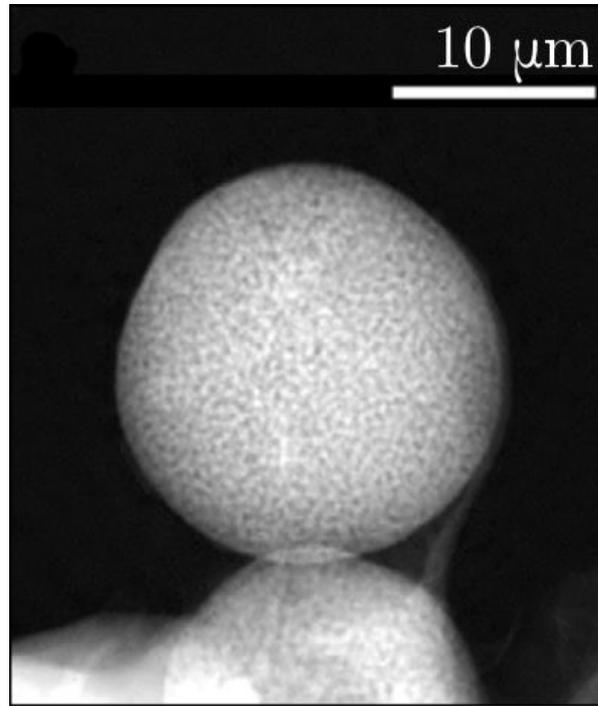
## A few examples

Visible light



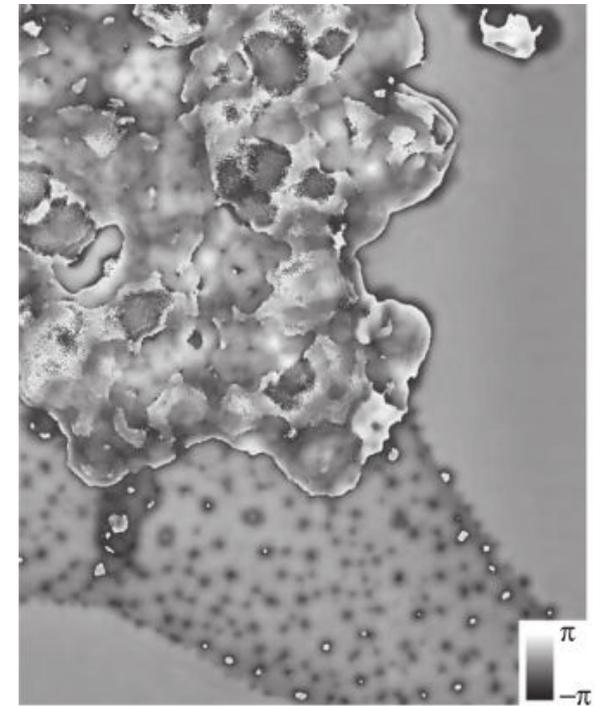
A. Maiden *et al.*, *Opt. Lett.* **35**,  
2585-2587 (2010).

X-rays



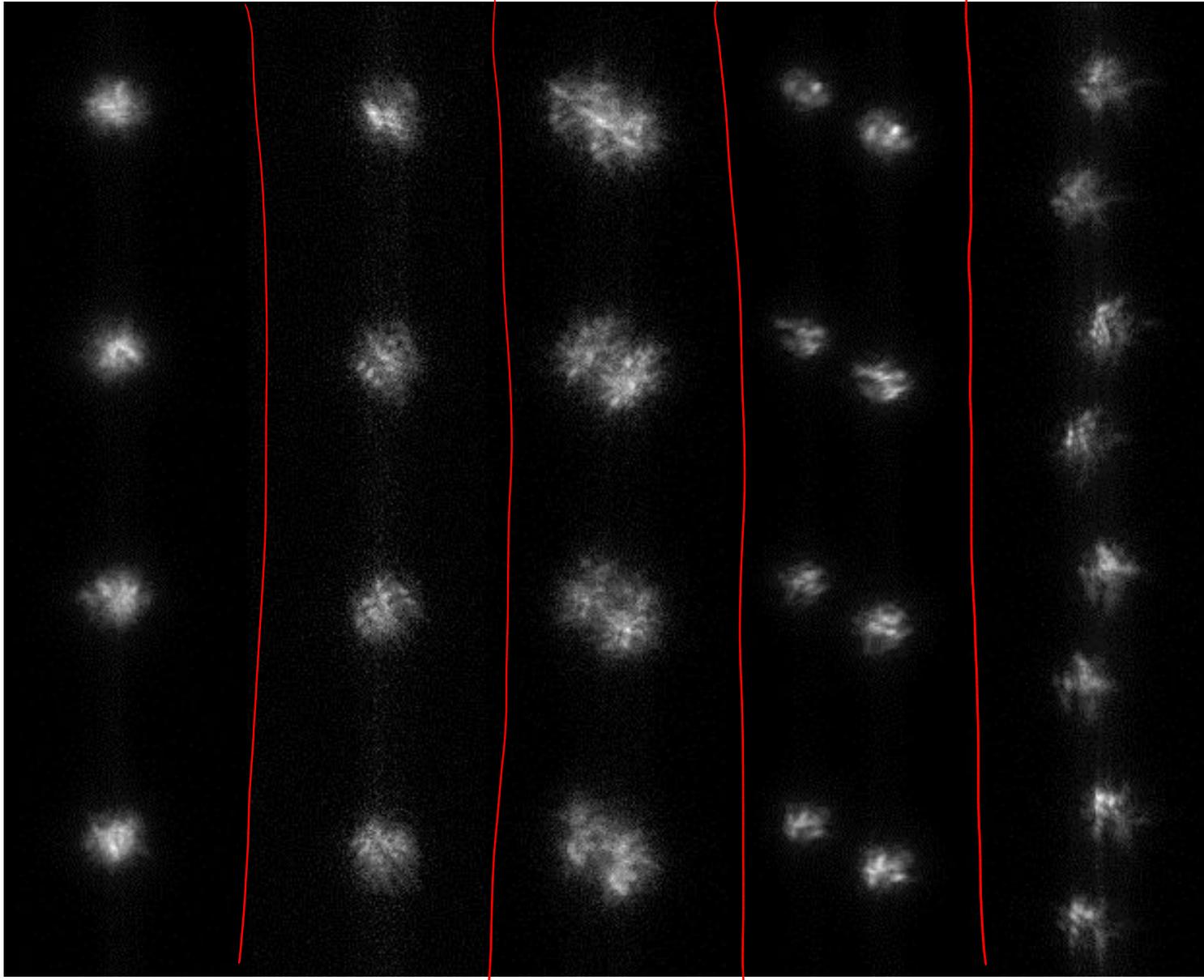
P. Thibault *et al.*, *New J. Phys* **14**,  
063004 (2012).

electrons



M. Humphry *et al.*,  
*Nat. Comm.* **3**, 730 (2012).

# Speckle imaging in astronomy



different  
measurements  
of the same  
objects.

Difference  
is caused  
by turbulence

Source: <http://www.cis.rit.edu/research/thesis/bs/2000/hoffmann/thesis.html>

# Speckle imaging in astronomy

## Model

One measurement

$$I(\vec{r}) = O * P$$

instantaneous  
point-spread function

$$\tilde{I}(\vec{u}) = \tilde{O}(\vec{u}) \cdot M \leftarrow \text{MTF}$$

$$|\tilde{I}(\vec{u})|^2 = |\tilde{O}(\vec{u})|^2 |M|^2$$

can be modeled  
from fluid  
dynamics

many  
measurements

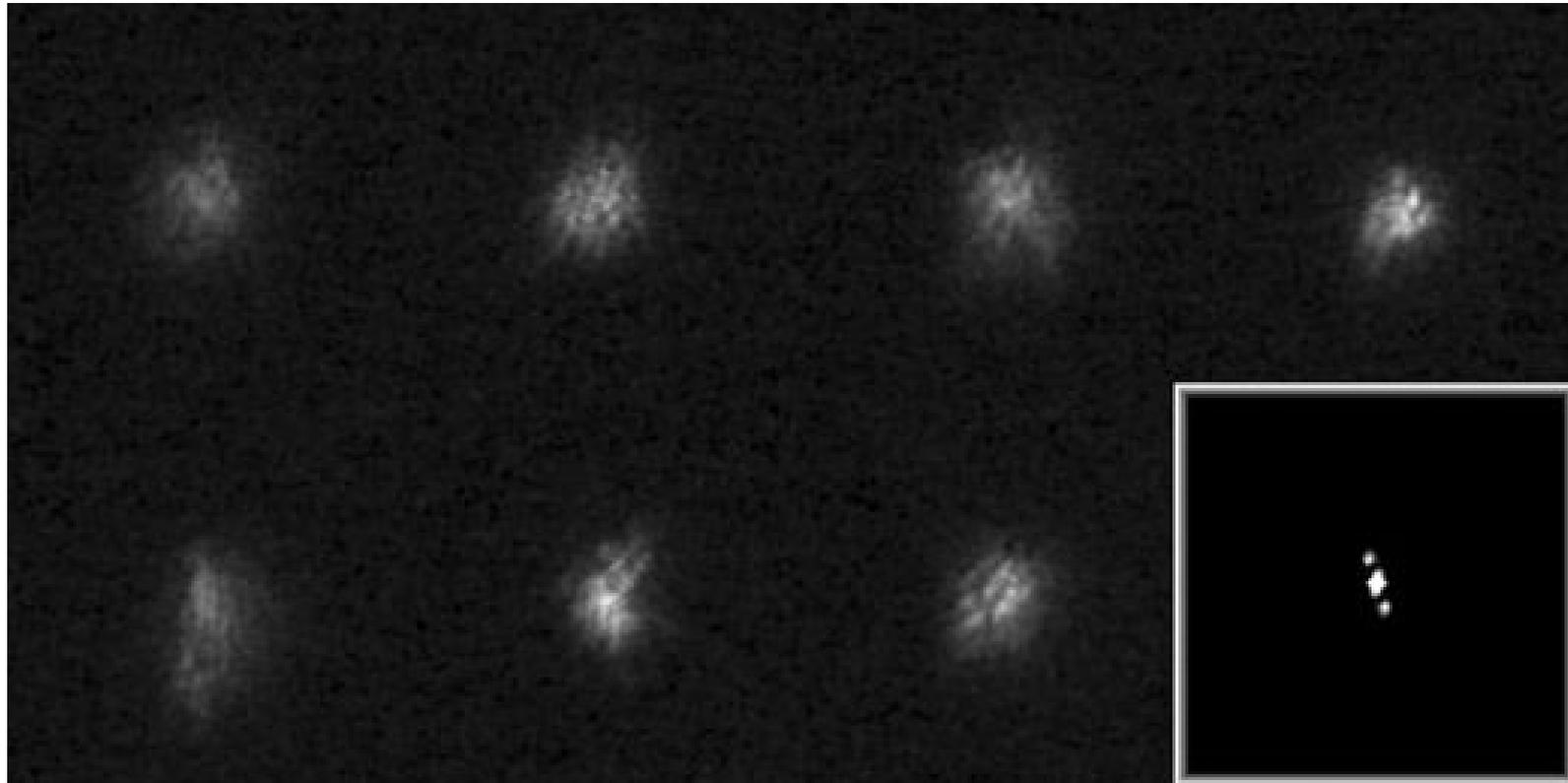
$$\langle |\tilde{I}(\vec{u})|^2 \rangle = |\tilde{O}|^2 \langle |M|^2 \rangle$$

$$|\tilde{O}|^2 \approx \frac{\langle |\tilde{I}|^2 \rangle}{\langle |M_{\text{model}}|^2 \rangle}$$

recovering  $O$  from  $|\tilde{O}|^2$  same as Coherent diffractive imaging!

# Speckle imaging in astronomy

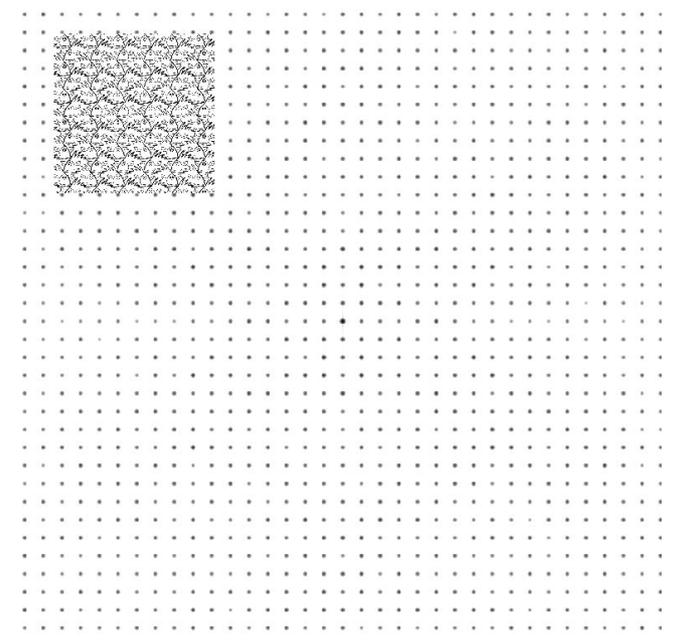
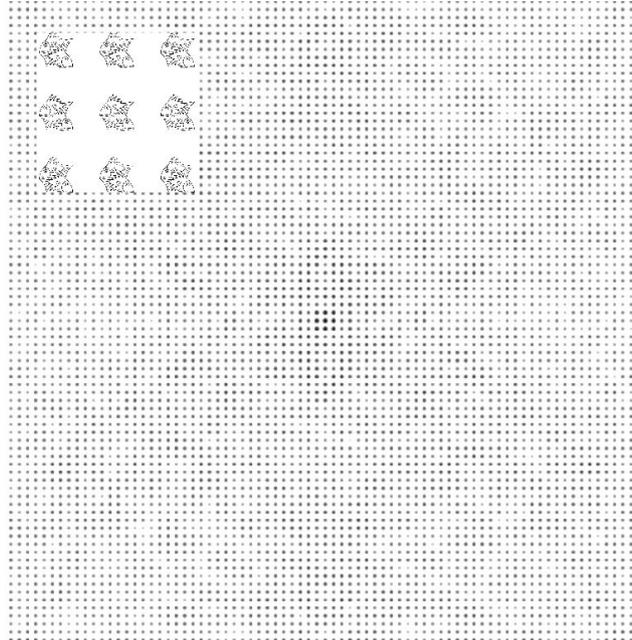
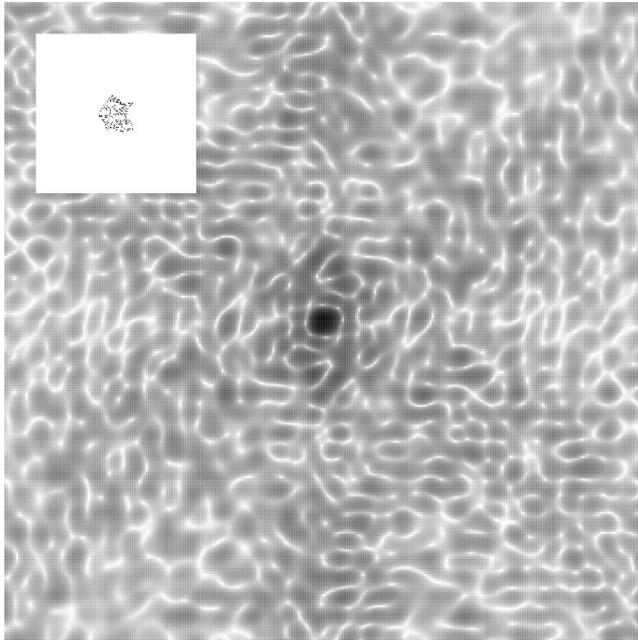
## Retrieval of the autocorrelation



Source: <http://www.astrosurf.com/hfosaf/uk/speckle10.htm>

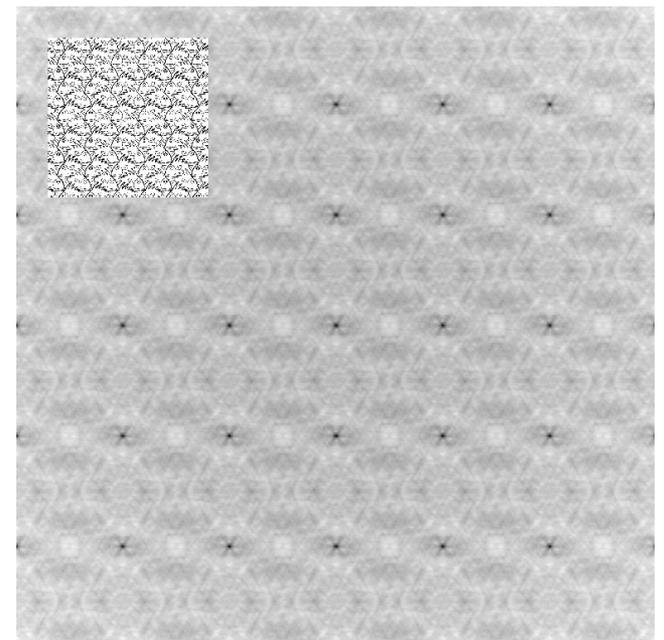
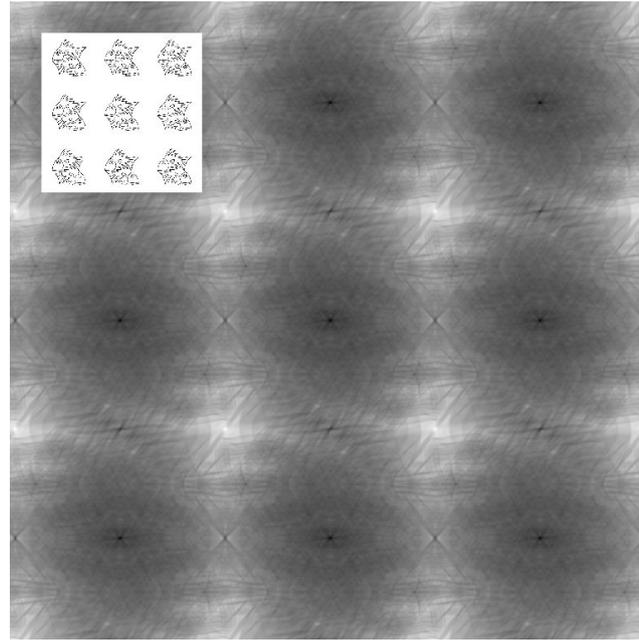
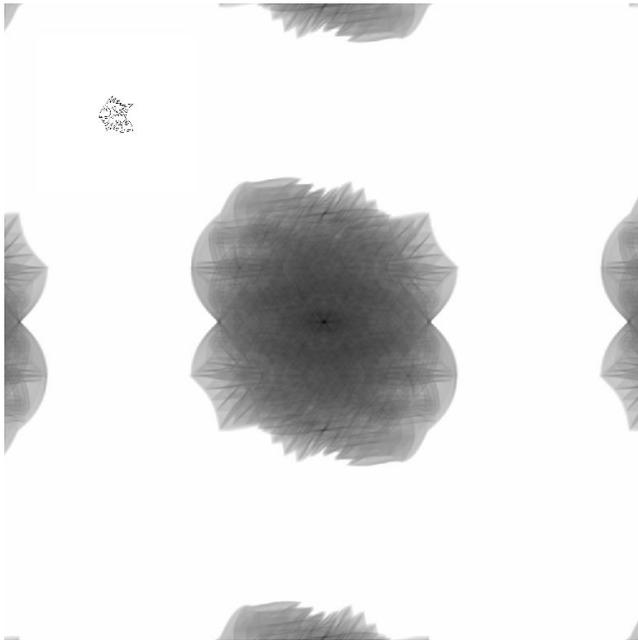
# Crystallography

## Diffraction by a crystal: Bragg peaks



# Crystallography

## Fourier transform of intensity: autocorrelation



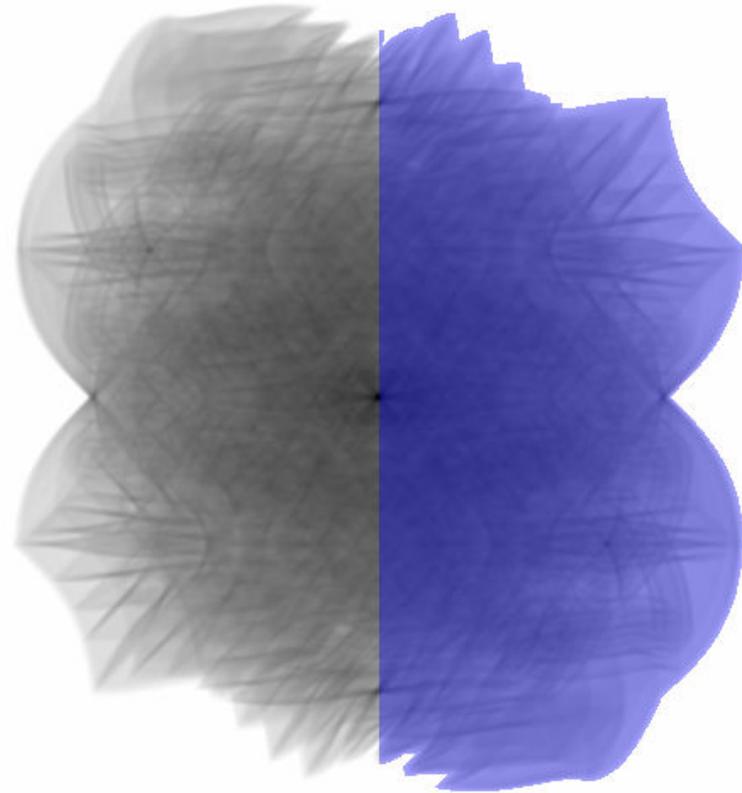
F.T. of crystal diffraction  
called Patterson map

# Crystallography

Problem is overconstrained with an isolated sample



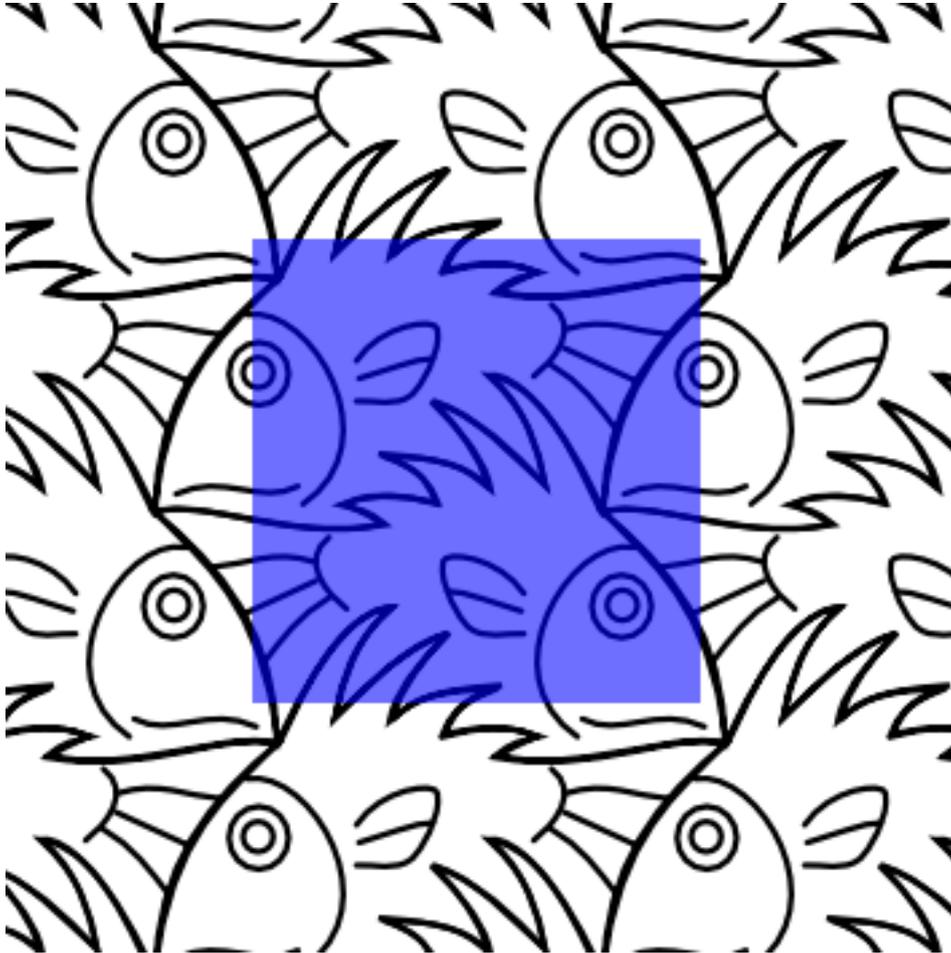
**unknowns =  $N$**



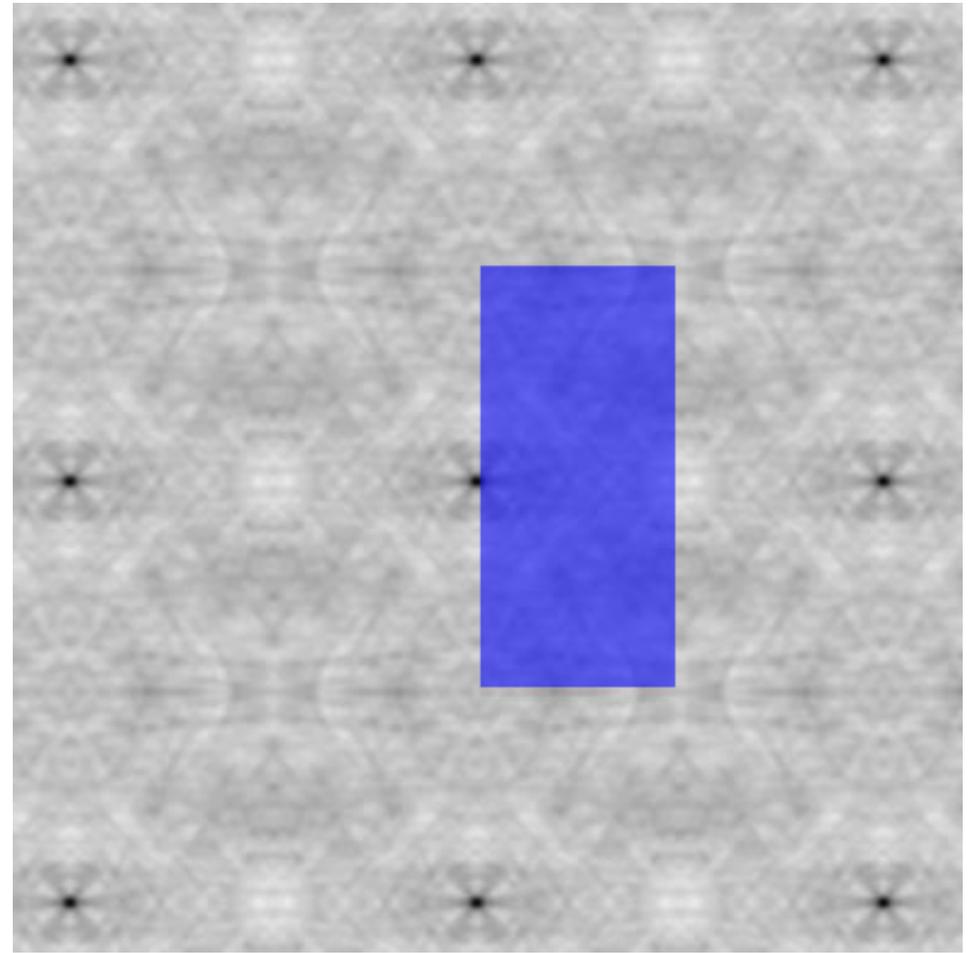
**constraints  $\geq 2N$**

# Crystallography

Problem is **underconstrained** with a crystal



**unknowns =  $N$**



**constraints =  $N/2$**

# Crystallography

## Structure determination

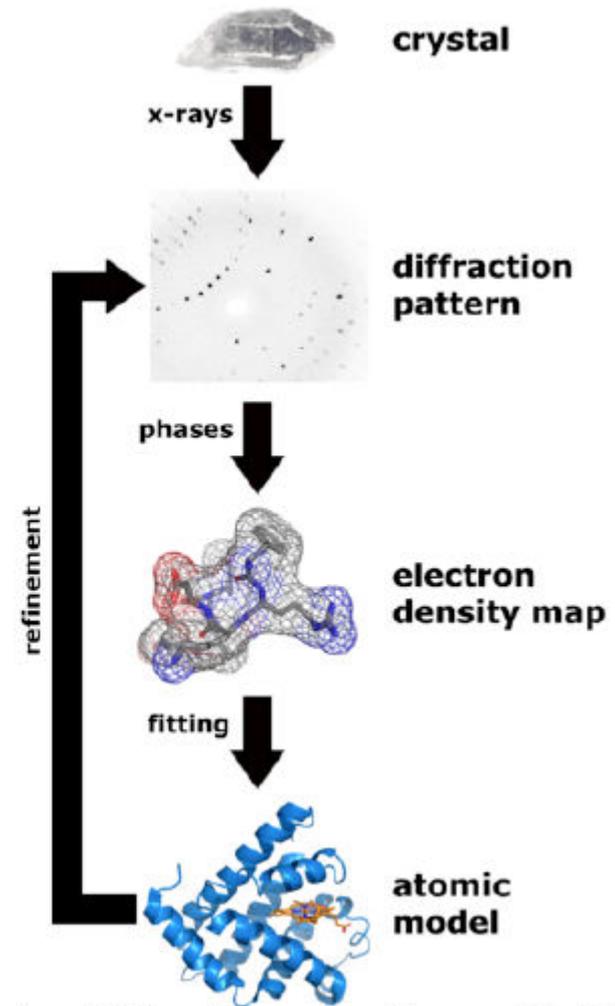
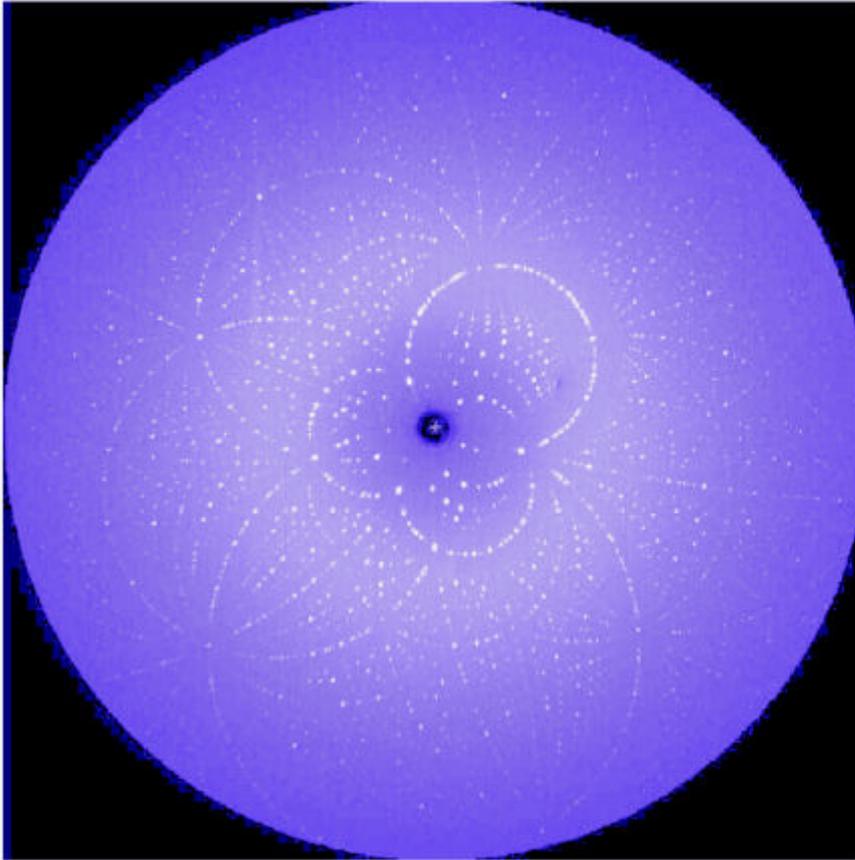


Image from Wikimedia courtesy Thomas Splettstoesser

# Crystallography

## Structure determination

- Hard problem: few measurements for the number of unknowns
- Luckily: crystals are made of atoms → strong constraint
- Also common: combining additional measurements (SAD, MAD, isomorphous replacement, ...)

# Summary

// lensless imaging

## Imaging from far-field amplitudes

- Used when image-forming lenses are unavailable (or unreliable) or to obtain more quantitative images.
- In general difficult because of the phase problem
- Solved with the help of additional information:
  - Strong *a priori* knowledge (e.g. CDI: support)
  - Multiple measurements (e.g. ptychography)